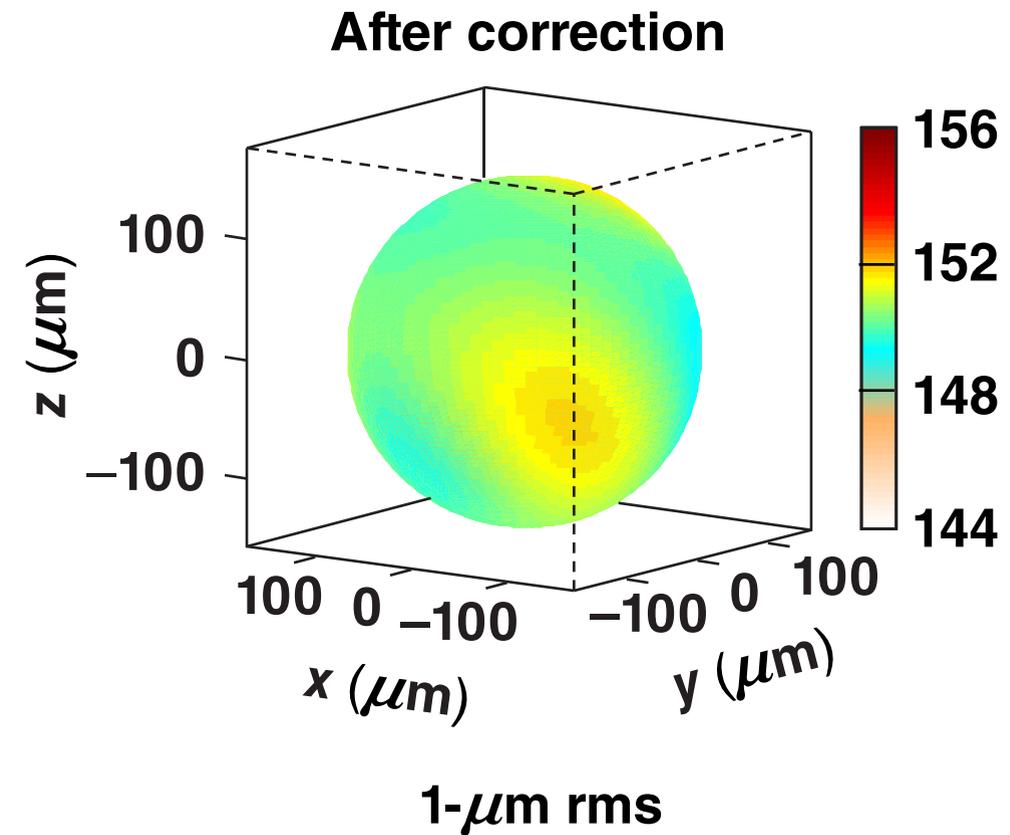
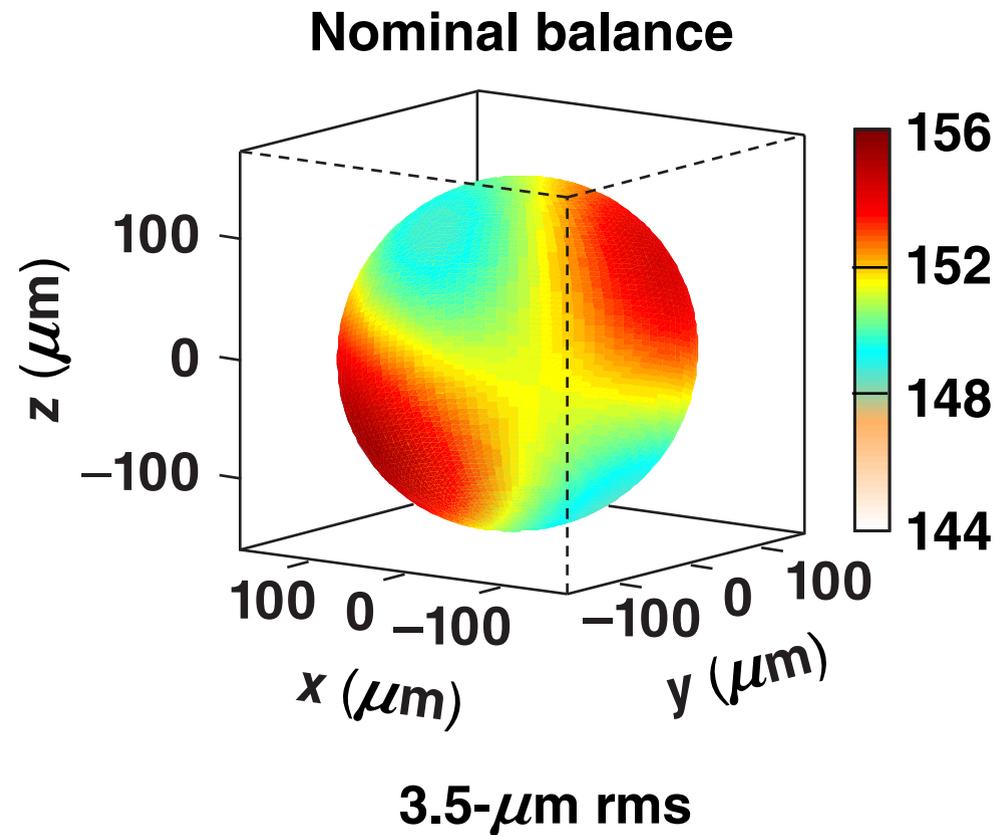


Subpercent Scale Control of 3-D Modes 1, 2, and 3 of Targets Imploded in Direct-Drive Configuration on OMEGA



D. T. Michel
University of Rochester
Laboratory for Laser Energetics

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Improved drive symmetry has been demonstrated on OMEGA

- In a series of direct-drive implosions, multiple self-emission x-ray images were used to tomographically measure their 3-D modes 1, 2, and 3 at a convergence ratio of ~ 3
- The target modes were shown to vary linearly with the laser modes from approximately constant static modes
- This demonstrated that the target modes can be mitigated by adjusting the laser beam-energy balance to compensate the static modes

This method was applied to low-adiabat shots and made it possible to reduce the low-mode nonuniformities from $3.5 \mu\text{m}$ to $1 \mu\text{m}$.

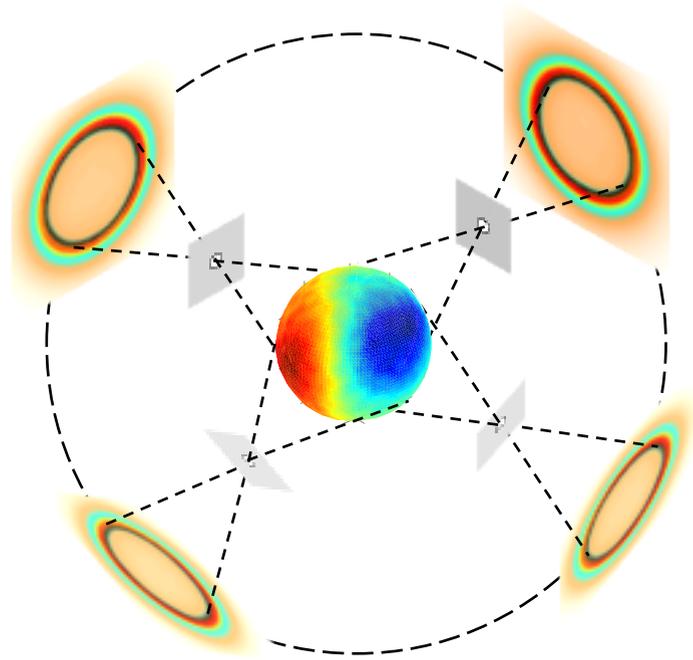
Collaborators



**I. V. Igumenshchev, A. K. Davis, D. H. Edgell,
D. H. Froula, D. W. Jacobs-Perkins, V. N. Goncharov,
S. P. Regan, R. Shah, A. Shvydky,
and E. M. Campbell**

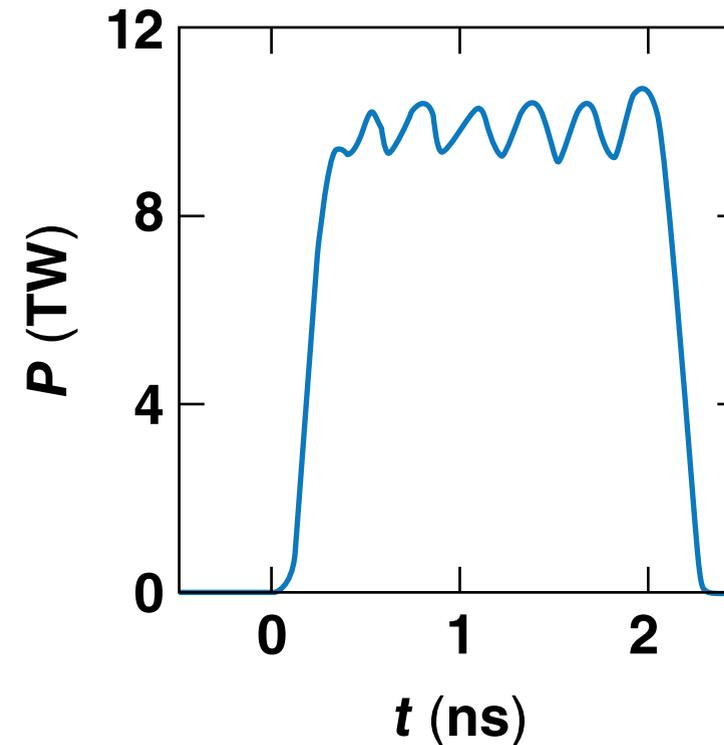
**University of Rochester
Laboratory for Laser Energetics**

Self-emission shadowgraphy* from multiple lines of sight was used to tomographically measure the 3-D modes $\ell = 1$, $\ell = 2$, and $\ell = 3$ of targets imploded on OMEGA



Synchronized observation $\langle R \rangle = 150 \mu\text{m}$

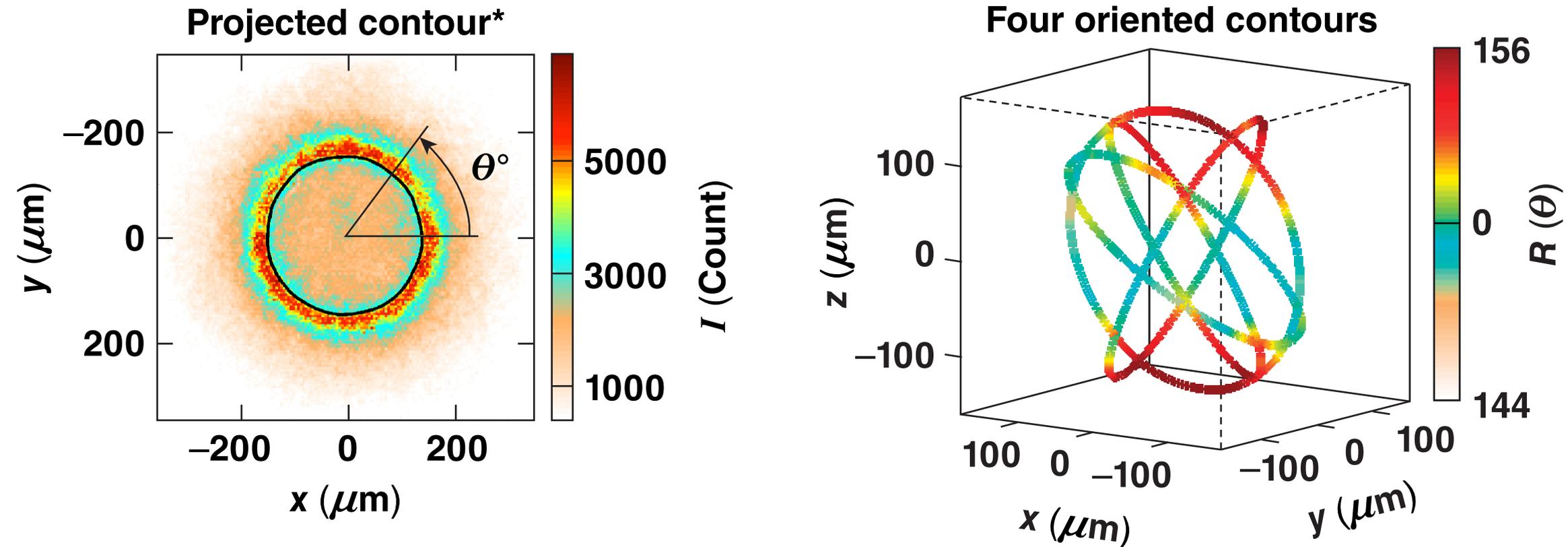
Each camera measures the projected ablation-front surface along the line of sight of the diagnostic.



$$R(\theta, \phi) = \sum_{\ell=0}^3 \sum_{m=-\ell}^{\ell} \sqrt{4\pi} r_{\ell}^m Y_{\ell}^m(\theta, \phi) **$$

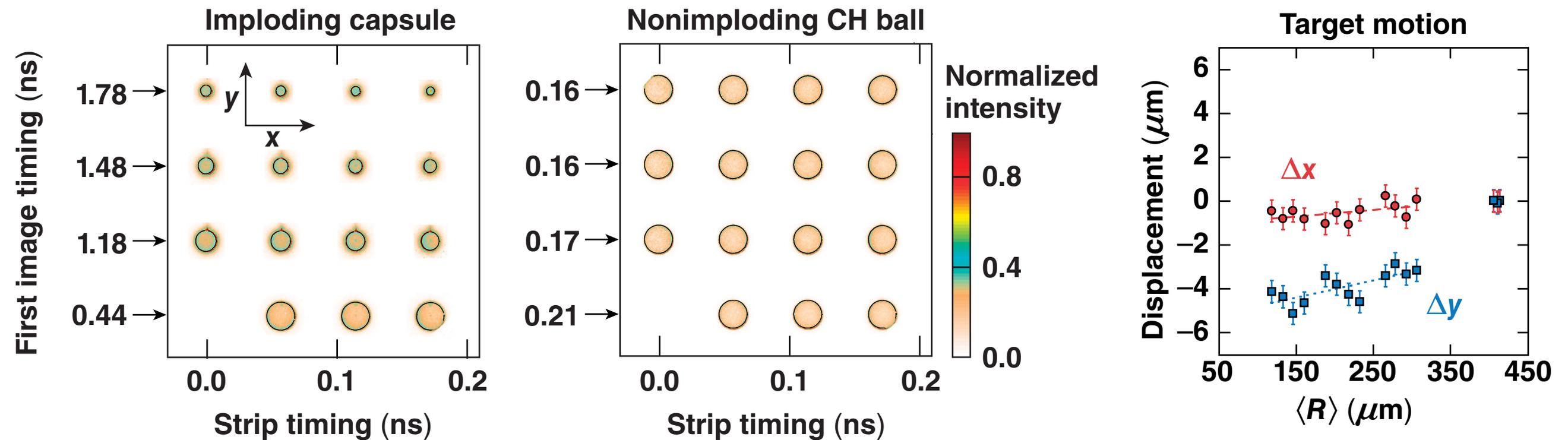
*D. T. Michel *et al.*, Rev. Sci. Instrum. **83**, 10E530 (2012); D. T. Michel *et al.*, High Power Laser Sci. Eng. **3**, e19 (2015).
 ** Y_{ℓ}^m are the tesseral spherical harmonics, E. T. Whittaker and G. N. Watson, *A Course of Modern Analysis* (Cambridge University Press, 1927), p. 392; $R(\theta, \phi)$ is normalized in percent ($r_0^0 = 100\%$).

On each camera, the angular variation of the projected ablation contour $R(\theta)$ was determined for an averaged radius of $150 \mu\text{m}$



The 3-D shape of the target was obtained by orienting the four contours perpendicular to the camera axis.

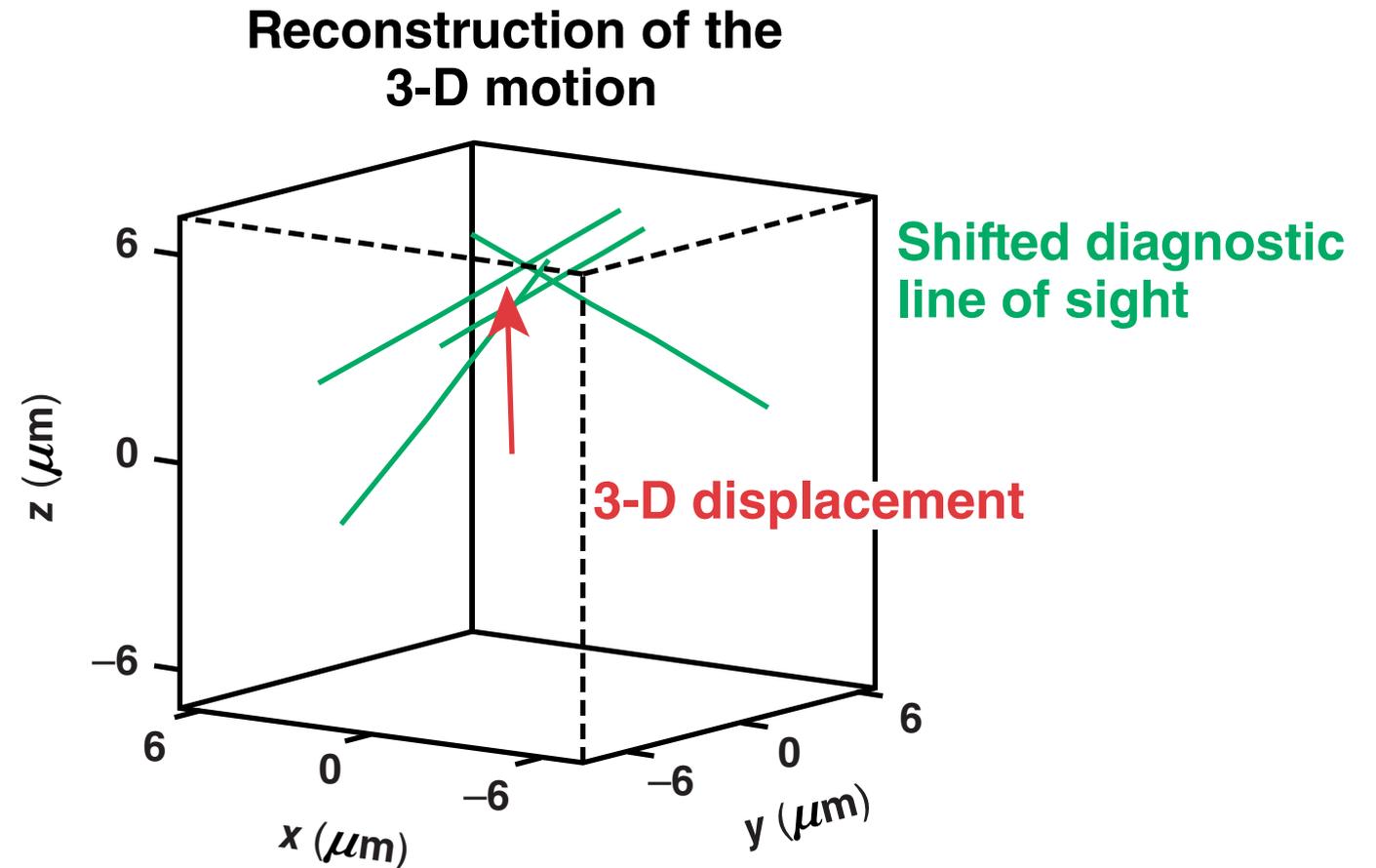
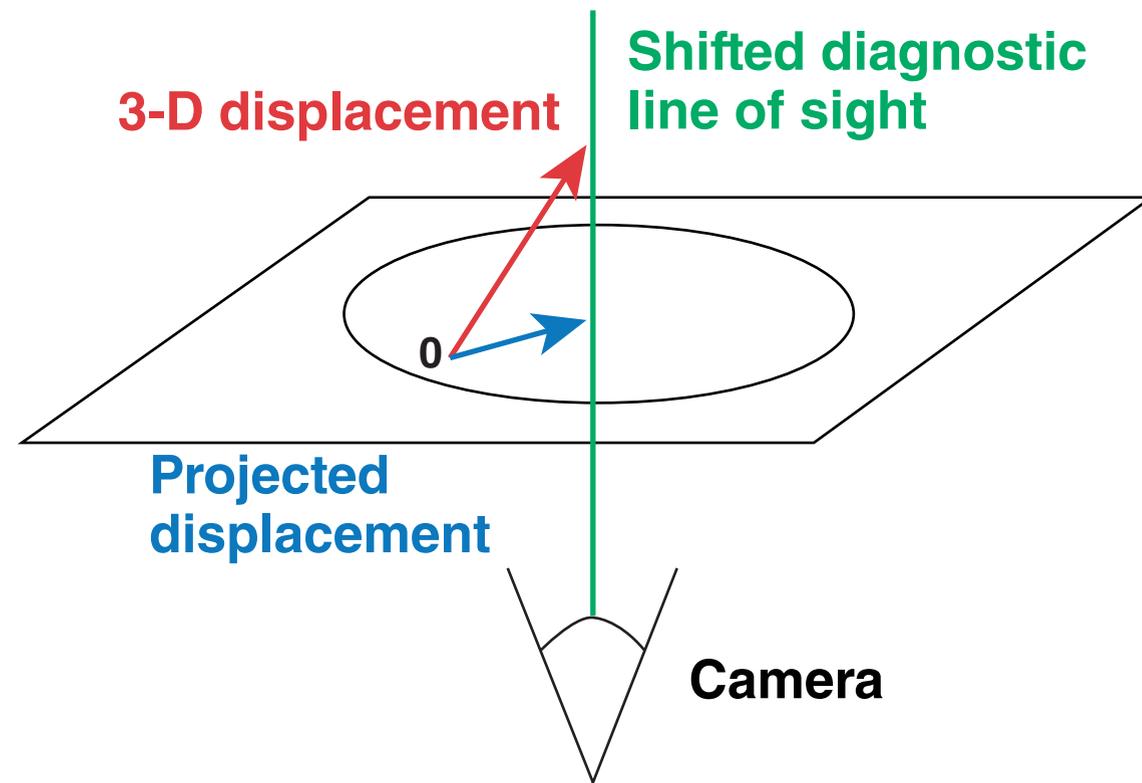
The target motion was obtained by comparing the positions of the contours centers with the corresponding contour centers measured on a nonimploding solid CH ball shot



The target motion at 150 μm was obtained using linear fits.*

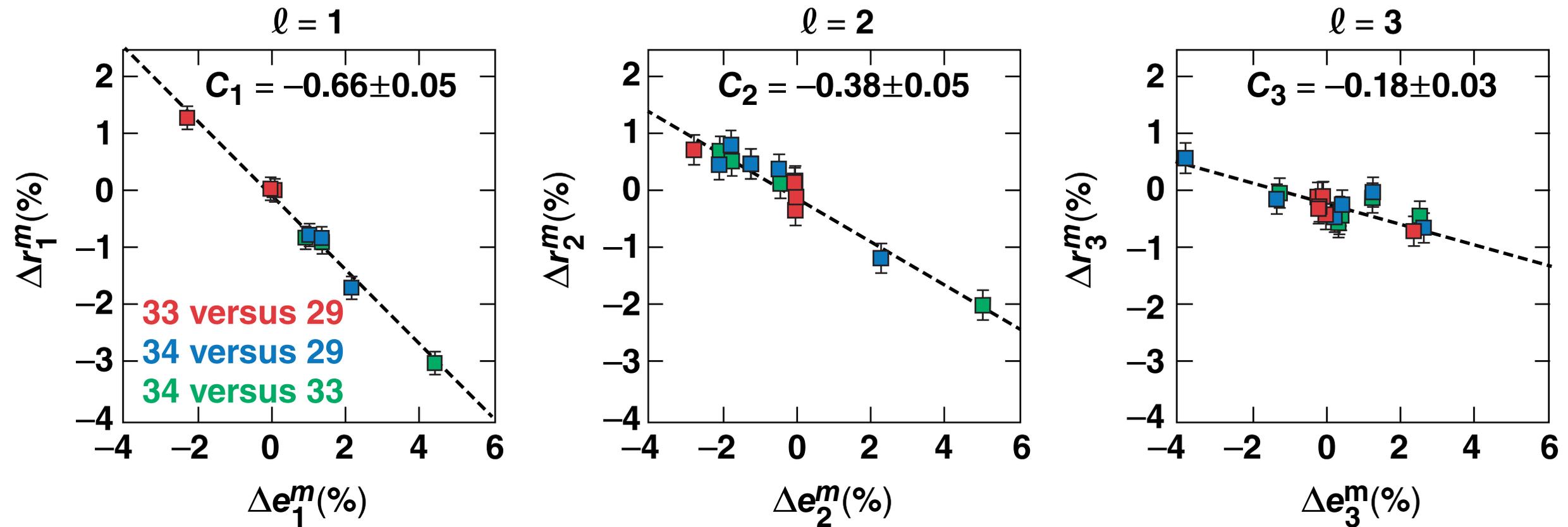
* $\delta(\Delta R_{\text{center}}) = \pm 1.3 \mu\text{m}$, resulting, in $\delta[(\Delta R_{\text{center}})_{150}] = \pm 0.6 \mu\text{m}$ at the 90th percentile of the student's t distribution

The 3-D target displacement is located at the intersection of the four lines defined by the camera axis, translated by the measured projected target motions



For each mode ℓ , a linear evolution of the target modes (Δr_ℓ^m) with the laser beam-energy balance (Δe_ℓ^m) was measured

$$C_\ell = \Delta r_\ell^m (150 \mu\text{m}) / \Delta e_\ell^m$$



The reduction of C_ℓ with ℓ was explained by the individual beam shape that modifies the laser illumination nonuniformity compared to the laser beam-energy balance.

*The laser modes are obtained by minimizing $\sum_{\ell=0}^3 \sum_{m=-1}^{\ell} \sqrt{4\pi} e_\ell^m \gamma_\ell^m(\theta_b, \phi_b) - E_b$, where E_b is the beam energy normalized to the averaged beam energy.

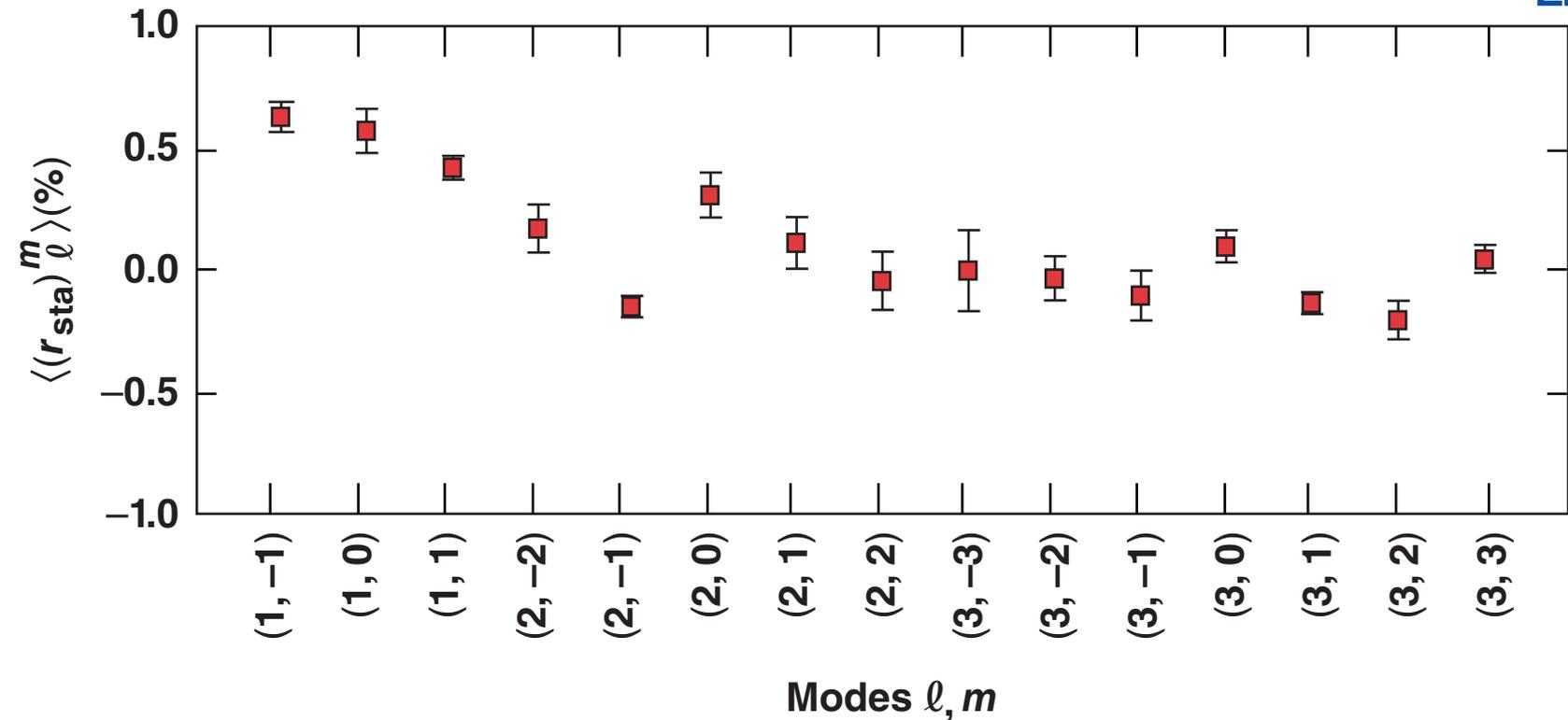
The linear variations show that the target modes are the result of two components: a part that varies linearly with the laser modes and a static part

The static modes are given by

$$\langle (r_{\text{sta}})_\ell^m \rangle_{\text{shots}} = \langle r_\ell^m - \mathbf{C}_\ell \mathbf{e}_\ell^m \rangle_{\text{shots}}$$

The optimized beam energy balance is given by

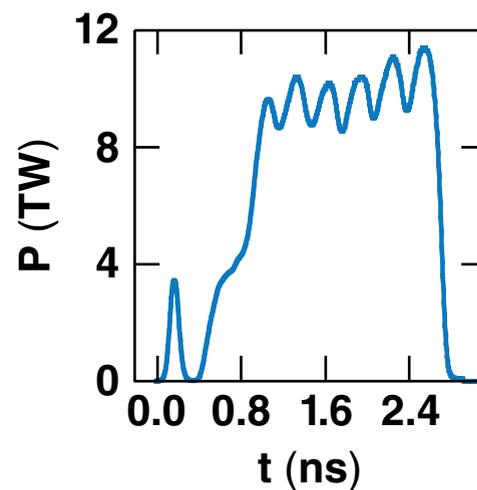
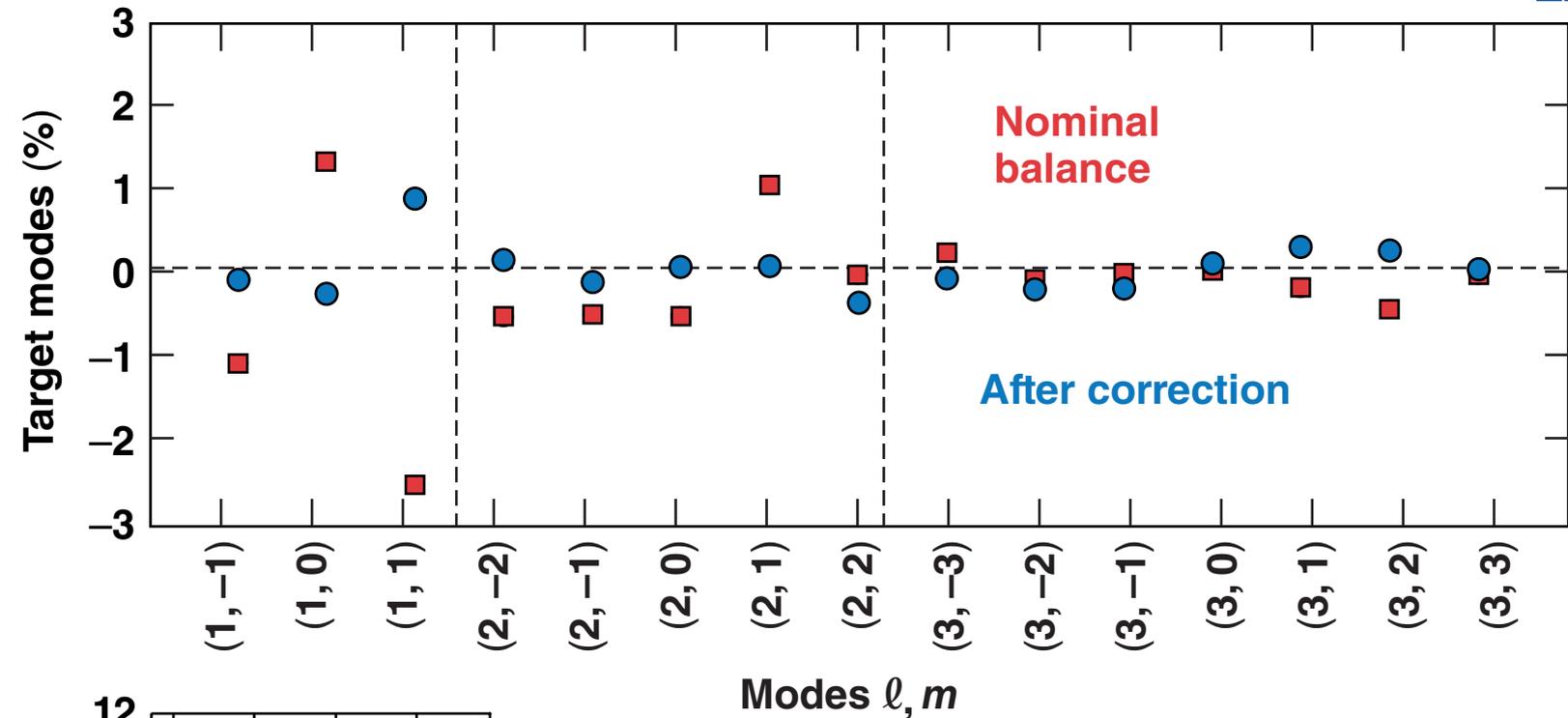
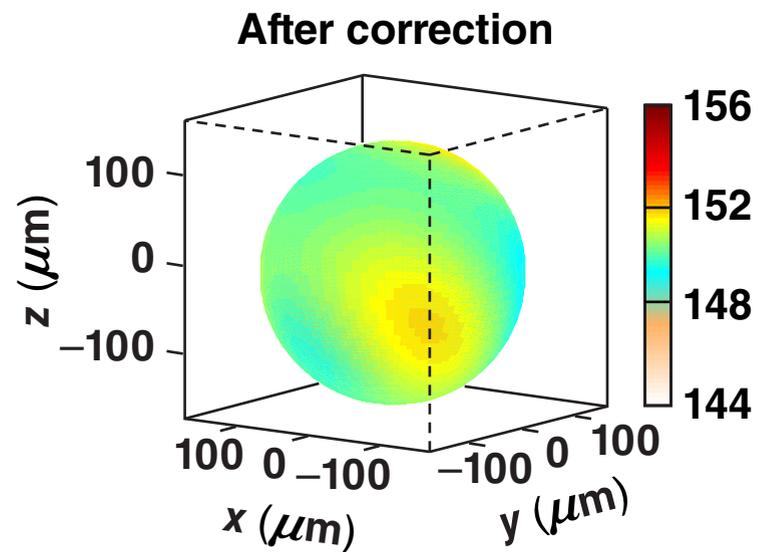
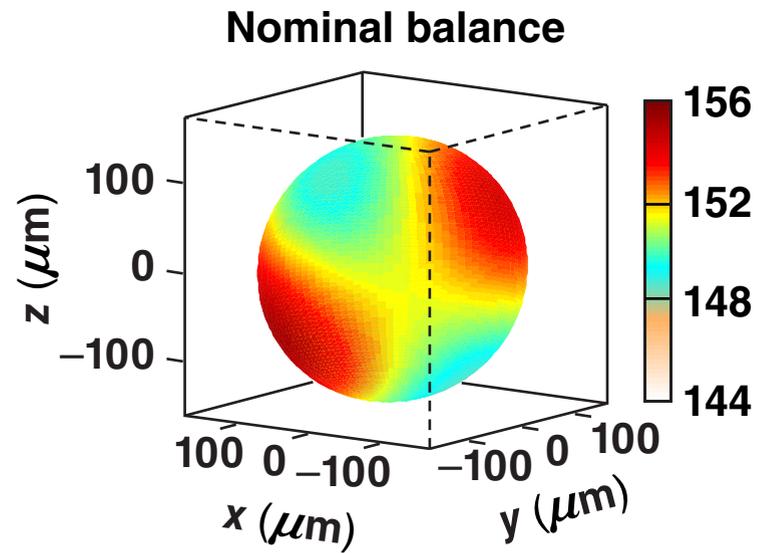
$$(\mathbf{e}_{\text{opt}})_\ell^m = - \frac{\langle (r_{\text{sta}})_\ell^m \rangle_{\text{shots}}}{\mathbf{C}_\ell}$$



This demonstrates that the target modes can be mitigated by adjusting the laser modes to compensate the static modes.

D. T. Michel *et al.*, "Subpercent-Scale Control of 3-D Modes 1, 2, and 3 in Direct-Drive Configuration on OMEGA," submitted to Physical Review Letters.

This method was successfully applied to mitigate the target nonuniformities on a low-adiabat warm implosion



The correction made it possible to reduce the standard deviation of the target modes 1, 2, and 3 from $\sim 3.5 \mu\text{m}$ to $\sim 1 \mu\text{m}$.

Improved drive symmetry has been demonstrated on OMEGA

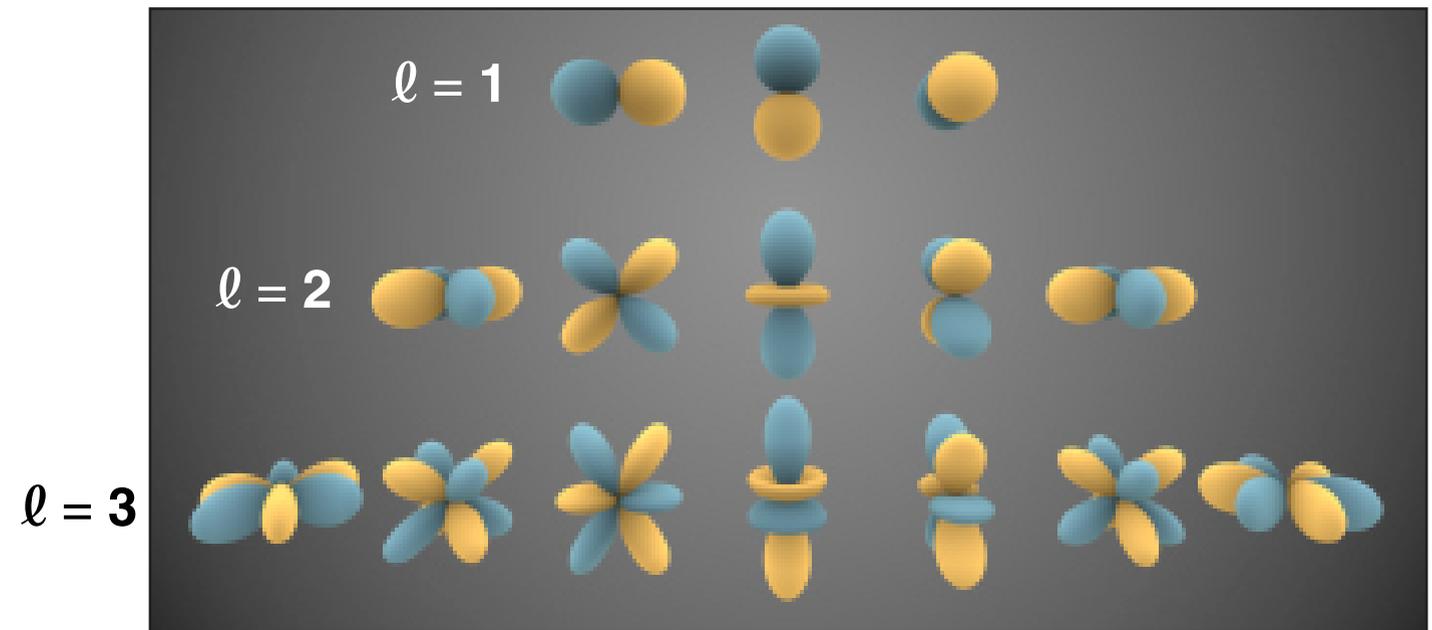
- In a series of direct-drive implosions, multiple self-emission x-ray images were used to tomographically measure their 3-D modes 1, 2, and 3 at a convergence ratio of ~ 3
- The target modes were shown to vary linearly with the laser modes from approximately constant static modes
- This demonstrated that the target modes can be mitigated by adjusting the laser beam-energy balance to compensate the static modes

This method was applied to low-adiabat shots and made it possible to reduce the low-mode nonuniformities from $3.5 \mu\text{m}$ to $1 \mu\text{m}$.

Over three shots, the beam-energy balance was changed to modify their modes $\ell = 1$, $\ell = 2$, $\ell = 3$, and for $m = 0$

$$\bar{E}(\theta_b, \phi_b) = \sum_{\ell=0}^3 \sum_{m=-1}^{\ell} \sqrt{4\pi} e_{\ell}^m Y_{\ell}^m(\theta_b, \phi_b)^*$$

Shots	Δ	Δ	Δ
84633 versus 84629	-2.2	-2.6	2.4
84634 versus 84629	2.2	2.2	-3.5



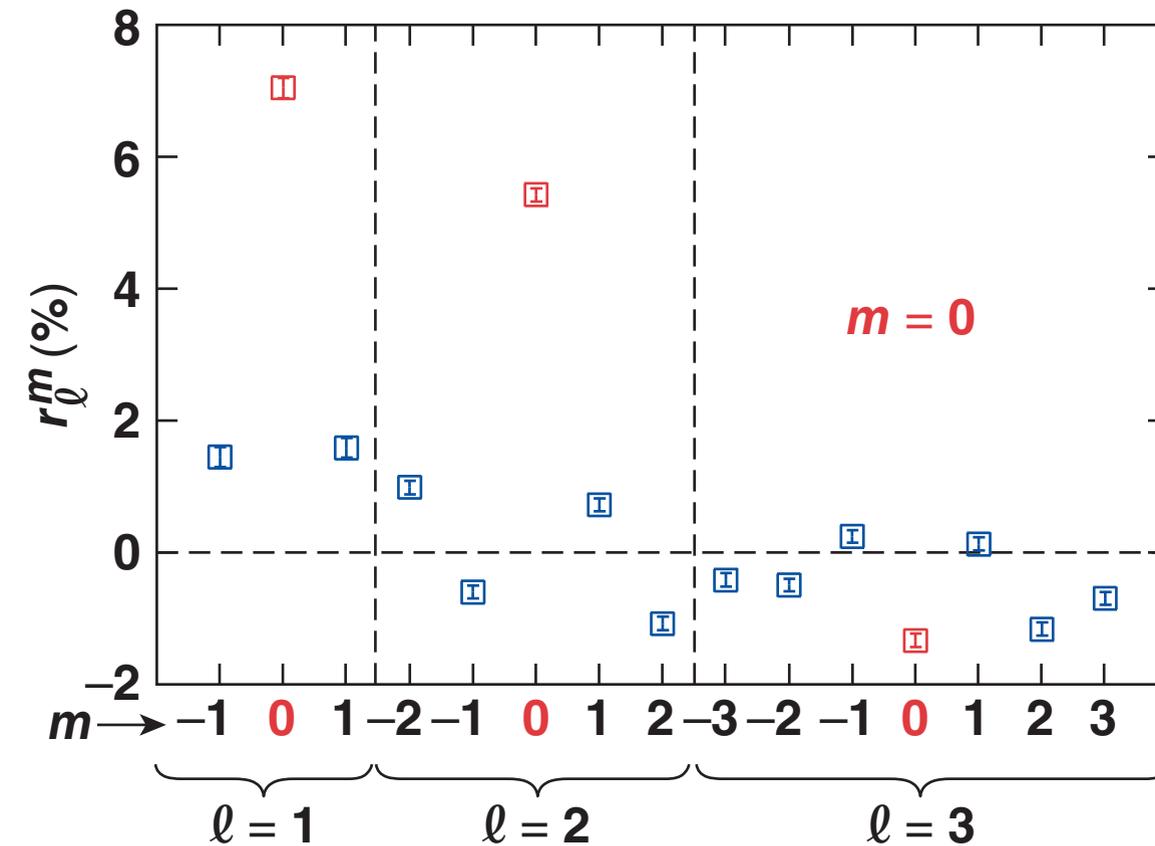
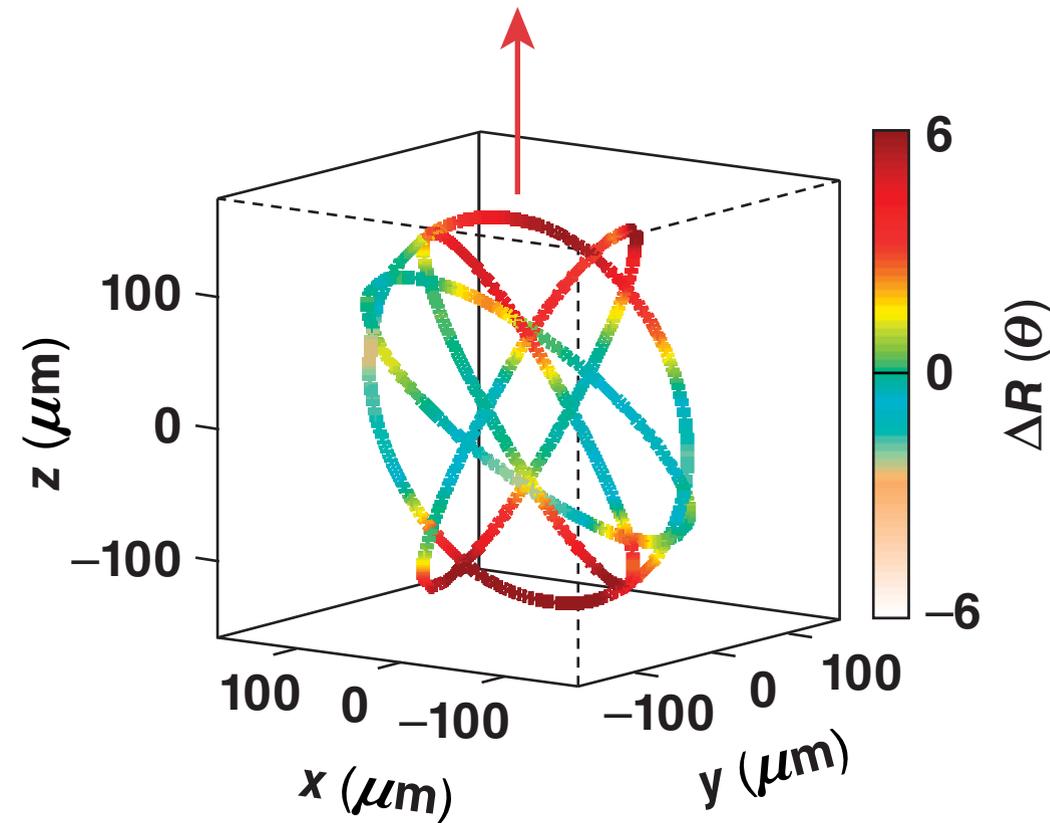
$$C_{\ell} = \Delta r_{\ell}^m (150 \mu m) / \Delta e_{\ell}^m$$

The low-mode coupling coefficients were determined by measuring the variation of the target modes (Δr_{ℓ}^m) as a function of the laser beam-energy balance (Δe_{ℓ}^m).

**The modes are obtained by minimizing $\sum_{\ell=0}^3 \sum_{m=-1}^{\ell} \sqrt{4\pi} e_{\ell}^m Y_{\ell}^m(\theta_b, \phi_b) - \bar{E}_b$, where \bar{E}_b is normalized to the averaged beam energy.

The target modes were obtained by decomposing the four contours translated by the target displacement over spherical harmonics

$$\bar{R}(\theta_c, \phi_c) = \sum_{\ell=0}^3 \sum_{m=0}^3 r_{\ell}^m Y_{\ell}^m(\theta_c, \phi_c)$$



Errors of $\delta(r_1^m) = \pm 0.15\%$, $\delta(r_2^m) = \pm 0.1\%$, and $\delta(r_3^m) = \pm 0.1\%$ were obtained by simulating the errors in $[\Delta R(\theta)]_{150}$ and $[\Delta R_{\text{center}}]_{150}$.

The decrease of C_ℓ with ℓ was a result of the beam profiles that modify the amplitude of the laser modes on target*

- The laser modes are described by minimizing ($\partial A / \partial \mathbf{e}_1^m = 0$):
$$A = \sum_{b=1}^{60} \left[\sum_{\ell=0}^{\infty} \sum_{m=-\ell}^{\ell'} \sqrt{4\pi} \mathbf{e}_\ell^m Y_\ell^m(\theta_b, \phi_b) - \bar{E}_b \right]^2$$

- This results in

$$\begin{aligned} \sum_{b=1}^{60} [\bar{E}_b Y_\ell^m(\theta_b, \phi_b)] &= \sum_{\ell'=0}^{\infty} \sum_{m'=-\ell'}^{\ell'} \sqrt{4\pi} \mathbf{e}_{\ell'}^m \sum_{b=1}^{60} [Y_{\ell'}^m(\theta_b, \phi_b) Y_{\ell'}^{m'}(\theta_b, \phi_b)] \\ &\approx \sum_{\ell'=0}^{\infty} \sum_{m'=-\ell'}^{\ell'} \sqrt{4\pi} \mathbf{e}_{\ell'}^m (60/4\pi) \int_{\Omega} [Y_{\ell'}^m(\theta_b, \phi_b) Y_{\ell'}^{m'}(\theta_b, \phi_b)] \approx \sum_{\ell'=0}^{\infty} \sum_{m'=-\ell'}^{\ell'} \sqrt{4\pi} \mathbf{e}_{\ell'}^m (60/4\pi) \delta_{\ell\ell'} \delta_{mm'} \approx \sqrt{4\pi} \mathbf{e}_\ell^m (60/4\pi) \end{aligned}$$

- The mode decomposition of a single beam energy per solid angle on target is given by

$$\tilde{E}_b(\theta, \phi) = \bar{E}_b \sum_{\ell=0}^{\infty} \frac{2\ell+1}{4\pi} a_\ell P_\ell(\cos \gamma) = \bar{E}_b \sum_{\ell=0}^{\infty} a_\ell \sum_{m'=-\ell}^{\ell} Y_\ell^m(\theta_b, \phi_b) Y_\ell^m(\theta, \phi)$$

- The mode decomposition of the total energy per solid angle on target is given by

$$\tilde{E}_{\text{tot}}(\theta, \phi) = \frac{60}{4\pi} \sum_{b=1}^{60} \tilde{E}_b(\theta, \phi) = \sum_{\ell=0}^{\infty} \sum_{m=-\ell}^{\ell} \sqrt{4\pi} \tilde{\mathbf{e}}_\ell^m Y_\ell^m(\theta, \phi) \text{ with: } \tilde{\mathbf{e}}_\ell^m = \frac{a_\ell}{\sqrt{4\pi}} \frac{4\pi}{60} \sum_{b=1}^{60} \bar{E}_b Y_\ell^m(\theta_b, \phi_b) = a_\ell \mathbf{e}_\ell^m,$$

where $a_\ell = 2\pi \int_{-1}^1 \bar{E}_b(\theta, \phi) P_\ell(\cos \gamma) d(\cos \gamma)$ are normalized to have $\mathbf{e}_0^0 = 100\%$

The decrease of C_ℓ with ℓ was a result of the beam profiles that modify the amplitude of the laser modes on target*

- The modes of the laser beam energy balance are described by minimizing

$$A = \sum_{b=1}^{60} \left[\sum_{\ell=0}^3 \sum_{m=-1}^{\ell} \sqrt{4\pi} e_\ell^m Y_\ell^m(\theta_b, \phi_b) - \bar{E}_b \right]^2$$

- Accounting for the beam profile, the mode decomposition of the total energy per solid angle on target is given by

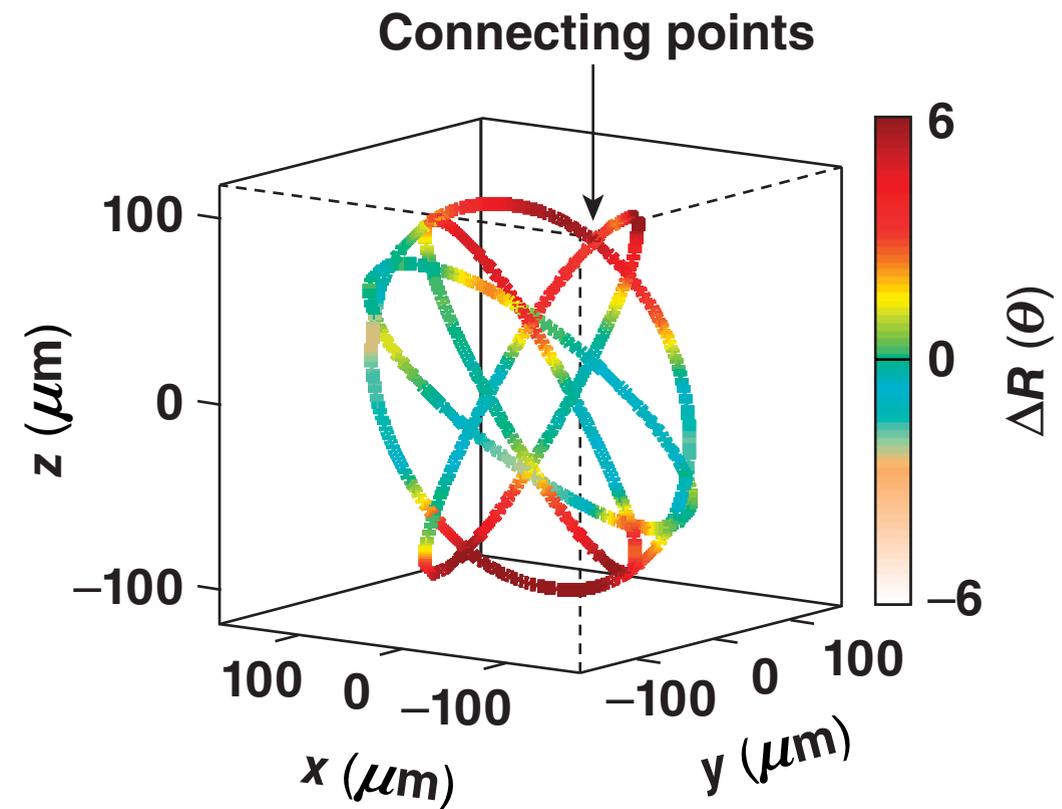
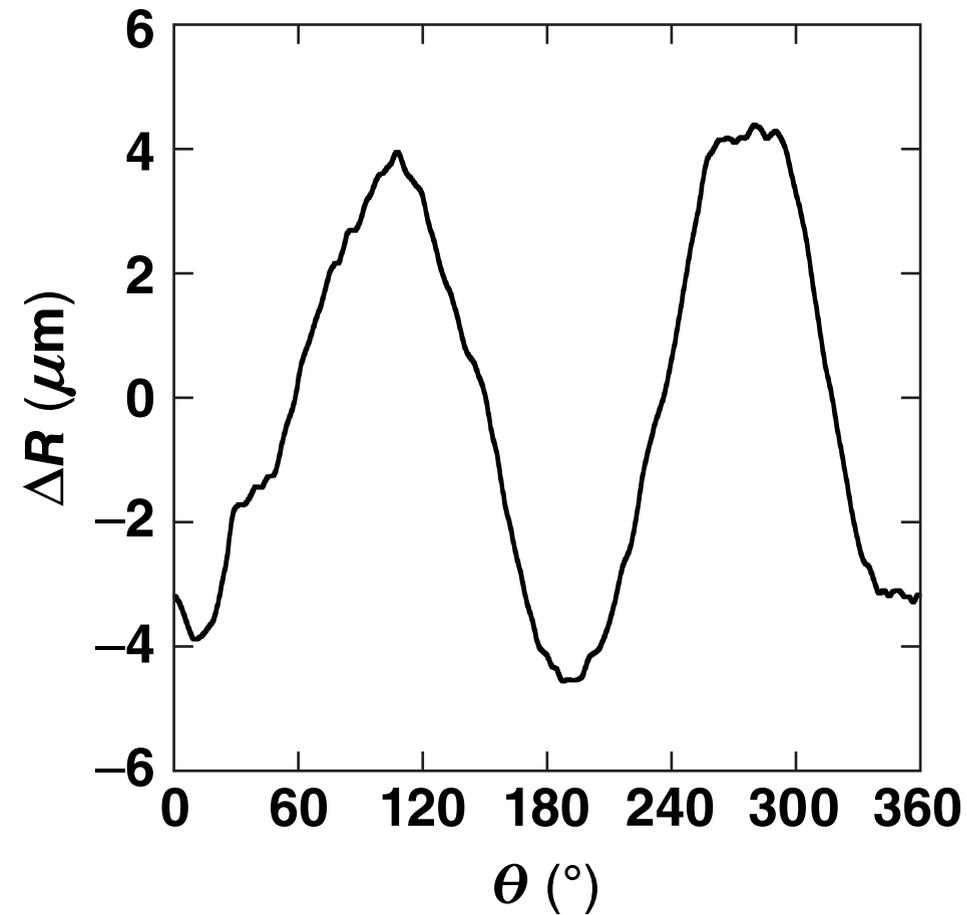
$$E_{\text{tot}}(\theta, \phi) = \sum_{\ell=0}^3 \sum_{m=-1}^{\ell} \sqrt{4\pi} e_\ell^m Y_\ell^m(\theta, \phi)$$

with $\tilde{e}_\ell^m = a_\ell$, where $a_\ell = 2\pi \int_{-1}^1 \bar{E}_b(\theta, \phi) P_\ell(\cos \gamma) d(\cos \gamma)$

a_1	0.79
a_2	0.47
a_3	0.20

A constant coupling of the modes of the target irradiation pattern to the target modes is obtained of $C_\ell/a_\ell = -0.85 \pm 0.07$.

The 3-D shape of the target was obtained by orienting each contour perpendicular to the camera axis



An error of $\delta[\Delta R(\theta)]_{150}$ of 1 μm was evaluated by comparing the contours at the connecting points; this error is comparable to the error of $\pm 0.4 \mu\text{m}$ estimated previously.

This method was applied to correct the target nonuniformities on a low-adiabat warm implosion

