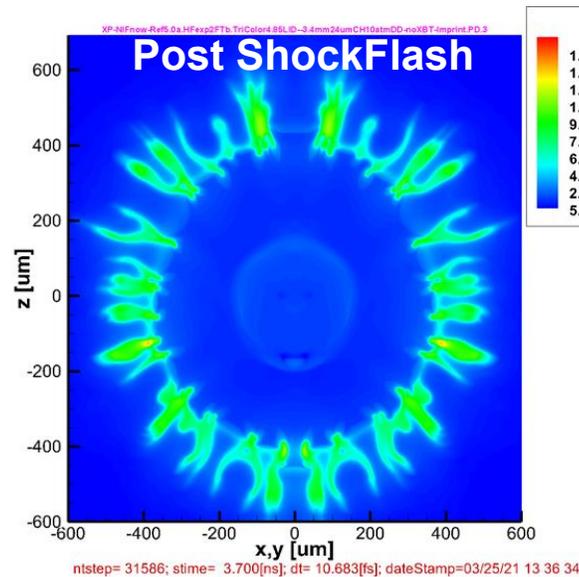
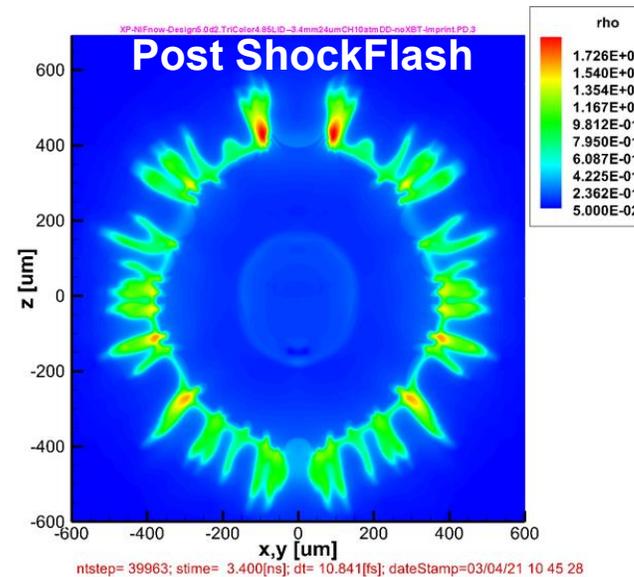


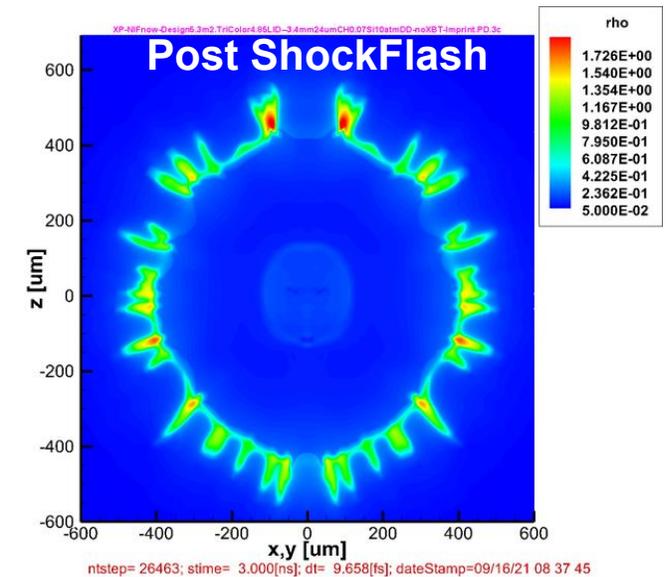
National Ignition Facility (NIF) Polar-Direct-Drive Exploding-Pusher Experiments — Improving Performance via Imprint Mitigation



Reference case



Picketed pulse reduces imprint



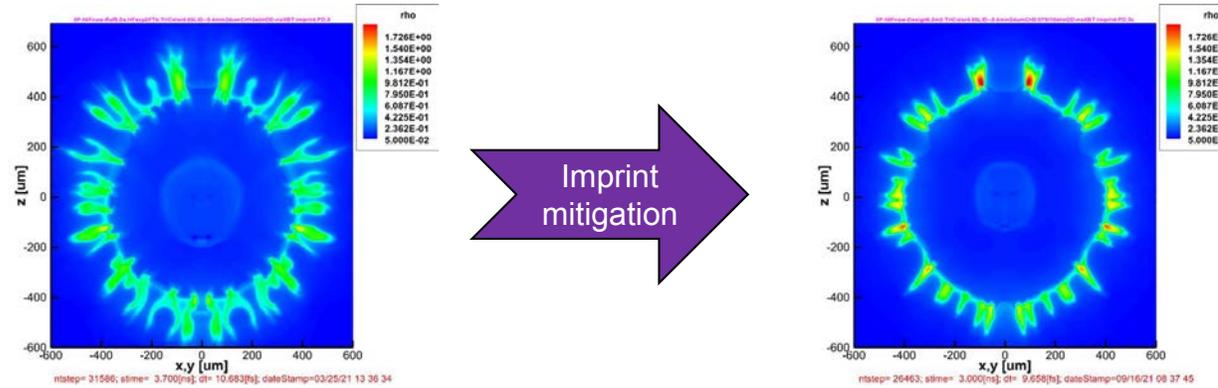
Picketed pulse plus Si-doped CH-shell
~equivalent no-Imprint yield

J. A. Marozas
University of Rochester
Laboratory for Laser Energetics

63rd Annual Meeting of the American Physical Society
Division of Plasma Physics
8-12 November 2021

Imprint mitigation schemes alleviate Exploding Pusher (XP) susceptibility during the final compression phase

- XPs require imprint mitigation to realize the benefits of shape control
 - Current NIF allows pulse-shape (pickets) and target solutions



- Target solutions offer ample imprint mitigation for XPs
 - Si-doped CH-shell*; split layer [Q2Q3FY22]
 - CH-foams; low density overcoat [future dev.]

* Suxing *et al.*, LLE (202?)

Collaborators

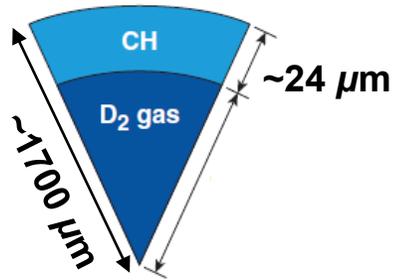


**P. W. McKenty, T. J. B. Collins, M. J. Rosenberg,
H. G. Rinderknecht, S. P. Regan and E. M. Campbell**
University of Rochester, Laboratory for Laser Energetics

**C. B. Yeamans, B. Blue, L. Divol,
G. E. Kemp, and H. D. Whitley**
Lawrence Livermore National Laboratory

FY21 XP shots assessed symmetry control and energy coupling through beam pointing and pulse shaping

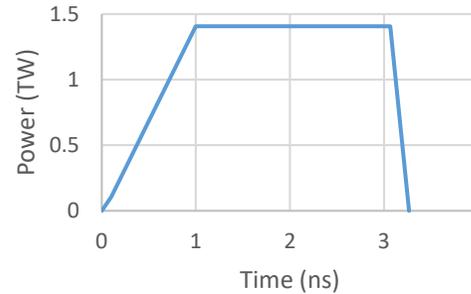
Exploding pusher
 ~700 kJ
 ~280 TW



Old-pointing & pulse Reference

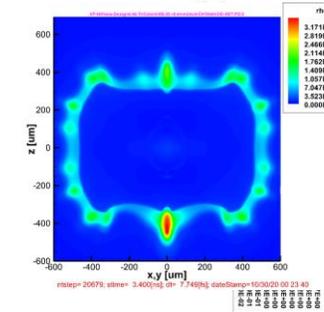
N201220-002

Per-Beam Laser Power

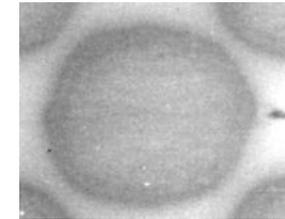


End of Acceleration

DRACO pre-shot
 3.4 ns



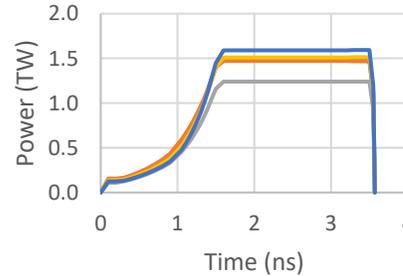
t~3.4 ns



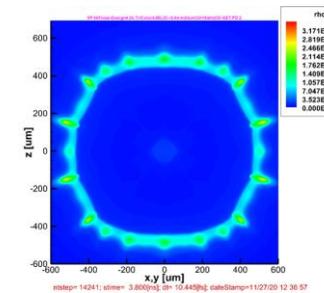
DRACO (new) pointing & pulse
 • Designed for improved symmetry and yield

N201220-001

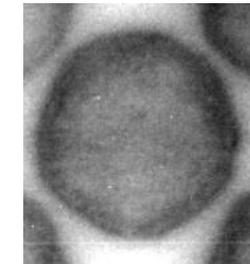
Per-Beam Laser Power



DRACO pre-shot
 3.8 ns



t~3.7 ns



Draco preShots predict expected shape

Self-emission images {equatorial view} agree with postshot simulations, achieving the expected control of the shape, however, sans the expected yield improvement

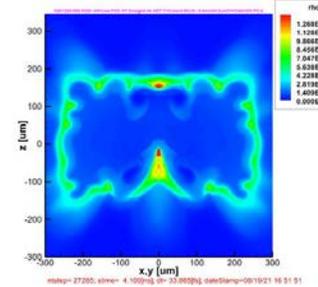
N201220-002

lowRes ~BT

t~4.2 ns



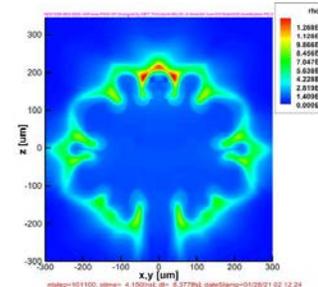
DRACO no-Imprint
postShot at Bangtime



Reference config., pointing & pulse
 $Y_{DD} = 4.8e13$ {exp.}

N201220-003

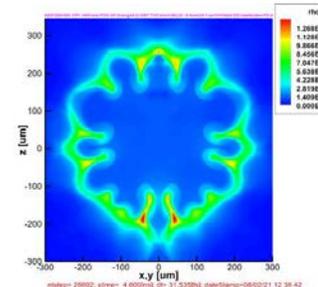
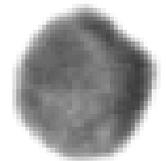
t~4.25 ns



Improved pointing and ref. pulse
 $Y_{DD} = 4.5e13$ {exp.}

N201220-001

t~4.7 ns



Improved pointing and pulse-shapes
 $Y_{DD} = 4.4e13$ {exp.}

Experimental pulse shapes lead to additional distortion

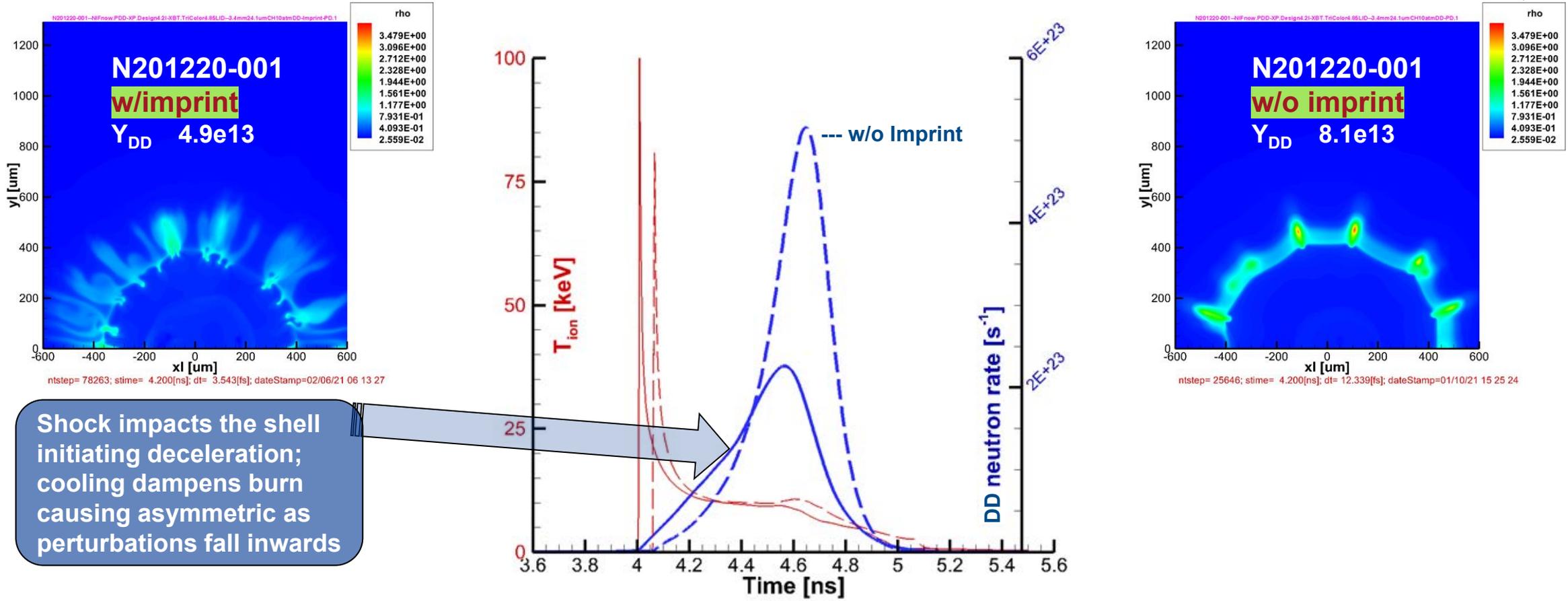
XP PostShots indicate that yield fell short of expectations due to imprint

Parameter	N201220-002	N201220-003	N201220-001
Pointing & PulseShape	Old&ramp2FT	New&ramp2FT	New&exp2FT
Exp. DD-n	4.8e13	4.5e13	4.4e13
PreShot (scaled), DD-n Draco (XBT, lowMode)	5.36e13	8.80e13	9.72e13
PostShot, DD-n Draco (XBT, lowMode)	5.97e13 {>preshot}	7.51e13 {~15% drop}	8.12e13 {~20% drop}
PostShot, DD-n Draco (XBT, Imprint)	4.58e13 {preliminary}		4.88e13 {preliminary}

XPs require imprint mitigation to realize the benefits of shape control

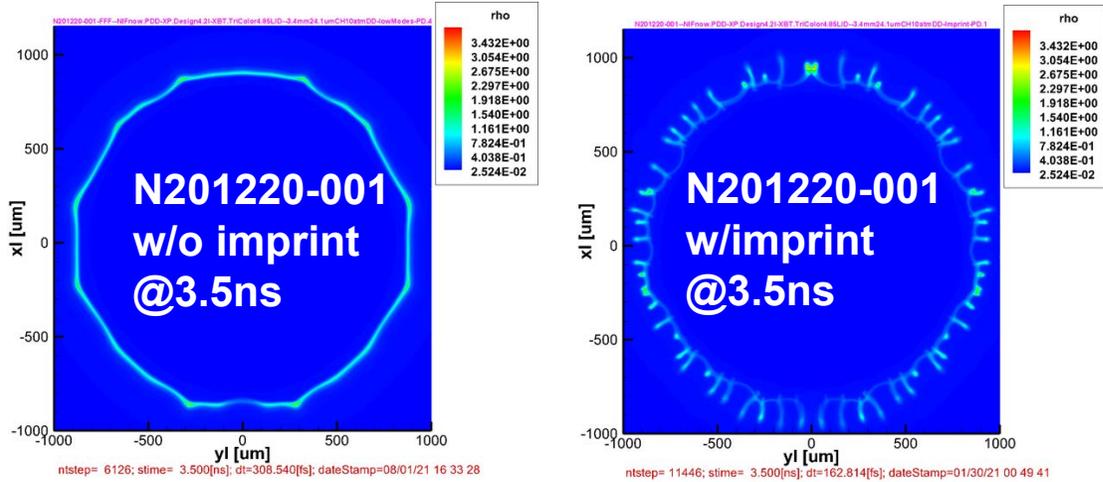
- “Improved” designs inadvertently had stronger imprint (beam overlap and defocusing)
 - during the design phase, imprint was not expected to effect XP

Recent XP shots have been showing evidence of imprint; reduced yield with dampened and wider, asymmetric burn curves (also observed in DT experiments)



- Postshot simulations illustrate the detrimental effect of imprint on yield
 - Decreased burn rate, earlier peak, asymmetric shape and wider burn-width

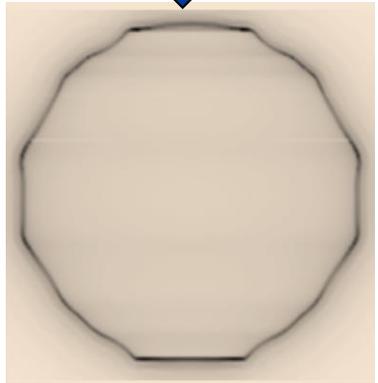
Self-emission images suggest signs of imprint when compared to simulations



Incentive for imprint mitigation shots

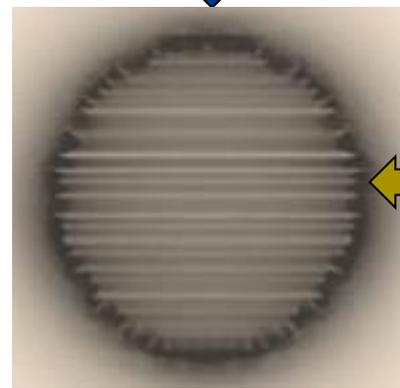
- July 2021: picketed pulse-shapes
- Q2Q3FY22: Si-doped CH-shells

Synthetic \downarrow x-ray image



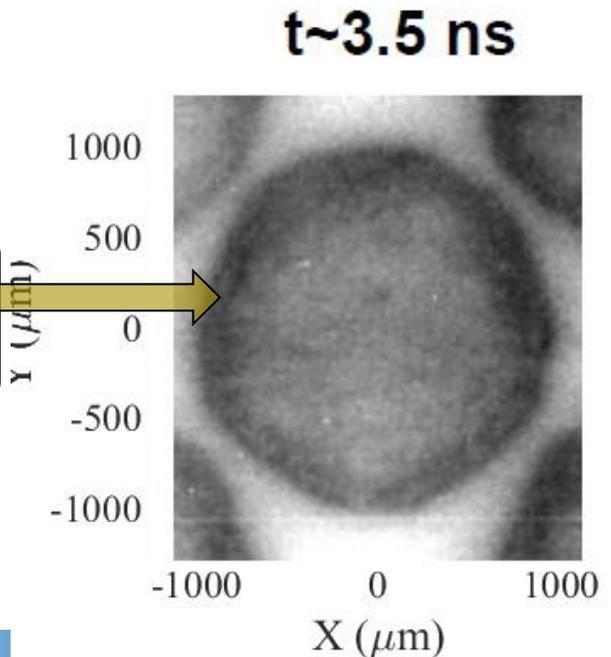
Spect3D @3.5ns

Synthetic \downarrow x-ray image



Spect3D @3.5ns

Thickened shell from imprint corroborates experimental images



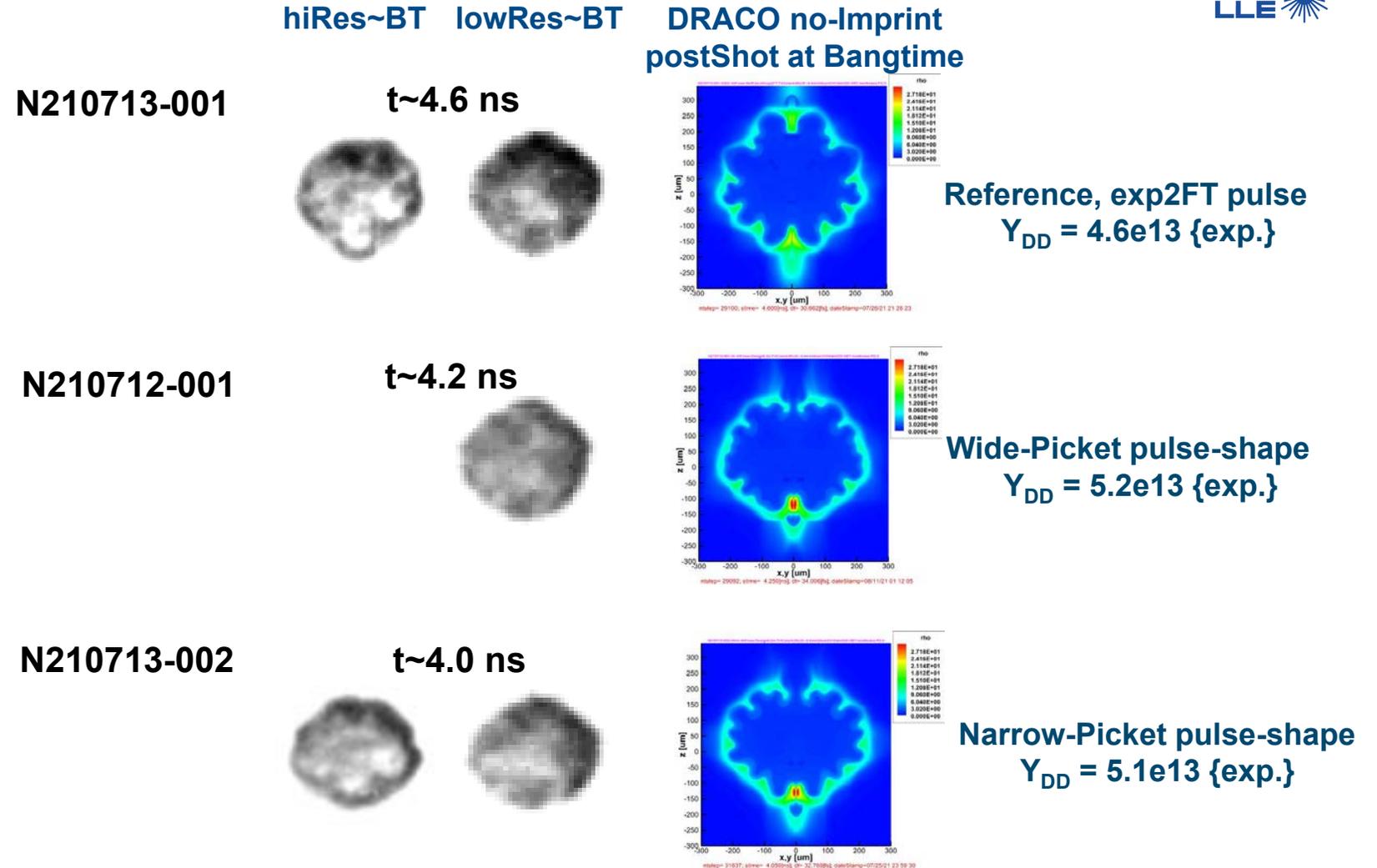
Note: Horizontal structure won't exist true 3D

The July 2021 XP shots attempted to demonstrate imprint mitigation by comparing a reference shot to two different picketed pulse shapes

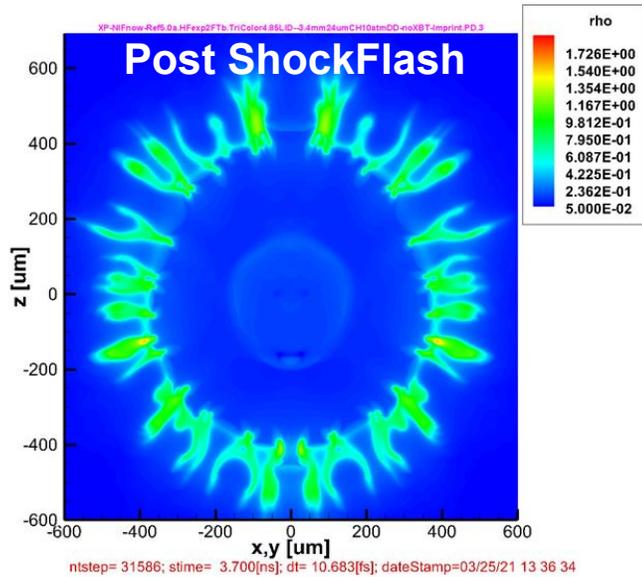
Gated x-ray self-emission images {equatorial view} show North/South asymmetry due to power imbalance

Initial imprint mitigation experiments only achieved ~10% increase due to inadequate picket formation.

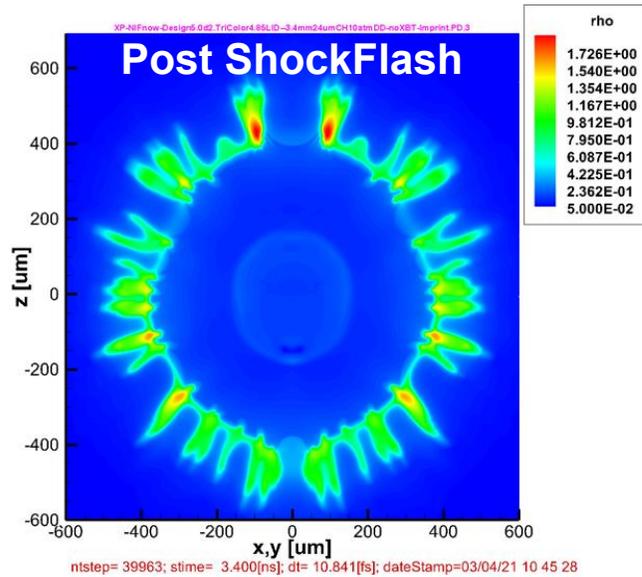
- Retuned pickets for reliable shape in November; hoping to achieve ~40% increase



Imprint mitigation schemes are predicted to improve XP performance, thereby realizing the benefit from shape control

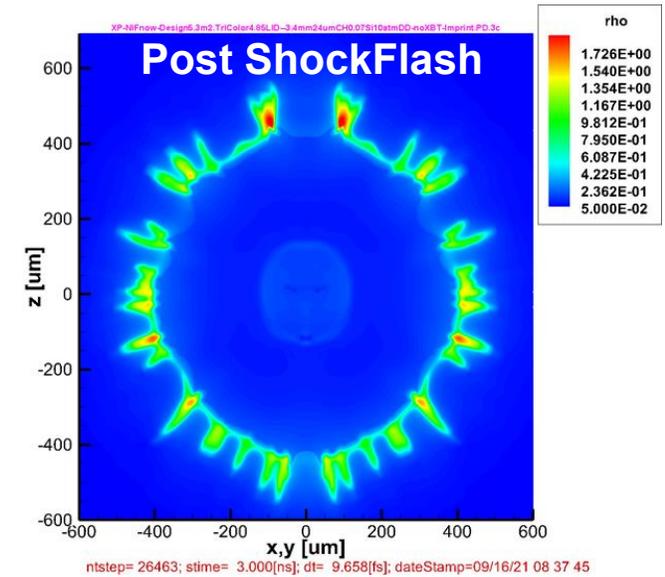


**Reference case
Without mitigation**



**Picketed pulse predicts
yield increase ~50%**

Initial experiments July 2021



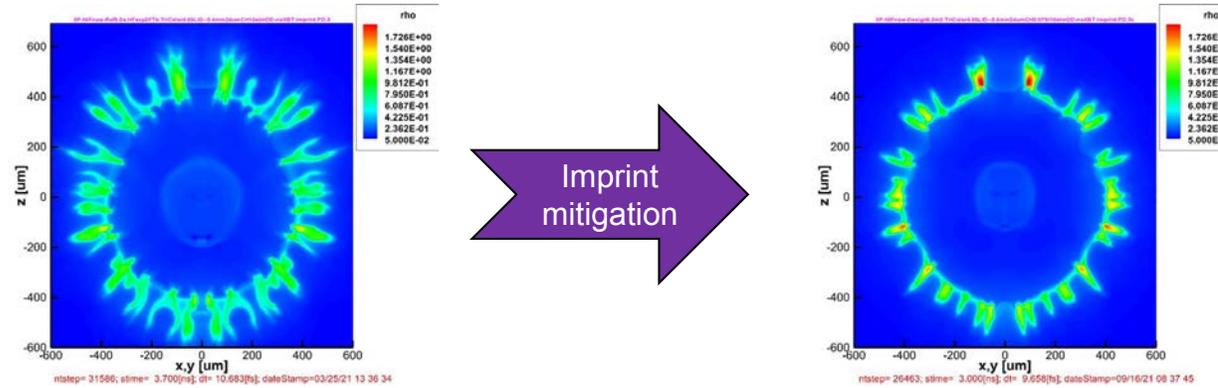
**Picketed pulse plus Si-doped CH-shell
~equivalent no-Imprint yield**

Experiments in Q2Q3FY22 scheduled to shoot Si doped CH shells

- Future experiments planned to use CH foam overcoats

Imprint mitigation schemes alleviate Exploding Pusher (XP) susceptibility during the final compression phase

- XPs require imprint mitigation to realize the benefits of shape control
 - Current NIF allows pulse-shape (pickets) and target solutions



- Target solutions offer ample imprint mitigation for XPs
 - Si-doped CH-shell*; split layer [Q2Q3FY22]
 - CH-foams; low density overcoat [future dev.]

* Suxing *et al.*, LLE (202?)

Backup --- NIF XP Shots data

XP PostShots indicate that yield fell short of expectations due to imprint

Parameter	N201220-002	N201220-003	N201220-001
Pointing & PulseShape	Old&ramp2FT	New&ramp2FT	New&exp2FT
Exp. DD-n	4.8e13	4.5e13	4.4e13
PreShot (scaled), DD-n Draco (XBT, lowMode)	5.36e13	8.80e13	9.72e13
PostShot, DD-n Draco (XBT, lowMode)	5.97e13 {>preshot}	7.51e13 {~15% drop}	8.12e13 {~25% drop}
PostShot, DD-n Draco (XBT, Imprint)	4.58e13 {preliminary}		4.88e13 {preliminary}

XPs require imprint mitigation to realize the benefits of shape control

- “Improved” designs inadvertently had stronger imprint (beam overlap and defocusing)
 - during the design phase, imprint was not expected to effect XP

MJDD NIFnow.PDD XP shots --- Q1FY21 December 2020



- Bangtimes have mixed agreement

Parameter	N201220-002	N201220-003	N201220-001
Pointing & PulseShape	Old&ramp2FT	New&ramp2FT	New&exp2FT
		4.2	4.65
		4.46	4.41
PreShot, BT [ns] Draco (XBT, lowMode)	3.83	3.87	4.32
PostShot, BT [ns] Draco (XBT, lowMode)			
			4.58 cf. Spider

pTOF data inconsistent with expected 450-500ps pre- & post-shot delta

“Early” due to delayed exp. rise

MJDD NIFnow.PDD XP shots --- Q1FY21 December 2020



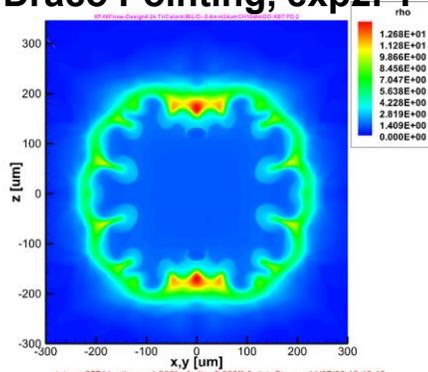
- Tion have mixed agreement depending on sim-config.

Parameter	N201220-002	N201220-003	N201220-001
Pointing & PulseShape	Old&ramp2FT	New&ramp2FT	New&exp2FT
Tion {exp.} [keV]	7.4	7.1	7.3
PreShot, Tion [keV] Draco (XBT, lowMode)	6.87	7.38	7.60
PostShot, Tion [keV] Draco (XBT, lowMode)	6.42	6.92	7.10
PostShot, Tion [keV] Draco (XBT, Imprint)	6.50		6.79

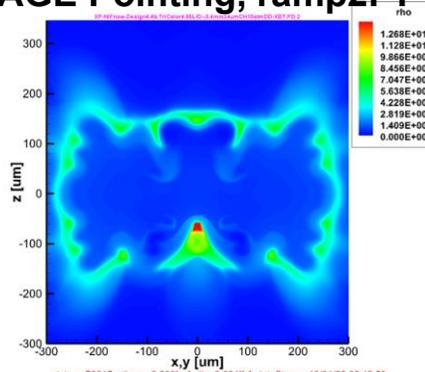
The pulse shapes and imprint alter the shapes compared to designs at bangtime; only measured target and pulse-shapes changed; N201220

PreShot Designs

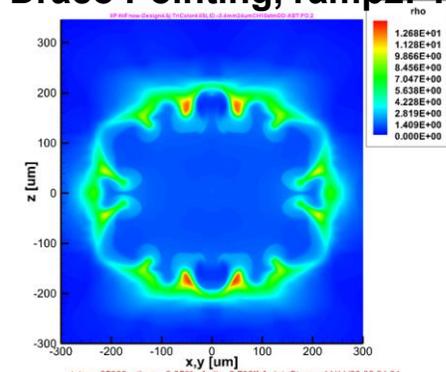
N201220-001 [FFF]
Draco Pointing, exp2FT



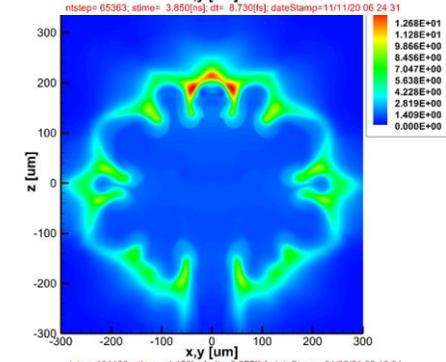
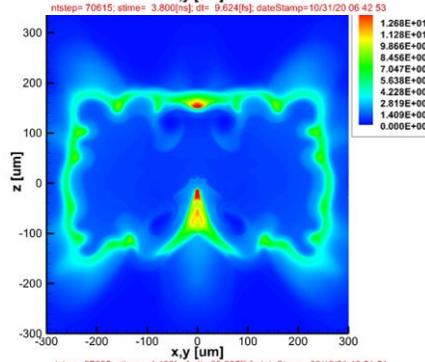
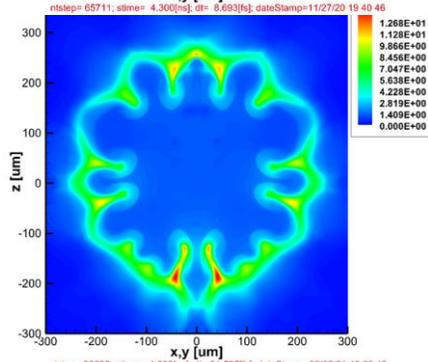
N201220-002 [DDD]
SAGE Pointing, ramp2FT



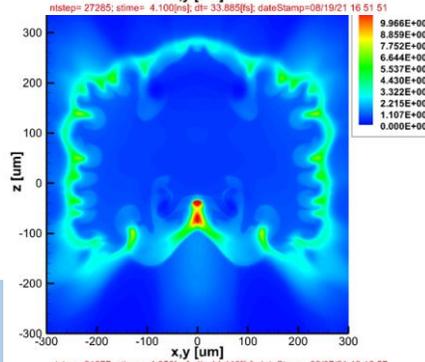
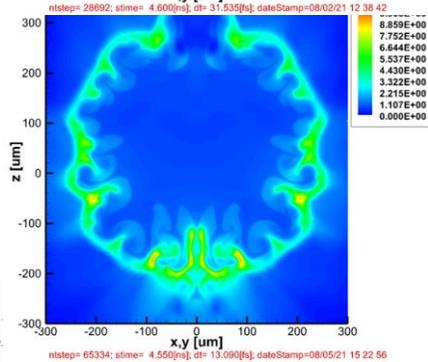
N201220-003 [EEE]
Draco Pointing, ramp2FT



PostShots w/
Experimental
PulseShapes



PostShots w/
Experimental
PulseShapes &
Imprint



Shot	$Y_{DT} (10^{13})$		T_i (keV)		t_{bang} (ns)	
	Exp.	Sim.	Exp.	Sim.	Exp.	Sim.
N201220-001	4.4	4.88	7.3	6.79	4.65*	4.58
N201220-002	4.8	4.58	7.4	6.50	4.21	4.08
N201220-003	4.5		7.1		4.2*	

MJDD NIFnow.PDD XP shots --- Q4FY21 July 2021

Postshots show yield fell short of expectations due to inadequate picket shapes



- Nominal XP target: 24um CH shell, 3.4mm diameter, 10atm DD-fill
- Nominal laser specs: ~290TW total peak power, ~710kJ

Parameter	N210713-001	N210712-001	N210713-002
Config	exp2FT	Wide picket	Narrow picket
Exp. DD-n	4.64e13	5.16e13	5.13e13
PreShot, DD-n Draco (XBT, lowMode)	9.77e13	9.47e13	10.1e13
PreShot, DD-n Draco (XBT, Imprint)	5.76e13		8.43e13
PostShot, DD-n Draco (XBT, lowMode)	10.7e13	10.2e13	10.9e13
PostShot, DD-n Draco (XBT, Imprint)			

Imprint simulations are underway including the effects to the imprint spectrum, both SSD modulators (important for SSDMPE) and non-zero m-modes.

- Picketed shapes were compromised and large North/South imbalances.

MJDD NIFnow.PDD XP shots --- Q4FY21 July 2021

- Bangtimes have mixed agreement

Parameter	N210713-001	N210712-001	N210713-002
Config	exp2FT	Wide picket	Narrow picket
X-ray BT (SPIDER) [ns]	~4.33	~4.05	~3.8
DD-n BT (pTOF) [ns]	4.72		4.19
PreShot, BT [ns] Draco (XBT, lowMode)	4.42	4.13	
PreShot, DD-n Draco (XBT, Imprint)	4.36		3.88
PostShot, BT [ns] Draco (XBT, lowMode)	4.61	4.27	4.04
PostShot, BT [ns] Draco (XBT, Imprint)			

MJDD NIFnow.PDD XP shots --- Q4FY21 July 2021

- Tion have mixed agreement depending on sim-config.

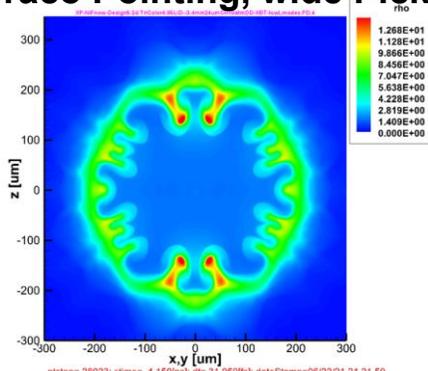
Parameter	N210713-001	N210712-001	N210713-002
Config	exp2FT	Wide picket	Narrow picket
Tion {exp.} [keV]	7.47	7.56	7.90
PreShot, Tion [keV] Draco (XBT, lowMode)	7.07	6.84	6.95
PreShot, DD-n Draco (XBT, Imprint)	6.79		6.70
PostShot, Tion [keV] Draco (XBT, lowMode)	6.91	6.75	6.92
PostShot, Tion [keV] Draco (XBT, Imprint)			

The pulse shapes and imprint alter the shapes compared to designs; the only changed is measured target and pulse-shapes; N210712-13



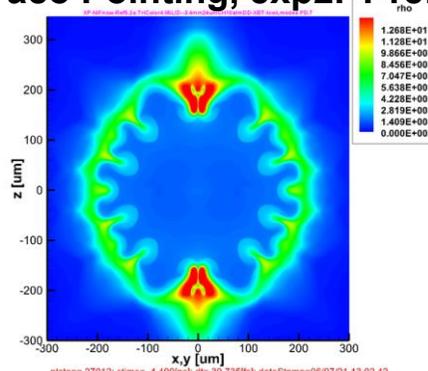
PreShot
Designs

N210712-001 [III]
Draco Pointing, wide Picket



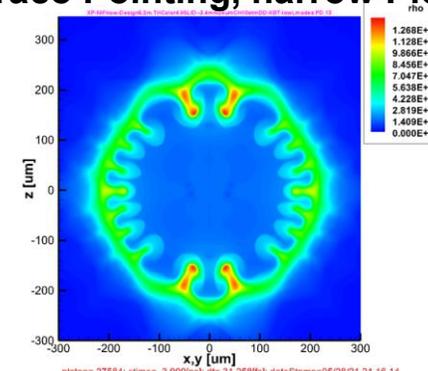
ntstep= 28023; stime= 4.150[ns]; dt= 31.959[fs]; dateStamp=06/22/21 21 21 50

N210713-001 [GGG]
Draco Pointing, exp2FT ref.



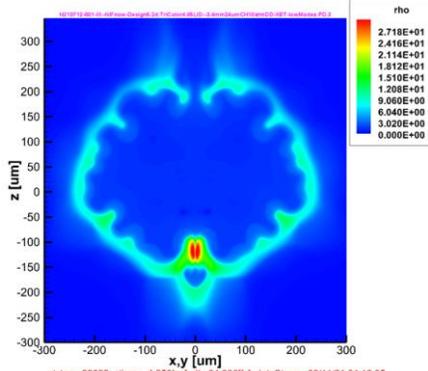
ntstep= 27012; stime= 4.400[ns]; dt= 30.735[fs]; dateStamp=06/07/21 13 02 42

N210713-002 [HHH]
Draco Pointing, narrow Picket

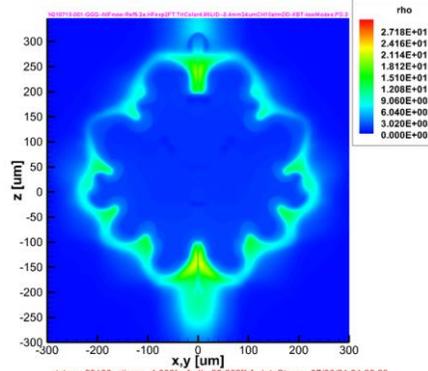


ntstep= 27584; stime= 3.900[ns]; dt= 31.258[fs]; dateStamp=05/26/21 21 16 14

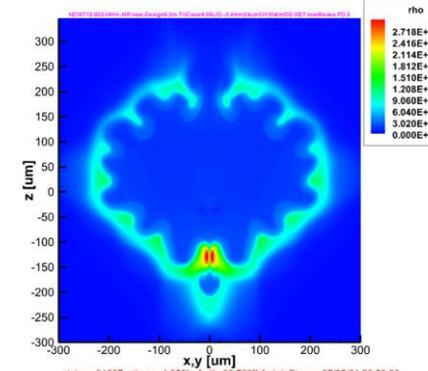
PostShots w/
Experimental
PulseShapes



ntstep= 29092; stime= 4.250[ns]; dt= 34.006[fs]; dateStamp=08/11/21 01 12 05



ntstep= 29100; stime= 4.600[ns]; dt= 30.662[fs]; dateStamp=07/26/21 21 28 23



ntstep= 31637; stime= 4.050[ns]; dt= 32.769[fs]; dateStamp=07/25/21 23 59 30

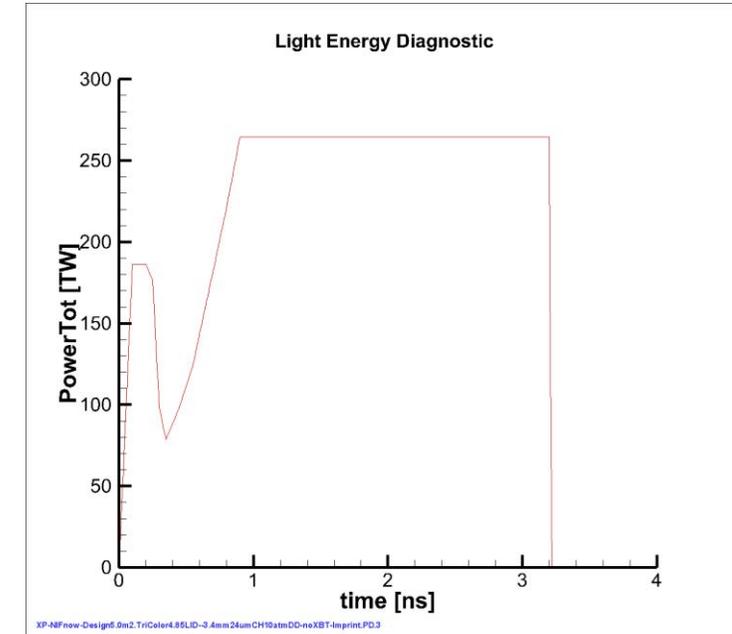
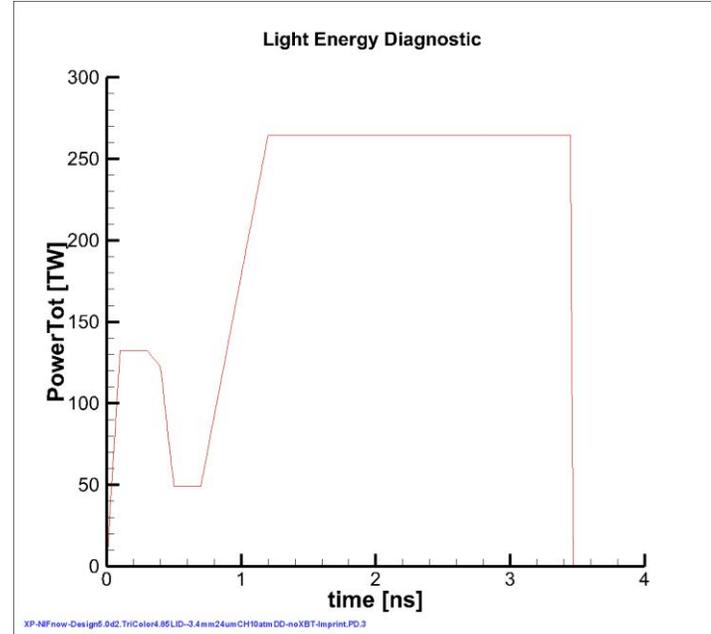
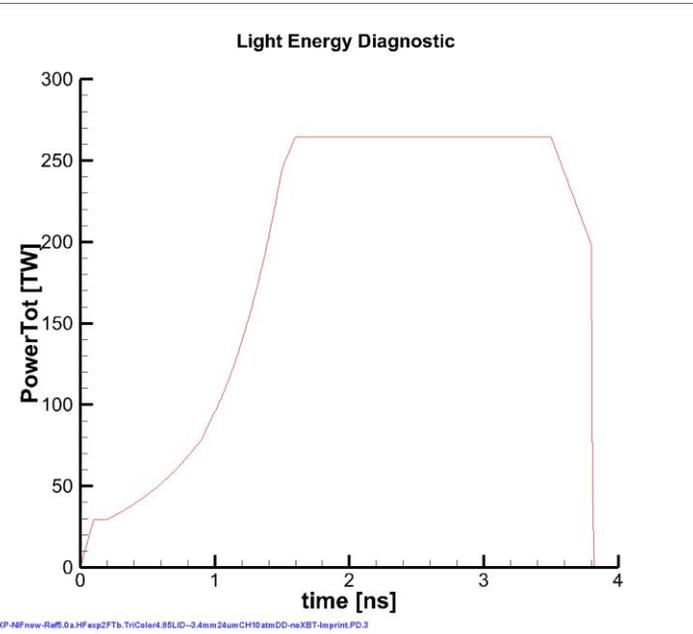
Backup --- NIF PulseShaping

MJDD NIFnow.PDD XP shots --- Q4FY21 July 2021

The three proposed pulse-shapes



- Reference, Design5.0d2, Design5.0m2

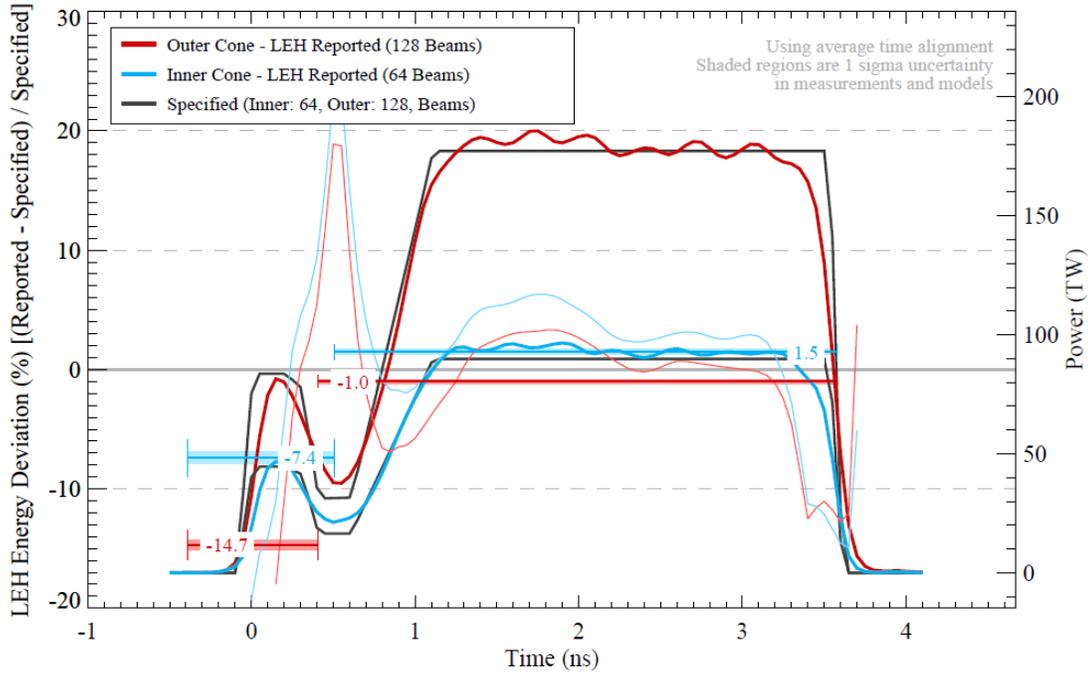


- Picket goals to achieve imprint mitigation:
 - ~50-60kJ; exact shape not so important
 - Narrow as possible; an impulse function is ideal but not necessary
 - Shortest dwell, and lowest power (return to zero ideal); need to rise to FT ~1ns

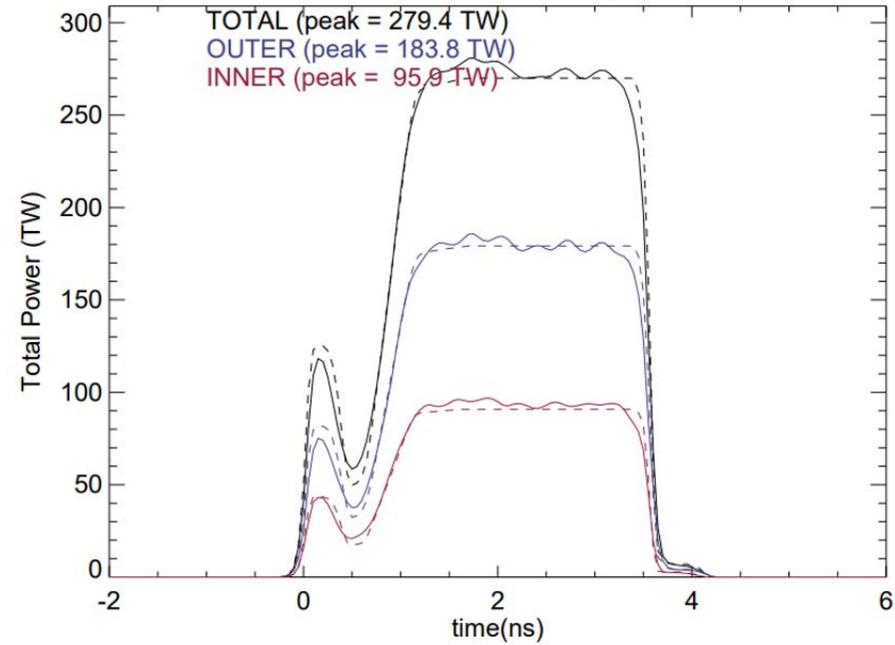
Pulse shapes from July 2021 XP shots

- Wide picket not well matched, plus strong imbalances

Power Accuracy: LEH Reported vs. Specified
N210712-001-999, I_MJDD_PDD_DDExpPush_S09a

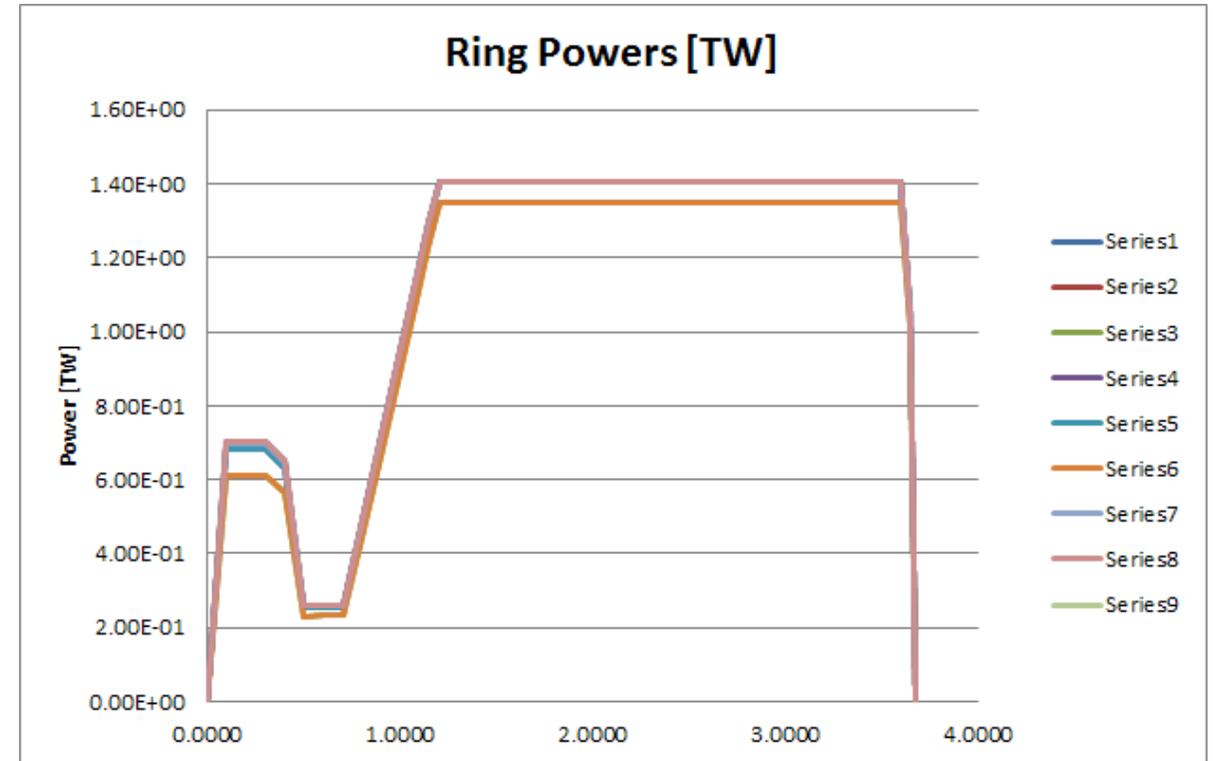
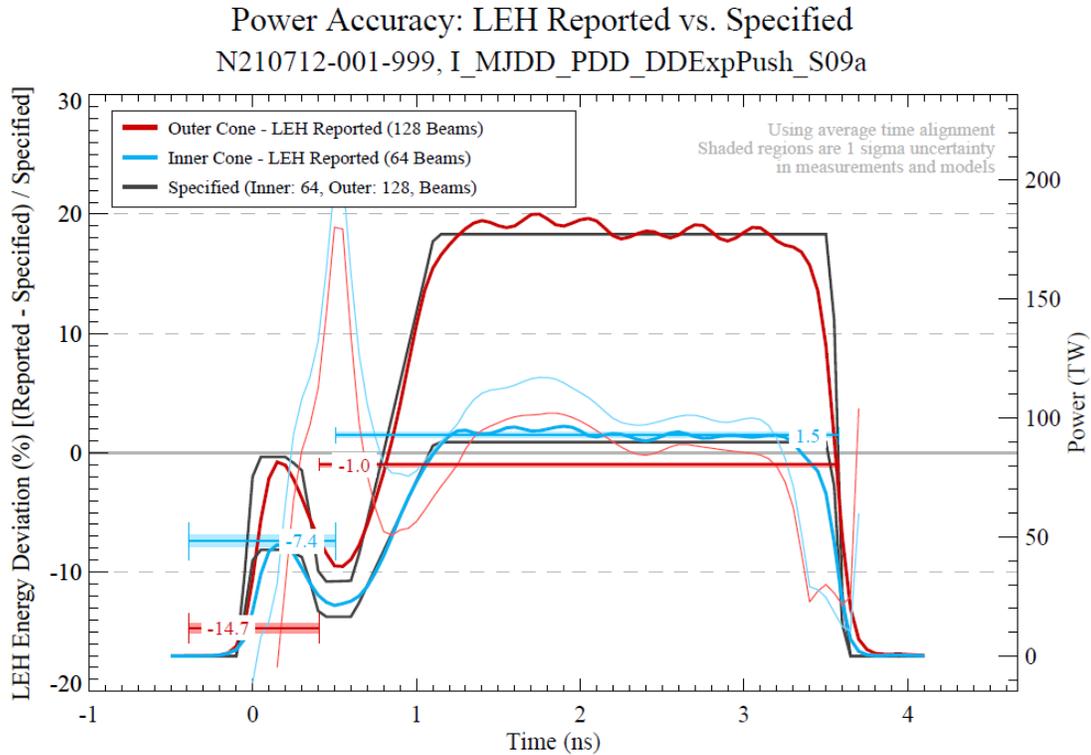


N210712-001-999
EXP : I_MJDD_PDD_DDExpPush_S09a



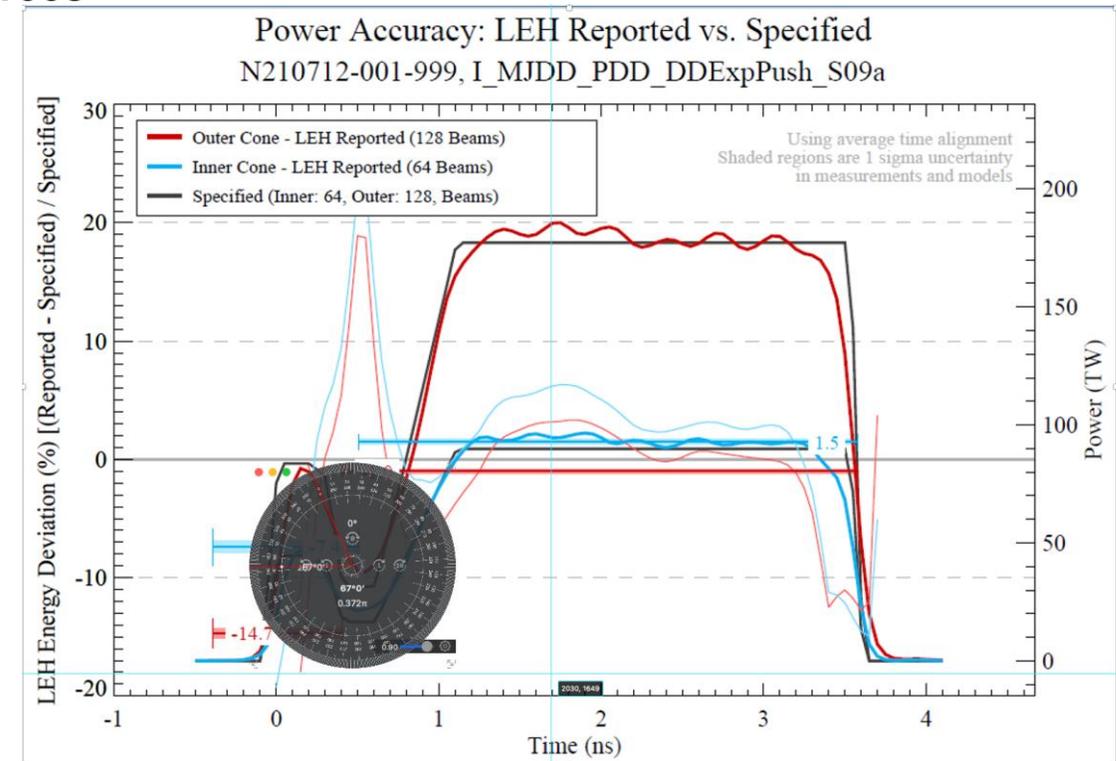
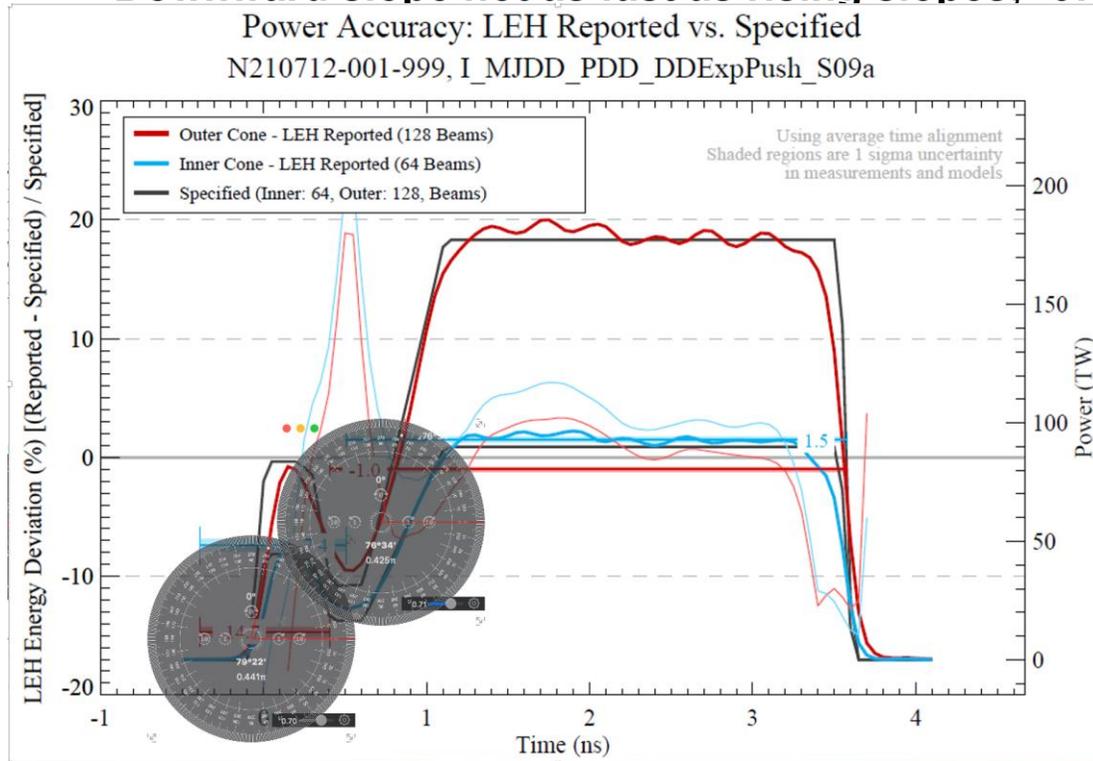
Pulse shapes from July 2021 XP shots

- **Black trace shown as “specified” is not exactly as-specified; 3 added ‘corner’ features; where did those come from?**



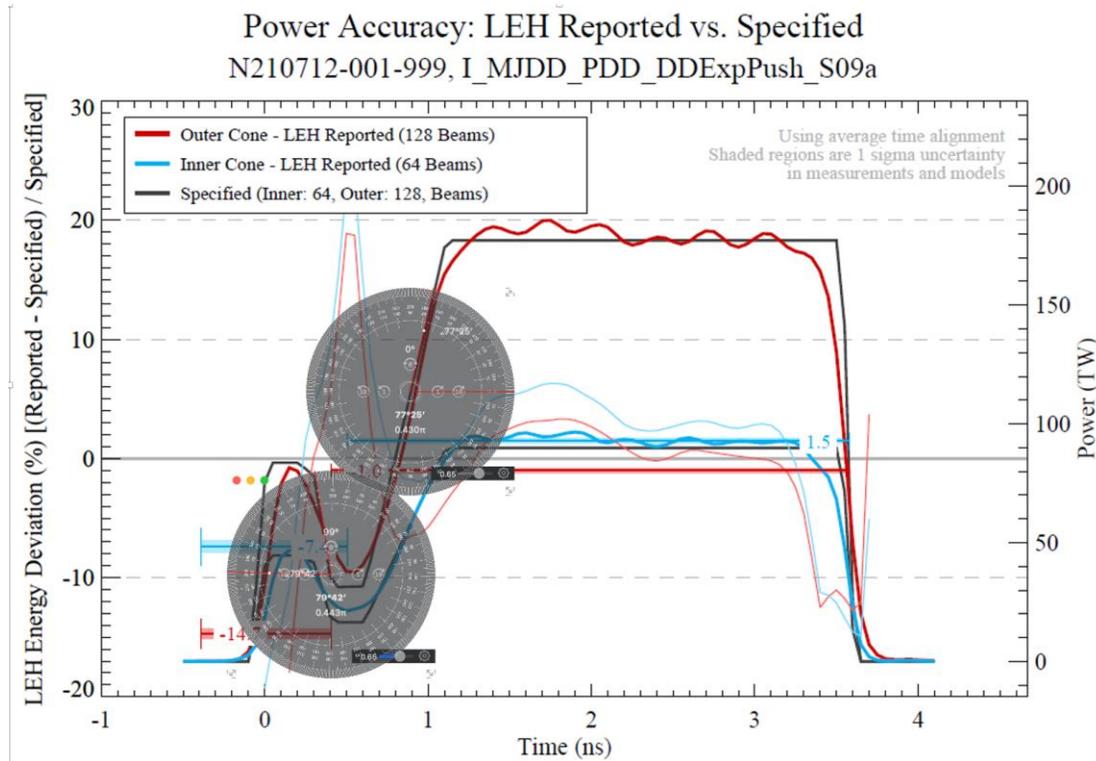
Pulse shapes from July 2021 XP shots

- Slopes....
- Takes ~200ps to rise to 50% of peak power; half the rate in the guide
- Initial slope not much faster than “internal” feature rise2main; 76- vs. 79-degrees
- Downward slope not as fast as rising slopes; -67-degrees



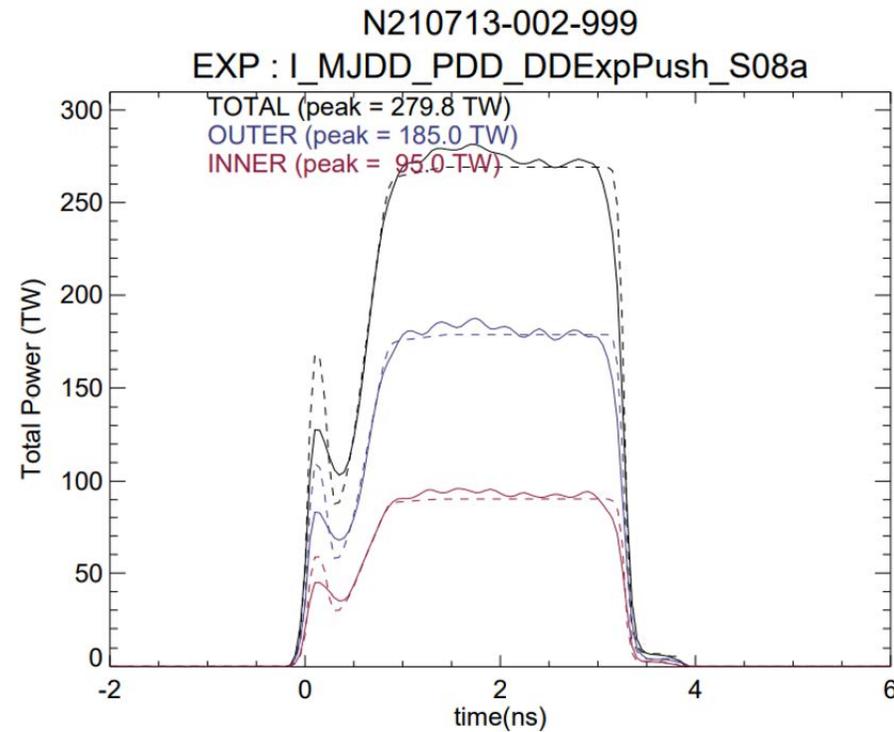
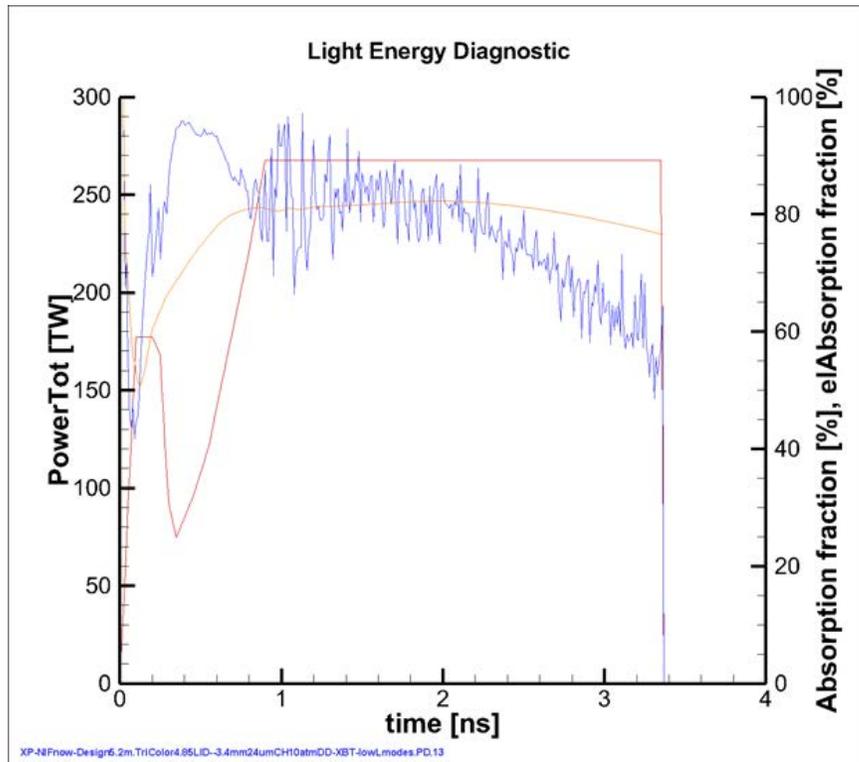
Pulse shapes from July 2021 XP shots

- Slopes of the falling2dwell and rise2main are nearly the same; 77- vs. 79-degrees. Shouldn't the NIF pulse have similar slopes? Where in the guide does it say why?



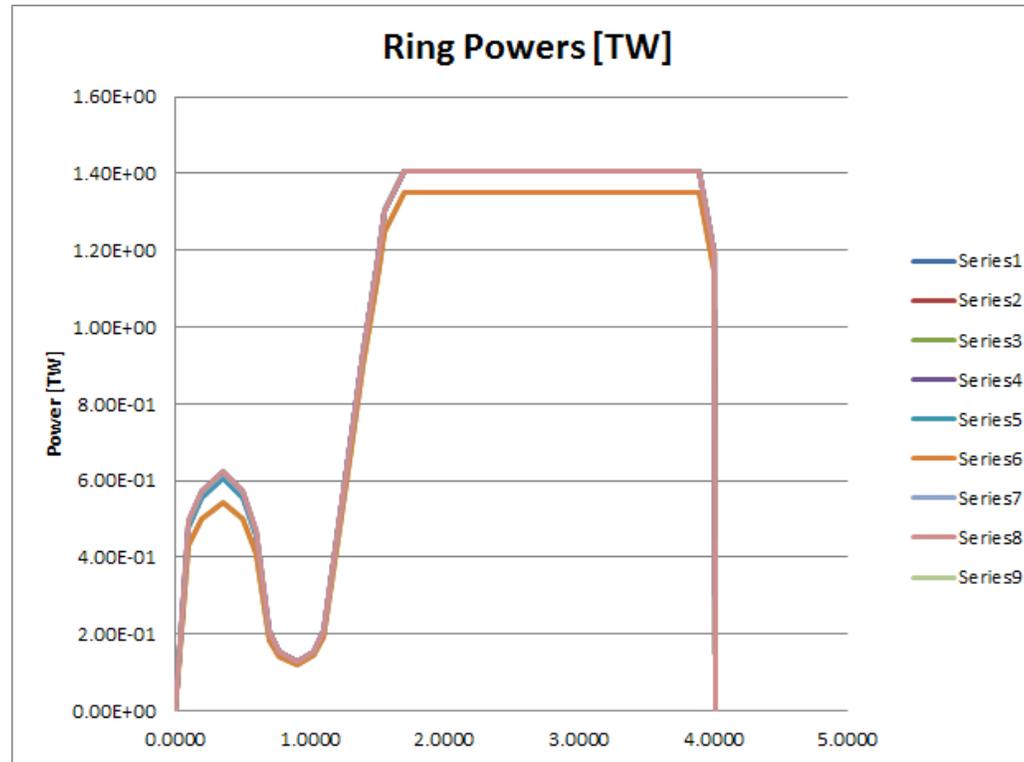
Pulse shapes from July 2021 XP shots

- **Narrow picket not well matched --- more like a step than a picket, plus strong imbalances**

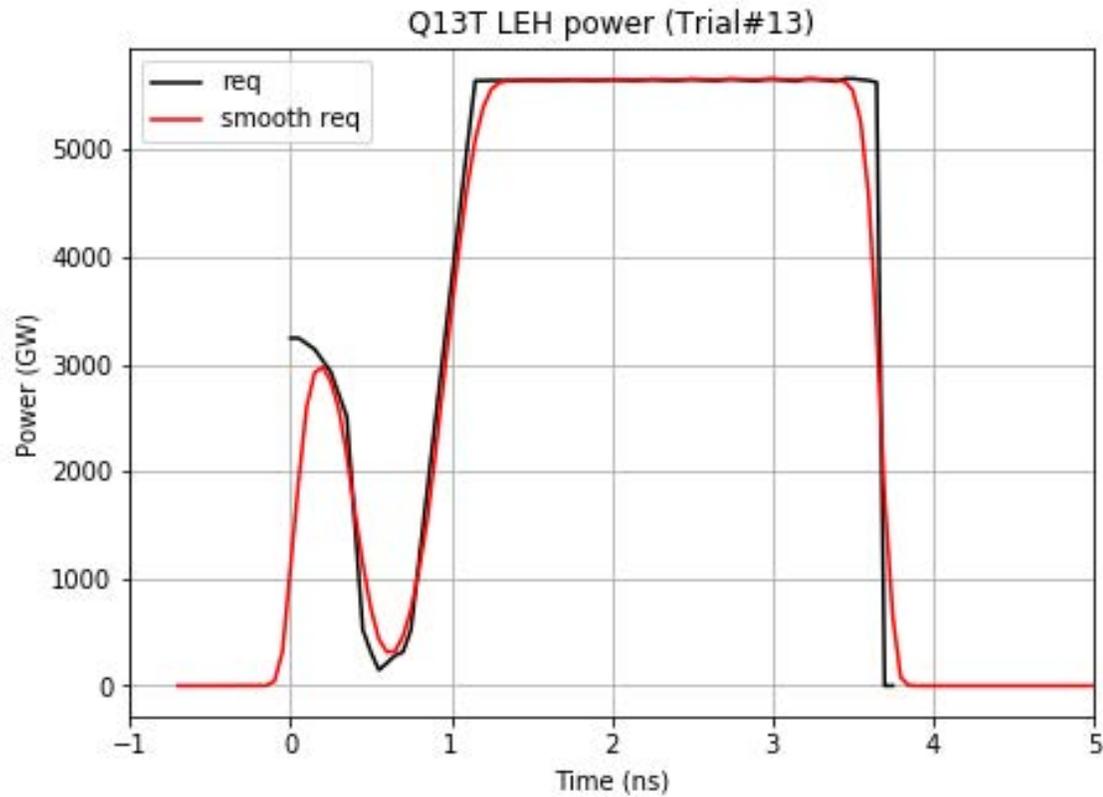


Attempted this pulse, but it failed in simulations w.r.t. imprint mitigation – dwell too long and imprinted too long in picket

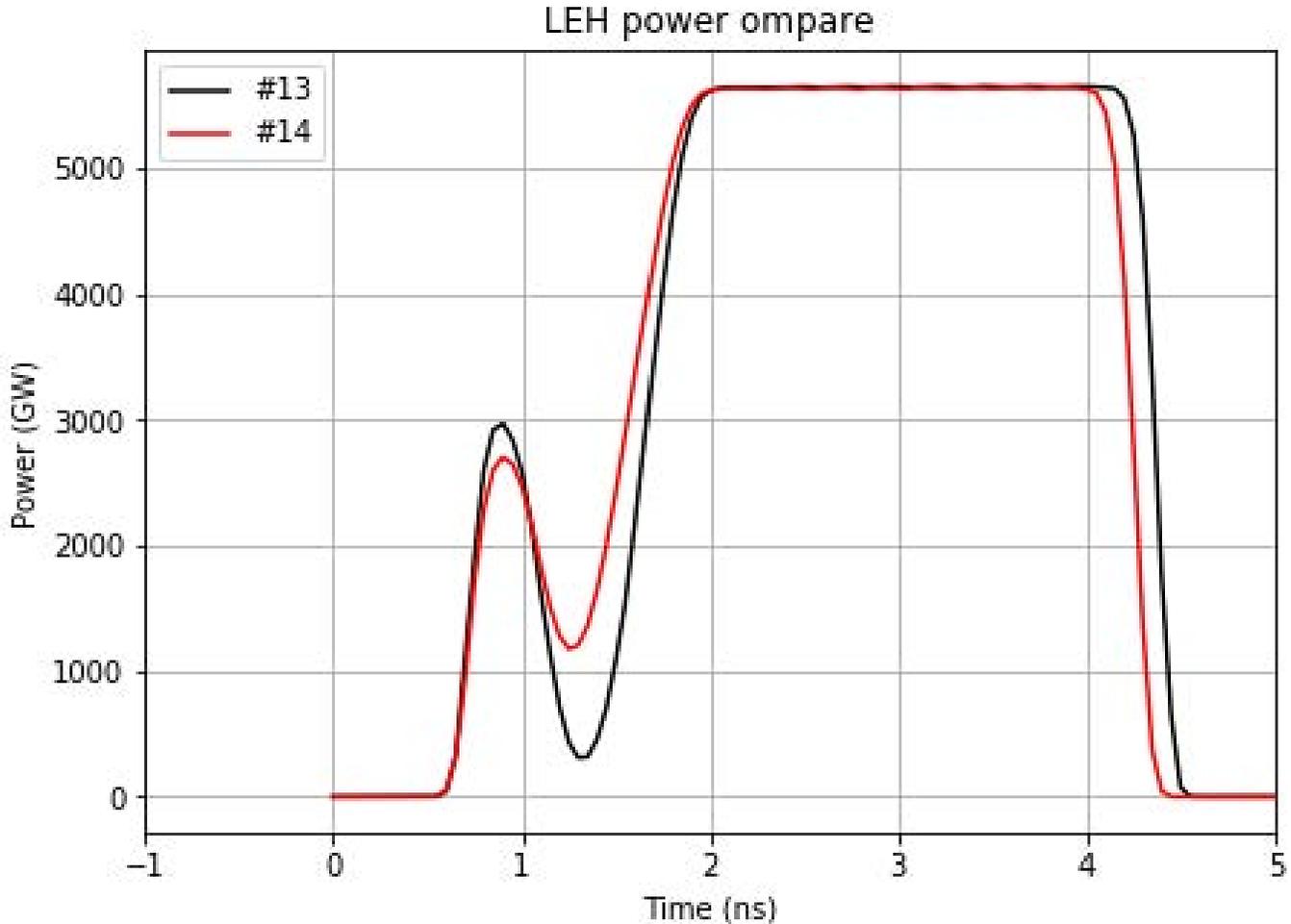
- Increase picket & dwell duration by 200ps, rounding the picket and dwell



#13 & #14 pulse shapes

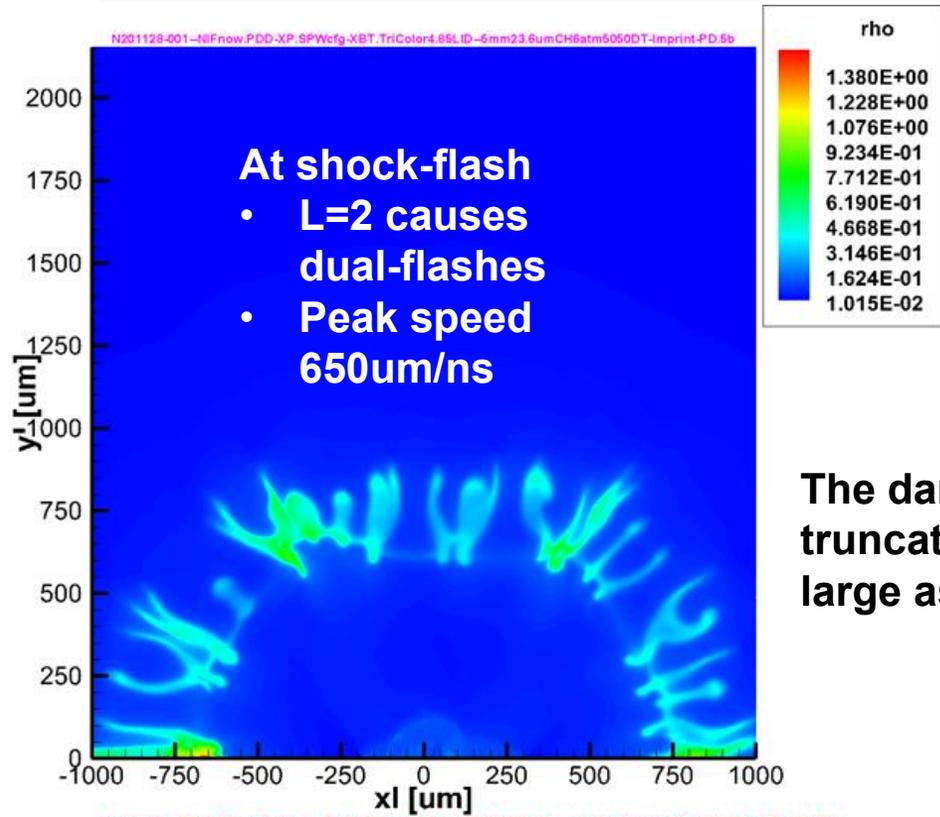


#13 & #14 comparison

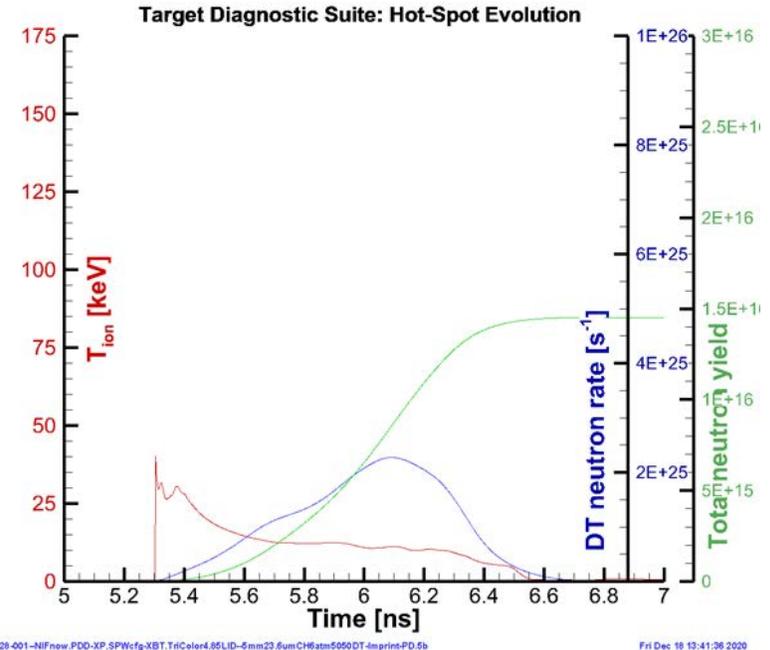


Backup

Including some imprint and CBET brings the simulation closer to experiment and hints at potential elimination of compression yield in experiment



The dampened (not full truncation) burn-curve shows large asymmetry



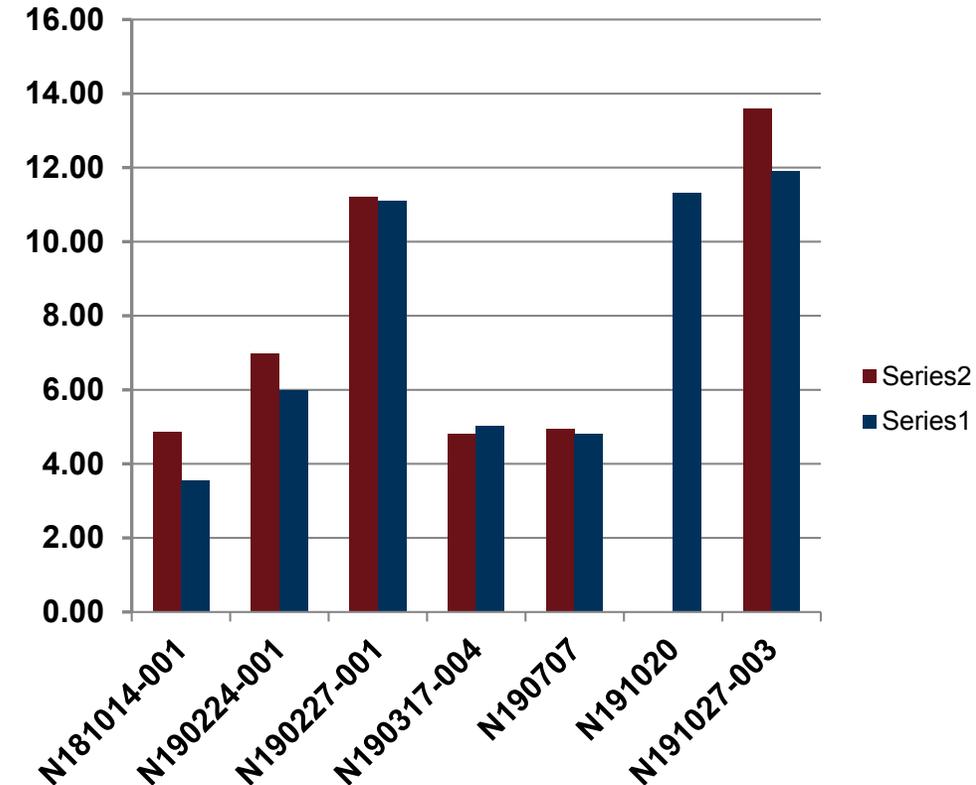
- Imprint barely effects the absorption
- Slows the average shell speed but advances the the distorted inner shock-front (see next slide)
- Similar burn-curve base (cf. no imprint), but with dampened compression phase leads to wider burn-width

N201128 001 Exp. vs. Sim	Y_{DT} (10^{15})	T_i (keV)	t_{bang} (ns)	t_{burn} (ps)
Exp.	6.35	8.17	6.29	505
noCBET ref.	51.3	12.4	5.81	250
CBET	28.9	9.39	6.16	360
Imprint&CBET	14.5	9.22	6.09	634

Backup --- XP and Draco

The LLE 2-D radiation-hydrodynamics code *DRACO* has proven reliable in predicting performance metrics for XP targets

Shot	Y_{DT} (10^{15})		T_i (keV)		t_{bang} (ns)		t_{burn} (ps)	
	Exp.	Sim.	Exp.	Sim.	Exp.	Sim.	Exp.	Sim.
N181014	3.56	4.85	6.30	6.49	5.52	5.54	618	430
N190224	5.97	6.97	7.65	10.29	4.88	4.41	502	322
N190227	11.1	11.2	8.94	9.65	4.22	4.25	452	300
N190317	5.01	4.83	7.37	8.20	4.58	4.38	607	402
N190707	4.81	4.94	11.14	12.0	2.71	2.84	311	202



- NIF experiments will help further code validation/verification to provide a means to design targets, reaching >100-kJ DT fusion yields at low convergence (~10)
- The XP platform allows the study of some required NIF improvements to achieve high-convergence implosions such as CBET mitigation and far-field spot profile, as well as our modeling of heat conduction*

* H. D. Whitley *et al.*, Physics Archive, <https://arxiv.org/abs/2006.15635> (2020).

The XP targets on the NIF provide a viable direct-drive platform while under non-ideal illumination conditions



- The XP platform generates a high neutron flux*, albeit at low areal density (ρR), using simple targets
 - ideal “open” platform to study implosion physics modeling†, NIF improvements, and neutron-survivability device testing*
- Moderately robust to the NIF current laser configuration
 - Spot-shape nonuniformity; size & shape
 - Modest laser smoothing; intended for SBSS
 - Modest wavelength separation (CBET mitigation)

The LLE 2-D radiation-hydrodynamics code *DRACO* has proven reliable in predicting performance metrics for XP targets.

Saturation of the performance at $Y \sim 10^{16}$ is likely caused by limited implosion symmetry and imprint growth. The path forward includes mitigation strategies.

- Mitigation
 - Symmetry; Draco pointing and pulse-shape strategy
 - Imprint: picketed pulses and target solutions

Experiments by C. Yeamans (LLNL)

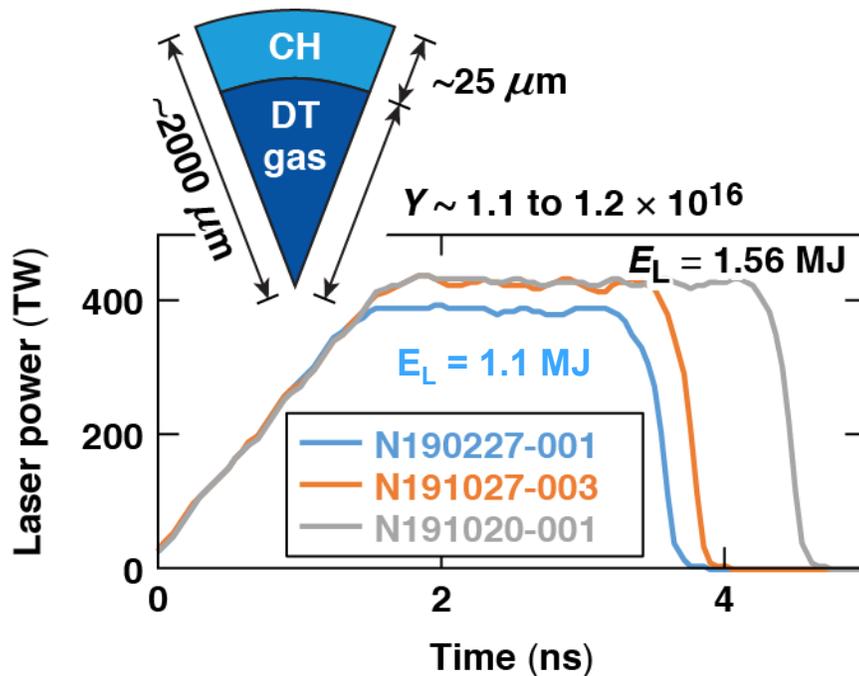
*C. B. Yeamans *et al.*, “High Yield Polar Direct Drive Fusion Neutron Sources at the National Ignition Facility,” submitted to Nuclear Fusion.

†J. A. Marozas *et al.*, Phys. Rev. Lett. **120**, 085001 (2018);

NIF: National Ignition Facility

CBET: Crossed-Beam Energy Transfer

XP are high v_{imp} PDD warm implosions on the NIF that couple up to $E_k = 100$ kJ and $E_{\text{hs}} = 30$ kJ, producing $Y_{DT} > 1 \times 10^{16}$



E29180a

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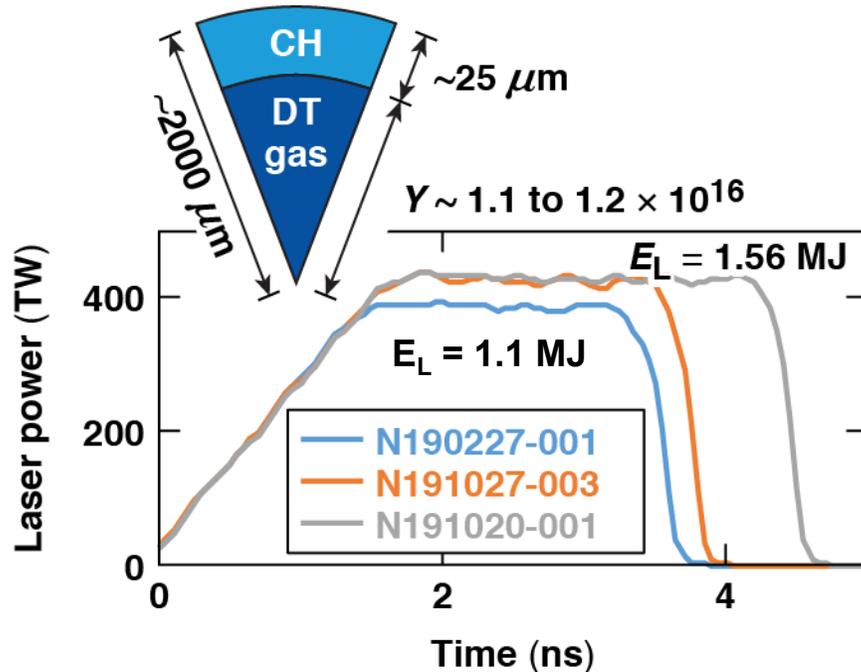
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E29180a

$$\eta_H \sim \frac{p_{\text{abl}} v_{\text{imp}}}{I}; \text{IFAR} \sim v_{\text{imp}}^2 / (p_{\text{abl}}^{2/5} \alpha^{3/5}); p_{\text{abl}} \sim I^{2/3} (A/Z)^{1/3}$$

- High α to improve stability while high v_{imp} increases hydroefficiency and decreases stability
- $v_{\text{imp}} > 800$ um/ns
- CR ~ 6 to 10
- Incident drive intensity $I = 8$ to 9×10^{14} W/cm²
- Hot-electron preheat does not affect performance (α is too high)

Saturation of the performance at $Y \sim 10^{16}$ is likely caused by limited implosion symmetry and imprint growth. The path forward includes mitigation strategies.

The LLE 2-D radiation-hydrodynamics code *DRACO* has proven reliable in predicting performance metrics for XP targets.

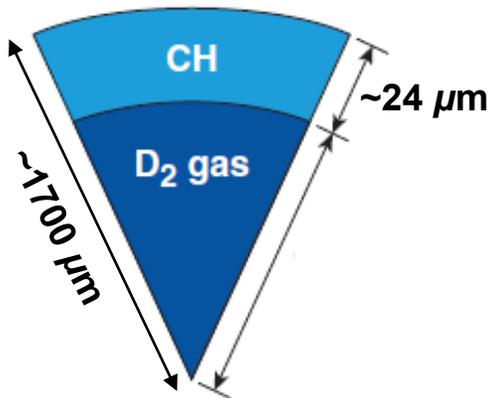
Experiments by C. Yeaman (LLNL)
CR: convergence ratio

Recent FY21 exploding pusher experiments assessed control of symmetry and energy coupling through laser pointing and pulse shaping

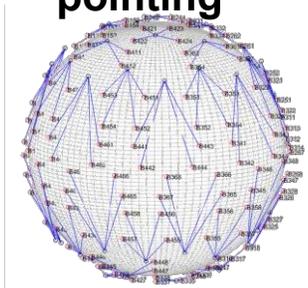


Exploding pusher

~700 kJ
~280 TW

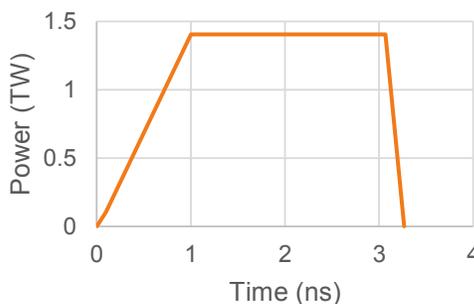


SAGE (old) pointing

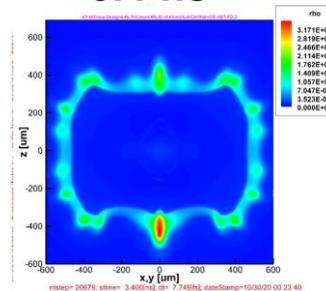


N201220-002

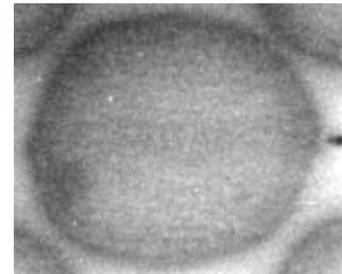
Per-Beam Laser Power



DRACO pre-shot
3.4 ns



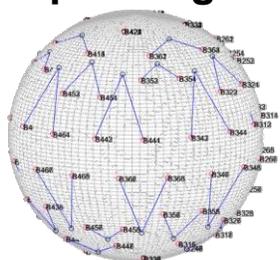
t~3.0 ns



t~3.8 ns

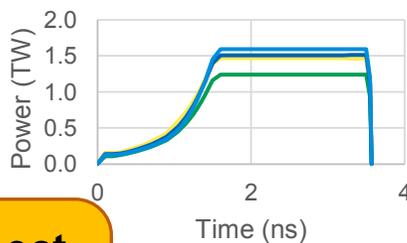


DRACO (new) pointing

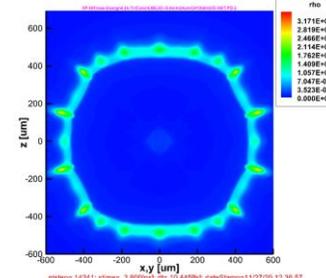


N201220-001

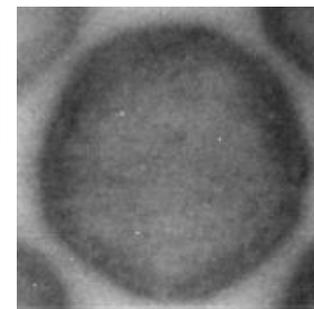
Per-Beam Laser Power



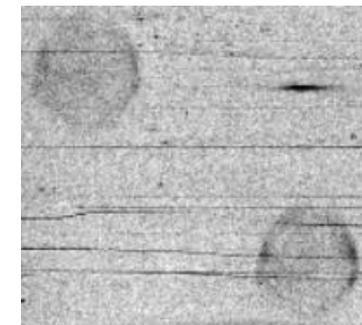
DRACO pre-shot
3.8 ns



t~3.5 ns



t~4.3 ns



2 mm

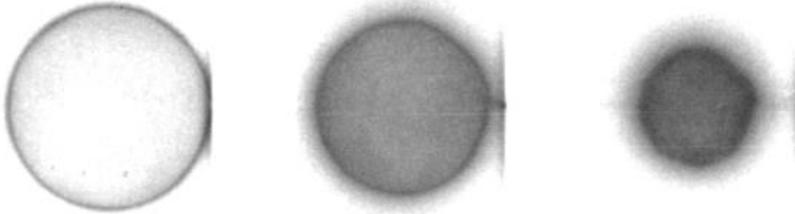
Design: J. Marozas

Design-001 maintains shape aspect ratio while 002 distorts in time implying a velocity nonuniformity

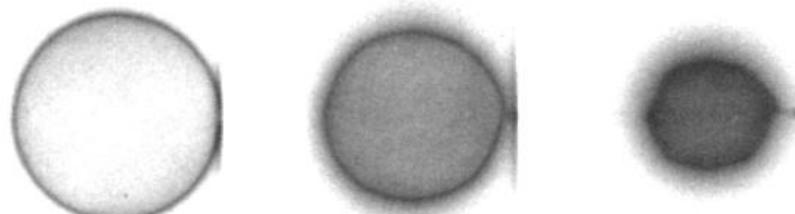
Qualitative 2D shape improvement close to pre-shot predictions

Gated x-ray self-emission images show agreement with postshot simulations, achieving the expected control of the shape

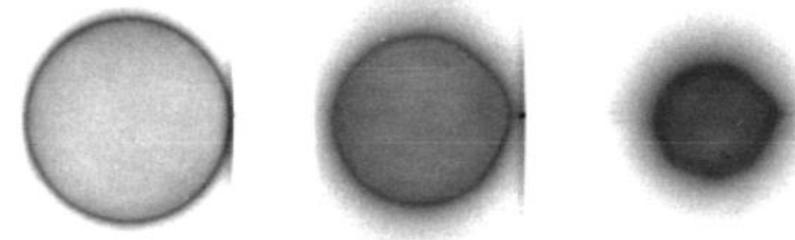
N201220-001 $t \sim 1.5$ ns $t \sim 2.5$ ns $t \sim 3.5$ ns $t \sim 4.3$ ns



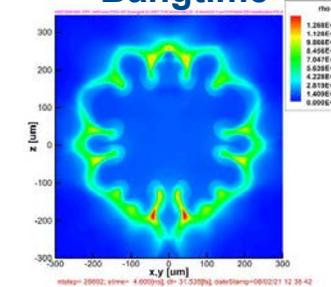
N201220-002 $t \sim 1.0$ ns $t \sim 2.0$ ns $t \sim 3.0$ ns $t \sim 3.8$ ns



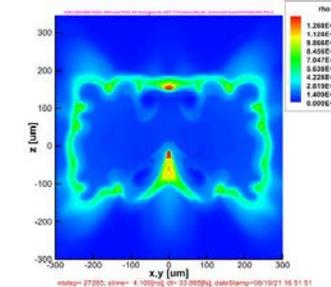
N201220-003 $t \sim 1.0$ ns $t \sim 2.0$ ns $t \sim 3.0$ ns $t \sim 3.85$ ns



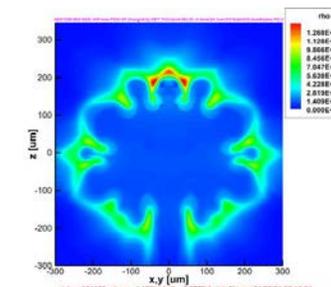
DRACO postShot at Bangtime



Improved pointing and pulse-shapes



Reference config.

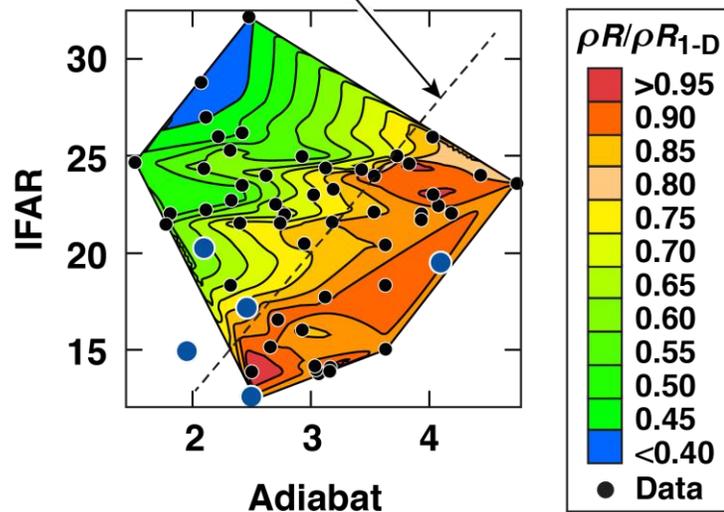


Improved pointing and ref. pulse

Simulations are employed to study the effect of imprint on target performance

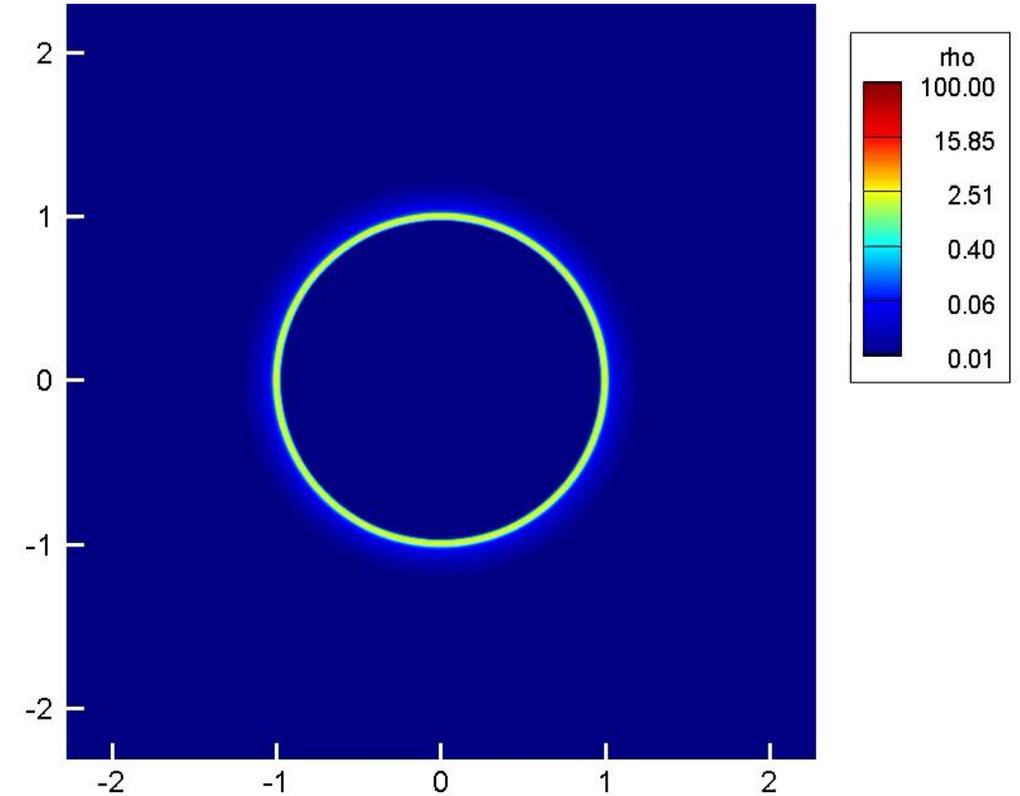
- Imprint has a strong impact on target performance
 - Compromises the shell integrity and lowers density and temperature and increases surface area
- Laser smoothing techniques mitigate imprint
- Simulations help gauge the laser smoothing requirements and guide target designs

Stability boundary $IFAR_s = 20 (\alpha/3)^{1.1}$



TC12040d

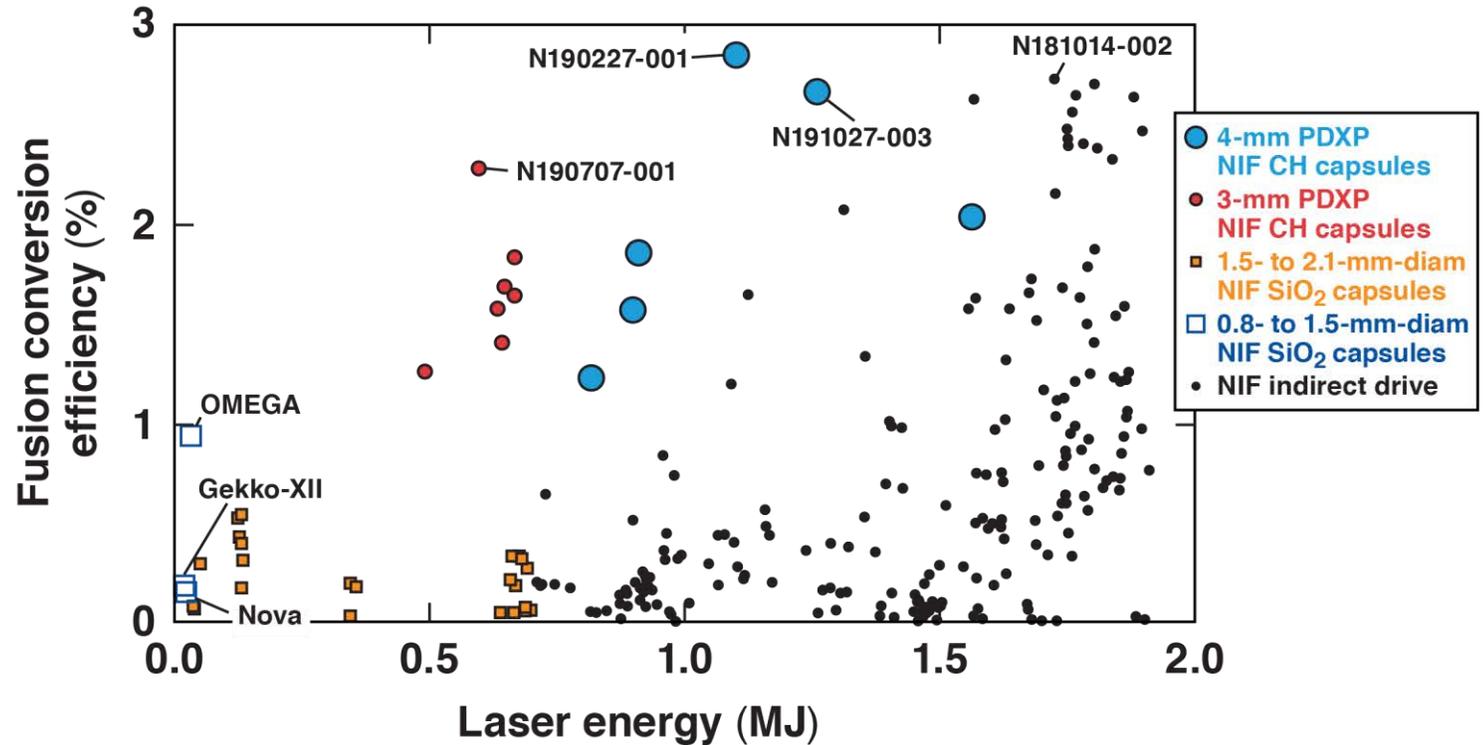
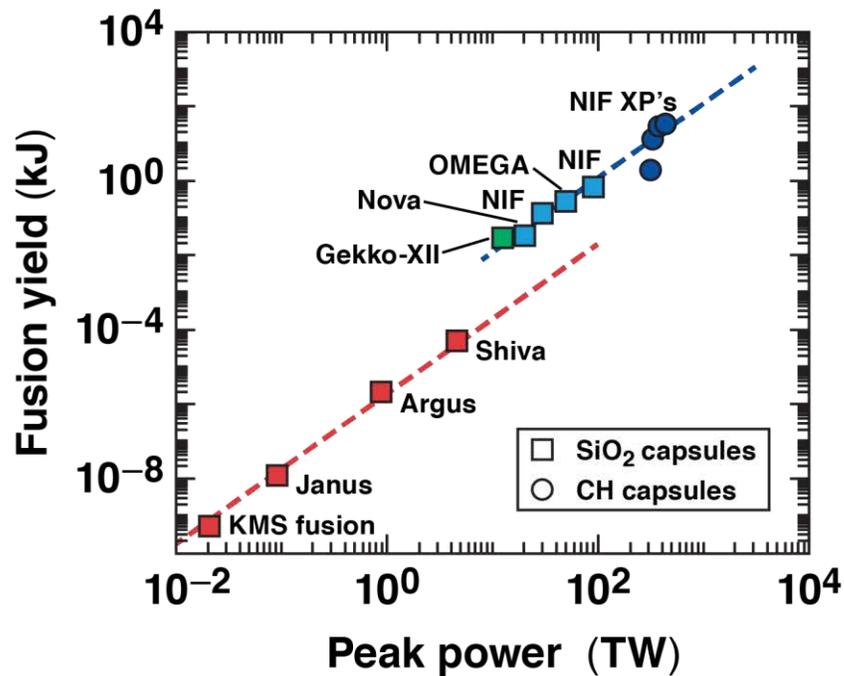
Imprint simulation* follows the shell in flight



ntstep= 6061; stime= 1.800[ns]; dt=191.709[fs]; dateStamp=03/06/19 17 16 57

*2-D DRACO Imprint Simulations with LLE Ray Trace, performed by T. Collins (2019).

XP targets on the NIF are being employed to study laser–energy coupling and implosion dynamics in an effort to push toward high yields to prepare the pathway toward ignition

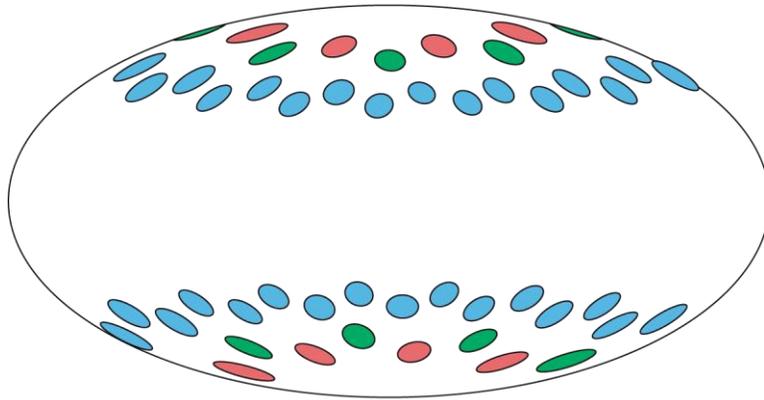


E29393

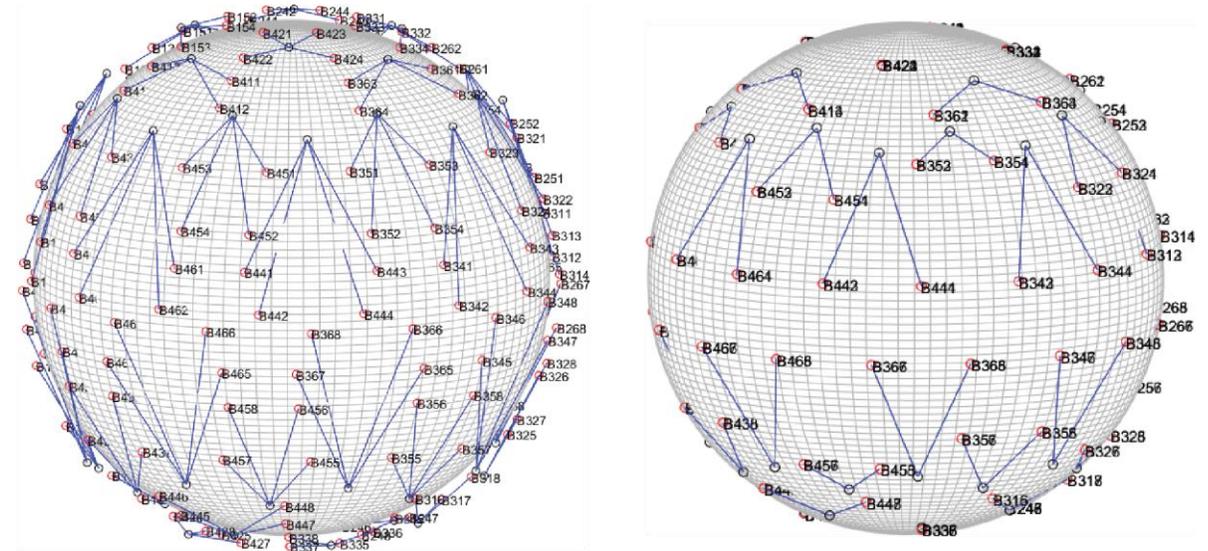
XP targets can achieve large DT_n yields owing to their high implosion speed ~ 700 to $900 \mu\text{m/ns}$, but their high electron preheat, low convergence and low ρR all prevent a propagating burn wave, i.e., ignition.

The NIF beams must be repointed for direct-drive experiments owing to its clustered polar port arrangement; referred to as polar direct drive (PDD)

LID port-color arrangement

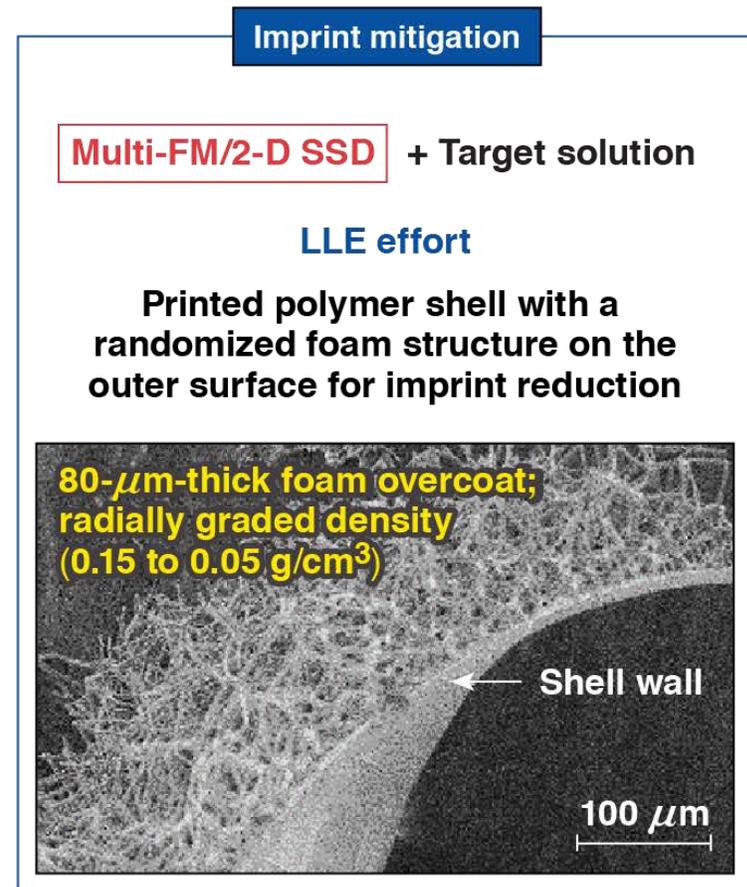
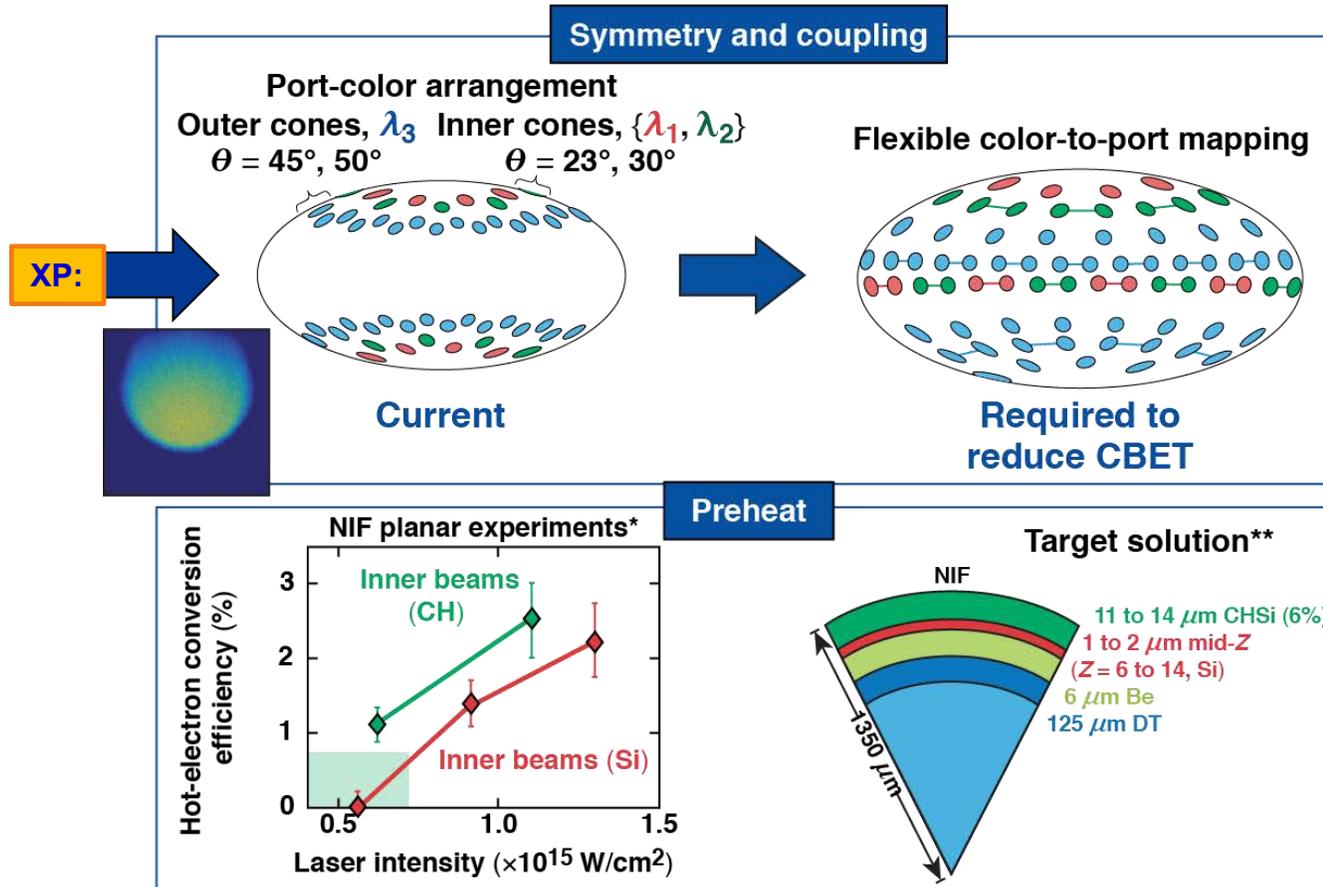


LDD-PDD port repointing designs



TC15447

NIF facility enhancements are required to perform $CR > 15$ PDD experiments



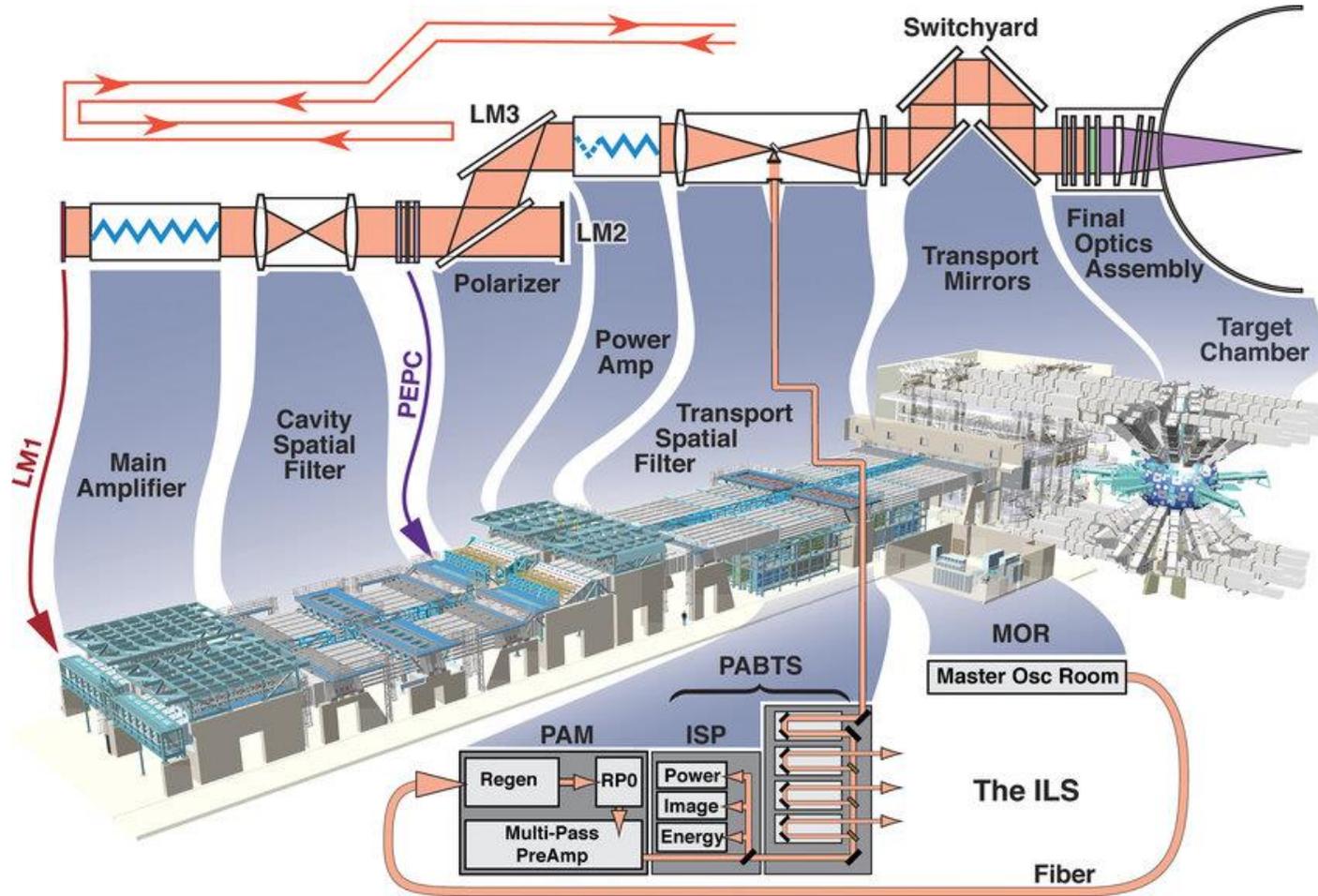
Experiments are underway to understand acceptable laser-target parameter space.

* A. A. Solodov *et al.*, Phys. Plasmas **27**, 052706 (2020).

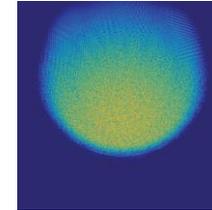
** V. N. Goncharov *et al.*, Phys. Plasmas **21**, 056315 (2014).

FM: frequency modulation
SSD: smoothing by spectral dispersion

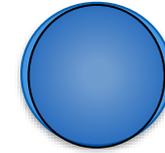
The XP platform will evaluate optimal pulse shape, spot-shapes (DPP), contoured shells, flexible color-to-port mapping (fC2Pm), and larger color separation ($\Delta\lambda_0$ detuning)



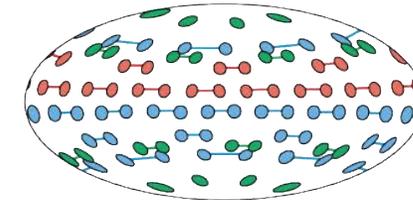
PDD-DPP



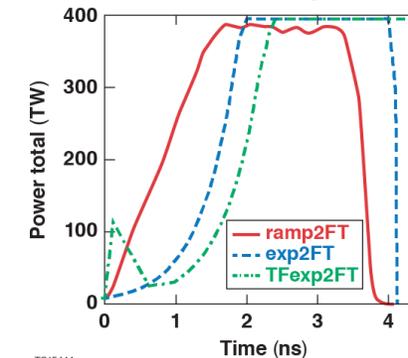
Contoured shell



fC2Pm with large $\Delta\lambda_0$



Pulse shapes



TC15444

Backup --- ICF Physics

The DT-plasma fuel assembly areal density (ρR) is a key ICF performance parameter



- The Lawson criterion defines a simple condition required for ignition:
 - ignition: (fusion products power density) \times (containment-time) $>$ (plasma energy density)

for 50-50 DT plasma:
$$n^2/4 \langle \sigma v \rangle Q \tau > 3nk_B T$$

solving for $n\tau$:*
$$n\tau > \frac{12k_B T}{\langle \sigma v \rangle Q} \approx 10^{14} : 10^{15} [\text{s/cm}^3]$$
 Lawson Criterion

- Containment time can be approximated by disassembly time; i.e., related to the ion sound speed traversal across the fuel assembly as

$$\langle \tau \rangle \cong \frac{R_f}{4c_s}, \quad n \equiv \rho/m \Rightarrow n\tau \cong \frac{\rho R_f}{4\langle M_i \rangle c_s}$$

$$\rho R_f \gtrsim 3 [\text{g/cm}^2] \text{ for efficient burn, i.e., large gains}$$

The areal-density, ρR , can be expanded to express other important ignition criteria...

* $Q_{DT} = 17.6 \text{ MeV}$

Maximizing hydroefficiency is key for successful ICF implosions

- Hydroefficiency $\eta_H = E_k/E_L$

$$\eta_H \sim \frac{p_{abl} V_{imp}}{I}$$

$$p_{abl} = f_{abs} I^{2/3}$$

To maximize hydroefficiency

Maximize laser coupling

- Mitigate LPI losses
- Use mid-Z ablator (balance between κ_B and radiation losses)

Increase implosion velocity

- Reduce shell mass

Reduce laser intensity

- Good for LPI but stability is a concern
- Advantage in multi-shell designs

Stability must be addressed:

$$IFAR \equiv R/\Delta R \sim v_{imp}^2 / p_{abl}^{2/5} \alpha^{3/5}$$

While working on mitigating the deleterious effects of LPI, the LDD program on the NIF and OMEGA leverages on v_{imp} and adiabat to maximize the performance.

LPI: laser-plasma instability
IFAR: in-flight aspect ratio

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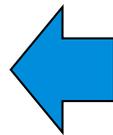
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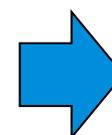
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$\rho R_f \gtrsim 3 [\text{g/cm}^2]$ for efficient burn, i.e., large gains

$$\rho R \sim \frac{(P_{\text{abl}}^2/I)^{1/3} E_L^{1/3}}{\sqrt{\alpha}}$$



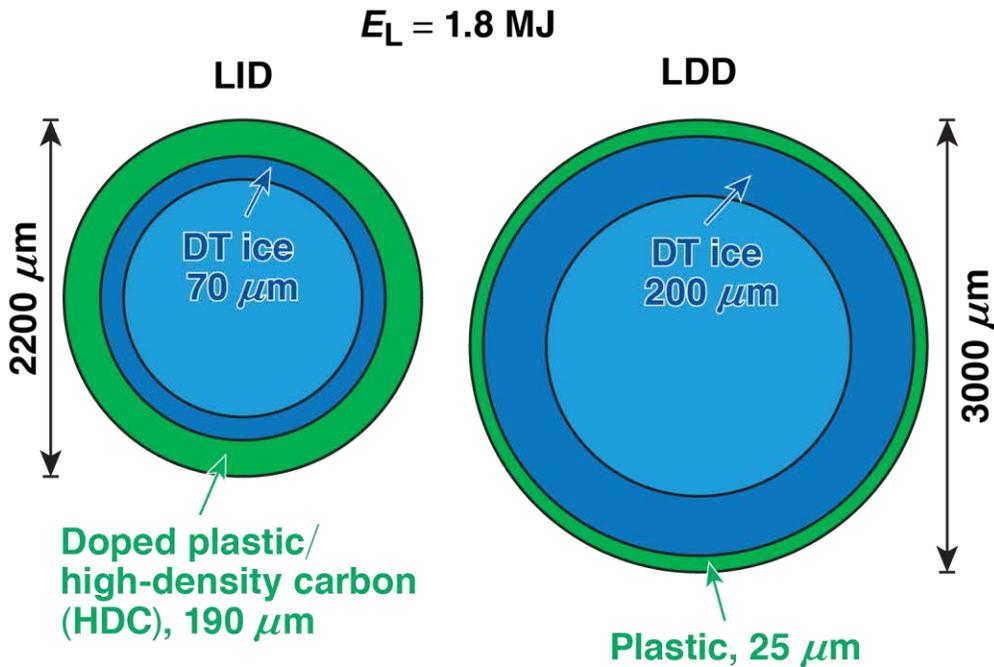
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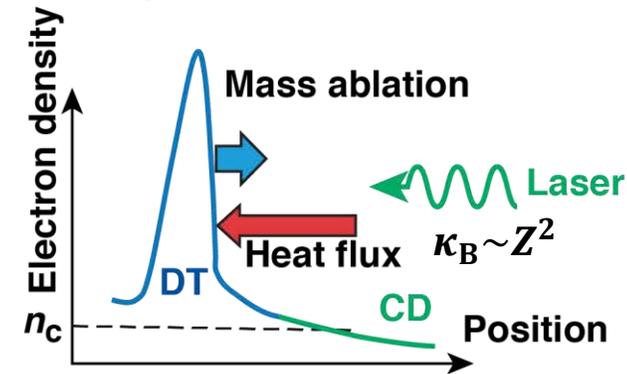
$$P_{\text{hs}} > 250 \text{ Gbar} \left(\frac{E_{\text{hs}}}{10 \text{ kJ}} \right)^{-1/2} = P_{\text{th}}$$

* $Q_{\text{DT}} = 17.6 \text{ MeV}$

LDD couples more energy to the hot spot compared to LID, which means more fuel mass and thinner ablators



Thinner plastic maximizes drive efficiency



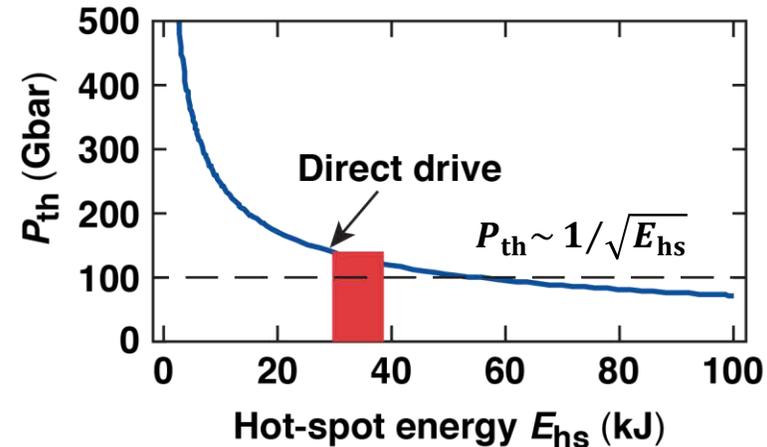
Atomic number \downarrow

$$m_{abl} \sim I^{1/3} \left(\frac{A}{Z}\right)^{2/3}$$

Ion charge \uparrow

$$p_{abl} \sim I^{2/3} \left(\frac{A}{Z}\right)^{1/3}$$

Threshold hot-spot pressure for ignition



E29127a

LDD ignition: $CR > 22$, $P_{hs} > 120 \text{ Gbar}$, $P_{hs} > 350 \text{ Gbar}$ has already been demonstrated in LID implosions.

Two critical parameters for ignition are hot-spot pressure and internal energy

From the Lawson Criterion:*

$$(\rho R)_{\text{hs}} \times T_i > 0.3 \text{ g/cm}^2 \times 5 \text{ keV}$$

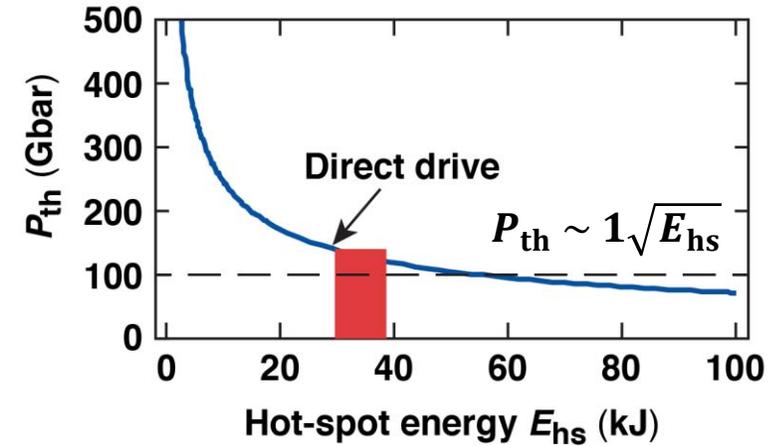
$$P_{\text{DT}} = 2\rho T / 2.5m_p$$

$$P_{\text{hs}} = 100 \text{ Gbar} \left(\frac{100 \mu\text{m}}{R_{\text{hs}}} \right)$$

$$E_{\text{hs}} = 3/2 P_{\text{hs}} V_{\text{hs}}$$

$$P_{\text{hs}} > 250 \text{ Gbar} \left(\frac{E_{\text{hs}}}{10 \text{ kJ}} \right)^{-1/2} = P_{\text{th}}$$

Threshold hot-spot pressure for ignition



E29127b

To burn the main fuel: Gain > 1, $\rho R > \text{g/cm}^2$

$$\rho R \sim \frac{(P_{\text{abl}}^2 / I)^{1/3} E_L^{1/3}}{\sqrt{\alpha}}$$

* R. Betti et al., Phys. Plasmas **17**, 058102 (2010).

Maximizing hydroefficiency is key for successful ICF implosions

Hydroefficiency $\eta_H = E_k/E_L$

$$\eta_H \sim \frac{p_{abl} v_{imp}}{I}$$

$$p_{abl} \sim f_{abs} I^{2/3}$$

Shell aspect ratio increases as I decreases (same V_{imp})

To maximize hydroefficiency

Maximize laser coupling

- Mitigate LPI losses
- Use mid-Z ablator (balance between κ_B and radiation losses)

Increase implosion velocity

- Reduce shell mass

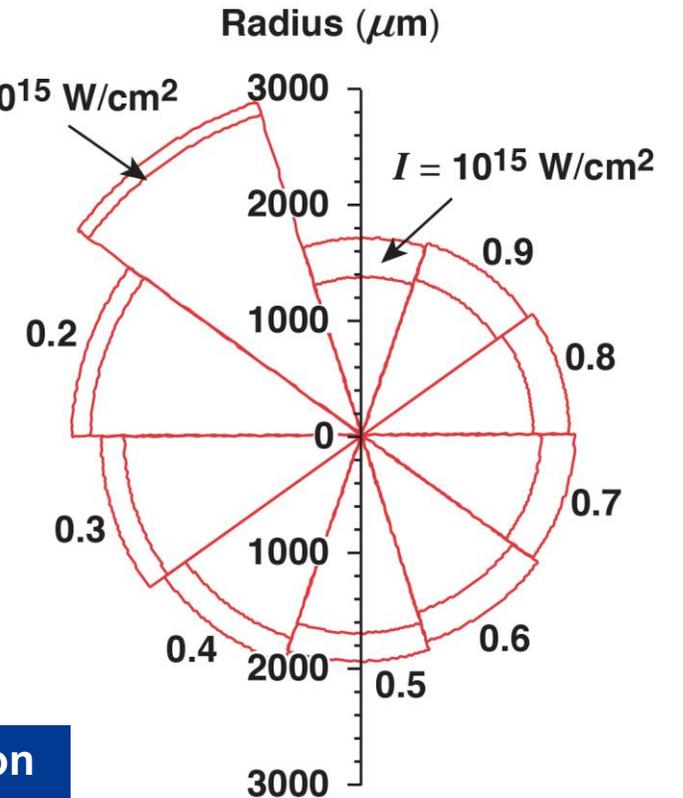
Reduce laser intensity

- Good for LPI but stability is a concern
- Advantage in multi-shell designs

Stability must be addressed:

$$IFAR = \frac{R}{\Delta R} \sim v_{imp}^2 / p_{abl}^{2/5} \alpha^{3/5}$$

$I = 0.1 \times 10^{15} \text{ W/cm}^2$

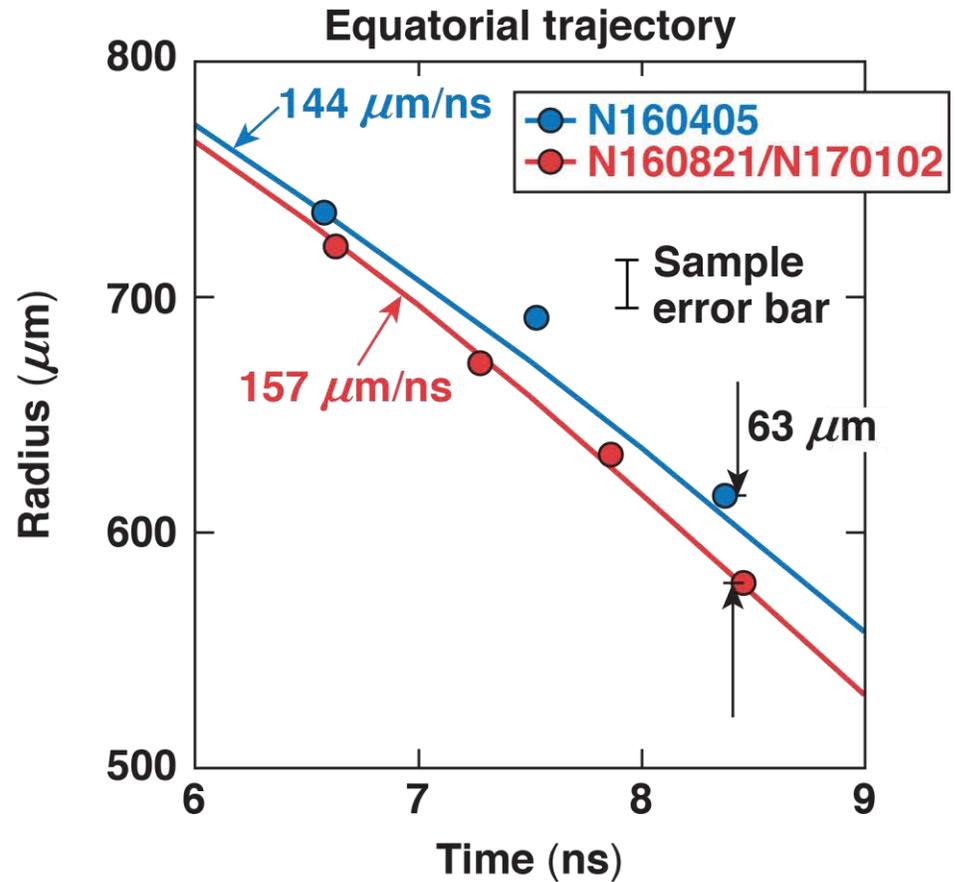
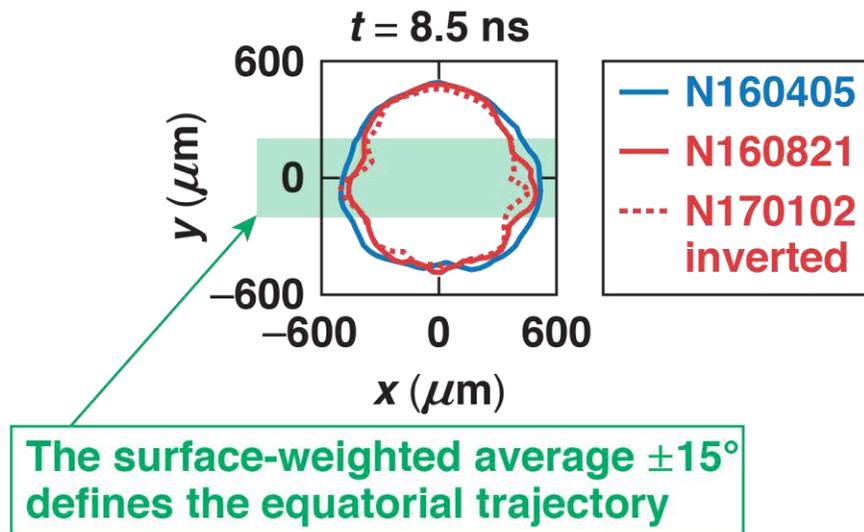


While working on mitigating the deleterious effects of LPI, the LDD program on the NIF and OMEGA leverages on v_{imp} and adiabat to maximize the performance.

LPI: laser-plasma instability
IFAR: in-flight aspect ratio

Improved equatorial coupling from wavelength detuning is inferred from gated radiographs

- The predicted and measured trajectories* show the expected faster implosion speeds near the equator



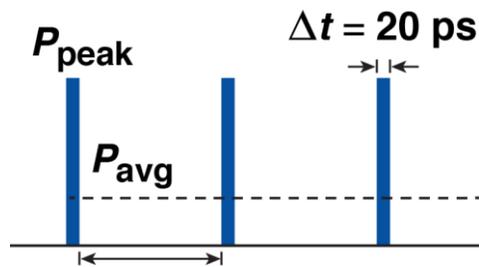
TC13351b

* Analysis by D. Turnbull

Laser energy coupling loss caused by CBET can be mitigated in different domains that can be combined; **temporal, spatial, and spectral**

- Temporal domain**

- multiplexing the beams reduces interaction



- STUD* pulses

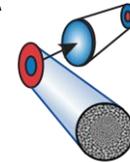
* STUD: spike trains of uneven duration and delay; B. Afeyan and S. Hüller, EPJ Web Conf. **59**, 05009 (2013).

- Spatial domain**

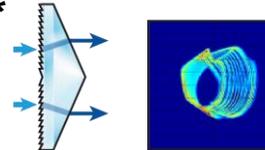
(reduce interaction volume)

- dynamic spot shape

- two stage**



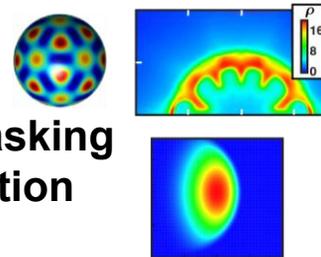
- Graxicon***



- KrF lasers (NRL)

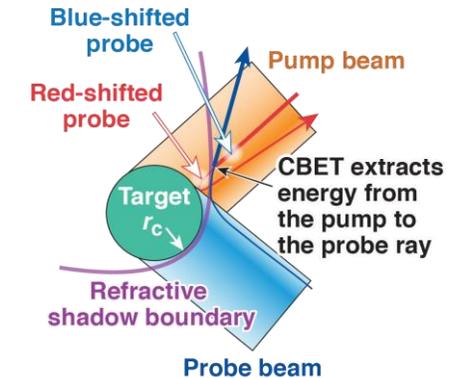
- spots smaller than target (e.g., R_{75})†

- spot-masking apodization (SMA)



- Spatial domain**

- wavelength detuning; Dm_0



- wide bandwidth within each beam (e.g., SRRS‡)
- lower intensity per band and incoherence disrupts growth

** D. H. Froula *et al.*, Phys. Plasmas **20**, 082704 (2013).

*** T. J. Kessler and H. Huang, presented at the Ninth International Conference on Inertial Fusion Sciences and Applications (IFSA 2015), Seattle, WA, 20–25 September 2015 (Abstract Mo.Po.61).

† S. P. Regan *et al.*, Phys. Rev. Lett. **117**, 025001 (2016); 059903(E) (2016).

‡ SRRS: stimulated rotational Raman scattering;

J. Weaver *et al.*, “Spectral and Far-Field Broadening due to Stimulated Rotational Raman Scattering Driven by the Nike Krypton Fluoride Laser,” to be published in Applied Optics.

The XP platform has a long history dating back to the early days of ICF, circa 1970's

- Thin-shelled glass microballoons filled with DT gas were shot with an intense short laser burst

A THEORETICAL INTERPRETATION OF EXPLODING PUSHER LASER FUSION EXPERIMENTS

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ABSTRACT. The implosion of neon-filled glass microballoons has been examined via numerical simulations using a one-dimensional Lagrangian hydrodynamic code. Comparison has been made to experiments performed on a four-beam neodymium glass laser system using an array of diagnostics including X-ray pinhole

Exploding pusher performance – A theoretical model

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University of California, Lawrence Livermore Laboratory, Livermore, California 94550
(Received 24 July 1978; final manuscript received 8 March 1979)

A simple model of "exploding-pusher" laser fusion targets is presented. The peak ion temperature is calculated by assuming that the ion heating mechanism is a shock followed by isentropic compression. The compression is calculated by conservation of mass and reasonable assumptions on the density profile at peak compression time. The neutron yield predictions of this model agree over a broad range of parameter space with those of complex one-dimensional numerical simulation, which in two dimensions have tracked experimentally measured yields quite accurately.

I. INTRODUCTION

Exploding pusher targets¹ have been the most common targets of early laser fusion experiments, and are the first targets to be shot on the 25 TW SHIVA laser system at Lawrence Livermore Laboratory. Typically, they consist of a glass microballoon filled with deuterium-tritium (DT) gas. For SHIVA, typical dimensions are radius 200 μ , thickness 2 μ , and DT fill 2 mg/cc.

conduction and by expansion.

Whereas the high-compression, isentropic implosion targets are sensitive to electron preheat and to Rayleigh Taylor instability, exploding pushers, by virtue of their rapid thermal wave early heating and by their nonablative implosion dynamics, are not sensitive to the aforementioned problems. On the other hand, the preheat levels of these targets and exploding pusher



A NIF XP*

How do they work?...

* M. J. Rosenberg *et al.*, Phys. Plasmas **21**, 122712 (2014).

The XP platform has a long history dating back to the early days of ICF, circa 1970's

- Thin-shelled glass microballoons filled with DT gas were shot with an intense, short laser burst
 - the laser energy rapidly heats the electrons
 - preheats the DT electrons (**reduces convergence, low ρR**)
 - causes **$\sim 1/2$** the shell-mass to explosively ablate (**high stability**)
 - the imploding inner **$\sim 1/2$** -shell acts as a piston, driving a strong inward shock (**high speeds**)
 - heats the ions to thermonuclear temperatures, producing fusion reactions
 - the exploding outer shell and imploding/compressing inner shell led to these targets being called **exploding pushers**



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Current XP experiments use similar concepts, scaled to higher energy and similar aspect ratio targets but leveraging shaped pulses to drive the remaining shell during compression, further increasing the DT_n yield.

The XP platform has a long history dating back to the early days of ICF, circa 1970's

- **Thin-shelled glass microballoons filled with DT gas were shot with an intense short laser burst**
 - the laser energy rapidly heats the electrons, driving an inward heat wave into the shell and DT gas preheating the DT electrons (**reduces convergence, low ρR**)
 - the thermal wave heats the shell ions (strong coupling) but not the DT ions (weak coupling)
 - the rapid deposition causes $\sim 1/2$ the shell mass to explosively ablate (**high stability**)
 - the imploding inner 1/2-shell acts as a piston driving a strong inward shock (**high speeds**)
 - this shock wave principally heats the ions
 - the inner shell continues to move inward behind the shock front, compressing the post-shock DT gas, further heating the ions to thermonuclear temperatures, producing fusion reactions
 - the shock and compression cause $T_i \gg T_e$, while the weak coupling preserves the imbalance
 - the exploding outer shell and imploding/compressing inner shell led to the moniker **exploding pushers**
 - eventually the thermal pressure builds and the return shock impacts the shell, both contributing to deceleration and halting compression
 - fusion reactions continue until the DT is cooled by thermal conduction and expansion
- **Current XP experiments use similar concepts but with larger targets with similar aspect ratios and shaped pulses that continue to drive the remaining shell during compression, further increasing the DT_n yield**

The XP platform has a long history dating back to the early days of ICF, circa 1970's

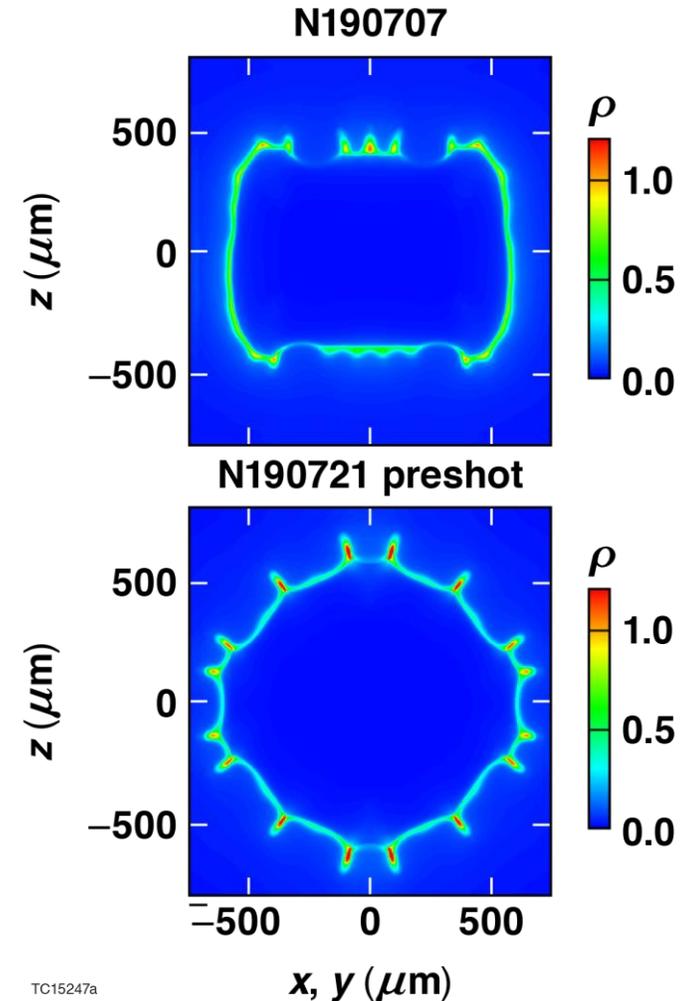
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 - the exploding outer-shell and imploding/compressing inner-shell led to these targets being called **exploding pushers**

Current XP experiments use similar concepts scaled to higher energy and similar aspect ratio targets, but leveraging shaped pulses to drive the remaining shell during compression, further increasing the DT_n yield.

Upcoming XP experiments on the NIF hope to increase neutron yield using improved pulse shapes and pointings based on 2-D simulations that include necessary physics

Elements of the new design

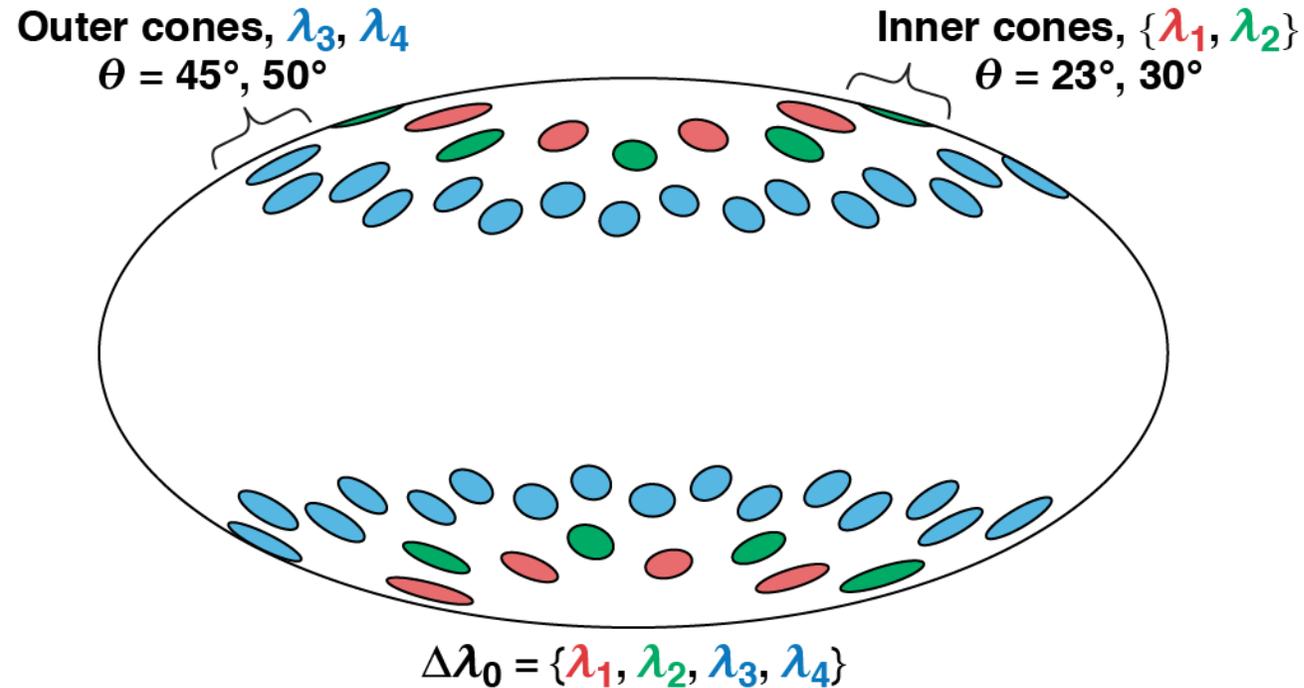
- Dual shocks improve yield while remaining high adiabat and predominantly “shock yield”
- The pulse shape improves separation of shell from shock, improving yield
- Steep main pulse rise improves coupling
- Simpler quad splitting improves power imbalance
- Repointing and pulse shapes yield rounder implosions
- Extensible
 - large targets with optimized DPP’s >100 kJ



TC15247a

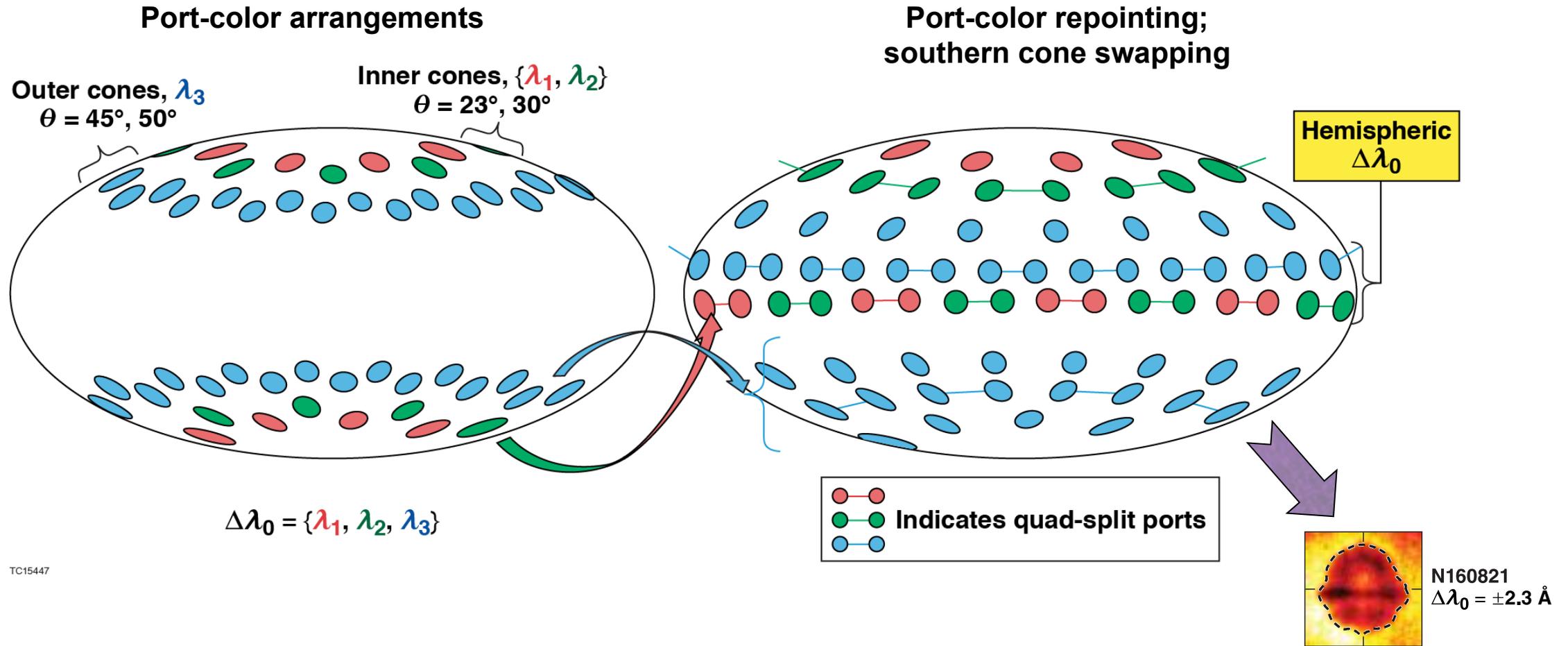
The NIF LID four-color laser sources are currently fixed into a symmetric pattern

LID port-color arrangement



TC15447

Early NIF PDD experiments employing wavelength detuning to study CBET mitigation used extreme repointing to induce a hemispheric color change

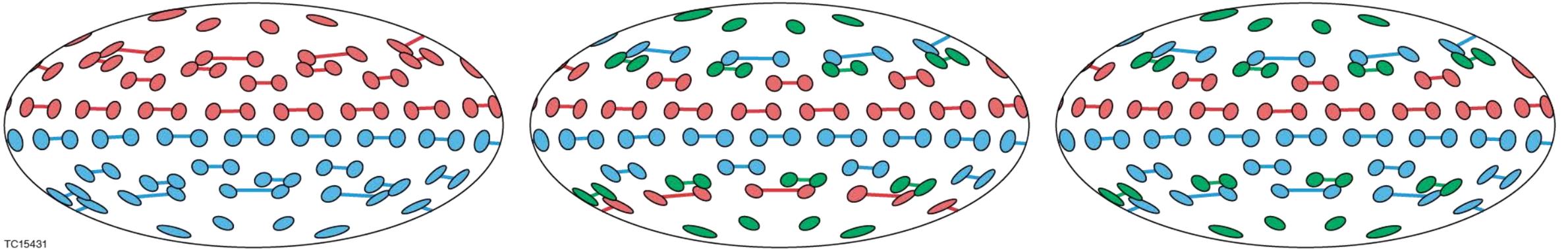


TC15447

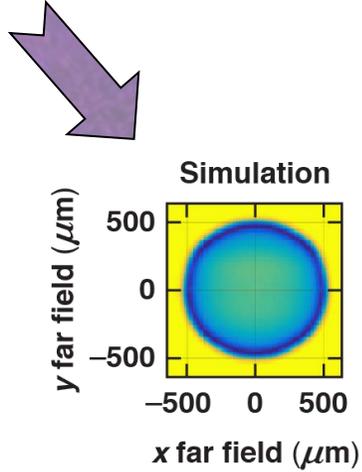
* J. A. Marozas et al., Phys. Rev. Lett. **120**, 085001 (2018).

The NIF fiber-optic front end will be updated to flexibly remap the four-color sources to a variety of port configurations to help mitigate CBET

Flexible color-to-port mapping (fC2Pm) with large $\Delta\lambda_0$

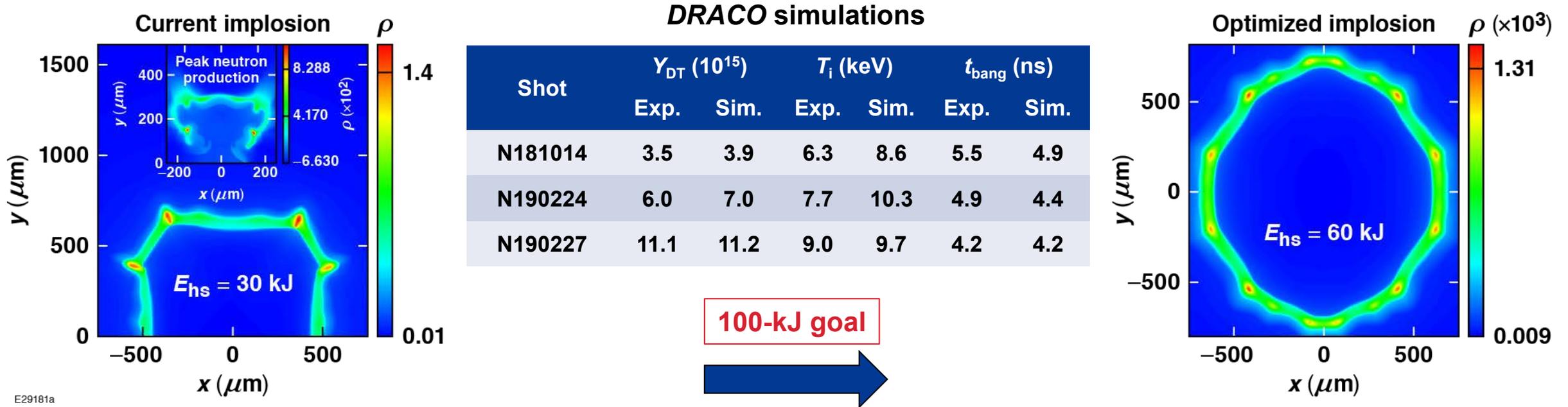


The new front-end remapping eliminates the need for extreme repointing to induce a hemispheric color change.



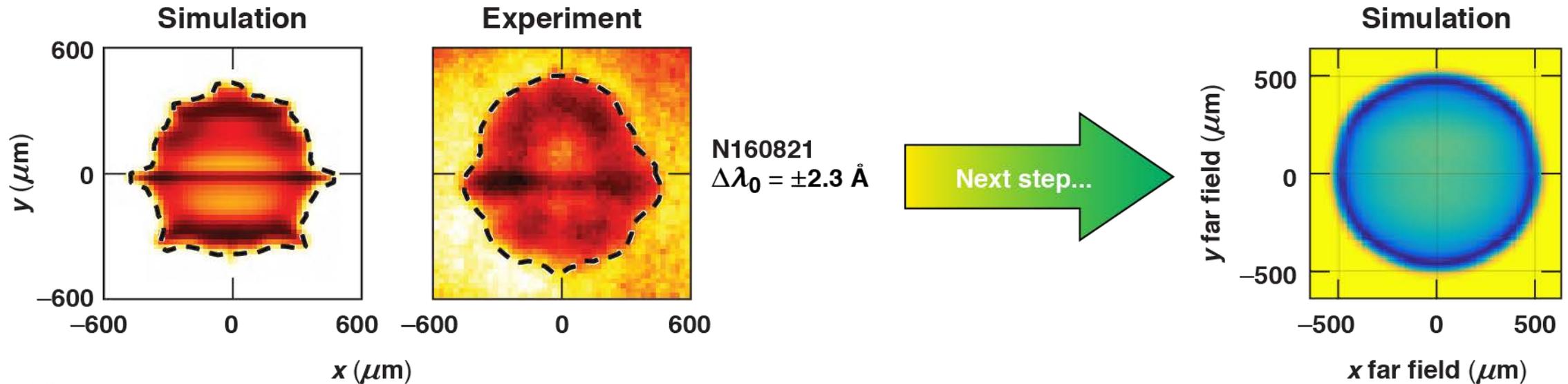
The performance of current implosions can be improved by better symmetry control and enhanced laser coupling

- **CBET mitigation strategies** – smaller SMA spots, $\Delta\lambda = \pm 6 \text{ \AA}$ $\Rightarrow E_k = 100 \rightarrow 120 \text{ kJ}$
- **Symmetry control strategies** – pulse shaping, repointing, PDD phase plates $\Rightarrow E_{hs}/E_k = 30\% \rightarrow 50\%$



E29181a

Application of both large wavelength detuning, fC2Pm, and optimized PDD spot shapes promises large returns for high-convergence PDD experiments on the NIF



TC13885e

Low convergence XP: \rightarrow >100kJ yield
High Convergence: alpha burner, ignition