Backlighting of Polar Driven implosions on OMEGA have led to precise comparisons with 2-D simulations

- Low mode perturbations of the shell, due principally to illumination nonuniformity, are easily measured with time-resolved x-ray backlighting.
- 2-D DRACO simulations match the observed shell perturbations in time and shape with some small differences in mass distribution.
- Using beam pointing alone the L=2 perturbations have been minimized leaving only L=4 and higher harmonics.

40 of the OMEGA beams are used to emulate the NIF 48 beam indirect-drive configuration



from 21° to 59°, are used to emulate the NIF geometry.

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Additional OMEGA beams are used for x-ray backlighting.

Abel inversion can be used to determine the plasma density from x-ray radiographs*

Absorption of backlighter x rays along a path follows the relation

$$I = I_0 \exp\left[-\int \kappa(\boldsymbol{E}, \boldsymbol{r}) \, \mathrm{d}\boldsymbol{z}\right].$$

The inverse Abel transform gives the radially dependent opacity

$$\kappa(\boldsymbol{E},\boldsymbol{r}) = \frac{1}{\pi} \int_{\boldsymbol{r}}^{\infty} \frac{\mathrm{d}}{\mathrm{d}\boldsymbol{y}} \left\{ \ln\left[\frac{\boldsymbol{I}(\boldsymbol{y})}{\boldsymbol{I}_0}\right] \right\} \frac{\mathrm{d}\boldsymbol{y}}{\sqrt{\boldsymbol{y}^2 - \boldsymbol{r}^2}}.$$

If the mass absorption coefficient is approximately constant through the plasma, as is the case for bound-free absorption by inner-shell electrons, then O(r) = r(F, r)/r(r - F)

$$\rho(\mathbf{r}) = \kappa(\mathbf{E}, \mathbf{r}) / \mu_{\text{eff}}(\mathbf{E}),$$

where $\mu_{\text{eff}}(E)$ is the mass absorption coefficient averaged over the effective energy band of the radiograph, and can be determined as follows:

$$M_{\text{shell}} = \int \rho(r) dV = 4\pi \int \frac{\kappa(E, r)}{\mu_{\text{eff}}(E)} r^2 dr.$$

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^{*}F. J. Marshall et al., Phys. Rev. Lett. 102, 185004 (2009).

Abel inversion is used to compute the density profiles from framed x-ray radiographs of a polar-driven, direct-drive implosion



The density distributions are determined for each time from the radiographs and can be compared to simulation



^{*}F. J. Marshall et al., Phys. Rev. Lett. 102, 185004 (2009).

The measured areal-density time history is consistent with 1-D simulations up to bang time



Low convergence ratio c_r~10

For the pointing cases with no azimuthal variation the polar-drive implosions varied from oblate to prolate



LLE

Low convergence ratio c_r~10

Framed backlit images and DRACO simulations agree well for oblate- to prolate-shaped PD implosions



The observed shell perturbations are accurately reproduced by DRACO 2-D simulations



F. J. Marshall et al., J. Phys. IV France 133, 153 (2006).

 $-\ell n (I/I_0)$

Low-mode perturbations of the hot spot are determined from framed x-ray images of target self-emission



The observed hot-spot perturbations deviate significantly from *DRACO* 2-D simulations

100 µm

OMEGA shot 49331, PD implosion Framed x-ray images 2- to 4-keV emission 2.46 ns 2.51 ns 2.40 ns Stalk direction indicated with arrows The differences in measured and simulated shapes may be due to the stalk, whose effect is not included in the simulations. Viewed from $\theta = 101^{\circ}$ DRACO/Spect3D simulation

LR

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High convergence ratio c_r~18, slight L=2 (prolate)

First experiments and simulations indicated that an L=2 component was present in the implosions



High convergence ratio c_r~18, slight L=2 (prolate)

Modal decomposition of the results and simulations indicated both a small L=2 and a larger L=4 component



High convergence ratio c_r~18, minimum L=2

The measured and simulated radiographs exhibit nearly identical shapes for the 30,150,150 beam pointing case



High convergence ratio $c_r \sim 18$, minimum L=2

The modal decomposition of 30,150,150 pointing case agrees well with the DRACO simulation



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