Effect of Flow on Laser–Plasma Interactions near the Quarter-Critical Density in the Plasmas of Inertial Confinement Fusion





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In the modeling of direct-drive inertial confinement fusion plasmas, the reduction of two-plasmon decay instability and of the laser absorption due to plasma flow is observed



- Two LPI processes: crossed-beam energy transfer (CBET) and two-plasmon decay (TPD) have a large influence on the coupling of laser power to plasmas, and these two processes coexist in the plasma region near the quarter-critical density.
- Three-dimensional modeling using the LPSE platform has confirmed that plasma flow can limit the growth of TPD and increase CBET for the conditions relevant to the experiments on the OMEGA Laser System.
- At the optimum magnitude of plasma flow close to Mach-1 the intensity of plasma waves and the laser absorption is decreased by about one third compared to no flow case.



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University of Rochester Laboratory for Laser Energetics Laser-plasma interaction has been modeled in OMEGA plasmas near the quarter-critical density where the laser intensity profile is produced by multiple laser beams

Ray tracing of



2 LPI processes:



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A. Simon *et al*., Phys. Fluids <u>26</u>, 3107 (1983).

Computational platform *LPSE* can model nonlinear LPI, including TPD and CBET



• The model describes the evolution of laser light (enveloped near frequency ω_0), plasma-wave field (near ω_p), and ion-acoustic perturbation δN

Laser light
$$i\frac{\partial V_{0}}{\partial t} + i\gamma_{0} \circ V_{0} + \frac{c^{2}}{2\omega_{0}} \nabla^{2}V_{0} + \frac{\omega_{0}^{2} - \omega_{p}^{2}\left(1 + N_{0} + \delta N\right)}{2\omega_{0}}V_{0} = \frac{i\omega_{p}}{4\omega_{0}}\left[\left(\nabla \cdot V_{p}\right)V_{p}\right]_{T} \cdot e^{-i\delta\omega t}$$
Plasma wave
$$i\frac{\partial V_{p}}{\partial t} + i\gamma_{p} \circ V_{p} + \frac{3V_{Te}^{2}c^{2}}{2\omega_{1}} \nabla^{2}V_{p} + \frac{\omega_{p}\left(1 + N_{0} + \delta N\right)}{2}V_{p} = \frac{1}{\omega_{p}}\left[\nabla\left(v_{p}^{*} \cdot v_{0}\right) - \left(\nabla \cdot V_{p}^{*}\right)V_{0}\right] \cdot e^{i\delta\omega t}$$

$$\text{Ion acoustic} \qquad \frac{\partial^2 \delta N}{\partial \tau^2} + 2\gamma_{\text{ia}} \times \frac{\partial \delta N}{\partial \tau} - \frac{c^2}{2\omega_0} c_s^2 \nabla^2 \delta N = \frac{1}{16\pi n_0 m_{\text{i}}} \nabla^2 \left[\left| E_p \right|^2 + \frac{n_0}{n_c} \left| E_0 \right|^2 \right]$$

where
$$V_j = \frac{ieE_j}{m_e\omega_j} (j = 0, p) \frac{\partial}{\partial \tau} = \frac{\partial}{\partial t} + U_0 \times \nabla$$
,

U₀: flow; N₀: background density profile; $\delta \omega = \omega_0 - 2\omega_p$: frequency mismatch

It is possible to study the relative importance of different wave-coupling processes.

Modeling included all OMEGA laser beams – going into the target and also going out from the target





A. Simon et al., Phys. Fluids 26, 3107 (1983).

Simulations with LPSE found that the intensity of plasma waves in nonlinear TPD depends on the flow velocity



$I_{\rm av} = 4.8 \times 10^{14} \ {\rm W/cm^2}$

 $\eta = 1.4$

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The instability region is clearly localized near the quarter-critical density, It is more narrow for M=1.1 than for M=0 ROCHESTER

Two mechanisms for the coupling of laser energy to plasma are the collisional damping and the Landau damping of plasma waves





The nonlinear saturation of plasma waves depends on the flow magnitude.

The growth of CBET and TPD can be characterized by the magnitude of low-frequency density perturbations involved in CBET and in saturation of TPD





The laser absorption in the region near quarter-critical density is a result of both TPD and CBET, and depends on the flow velocity





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The largest reduction of laser absorption is found at when flow velocity is close to Mach 1.

Conclusions

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