Magnetized Collisionless Shock Formation Mediated by the Modified Two-Stream Instability



1.5 $t = 0.10 T_{ci1} (0.126 \text{ ns})$ 1.0 0.5 $k_y [\omega_{pe}/c]$ 0.0 -0.5-1.0 $E_{\rm x}$ Fourier spectrum -1.5+-1.5 -1.0-0.50.0 0.5 1.0 1.5 $k_x \left[\omega_{pe} / c \right]$ 20 40 60 80 100 0 a.u.

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Summary

A laboratory achievable perpendicular magnetized collisionless shock is mediated by a modified two-stream instability

- Supercritical perpendicular magnetized collisionless shock can be readily formed on TeraWatt laser systems (e.g. OMEGA-EP)
- Mode analysis in the shock transition region indicates a modified two-stream instability (MTSI) provides the main dissipation for shock formation
- The growth rate of MTSI (γ_{MTSI}) is much larger than (>10x) the ion gyro-frequency (Ω_{ci}); realistic ion/electron mass ratio substantially separates the shock formation time from T_{ci} (= $2\pi/\Omega_{ci}$)
- Shock reflected ions gyrate in the upstream and participate in shock front reformation



Collaborators



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High-power lasers and computational advances enable study of collisionless shocks in the laboratory

Studies on magnetized collisionless shock

- Telescopes*
- In situ spacecraft missions**
- Laser experiments[†]
- Simulations[‡]



Bow shock near galactic center* (Gemini North, Oct. 17, 2000)



International Sun-Earth Explorer (ISEE) Program**



Drive Lasers



IIE

Quasi-1D particle-in-cell simulation[‡]

** UCLA, http://www-ssc.igpp.ucla.edu/ssc/isee.html; Wang et al., Geophys. Res. Lett. 46, 562 (2019)

plasma plumes experiment[†]

- [†] Woolsey *et al.*, Phys. Plasmas **8**, 2439 (2001); Schaeffer *et al.*, Phys. Plasmas **19**, 070702 (2012); Schaeffer *et al.*, Phys. Plasmas **24**, 041405 (2017) [‡] Matsukiyo *et al.*, J. Geophys. Res. **108**, 1459 (2003); Matsukiyo *et al.*, ApJ **742** 47 (2011);
- Park et al., Phys. Plasmas 19, 062904 (2012); Park et al., ApJ 765 147 (2013); Schaeffer et al., Phys. Plasmas 27, 042901 (2020)

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^{*} NASA, https://apod.nasa.gov/apod/ap001017.html

A collisionless shock can be launched by a laser-driven super-magnetosonic piston on TeraWatt laser systems



$$\mathbf{v}^{[\text{sim./wall/piston}]} = \mathbf{v}^{[\text{lab}]} + V_{\text{d}}\mathbf{e}_{x}$$



Modified two-stream instability provides dissipation for the collisionless shocks



The shock can form within an ion gyro-period T_{ci} (=2 π/Ω_{ci}), or, a few tenths of a nanosecond under $B \sim 50$ T



A supercritical hydrogen shock forms within 0.1 T_{ci}



• The signature of a supercritical shock: reflected ions



Mach #: $M_s \sim 5.62, M_A \sim 2.06$

Weibel instability is suppressed by the strong external magnetic field



Shock compression ratio: $r \sim 2.35$

lons are reflected by self-generated *E*-field, resulting in MTSI

• Low-energy tail of incoming ion is reflected to $v_{xi} \sim 2V_{\text{shock}}$ by an electrostatic E_x field*





Mode analysis indicates MTSI is the operating instability





Motional electric field also accelerates reflected ions in the shock tangential direction

• Reflected ion is co-accelerated by the motional E_z field to 6.5-9.7 keV (average ~8 keV)



Accelerated ions will participate in shock front reformation at later times



Longer time 1D simulations shows shock front reformation as shock evolves

- Shock front reformation happens on length scale of $\sim \rho_{i1}$, with a period of $\sim 0.55 T_{ci1}$
- Rankine–Hugoniot jump conditions are satisfied for a portion of the duration in each reformation period (white dashed line)





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