Temperature distribution and magnetic field impact on heat flow in laser-heating experiments relevant to MagLIF

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Overview

- Laser heated plasma experiments relevant to Magnetized Liner Inertial Fusion (MagLIF)¹ have been conducted at Z
- Time-integrated spectrum and narrowband image data of Ar K-shell emission were collected during these experiments
- We have analyzed data from two experiments with magnetic field (H45, H50) and one without magnetic field (H51)
- Motivating questions:
 - 1. What can be extracted from the data and how can it be analyzed?
 - 2. What physics can we learn from the results?
 - 3. How can the results be used to plan for future experiments and investigation?

¹S.A. Slutz, M.C. Hermann, R.A. Vesey, A.B. Sefkow, D.B. Sinars, D.C. Rovang, K.J. Peterson, and M.E. Cuneo, Phys. Plasmas **17**, 056303 (2010)

Laser heating experiments at Z



- Be liner target
- Deuterium gas fill (60 psi) doped with 0.1% (atomic) Ar
- Laser pulse heats the fuel, 1mm diameter focal spot
- Experiments were performed at comparable laser pulse energies
- Time-integrated argon K-shell x-ray emission was observed and can be used for spectroscopy diagnosis
- Experiments with B (H45, H50) and without B (H51) have been conducted and analyzed

Data: B field impacts spectroscopic signatures



Image characteristics

- Time-integrated
- Axial resolution: 80µm
- Transverse resolution: 20µm
- hv₀ = 3124eV
- ∆hν = 7.5eV

Spectrum characteristics

- Time-integrated
- Observe 2-1 transitions in He-like and n=2 Li-like Ar
- Some Lyα emission in hotter regions of H50
- Axially resolved: ~200µm
- Spectral resolving power: ~1800

Spectra are sensitive to T_e

- Ar K-shell model was used to investigate spectrum sensitivity to electron temperature (T_e) and density (N_e)
- NLTE, steady-state simulations of D-Ar plasma produced synthetic spectra
- Important: radiation transport has small effect on the spectral range used for analysis
- PrismSpect¹ was used for the calculation
- Experimental conditions are below threshold of N_e dependence (~10²² cm⁻³)²
- A constant uniform N_e based on initial conditions and ionization of D is assumed in our analysis (2.0 x 10²⁰ cm⁻³)

Synthetic spectra



¹J. J. MacFarlane, I. E. Golovkin, P. Wang and P. R. Woodruff, High Energy Density Physics 3, 181 (2007) ²V.L. Jacobs, J.E. Rogerson, M.H. Chen, and R.D. Cowan, Physical Review A 32, 3382 (1985)

Analysis of axially resolved spectra produces $T_{e}(z)$



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Multi-objective data analysis enables $T_e(r,z)$ extraction



¹K. R. Carpenter et al, APS DPP contributed talk, October 2017, Milwaukee, WI ³L. A. Welser-Sherrill, R. C. Mancini, et al, Physical Review E **76**, 056403 (2007)

- Multi-objective data analysis permits extraction of spatial distribution of plasma conditions¹⁻⁴
- Simultaneous and self-consistent analysis of x-ray spectrum and narrow band image data¹
- Two alternative implementations produce consistent results
 - Forward reconstruction driven by a Pareto genetic algorithm
 - Quasi-analytic method in the optically thin approximation
- A collection of emissivity equations can be solved to get T_e(r,z)
- Emissivity weighted average of $T_e(r,z)$ is constrained with $T_e(z)$ from analysis of axially resolved spectra³

 ²I. Golovin, R. Mancini et al Physical Review Letters 88, 045002 (2002)
⁴T. Nagayama, R. Mancini et al Physics of Plasmas 19, 082705 (2012)

$T_e(r,z)$ was extracted for three experiments



• H50 was hotter and had a longer emission column than H45 and H51

- $T_e(r,z)$ is dependent on the window thickness and external magnetic field
- The results from H45 and H51 show qualitative similarities, though the experimental parameters were quite different

 $T_e(r,z)$ can be used to determine the experimental heat flux

Heat flux due to thermal conduction according to Braginskii¹:

Electrons:
$$\vec{q}_e = -\chi_{\parallel}^e \vec{\nabla}_{\parallel} T_e - \chi_{\perp}^e \vec{\nabla}_{\perp} T_e - \chi_{\Lambda}^e (\hat{B} \times \vec{\nabla} T_e)$$

lons:
$$\vec{q}_i = -\chi^i_{\parallel} \vec{\nabla}_{\parallel} T_i - \chi^i_{\perp} \vec{\nabla}_{\perp} T_i - \chi^i_{\Lambda} (\hat{B} \times \vec{\nabla} T_i)$$

 $\chi^e_{\parallel}, \chi^i_{\parallel}$ are functions of T and n

 $\chi^e_{\perp}, \chi^i_{\perp}, \chi^e_{\Lambda}, \chi^i_{\Lambda}$ are functions of T, n, and B

Details and Approximations:

- Assume no fluid motion: $V_e = V_i = 0$
- Assume $T_e = T_i$
- Used coefficients for a deuterium (Z=1) plasma
- Effect of Ar tracer has not yet been considered
- Azimuthally directed, cross-field component of flux not yet evaluated
- VT determined with finite differences approximation

Electron thermal conductivity varies with B

H51 (B = 0T)

H50 (B = 8.5T)

0

z (mm)



-0.5 6.5

0

r (mm)

6 0.5

Multiply conductivities by 1×10^{30} s⁻¹•m⁻¹ to get actual values

- Perpendicular thermal conductivity is 2-3 orders of magnitude lower than the parallel component
- Without B perpendicular and parallel conductivities are identical, i.e. thermal transport is isotropic

Magnetic field had similar effect in H45 (B=8.5T)



- Perpendicular conductivity in H45 is similar to H50 in both the magnitude and distribution of values
- Parallel conductivities are slightly smaller in H45 due to the lower temperature
- Thicker window used in H45 resulted in a shorter emission column and lower T_e (less laser energy coupling to fuel)

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Multiply conductivities by 1 \times 10^{30} s<sup>-1</sup>•m<sup>-1</sup> to get actual values
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Ion thermal conductivity is not as sensitive to B H50 (B = 8.5T)H51 (B = 0T)



0.5

r (mm)

z (mm)

Multiply conductivities by 1.0×10^{30} s⁻¹•m⁻¹ to get actual values

- Assume T_i=T_e
- Ion thermal conductivity values are one order of magnitude smaller than electrons
- Magnetic field has less of an effect on ions
- Values are similar between the magnetic field and no magnetic field cases

Electron heat flux: H50



- Gradient found with finite differences approximation
- Temperature gradient is predominantly radial
- Near the center, the effect of the B-field results in a flux that is mostly axial

Electron heat flux: H45



- Similar to H50, the temperature gradient is largely radial
- The magnetic field effect on the electrons produces an electron heat flux with large axial components

Conclusions

- We have analyzed time-integrated spectrum and narrowband image data from laser heating experiments relevant to MagLIF, done with/without B field
- Using multi-objective data analysis we have extracted 2-D resolved temperature distributions, T_e(r,z), for three experiments: H45(B=8.5T), H50(B=8.5T), H51(B=0T)
- LEH window thickness and magnetic field had prominent effects on the resultant distributions
- Using the formalism of Braginskii, we used T_e(r,z) to examine the experimental heat flux due to electron and ion thermal conduction
- The magnetic field significantly reduces the radial heat flow
- Thermal conductivity of electrons is more largely affected by B-field than that of ions
- It is important now to compare data analysis results with simulations
- The analysis can be performed on time-resolved data which would help unfold the time history of the distribution of temperature and thermal heat flow
- Still, time-integrated data analysis does show the B effects on heat flow

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Extra slides

Temporally and spatially averaged flux



0.5

0.6

lons

Time averaged energy flow





Two objective analysis produces $T_e(r,z)$



- Method is based on forward reconstruction
- Multi-objective optimization driven by a Pareto genetic algorithm (PGA) is performed
- Search for solutions that produce the best simultaneous and selfconsistent approximations to both pieces of data
- This technique has previously been implemented to extract the spatial structure of implosion cores^{1,2}

¹I. Golovkin, R. Mancini, S. Louis, Y. Ochi, K. Fujita, H. Nishimura, H. Shirga et al, Physical Review Letters **88**, 045002 (2002) ²R. Mancini, Ch. 15 in "*Applications of Multi-Objective Evolutionary Algorithms*", Eds. C. Coello-Coello and G. Lamont, World Scientific Pub. ISBN 981-256-106-4 (2004)