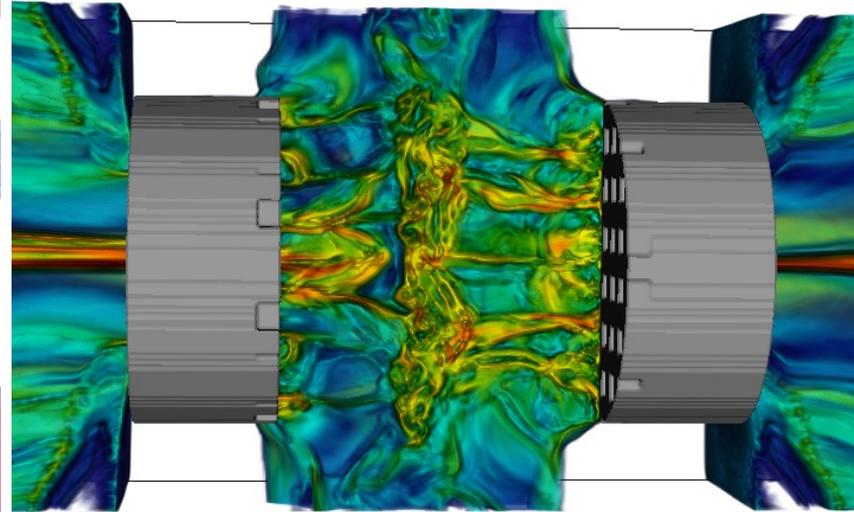
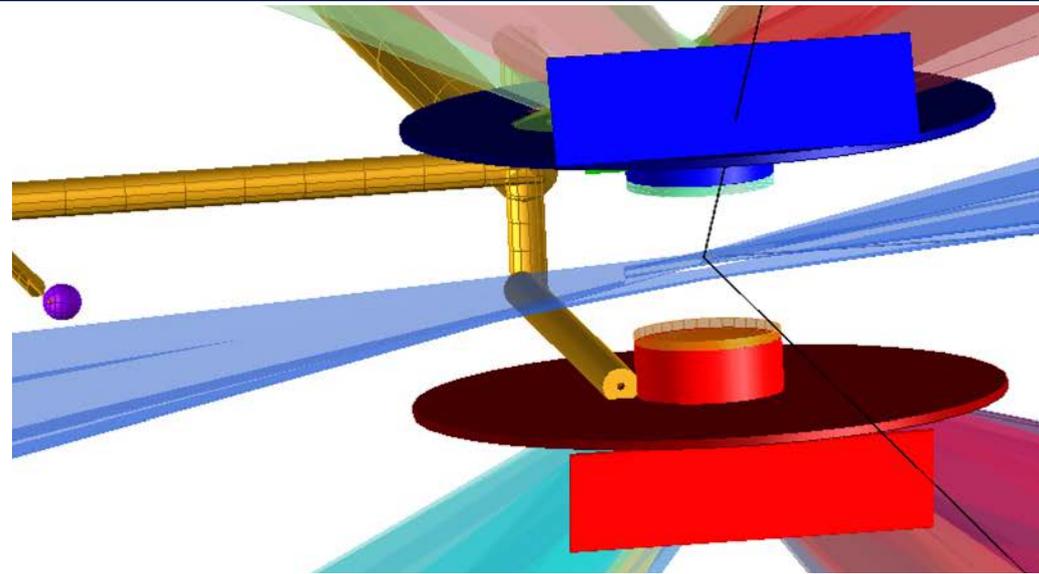


Numerical Modeling of Laser-Driven Plasma Experiments Aiming to Study Turbulent Dynamo and Thermal Conduction at the National Ignition Facility



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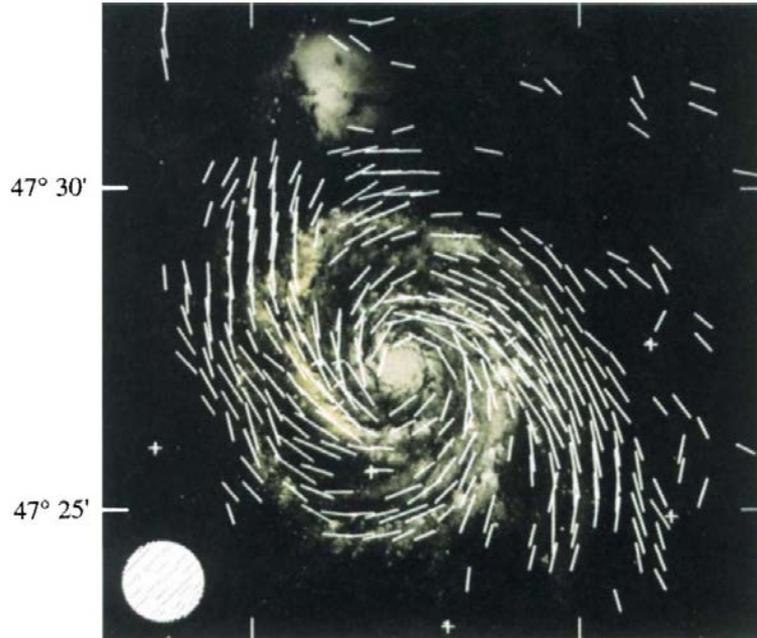
-  1. California State University Channel Islands, Camarillo, USA
-  2. University of Oxford, Oxford, UK
-  3. Princeton University, Princeton, USA
-  4. University of Rochester, Rochester, USA
-  5. Cornell University, Ithaca, USA
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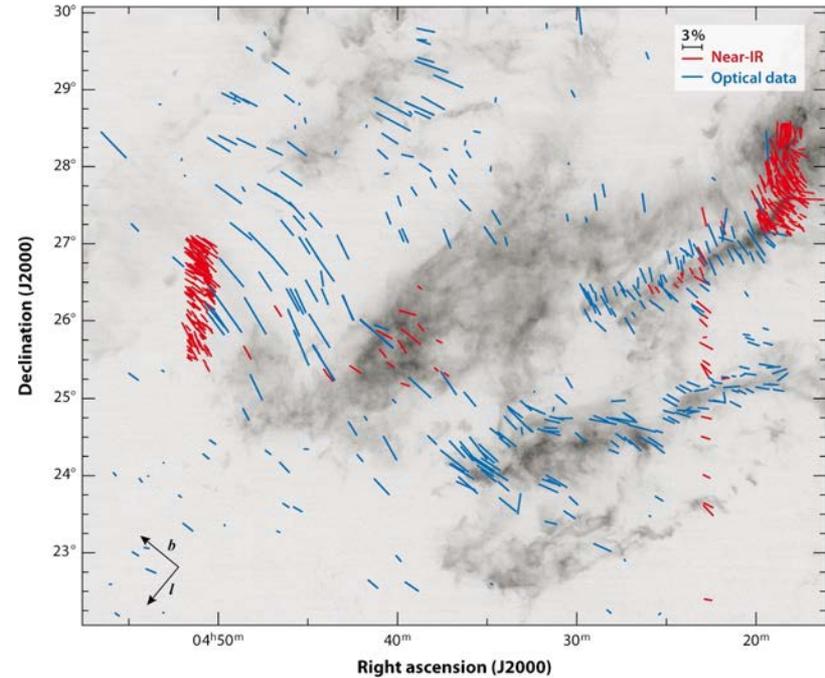


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- NSF Award PHY-2033925
- Subcontracts 536203 and 630138 with LANL and B632670 with LLNL.

- ❑ We design the experiment using FLASH simulations and create a turbulent plasma on NIF with magnetic field amplification under magnetic Reynolds number $R_m > 10^3$ and magnetic Prandtl number $P_m > 1$ condition.
- ❑ Fluctuation dynamo operates under these conditions and amplifies magnetic fields to the experimentally measured amplitude.
- ❑ Such a plasma with $P_m > 1$ is relevant to hot low-density plasmas found in astrophysical accretion disks and the intracluster medium (ICM).
- ❑ The magnetization is large enough for the electron Larmor radius to be less than mean free path, which can affect thermal conduction. Using FLASH simulations, we have created an experimental platform to study thermal conduction suppression as seen in active galactic nuclei (AGN) feedback.
- ❑ Synthetic proton radiography and X-ray self-emission images reveal characteristics of the turbulent plasma, e.g., electron temperature and magnetic field distributions, and enable direct comparison with the experiments.



13 h 28 min 00 s 13 h 27 min 30 s
B-field vectors of M51 (Zweibel 1997)



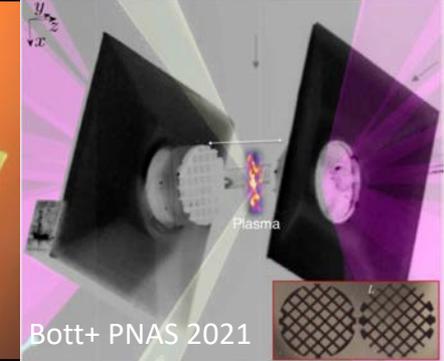
B-field in the Taurus dark-cloud complex (Crutcher 2012)

- ❑ Seed magnetic field generation by Biermann battery effect (Biermann 1950) or Weibel instability (Weibel 1959)
- ❑ Turbulent dynamo: small-value seed fields amplified to observed values (Parker 1979, Moffatt 1978, Kulsrud 1997)

Turbulent Dynamo at Omega Laser



Tzeferacos+
Nat Comm 2018

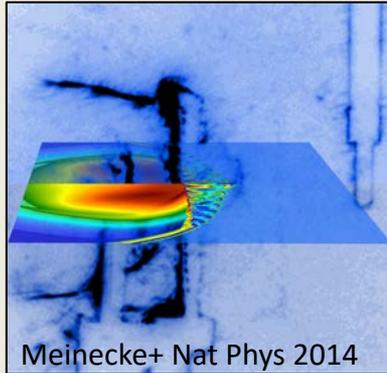


Bott+ PNAS 2021

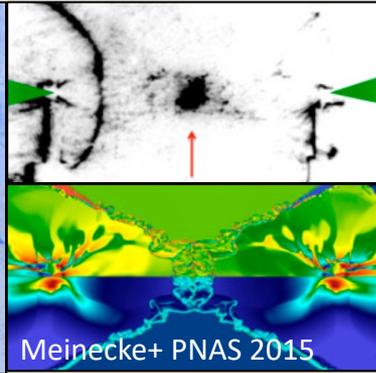


$Rm \sim 600$

Magnetic Field Amplification at Vulcan Laser



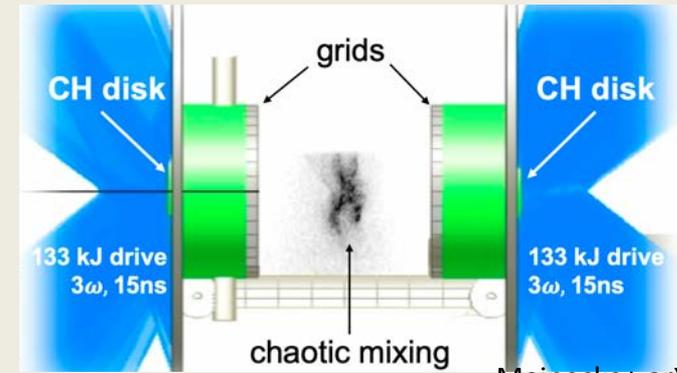
Meinecke+ Nat Phys 2014



Meinecke+ PNAS 2015

$Rm \sim 3 - 6$

NIF Dynamo Experiments and Effects on Heat Conduction



CH disk
133 kJ drive
 3ω , 15ns

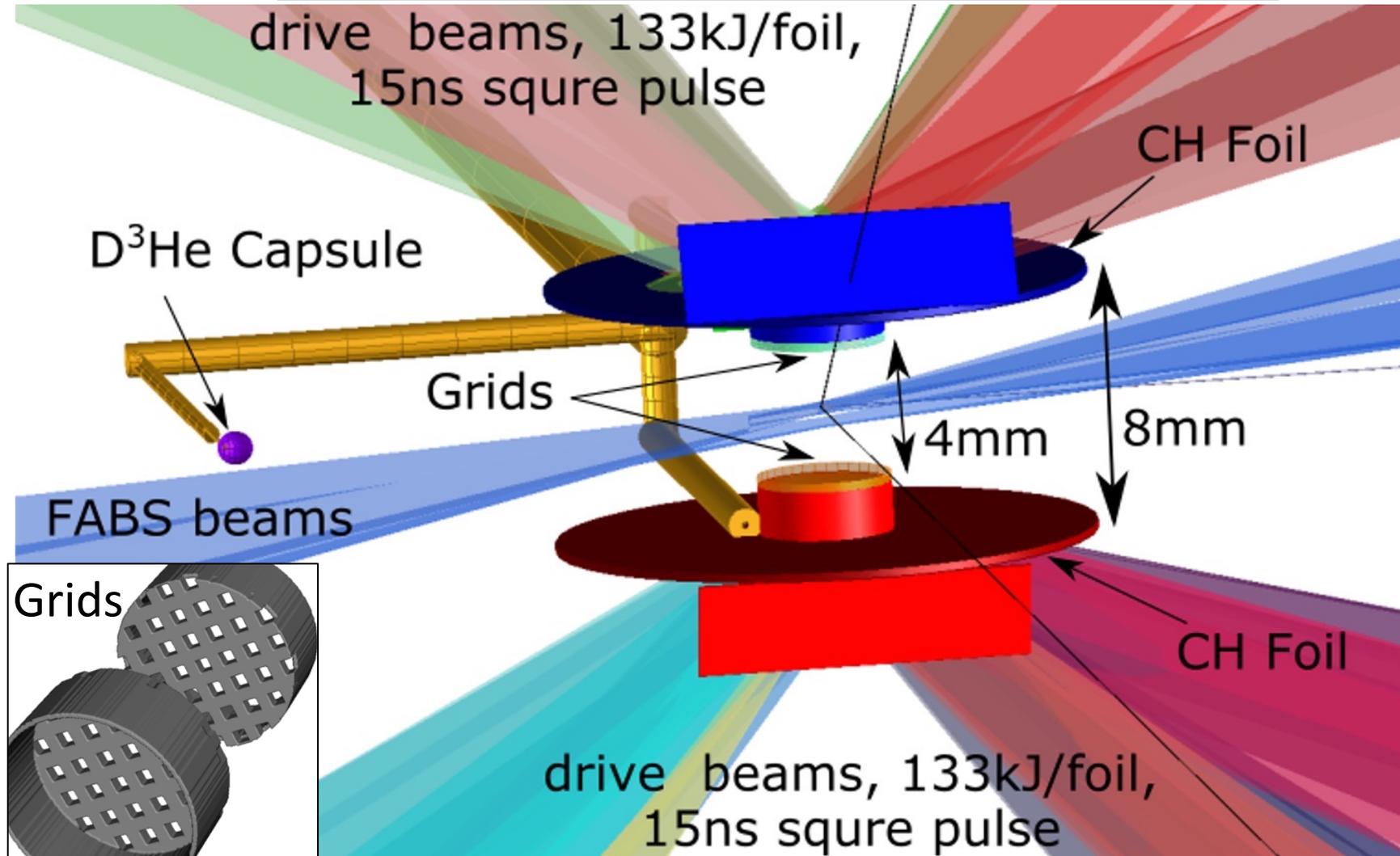
CH disk
133 kJ drive
 3ω , 15ns

chaotic mixing

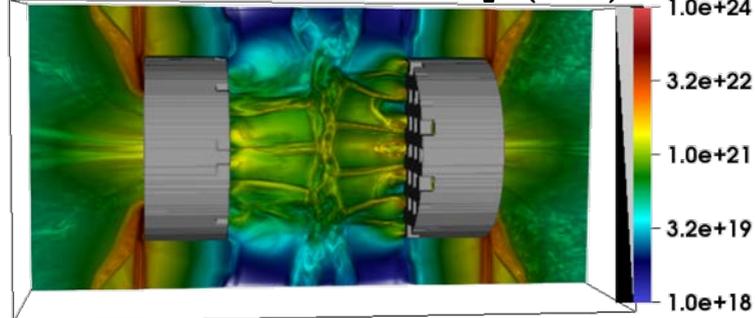
$Rm \sim 3500$

Meinecke+ arXiv:2105.08461

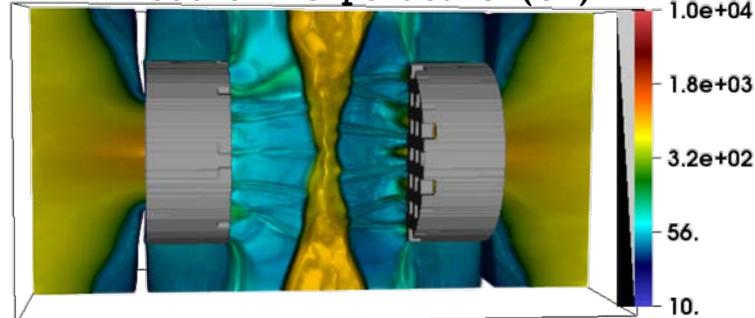
NIF Turbulent dynamo laser-target design



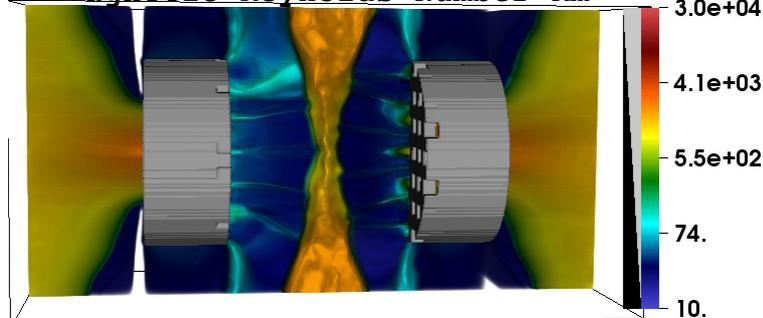
Electron Number Density (cm^{-3})



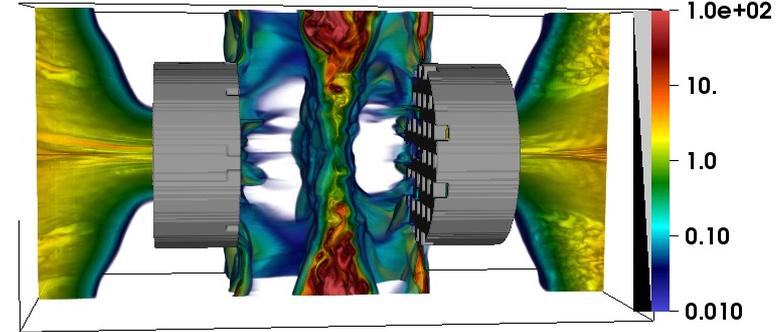
Electron Temperature (eV)



Magnetic Reynolds Number R_m



Magnetic Prandtl Number P_m



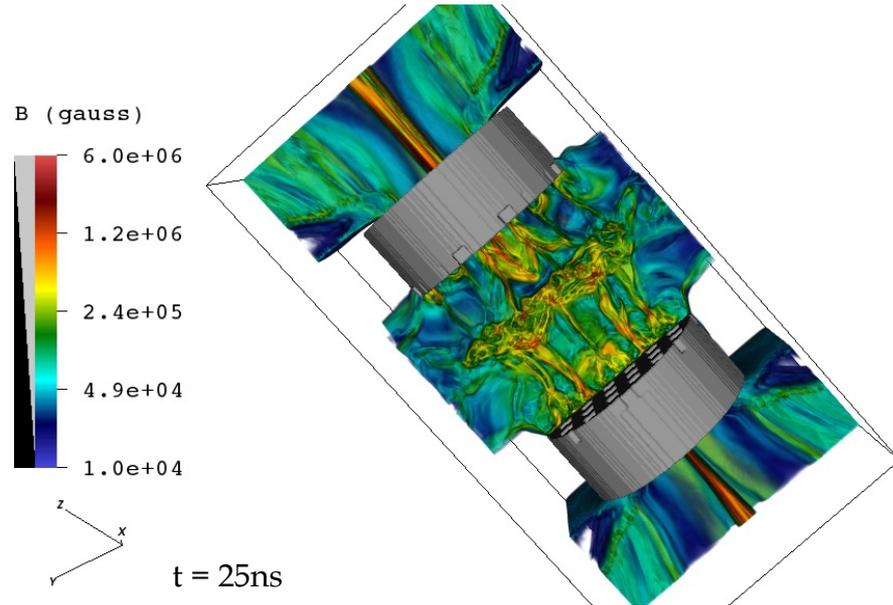
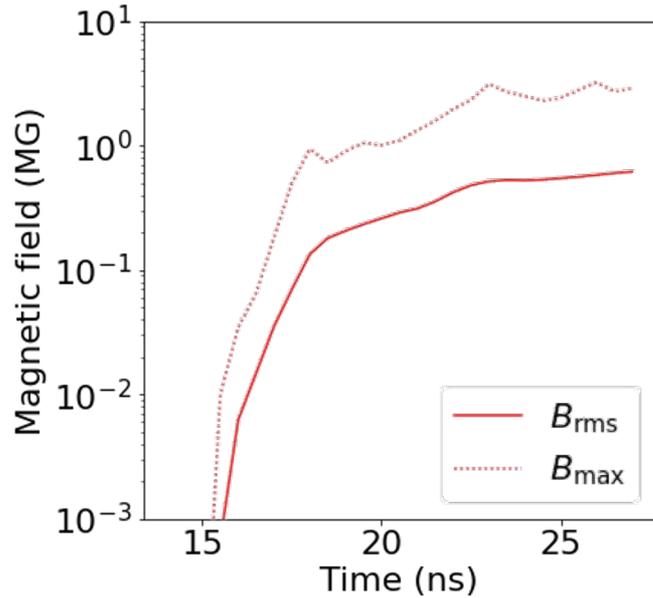
- The Interaction region turbulent plasma at 23ns has reached the following condition:

$$u_{\text{rms}} \sim 2 \times 10^7 \text{ cm/s},$$

$$n_e \sim 2.3 \times 10^{20} \text{ cm}^{-3}, T_e \sim 710 \text{ eV}$$

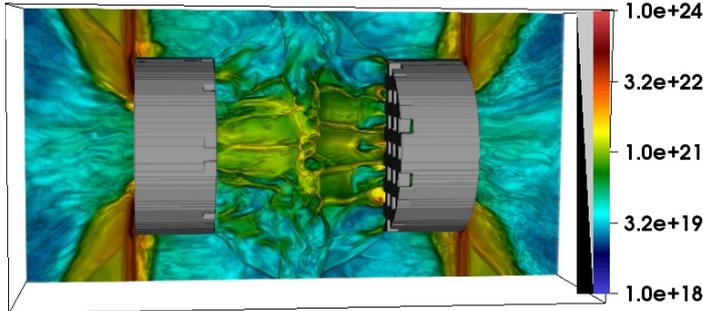
$$R_m \sim 1.8 \times 10^3, P_m \sim 4.$$

- Some locations in the interaction region can have higher R_m and P_m .
- We reach the condition for fluctuation dynamo to amplify the seed magnetic field.

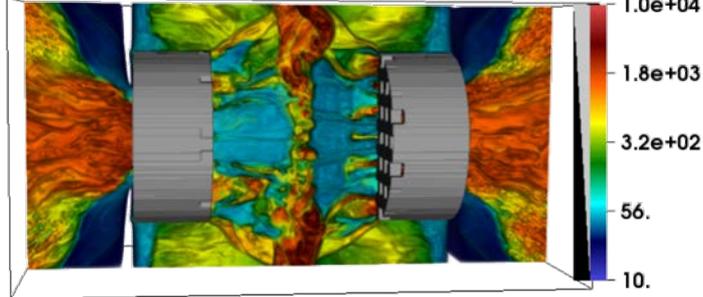


- ❑ Turbulent amplification of magnetic field results in $B_{rms} \sim 0.5$ MG and $B_{max} \sim 2.4$ MG.
- ❑ The plasma is in magnetized regime (electron Larmor radius $r_g <$ mean free path λ_e) with $r_g = 1.3 \times 10^{-4}$ cm and $\lambda_e = 1.1 \times 10^{-3}$ cm.
- ❑ Electron heat conduction is suppressed due to magnetization or other sources (e.g., kinetic instabilities, turbulence, non-local effects). We incorporate the suppression by switching off the electron thermal conductivity in FLASH.

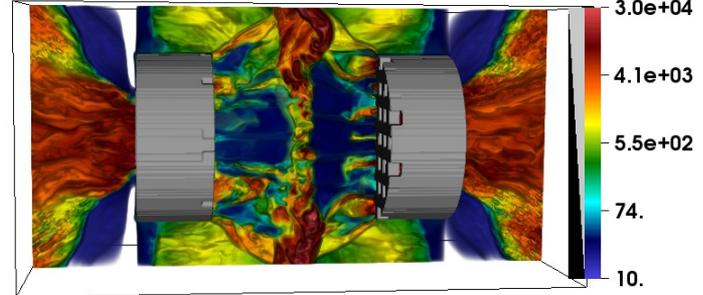
Electron Density (cm^{-3})



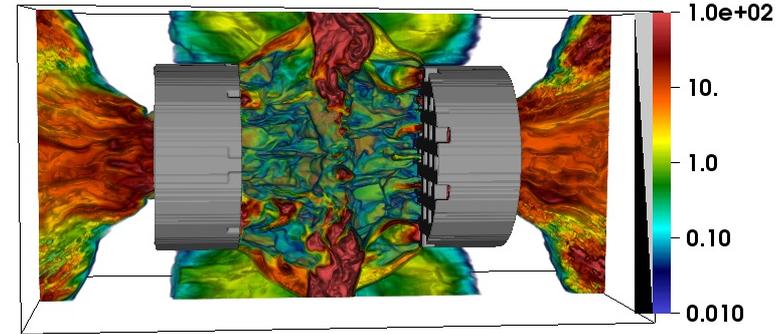
Electron Temperature (eV)



Magnetic Reynolds Number R_m



Magnetic Prandtl Number P_m

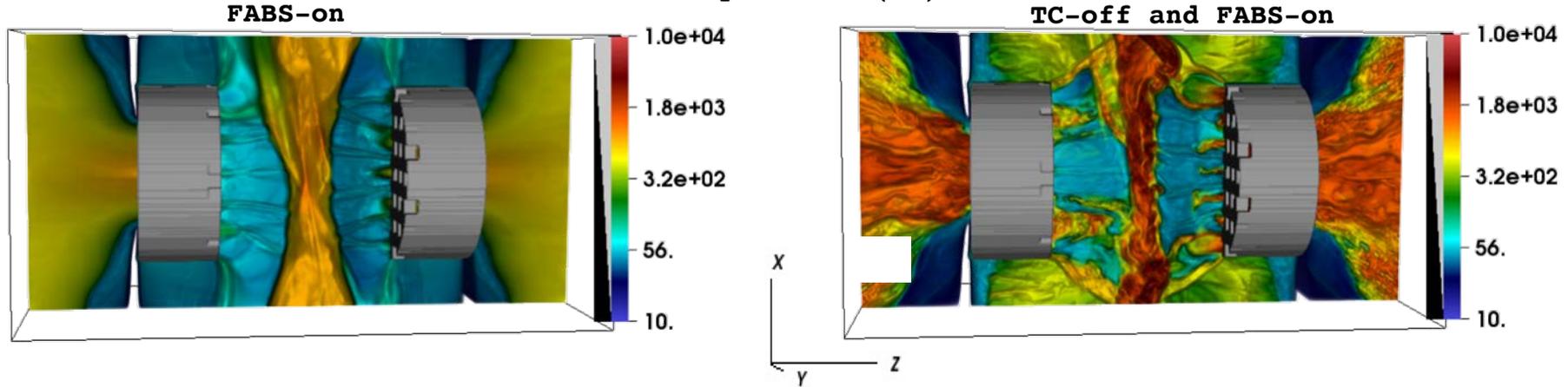


X
Y Z
t = 23ns

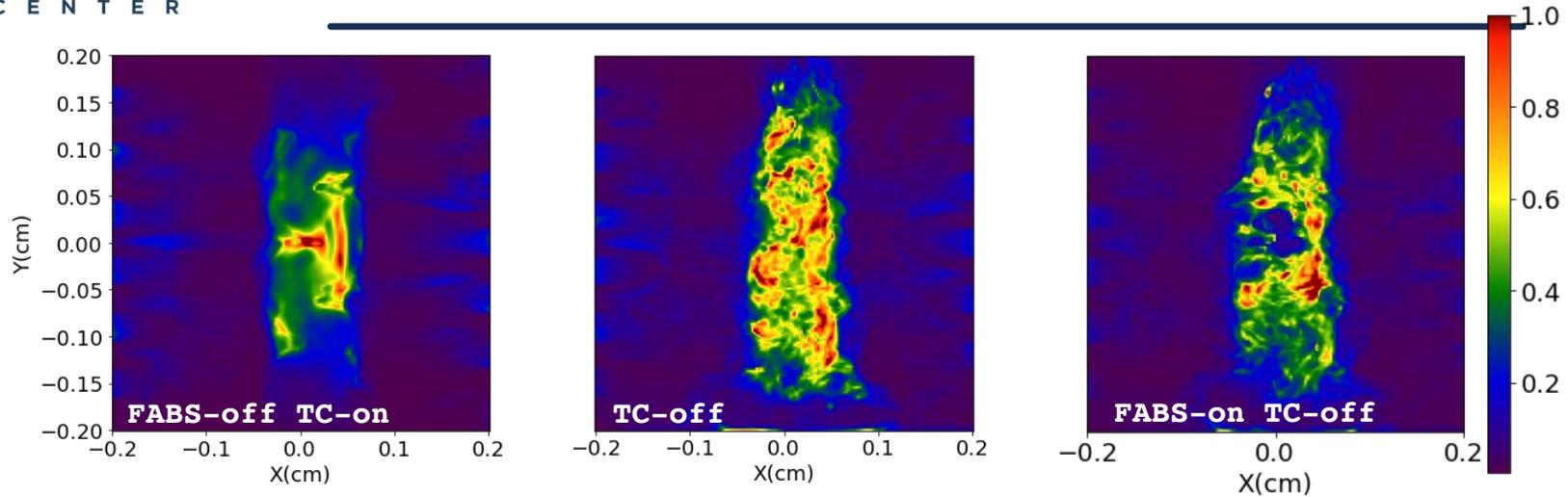
- ❑ For conduction off run:
 - ❑ The Interaction region at 23ns reaches
 - $n_e \sim 4 \times 10^{20} \text{ cm}^{-3}$, $T_e \sim 1.2 \times 10^3 \text{ eV}$
 - $R_m \sim 3.9 \times 10^3$, $P_m \sim 20$
 - ❑ We also see that high R_m and above-unity P_m condition last longer
 - ❑ Electron temperature profile is highly structured due to the suppression of electron thermal conduction

Electron temperature in conduction-off and FABS-on cases

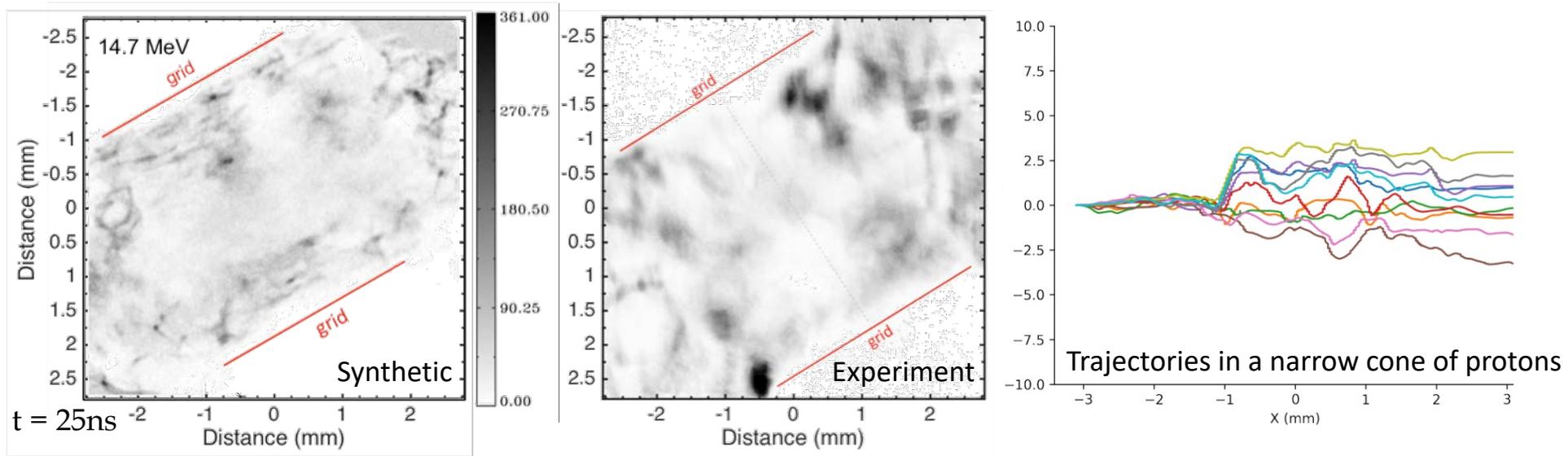
Electron Temperature (eV)



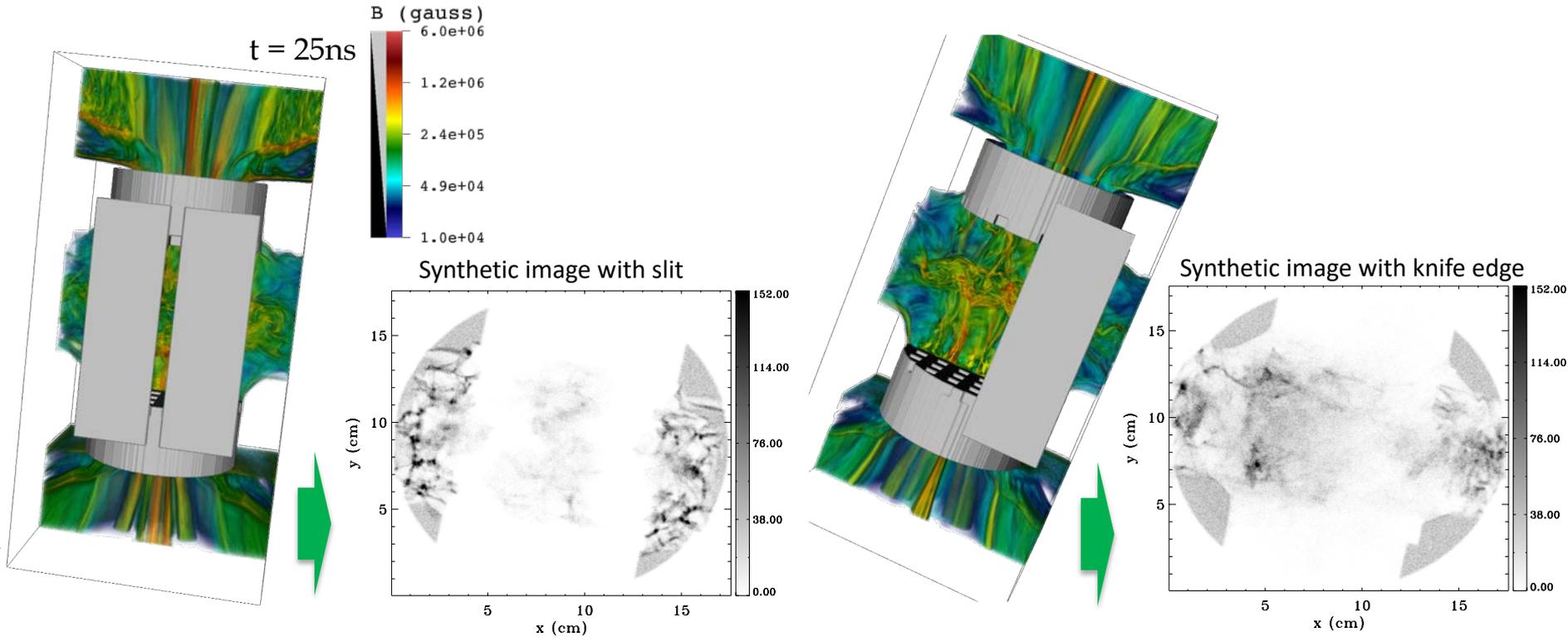
- ❑ Full-aperture backscatter system (FABS) being used to measure the turbulent velocity in the plasma can heat the plasma and affect the physical process. We compute the heating from FABS.
- ❑ In TC-on (conduction-on) case, we find moderate increase of temperature to $T_e \sim 750$ eV due to FABS.
- ❑ In TC-off, We have $T_e \sim 1.2 \times 10^3$ eV. FABS does not have a significant effect because we have thermal conduction turned off, and any heating is localized in the plane of FABS.



- ❑ For TC-off, electron temperature profile is highly structured for the conduction-off runs, resulting in highly structured X-ray self-emission intensity.
- ❑ Synthetic X-ray images enable quantitative comparison with the experimental data.
- ❑ The suppression of electron thermal conduction has been demonstrated in the experiment (Meinecke+ arXiv:2105.08461, also see the talk by Petros Tzeferacos in this session)



- ❑ Both the synthetic and experimental proton images are in the diffusive regime (Bott+ 2017) characterized as following:
 - ❑ Synthetic proton trajectories from a narrow cone source have crossings before they leave the plasma
 - ❑ The correlation length of the magnetic field is $l_B = 0.0036\text{cm}$, which is less than the RMS perpendicular displacement of protons going through the interaction regime plasma due to magnetic deflections $l_z \delta\theta = 0.014\text{cm}$.



- Using the image with slit or knife edge, we can measure the velocity diffusion coefficient and infer the range of RMS magnetic field strength.

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