

Design of magnetized, gas-filled capsule experiments for NIF

Meeting on Magnetic Fields in Laser Plasmas

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Magnetic fields in hohlraums and capsules: MHD Hydra simulations

Main effect of B field: reduced e- heat conduction perpendicular to B: $\omega_{ce}\tau_{ei} > 1$
Magnetic pressure \ll matter pressure: $\beta \gg 1$

New in this talk:

- Hohlraum sims of “bigfoot” NIF design
- Imposed axial field, “Biermann battery” fields and Nernst advection



No imposed B field: similar to W. Farmer, 2017¹

- Biermann fields \rightarrow hotter hohlraum fill
- Nernst advection reduces effect of B field
- Modest effect on implosion
- Small fields in capsule: < 50 T

- 1 W. A. Farmer et al., Phys. Plasmas 2017
- 2 D. J. Strozzi et al., J. Plasma Phys. 2015
- 3 L. J. Perkins et al., LLNL LDRD final report

Imposed axial field: similar to D. Strozzi, 2015^{2,3}

- Frozen-in law holds: B field compressed or rarified with plasma
- **Slightly** hotter hohlraum fill
- Improved inner-beam propagation: hotter, less-dense equator channel
- Capsule fields ~ 2 kT
- Gas-filled capsule yields increase up to 2x

“Bigfoot”¹ platform: starting point for warm magnetized design

“Bigfoot” campaign on NIF

- Robust hotspot: High $\rho \cdot R$, high velocity, high adiabat, lower convergence
- Shock overtaking in ablator
- Simple hohlraum: low gas fill, short laser pulse, low LPI
- HDC capsule: short laser pulse, smooth capsules
- Tied for highest yield on NIF

1 C. Thomas, APS-DPP invited talk, 2016

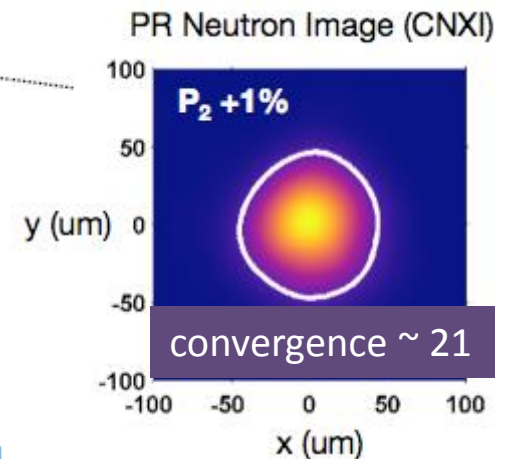
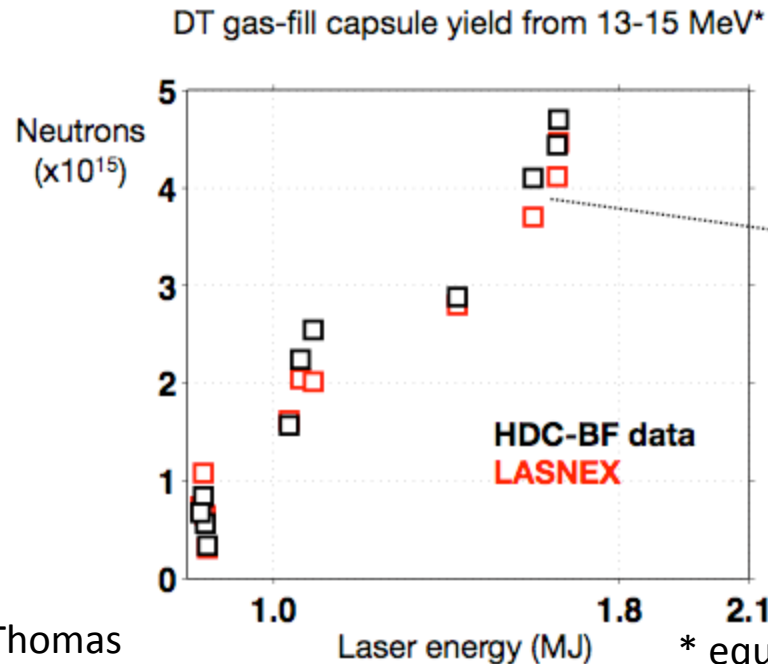


Figure courtesy C. Thomas

* equivalent DT yield from DD, D3He, ...

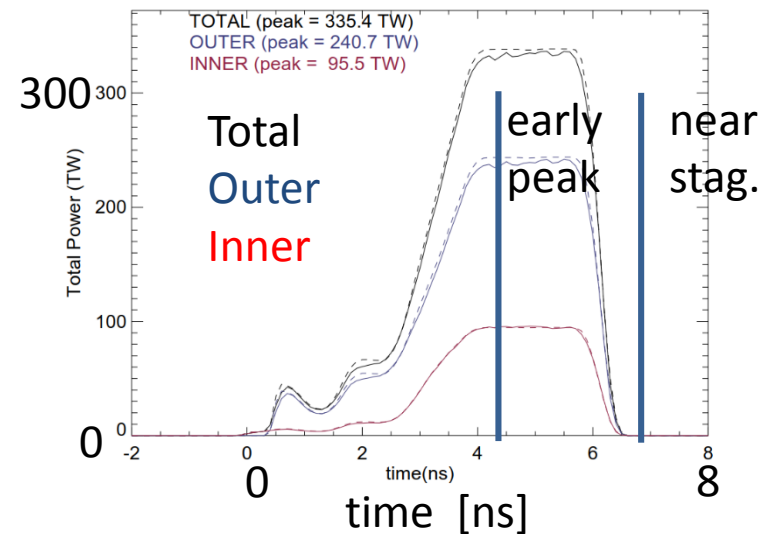
Why Bigfoot for warm magnetized design?

- Don't re-invent the wheel
- “Nice” features → predictable, easy to tune
 - Low LPI, low convergence
- But not so “nice” to be irrelevant!
 - Enough convergence to amplify B field, reduce e- conduction
 - Connection to existing, high-performance cryo platform

N161204: bigfoot NIF shot

- “Subscale” target: less taxing on laser:
 - 1.1 MJ, 340 TW
- Symcap: gas-filled capsule: D[30%]-He3[70%]
 - 5.5 mg/cc
 - no DT ice layer
- HDC capsule, W dopant
- Au hohlraum
- Low hohlraum gas fill density: 0.3 mg/cc He4

Laser power [TW]



HYDRA MHD model: Single-fluid Braginskii

Bulk momentum:

$$\rho \frac{D\vec{v}}{Dt} = -\nabla p + \vec{J} \times \vec{B}$$

Magnetic pressure:

$$\vec{J} \times \vec{B} = -\nabla \left(\frac{B^2}{2} \right) + \vec{B} \cdot \nabla \vec{B}$$

Maxwell:

$$\begin{aligned} \partial \vec{B} / \partial t &= -\nabla \times \vec{E} \\ \vec{J} &= \mu_0^{-1} \nabla \times \vec{B} \end{aligned}$$

Generalized
Ohm's law:

$$\vec{E} = \underbrace{-\vec{v} \times \vec{B} + \frac{1}{n_e e} \vec{J} \times \vec{B} - \frac{\nabla p_e}{n_e e}}_{\text{collisionless}} + \underbrace{\vec{\eta} \cdot \vec{J} - e^{-1} \vec{\beta} \cdot \nabla T_e}_{\text{collisional}}$$

advection / induction term
Hall term
Biermann battery
resistivity
thermal force*

- Plus analogs in electron energy equation
- Full Braginskii available in HYDRA
- No nonlocal limiting of Nernst: Brodrick, Sherlock

Just Nernst advection (draw B to lower T_e)
No Righi-Leduc in energy eq.

This talk:
$$\vec{E} = -\vec{v} \times \vec{B} - \frac{\nabla p_e}{n_e e} + \eta \vec{J} - e^{-1} \vec{\beta} \cdot \nabla T_e$$

→HYDRA Simulations: no imposed field

HYDRA Simulations: imposed axial field

N161204 “post-shot” sims: no imposed B field: Close on bangtime and yield



HYDRA methodology

- R-Z axisymmetric
- “HyPyD”: Pythonic framework:
J. Koning, J. Salmonson
- DCA non-LTE: Sept. 2017 model: H. Scott
- Electron heat flux limit $f = 0.15$ (high)
- X-rays on capsule artificially symmetrized

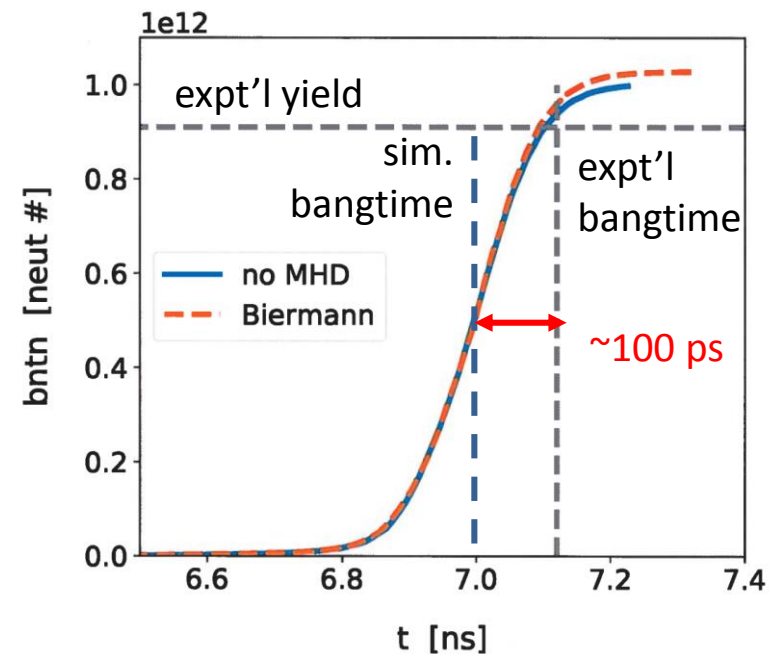
Ohm's law:
this slide

$$\vec{E} = -\vec{v} \times \vec{B} - \frac{\nabla p_e}{n_e e} + \eta \vec{J}$$

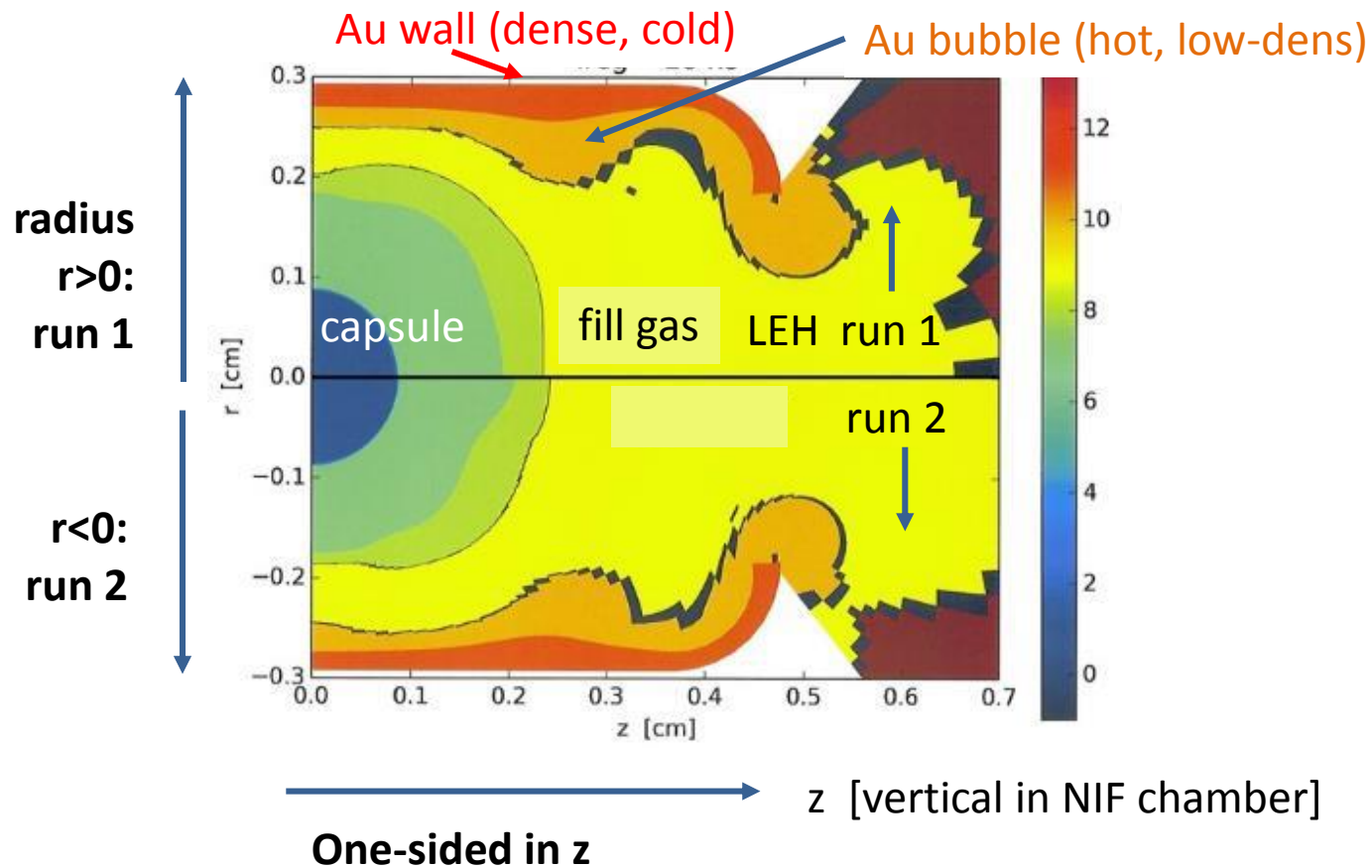
Without any hand tuning

- Sims' bangtime slightly early ~ 100 ps
- Sims 10% above measured yield
- Biermann fields have little effect

neutron yield



Hohlraum map legend

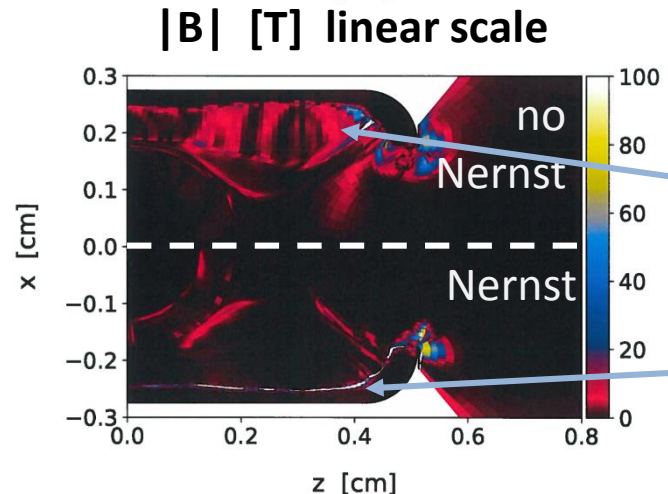
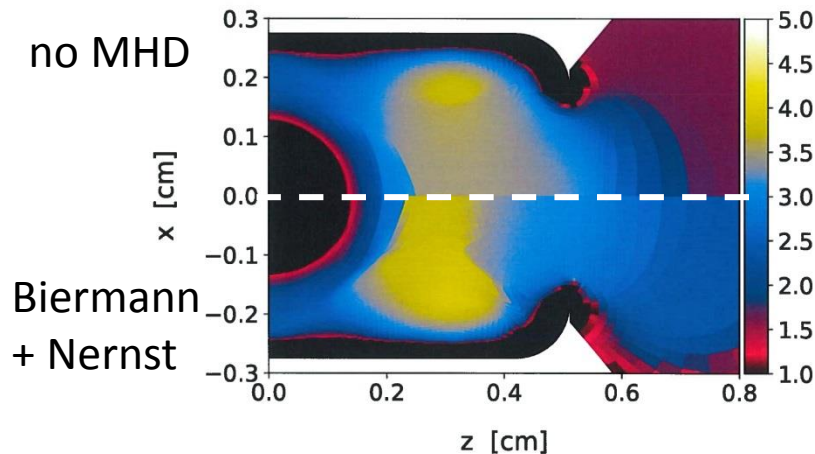
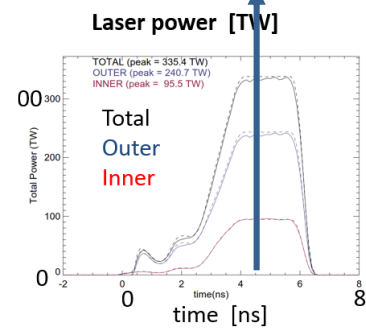
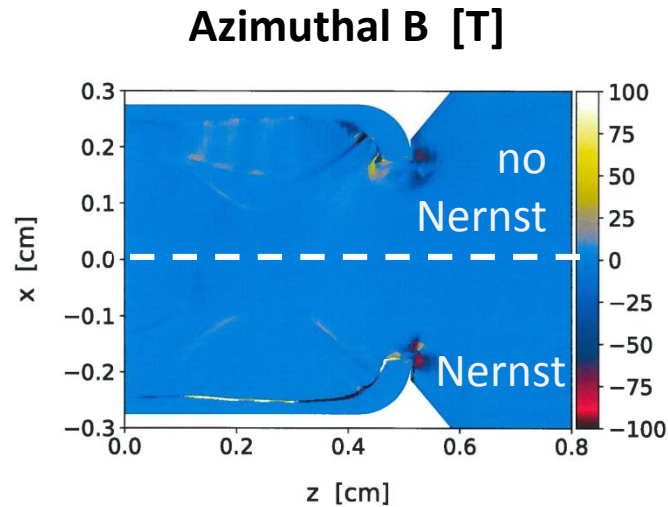
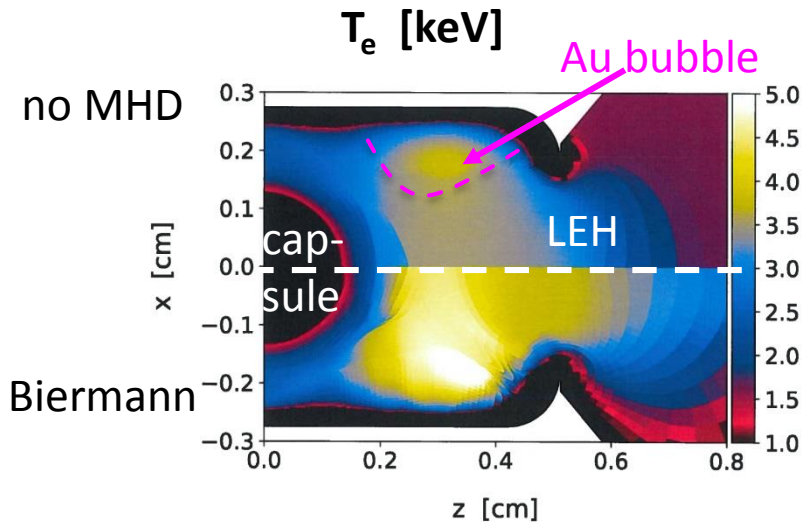


N161204: Biermann fields increase T_e , Nernst advection reduces the effect



Same story as Farmer 2017

Plots at 4.25 ns: early peak power



* Nernst advection
“erases” much of
Biermann field

* Draws field
deeper in to cold
Au wall

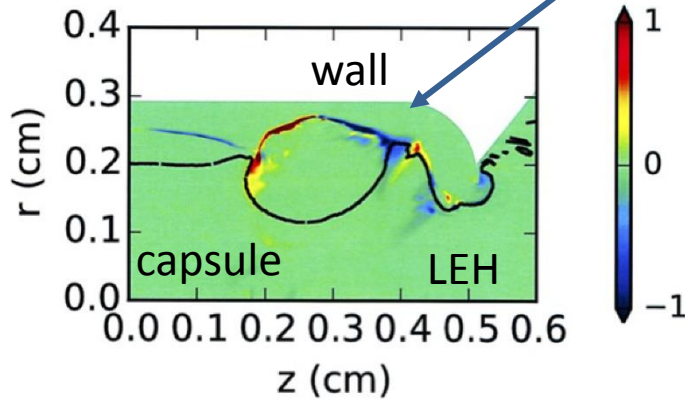
Hohlraums, no imposed field: Farmer PoP 2017

NIF shot N151122

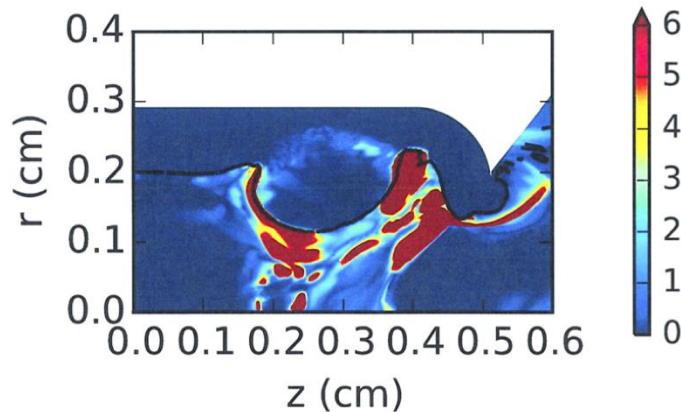
HDC capsule

0.3 mg/cc hohlraum gas fill

azimuthal B [MG]



Hall parameter $\omega_{ce}\tau_{ei}$



MHD: Biermann + Nernst

Highly localized
~ 100 T fields

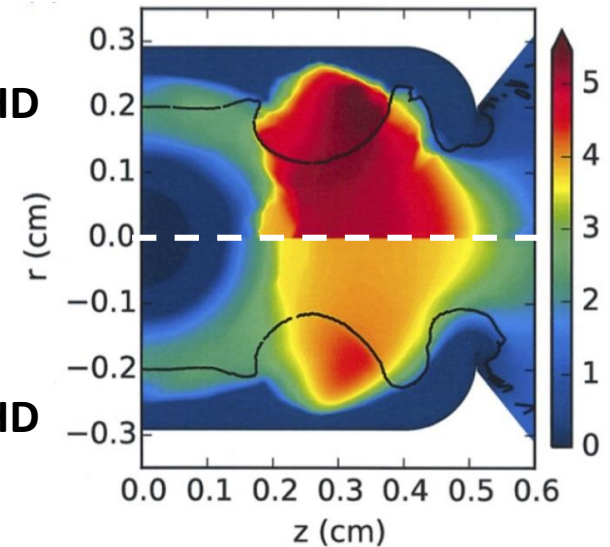
Plots at 5 ns:
late peak power

T_e [keV]

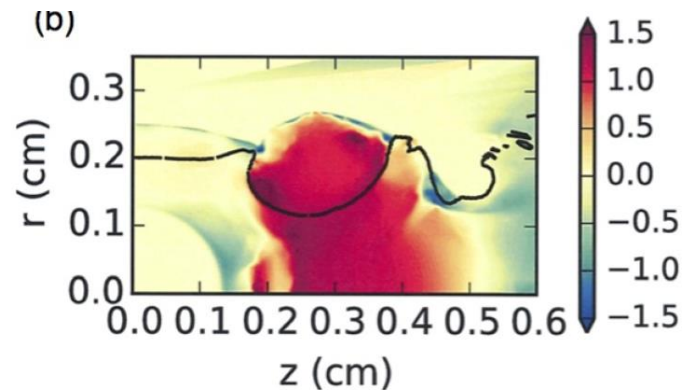
Farmer '17

MHD

**No
MHD**



T_e MHD – no MHD [keV]



Hohlraums, no imposed B: Nernst advection reduces effect of B field

Farmer '17

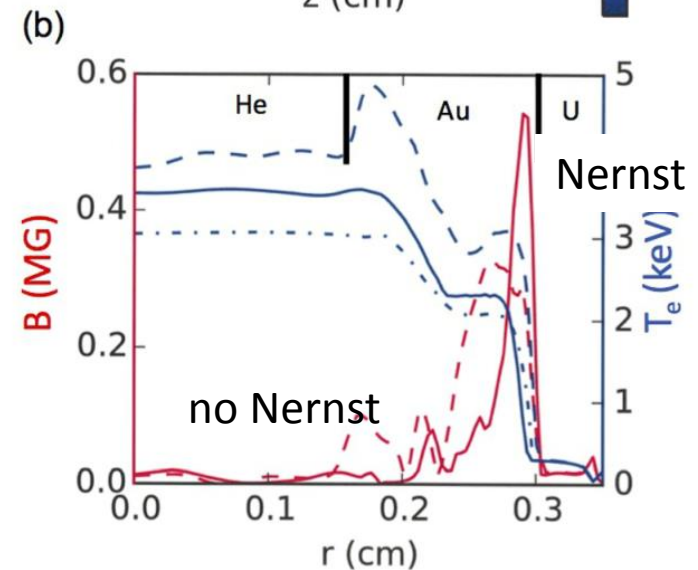
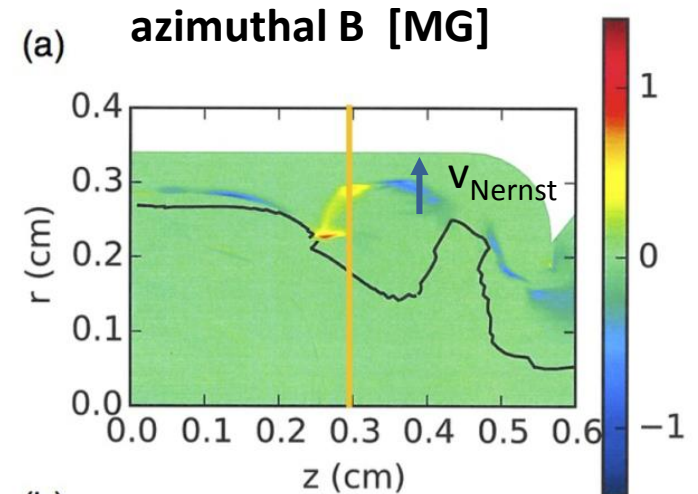
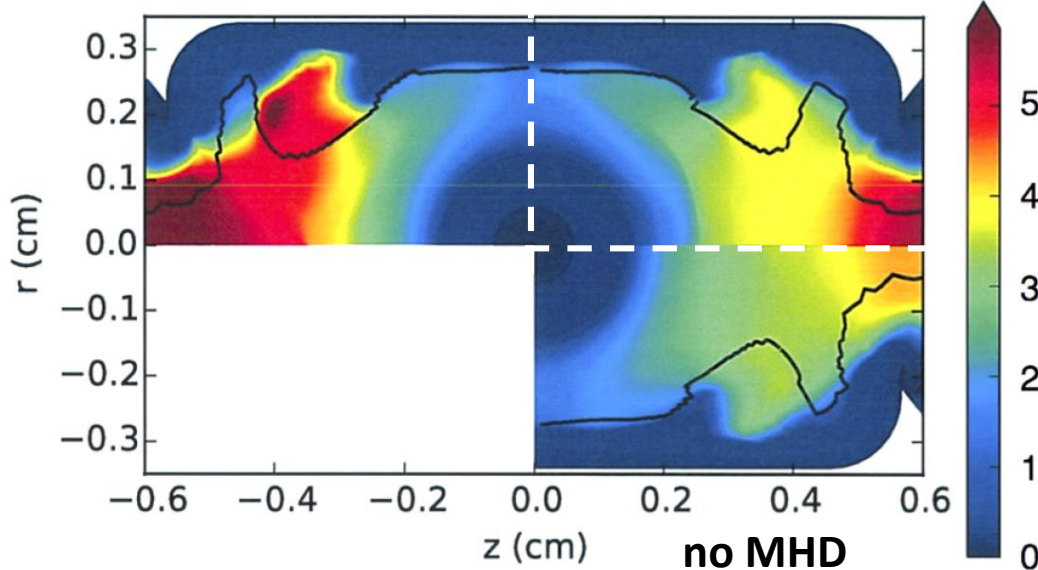
“High foot” design

CH capsule

0.6 mg/cc hohlraum gas fill

MHD no Nernst

MHD with Nernst



*“What Biermann giveth, Nernst taketh away”
– M. D. Rosen*

HYDRA Simulations: no imposed field

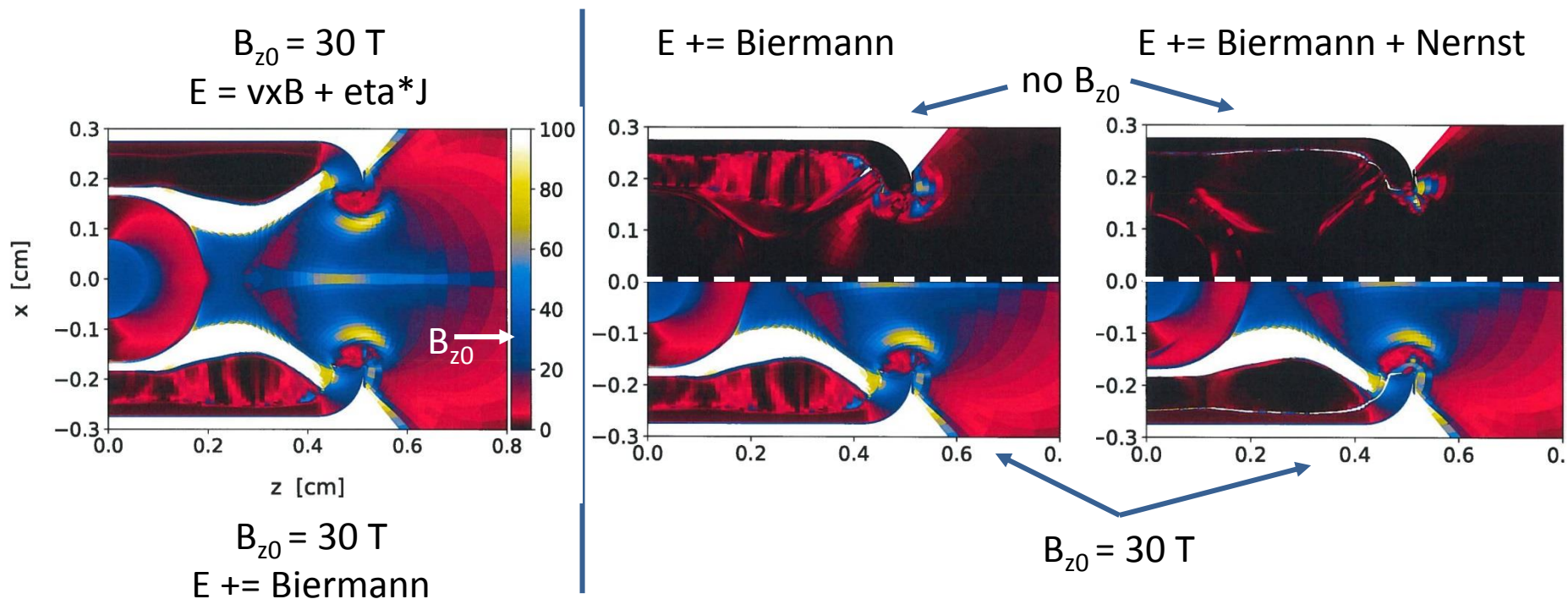
→ HYDRA Simulations: imposed axial field

N161204: Imposed $B_{z0} = 30$ T: field “adds” with Biermann in bubble / LEH



|B| [T]: same colormap

Plots at 4.25 ns: early peak power



- Imposed-field dynamics unchanged by Biermann or Nernst
- Biermann fields unchanged by imposed – at least by eye

N161204: Imposed $B_{z0} = 30$ T: little effect on hohlraum fill vs. Biermann



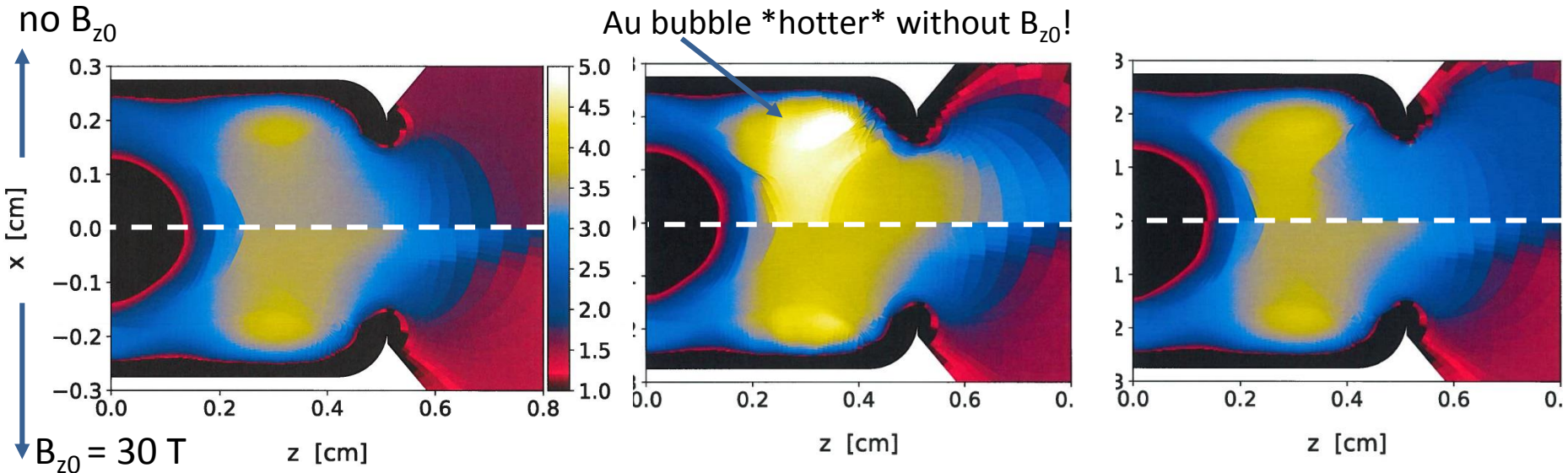
T_e [keV]: same colormap

Plots at 4.25 ns: early peak power

$E = v \times B + \eta * J$

$E += \text{Biermann}$

$E += \text{Biermann} + \text{Nernst}$



Why small effect from B_{z0} ?

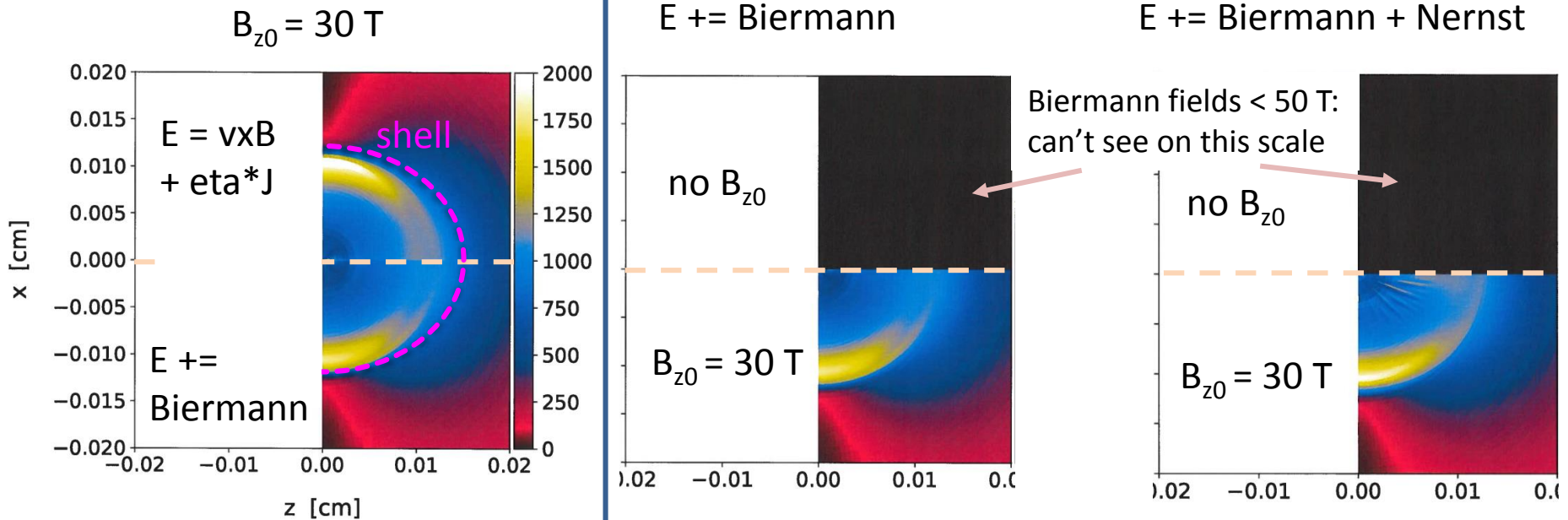
- Hall parameter > 1 in He gas fill with imposed field – not a “small field”
- imposed B reduced in Au bubble due to expansion: Frozen-in law
- Axial imposed field \rightarrow B in r-z plane: heat flow only reduced in 1 meaningful direction
- Biermann azimuthal field \rightarrow 2 directions reduced
- Seems we need B inside Au to increase T_e : Biermann does, imposed doesn't

N161204: Imposed $B_{z0} = 30$ T: capsule B field ~ 2 kT; Biermann fields small



$|B|$ [T]: same colormap

Plots at 6.75 ns: 0.25 ns before bangtime
x-ray flux on capsule artificially symmetrized



Frozen-in estimate of field increase

- Capsule initial radius 908 μm
- Radius at this time ~ 100 μm
- B increase $\sim (R_{\text{initial}} / R_{\text{final}})^2 = 81\times : 30 \text{ T} \rightarrow 2400 \text{ T}$

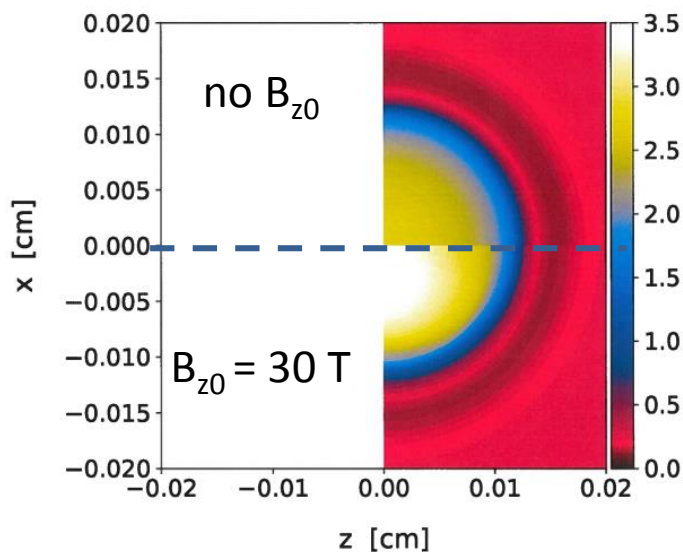
N161204: Imposed $B_{z0} = 30$ T: capsule hotter for all MHD models



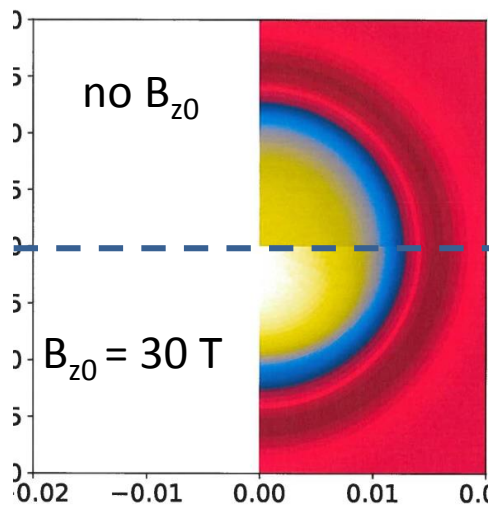
T_e [keV]: same colormap

Plots at 6.75 ns: 0.25 ns before bangtime
x-ray flux on capsule artificially symmetrized

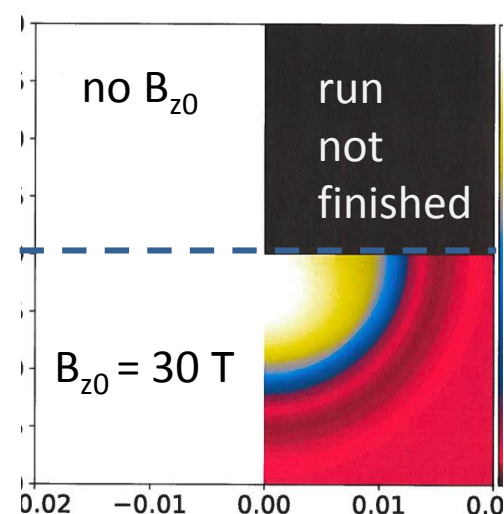
$E = vxB + \eta J$



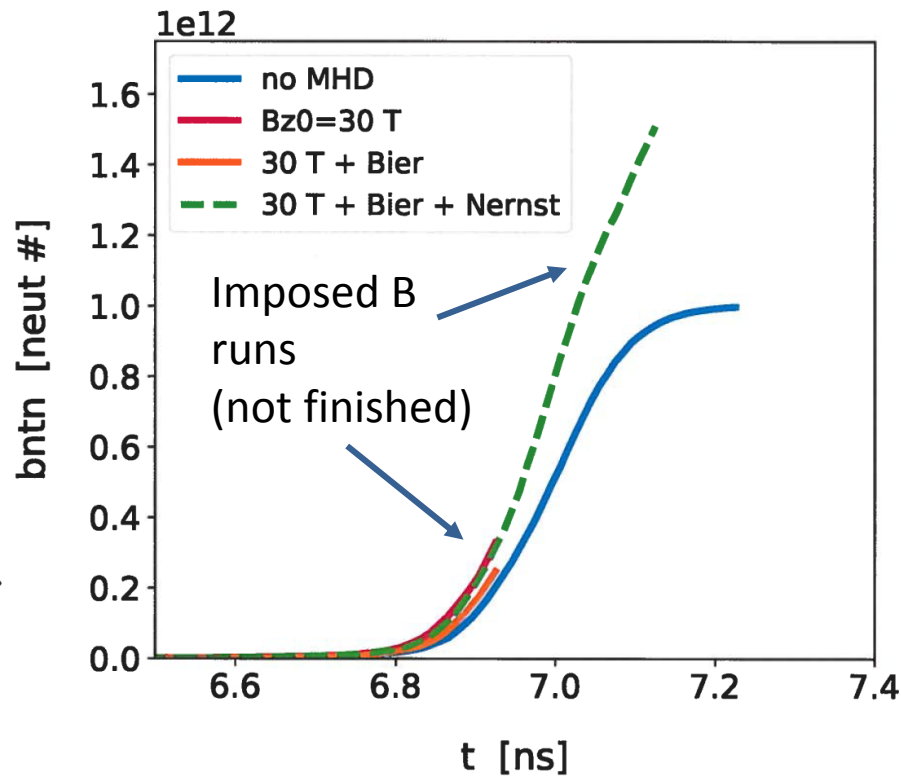
$E += \text{Biermann}$



$E += \text{Biermann} + \text{Nernst}$



N161204: Imposed $B_{z0} = 30$ T: bangtimes slightly earlier; yields higher

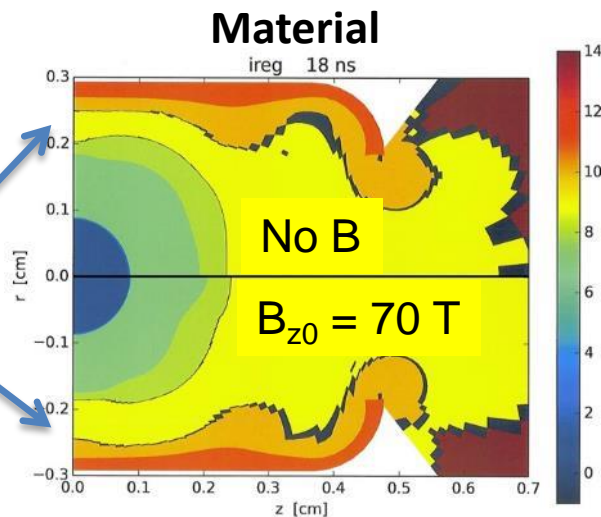


Imposed axial field (70 T) slightly raises T_e , improves inner-beam propagation

Strozzi '15
 $B_{z0} = 70$ T

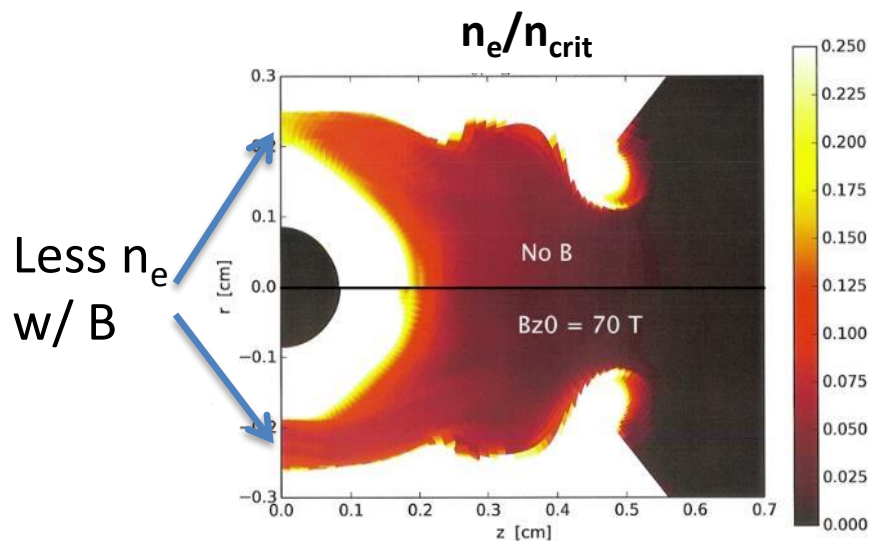
“Low-foot” shot N120321
CH capsule
18 ns: early peak power

Wider equator
channel with B



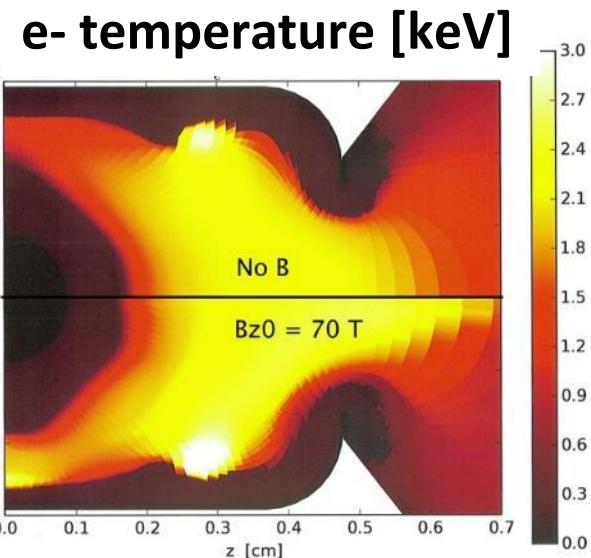
D J Strozzi, L J Perkins, et al.,
J. Plasma Phys. (2015)

Each figure: hohlraum
quadrants with initial $B_{z0} =$
70 T (top), and without MHD
(bottom)



Less n_e
w/ B

Higher T_e
w/ B, esp.
on equator



Imposed B field: 10 T similar effect in hohlraum as 70 T

Strozzi '15
 $B_{z0} = 10$ T

High-foot shot N121130
 $B_{z0} = 10$ T
15.2 ns: peak power

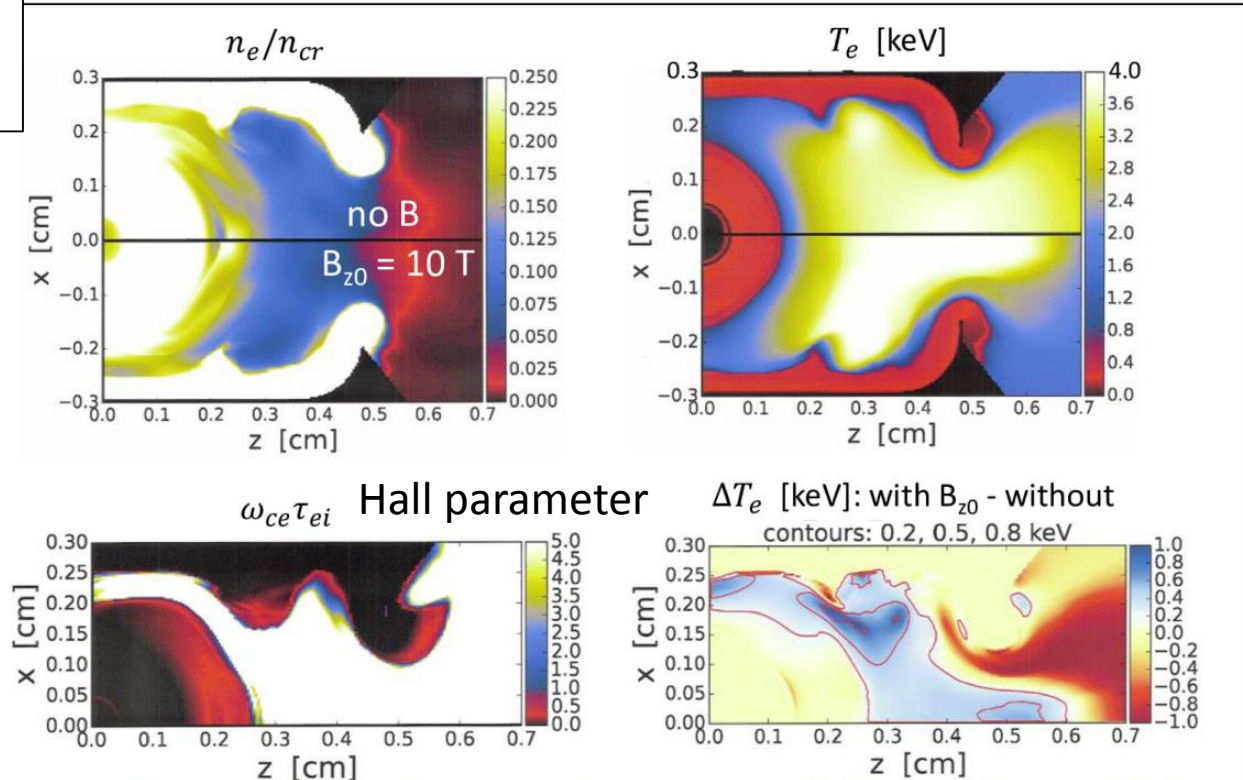


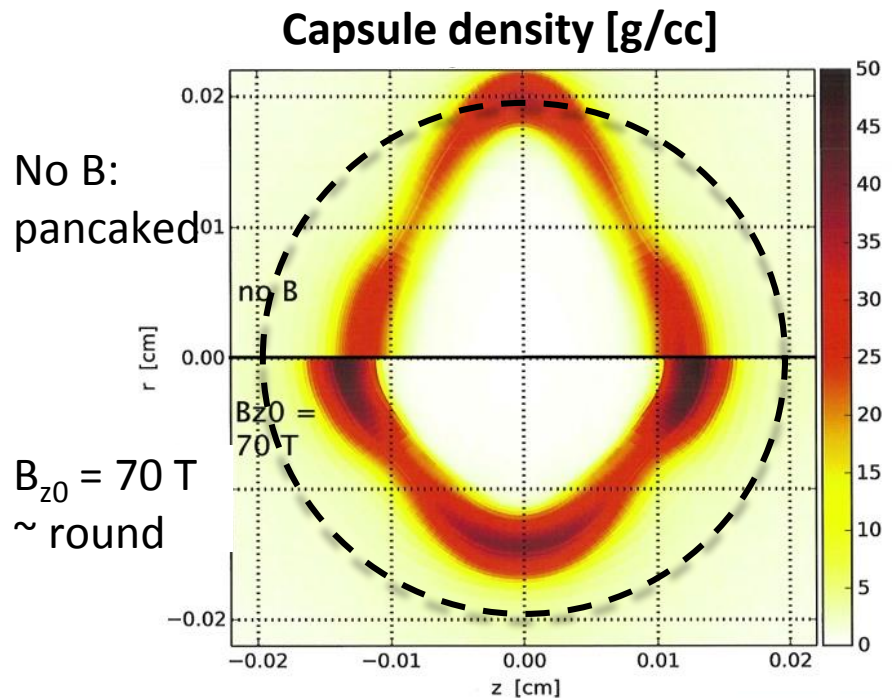
Figure 5. Plasma conditions at 14 ns (late peak power) from HYDRA simulations of NIF shot N121130. For n_e and T_e plots, top half ($x > 0$) has no field, and bottom half ($x < 0$) has $B_{z0} = 10$ T. The Hall parameter $\omega_{ce} \tau_{ei}$ is capped at 5 for clarity.

L. J. Perkins et al.,
LDRD final report

Imposed B: improved inner beam propagation, less pancaked implosion

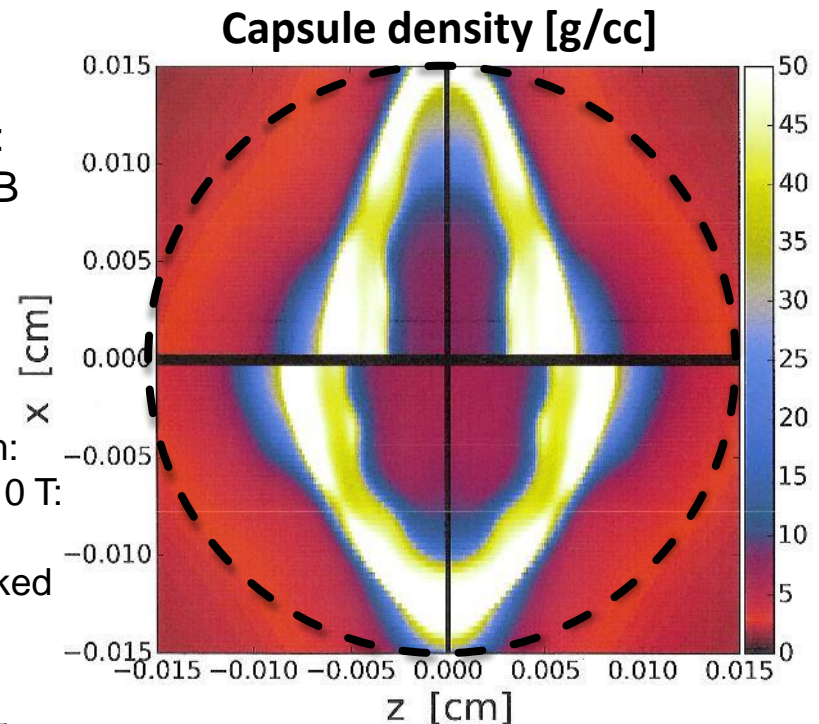
Strozzi '15
 $B_{z0} = 10, 70 \text{ T}$

Low-foot shot N120321¹
 $B_{z0} = 70 \text{ T}$
21.5 ns: end of pulse

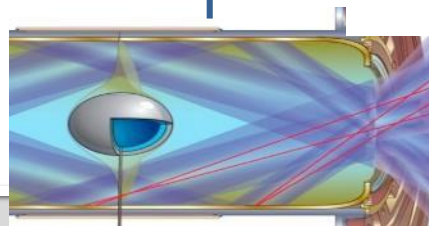


¹D. J. Strozzi, L. J. Perkins, et al.,
J. Plasma Phys. (2015)

High-foot shot N121130²
 $B_{z0} = 10 \text{ T}$
15.2 ns: peak power



²L. J. Perkins et al.,
LDRD final report



Magnetized “warm” (293 K) gas-filled capsules: established NIF process for cryo analogs

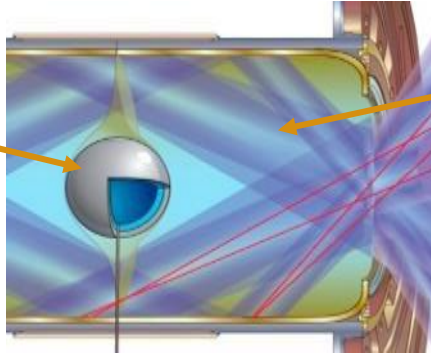
HDC capsule fill

cryo: 5.5 mg/cc D-He3

warm: pure D or D-He3

Magnetized shots from

TANDM, can't easily handle T



Hohlraum fill

cryo: 0.3 mg/cc He4

warm: C5H12, ~ same e- density

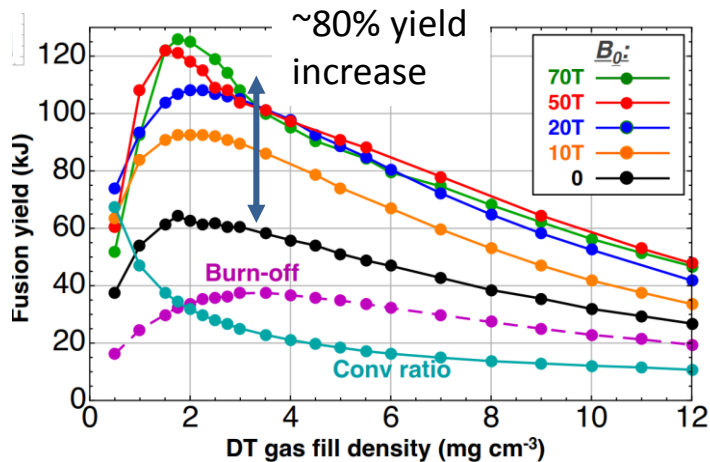
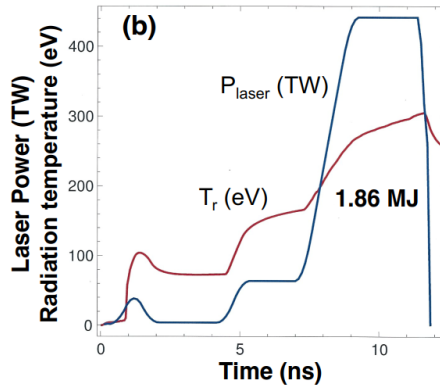
He4 → too much pressure on window

J. E. Ralph, D. J. Strozzi, et al., Phys. Plasmas 2016

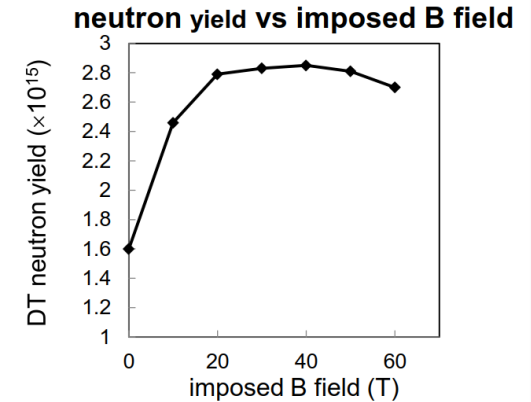
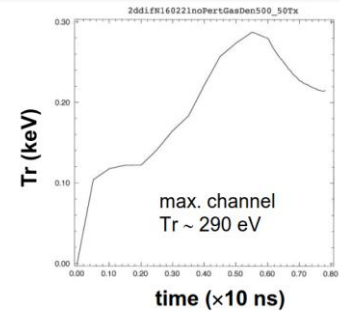
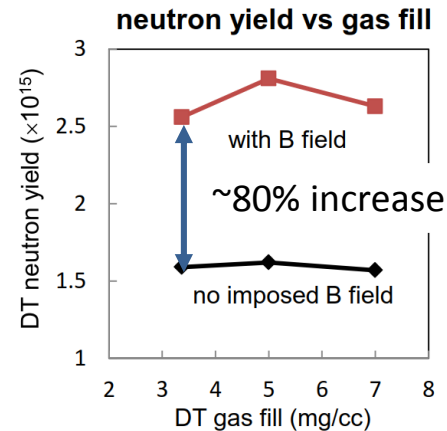
- Warm analogs of “low-foot” CH implosions
- Backscatter, x-ray drive, implosion shape similar
- Capsule gas: C3D8 – light species (H, D, ...) diffuse through CH –could aluminize
- HDC capsules should hold light species

Magnetized gas-filled capsules: up to 2x yield increase with imposed B field

L. J. Perkins [unpublished]:
HDC capsule, low adiabat



D. D. Ho [APS DPP 2016]:
HDC capsule, high adiabat



DT vs. DHe3 gas capsules

- Yield increased mainly by reduced e- conduction
- Not enough alpha's to matter
- Warm shots: D-He3 fill: e- conduction reduction should have similar effect

BACKUP BELOW



Hohlraums, no imposed field: MHD slightly reduces “drive deficit”, implosion less oblate

Farmer '17

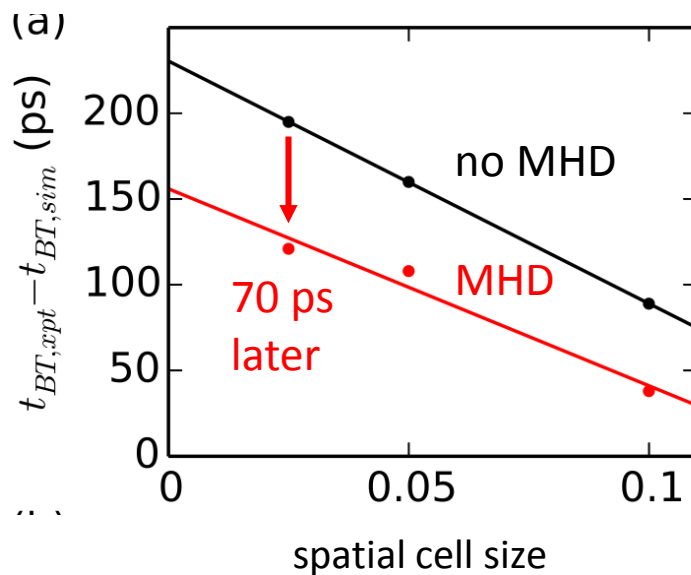
NIF shot N151122

HDC capsule

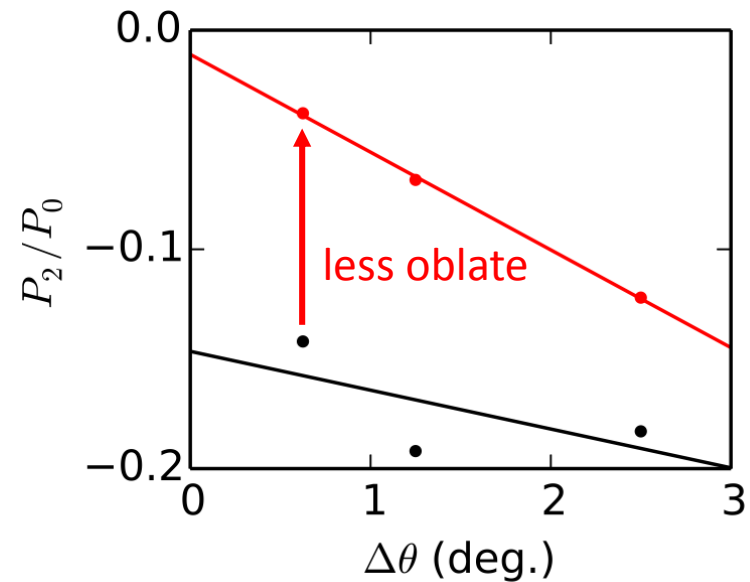
0.3 mg/cc hohlraum gas fill

W. A. Farmer, J. M. Koning, et al.,
Phys. Plasmas 2017

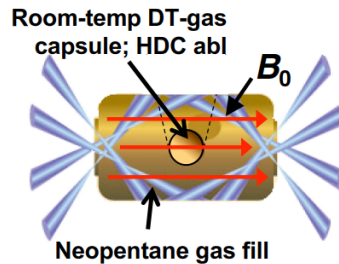
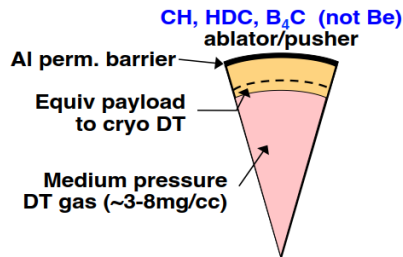
Bangtime: measured – simulated
reflects total x-ray drive



P_2/P_0 : hotspot emission shape



Room-temperature gas target performance, HDC shell – What's the most important role of the B-field?



Most important effect of B for (non-metal) gas targets is on electron heat conduction as there's few alphas.
⇒ Can get interesting results at low imposed B-fields (~20T) because $\omega\tau_e$ is still very high

