



Strong Suppression of Heat Conduction in Laser-driven Magnetized Turbulent Plasmas



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Flash Center for Computational Science Department of Physics and Astronomy Laboratory for Laser Energetics University of Rochester 63rd Annual Meeting of the APS Division of Plasma Physics November 8–12, 2021 Pittsburgh, PA Abstract: UO03.00010



TDYNO Collaboration





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- ear Security Administration DOE NNSA, DOE Office of Science, NSF, EPSRC
- DOE's INCITE & ALCC (@ANL), NLUF (@LLE) & DS (@LLNL)
- DOE NNSA NLUF: DE-NA0002724, DE-NA0003605, DE-NA0003934
- NSF/DOE Partnership in Basic Plasma & Engineering: PHY-2033925
- and General Atomics for target R&D and manufacturing!













- The advent of high-power lasers has opened new fields of plasma physics research and has enabled the reproduction of astrophysical environments in the laboratory.
- The TDYNO platform, which was designed with FLASH simulations to demonstrate fluctuation dynamo in the lab for the first time, at the Omega Laser Facility at the Laboratory for Laser Energetics (Tzeferacos+ 2018), has been successfully ported onto the National Ignition Facility (Meinecke et al. submitted, arXiv:2105.08461).
- The NIF experiments achieve fluctuation dynamo in the large magnetic Prandtl number regime, an experimental replica for the magnetized turbulence in the IGM.
- The strong magnetic fields generated in the experiments resulted in a reduction of local heat transport by two orders of magnitude or more, leading to strong temperature variations on small spatial scales, as is seen in cluster plasmas (Markevitch+ 2003).



Thermal transport in turbulent magnetized plasmas



- Galaxy clusters are diffuse, turbulent magnetized plasmas.
- In cluster cores, the temperatures remain anomalously high compared to what might be expected, given that the cooling time is short relative to the Hubble time.
- While feedback from the central active galactic nuclei is believed to provide most of the heating, there has been a long debate as to whether conduction of heat from the bulk to the core might help the core reach observed temperatures.
- Thermal conduction in magnetized, weakly collisional plasmas is a longstanding problem in ICF and fusion research.







TDYNO Collaboration









Magnetized turbulence & fluctuation dynamo
 Cosmic ray physics



TDYNO concerted effort



Astrophysical phenomena & astrophysical systems	Pm=Rm/Re	Cs	Facility	Dates	
Subsonic turbulence: SN shocks, Coma cluster, ISM	< 1	<1	Vulcan (UK)	2013-2014	
Subsonic dynamo: ISM, stars, planets, accretion disks	< 1, ≳ 1	<1	OMEGA (USA)	2015-present	
UHECR transport in turbulent magnetic fields: CR	< 1	<1	OMEGA (USA)	2016-2020	
Subsonic dynamo: IGM, galaxies, galaxy clusters	>1	<1	NIF (USA)	2016-present	
Supersonic turbulence: GMCs, ISM, star formation	< 1	>1	Vulcan (UK)	2016	
Supersonic dynamo: GMCs, ISM, star formation	< 1	>1	OMEGA (USA)		
Supersonic turbulence: GMCs, ISM, star formation	< 1	>1	LMJ (France)		
Second order Fermi: CR	< 1	<1	GSI (Germany)	Meinecke et al. submitted	
NRH instability: CR	< 1	<1	OMEGA+EP (USA	(arXiv:2105.08461)	



Planetary Science Planet interiors, equation of state, phase transitions, material properties









NIF Discovery Science and Laboratory Astrophysics

Laboratory Plasma Astrophysics Shocks, turbulence, particle acceleration and transport, hydro and MHD instabilities

















Nuclear Astrophysics Nuclear reaction rates and cross sections in stellar interiors, s- and r-processes









The FLASH code



- FLASH (Fryxell+ 2000) is a publicly available, high-performance computing (HPC), adaptive mesh refinement (AMR), finite-volume, radiation MHD code with extended physics capabilities (Tzeferacos+ 2015). Supported primarily by the U.S. DOE NNSA, LANL, and LLNL.
- FLASH is professionally managed software in continuous development for 20 years: coding standards; version control; daily automated regression testing; extensive documentation; user support; integration of contributions from external users.

> 3,500 users worldwide flash.rochester.edu >1,200 papers published with FLASH





NIF TDYNO platform





FLASH: plasma state and B-fields







- The FLASH simulations (Feister+ to be submitted) state with *Pm* >> 1 and an outer scale *L* ~ 600 μ characterized by *Rm* in the thousands, orders of (Schekochihin+ 2004)
- □ Within a few ns ($t_{eddy} \sim L / u \sim 3$ ns) the B-field r See also the talk by amplification saturates when the magnetic energy available kinetic energy (RMS)

X-ray imaging

X-ray imaging shows the formation of a wide (~ 4 mm), thick (~ 1-2 mm), turbulent interaction region after the interaction of the counter-propagating plasma flows

The TDYNO campaign fielded NIF's new OTS diagnostic to measure the electron density from the EPW wavelength shifts

Experimental plasma properties

xford

hysics

OTS Electron Density Fits 10^{21} Collection volume: 1,200 um x 50 um Electron Density (cm⁻³) 0 00 cylinder 4mm Analysis: Jena Meinecke and Steven Ross 16 18 20 22 24 26 Time (ns)

□ In the turbulent interaction region, we find N_e ~5x10²⁰ cm⁻³, which is corroborated by FABS; SRS: N_e ~5-8x10²⁰ cm⁻³, SBS: *u*~200 km/s

Proton radiography

PRAD shows protons being strongly deflected by the magnetic fields in the plasma.

Analysis suggests B_{rms}~0.8 MG and B_{max}~2-3 MG.

The magnetic field is strong enough to suppress thermal conduction: data shows hotspots surrounded by cooler plasma.

The magnetic field is generated via dynamo action in a turbulent plasma with with magnetic Prandtl numbers *Pm* ~ 12 and *Rm* ~ 3.5x10³.

Temperature distribution using two-color X-ray emission

- Plasma temperature is determined by comparing the X-ray emission in two different energy channels together with the detector response.
- The technique allows us to create 2D temperature maps of the density weighed path integrated temperature.
- We validated this technique with synthetic postprocessing using SPECT3d and FLASH simulations.

$T_{\rm e}$ profile controlled by magnetization of electrons

□ NIF data revealed **highly-structured temperature profile** when the Larmor radius, r_g , is smaller than the electron Coulomb mean free path, λ_e

The size of hot spots is limited by instrument resolution

The observed temperature is reproduced in simulations only if thermal conduction is off

- FLASH simulations show that such a structured temperature profile can be obtained if thermal conduction is uniformly suppressed.
- What is causing this is not understood (maybe micro-instabilities, e.g., Komarov+ 2018 and Roberg-Clark+ 2018, for whistler-unstable plasmas).

We estimate a significant reduction in heat conduction in the NIF experiments

Quantity	NIF	NIF	Omega
	t = 23ns	t = 25ns	t = 18ns
$n_e(\mathrm{cm}^{-3})$	7e20	5e20	1e20
$T_e \ (\mathrm{keV})$	1.4	1.6	0.4
$l_T(\mu m)$	50	50	200
$L \ (mm)$	1	1	0.5
$Z_{ m eff}$	5.7	5.7	5.3
$t_{\rm cond} \ ({\rm ps})$	27	14	1324
$t_{\rm age} \ ({\rm ns})$	1.7	1.6	1.6
$rac{t_{ m age}}{t_{ m cond}}$	63	114	1.2

- An estimate for heat conduction suppression is obtained by comparing the conduction timescale $t_{cond} \sim k_B n_e \ell_T^2 / \kappa_S$ with the time required for the turbulent structures to persist $t_{age} \sim L/c_s$.
- □ Our estimates suggest $(\kappa/\kappa_S)^{-1} \sim (t_{age}/t_{cond}) > 100-200x$ reduction of heat conduction in the NIF experiments.
- These conditions are, in fact, very similar to what we would expect to see in cluster of galaxies and they would potentially be relevant to the problem of cooling flows.

Projected temperature map of A754 overlaid with CHANDRA X-ray image at 0.8-5 keV (Markevitch+ 2003)

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Extra slides

Quantity	Experimental Value		
Composition of target	48.3% C, 49.6% H, 1.2% O,		
	0.6% N, $0.1%$ Mn, $0.1%$ Co		
Average atomic weight $(\langle M \rangle)$	6.7		
Mean ion charge $(\langle Z \rangle)$	3.6		
Effective ion charge (Z_{eff})	5.7		
Hydrogen effective ion charge $(Z_{\text{eff},B})$	5.5		
Electron temperature $(T_{\rm e})$	$1100 \mathrm{eV}$		
Ion temperature (T_i)	$1100 \mathrm{eV}$		
Electron number density $(n_{\rm e})$	$4.9 \times 10^{20} \text{ cm}^{-3}$		
Turbulent velocity (v_{turb})	$2 \times 10^7 \mathrm{~cm~s^{-1}}$		
Outer scale (L)	$0.06~\mathrm{cm}$		
RMS magnetic field $(B_{\rm RMS})$	$0.8 \ \mathrm{MG}$		
Maximum magnetic field (B_{max})	$3.0 \ \mathrm{MG}$		
Adiabatic index $(\gamma_{\rm I})$	5/3		

Table 1: Summary of target characteristics and direct experimental measurements of plasma parameters at t = 25 ns after the start of the drive laser pulse. The outer scale, L, represents the distance between two neighboring grid aperture centers. Values are reported in CGS, except for the temperature that is given in eV.

Extra slides

Quantity	Formula	Value
Coulomb logarithm $(\mathrm{log}\Lambda)$	$23.5 - \log n_{\rm c}^{1/2} T_{\rm c}^{-5/4} - \sqrt{10^{-5} + \frac{(\log T_{\rm c} - 2)^2}{16}}$	7.2
Mass density (ρ)	$1.7\times 10^{-24} (\sum_j M_j n_j)$	$1.6 \times 10^{-3} \mathrm{~g~cm^{-3}}$
Debye length (λ_D)	$7.4 \times 10^2 \frac{T_{\rm e}^{1/2}}{n_{\rm e}^{1/2}} \left[1 + \frac{T_{\rm e}}{T_{\rm i}} Z_{\rm eff}\right]^{-1/2}$	$4.2\times 10^{-7}~{\rm cm}$
Sound speed (c_s)	$9.8 \times 10^5 \frac{\left[(\langle Z \rangle + 1)\gamma_{\rm I} T_{\rm e}\right]^{1/2}}{\langle M \rangle^{1/2}}$	$3.5\times10^7~{\rm cm~s^{-1}}$
Mach number	$v_{ m turb}/c_{ m s}$	0.6
Plasma β	$4.0 \times 10^{-11} \frac{n_{\rm e} T_{\rm e} + \sum_j n_j T_{\rm i}}{B_{\rm RMS}^2}$	44
H-ion mean free path (λ_{Hion})	$2.1\times10^{13}\frac{T_{\rm i}^2}{\left(Z_{\rm H}^2 Z_{\rm eff,B} n_{\rm e}\log\Lambda\right)}$	$1.3\times 10^{-3}~{\rm cm}$
Electron-ion mean free path $(\lambda_{\rm e})$	$2.1 \times 10^{13} \frac{T_{\rm e}^2}{\left(Z_{\rm eff} n_{\rm e} \log \Lambda\right)}$	$1.2\times 10^{-3}~{\rm cm}$
Electron Larmor radius $(r_{\rm g})$	$2.4 \frac{T_c^{1/2}}{B_{\rm RMS}}$	$1.0\times 10^{-4}~{\rm cm}$
Hydrogen Larmor radius $(\rho_{\rm H})$	$1.0 \times 10^2 \frac{M_{\rm H}^{1/2} T_{\rm i}^{1/2}}{Z_{\rm H} B_{\rm RMS}}$	$4.1\times 10^{-3}~{\rm cm}$
Unmagnetized thermal diffusivity $(\chi_{\rm S})$	$3.0\times10^{21} \frac{T_{\rm e}^{5/2}}{Z_{\rm eff} n_{\rm e}\log\Lambda}$	$5.9 \times 10^6 \ {\rm cm^2 \ s^{-1}}$
Magnetized thermal diffusivity $(\chi_{\rm m})$	$\chi_{ m S} r_{ m g} / \lambda_{ m e}$	$4.7 \times 10^5 \ {\rm cm^2 \ s^{-1}}$
Suppressed thermal diffusivity (χ)	$\sim \chi_{\rm S}/150$	$3.9 \times 10^4 \ {\rm cm^2 \ s^{-1}}$
Turbulent Péclet number (Pe_{turb})	$v_{turb}L/\chi$	~ 30
Dynamic viscosity (μ)	$4.27 imes 10^{-5} rac{M_{ m H}^{1/2} T_{ m i}^{5/2} n_{ m H}}{\log \Lambda Z_{ m eff, B} n_{ m e}}$	$7 {\rm ~g} {\rm ~cm}^{-1} {\rm ~s}^{-1}$
Kinematic viscosity (ν)	μ/ ho	$4.3\times 10^3 \ {\rm cm^2 \ s^{-1}}$
Turbulent Reynolds number $(\mathrm{Re}_{\mathrm{turb}})$	${ m v_{turb}}L/ u$	280
Viscous dissipation scale (l_{ν})	$L/{ m Re}_{ m turb}^{3/4}$	$8.8\times 10^{-4}~{\rm cm}$
In-flow Resistivity $(\eta_{ })$	$3.1 \times 10^5 \frac{Z_{\rm eff} \log \Lambda}{T_{\rm e}^{3/2}}$	$350 \ {\rm cm^2 \ s^{-1}}$
Magnetic Reynolds number $(\mathrm{Rm}_{\mathrm{turb}})$	$\mathrm{v_{turb}}L/\eta_{ }$	3.5×10^3
Magnetic Prandtl number (Pm)	Rm / Re	12
Resistive dissipation scale (l_η)	$L/\mathrm{Pm}^{1/2}$	$2.5\times 10^{-4}~{\rm cm}$

Table 2: Summary of derived plasma parameters at t = 25 ns after the start of the drive laser pulse. The values are reported in CGS.

□ **Step 1**: We take FLASH simulations results for the interaction region.

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- □ Step 3: A relation between filter ratio and temperature is obtained by running PrismSPECT for a range of different plasma conditions.
- Step 4: The mass-weighted (along the line-of-sight) temperature map obtained from the filter ratio is compared to the mass-weighted temperature from FLASH simulations.

