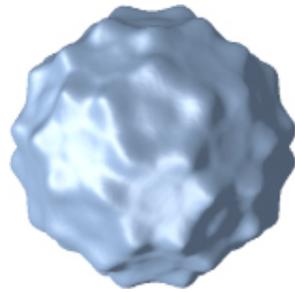
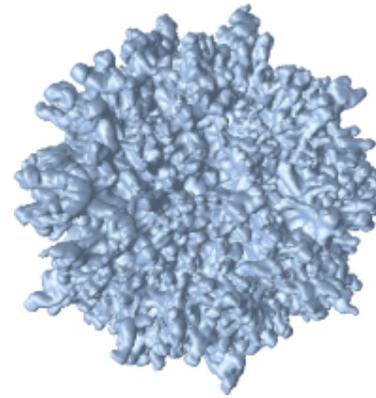


# Three-Dimensional Hydrodynamic Simulations of the Effects of Laser Imprint in OMEGA Implosions

**Fuel–ablator interface in room-temperature  
shot 84629 at neutron peak**

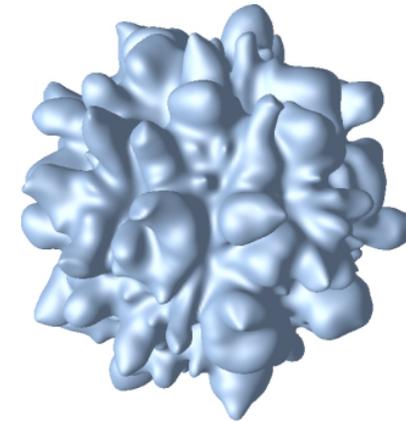


**No imprint**



**Imprint (SSD on)**

**Hot-spot shape (surface  $T_i = 1$  keV)  
in cryogenic shot 77066 at neutron peak**



**Imprint (SSD on)**

**I. V. Igumenshchev  
University of Rochester  
Laboratory for Laser Energetics**

**59th Annual Meeting of the  
American Physical Society  
Division of Plasma Physics  
Milwaukee, WI  
23–27 October 2017**

## Summary

# Simulations indicate that the effects of laser imprint alone are insufficient to explain the underperformance of cryogenic $\alpha \sim 4$ implosions on OMEGA



- The imprint model was developed and implemented in the 3-D hydrodynamics code *ASTER*\*
- Simulations reproduce observed improvement in implosion performance when polarization smoothing\*\* (PS) and smoothing by spectral dispersion† in two dimensions (2-D SSD) are applied
- Room-temperature targets suffer from imprint that introduces significant small-scale ( $\ell \sim 50$  to 150) modulations
- Imprint in cryogenic implosions develops broadband modulations with dominant  $\ell \sim 30$ ; these modulations have a moderate effect on the implosions

\*I. V. Igumenshchev *et al.*, Phys. Plasmas **23**, 052702 (2016).

\*\*T. R. Boehly *et al.*, J. Appl. Phys. **85**, 3444 (1999).

†S. Skupsky *et al.*, J. Appl. Phys. **66**, 3456 (1989).

# Collaborators

---



**E. M. Campbell, V. N. Goncharov, S. P. Regan, and A. Shvydky**

**University of Rochester  
Laboratory for Laser Energetics**

**A. J. Schmitt**

**Naval Research Laboratory**

# ASTER models imprint in OMEGA implosions by calculating far-field intensity modulations

- Evolution of the electric near field

$$E(x, y, t) = E_0(x, y) e^{i[\psi_{\text{DPP}}(x, y) + \psi_{\text{SSD}}(x, y, t)]}$$

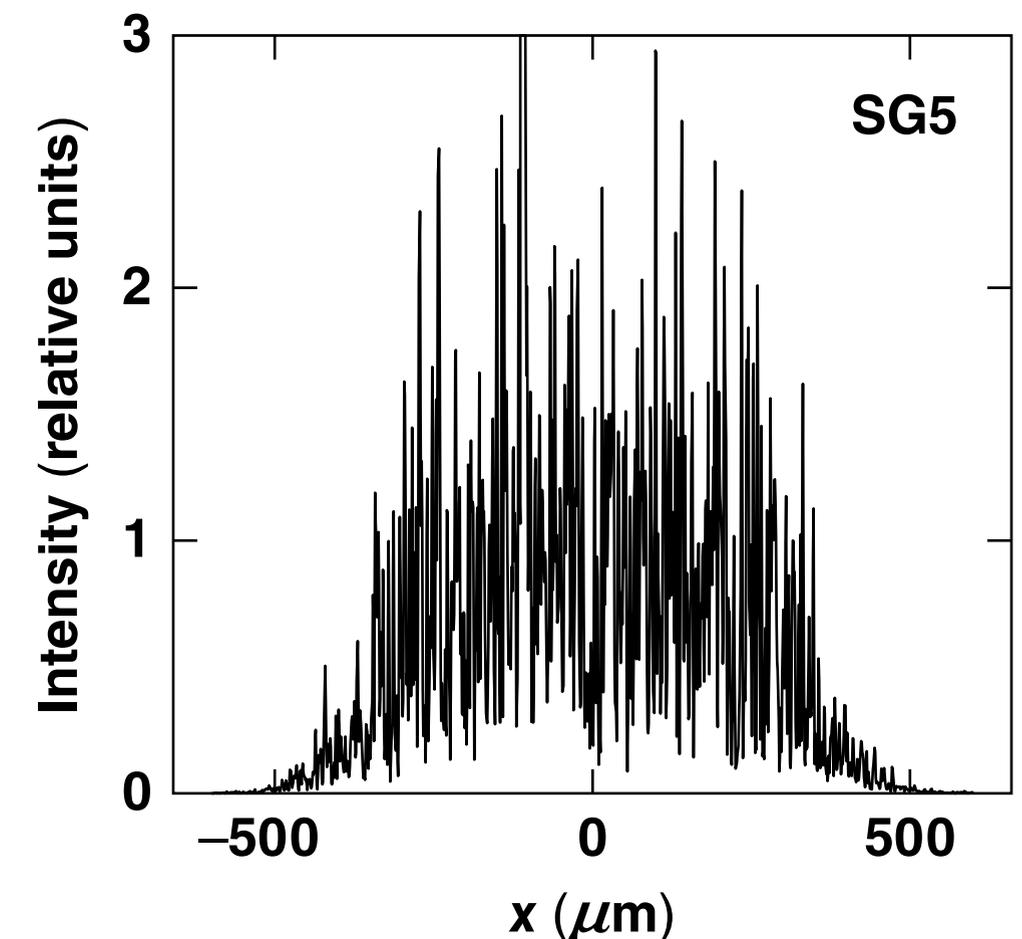
$$\psi_{\text{SSD}}(x, y, t) = \delta_x \sin(k_x x - 2\pi f_x t) + \delta_y \sin(k_y y - 2\pi f_y t)$$

- Evolution of the far-field (focal-spot) intensity

$$I_{\text{ff}}(x_{\text{ff}}, y_{\text{ff}}, t) \propto \left| \iint_{(x, y)} E(x, y, t) \cdot \exp\left[-i \frac{2\pi}{\lambda F} (x_{\text{ff}} \cdot x + y_{\text{ff}} \cdot y)\right] dx dy \right|^2$$

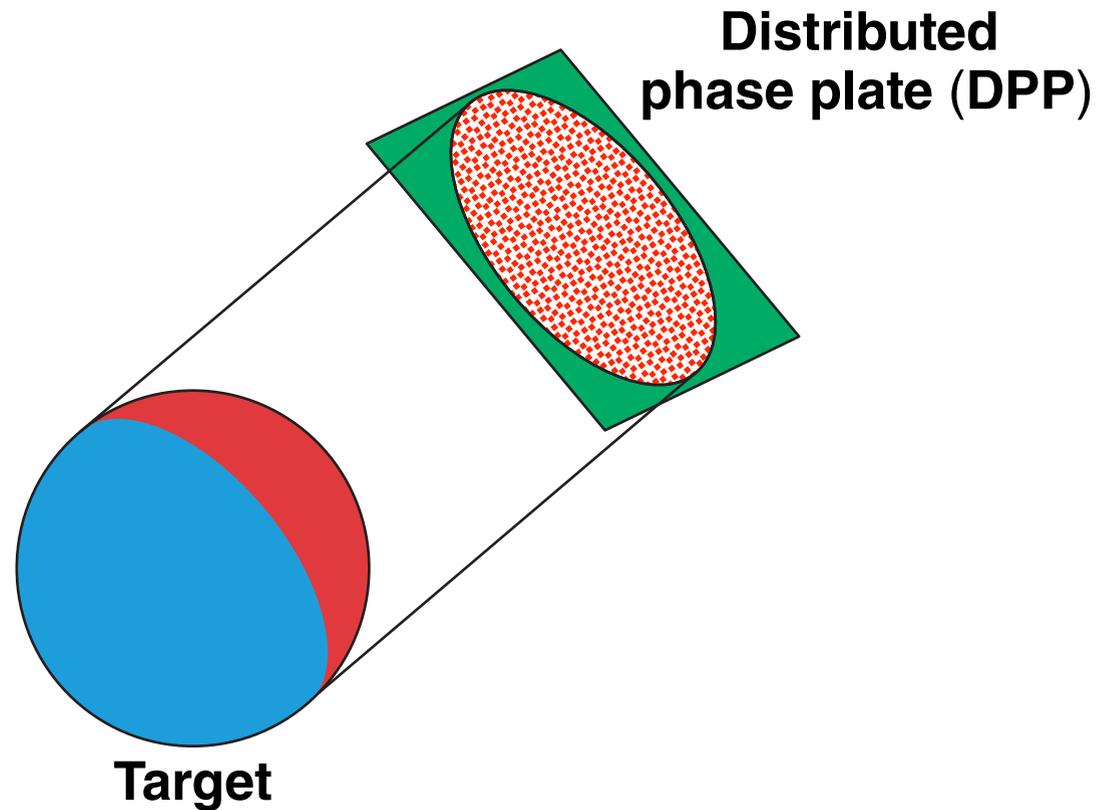
- Polarization smoothing: overlapping two copies of the intensity pattern separated by  $\Delta y = 86.4 \mu\text{m}$

Example beam-intensity profile

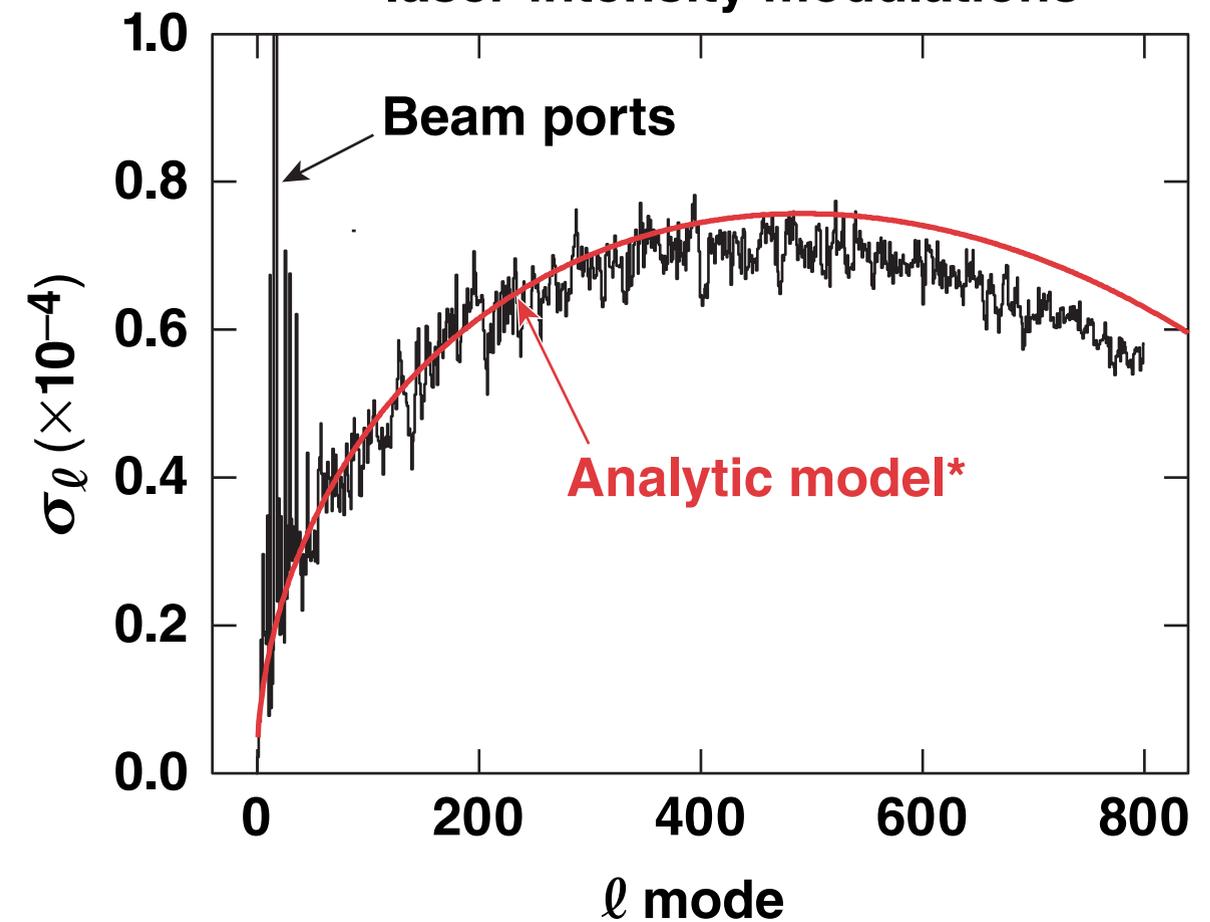


# ASTER models imprint in OMEGA implosions by calculating far-field intensity modulations (continued)

Calculated imprint patterns for all 60 OMEGA beams are projected onto a target

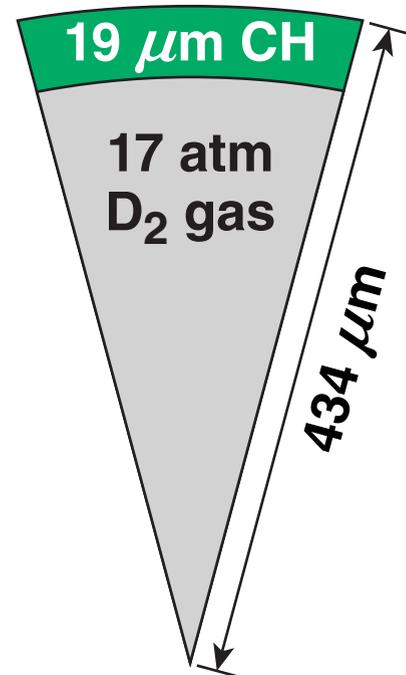


Spectrum of on-target laser-intensity modulations

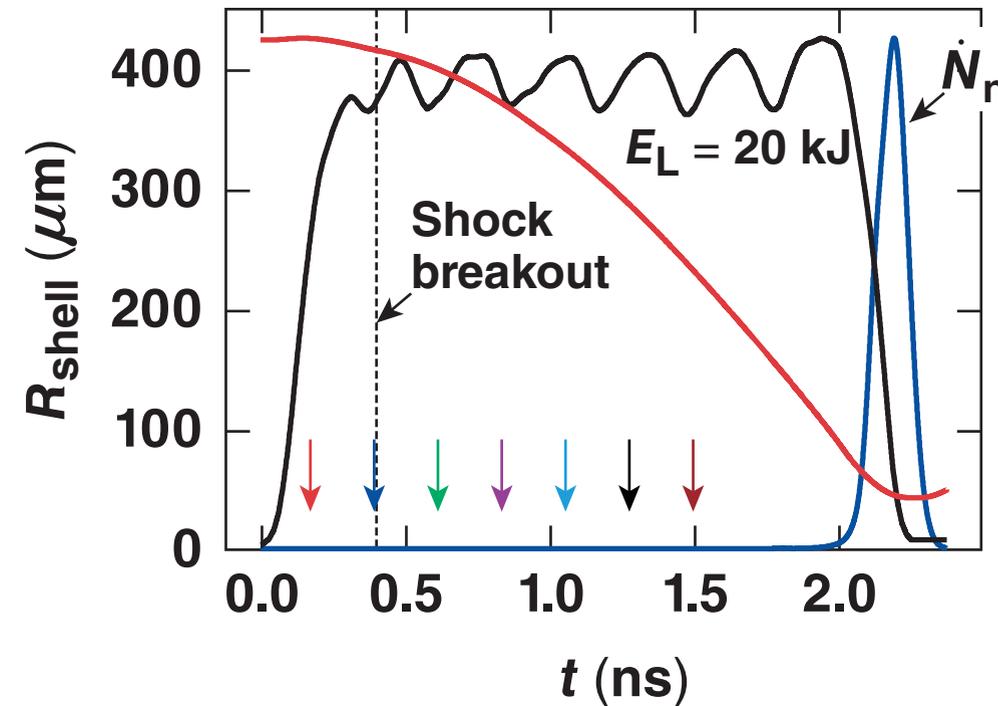


# Three-dimensional *ASTER* simulations of shot 84629 include imprint and beam overlapping effects

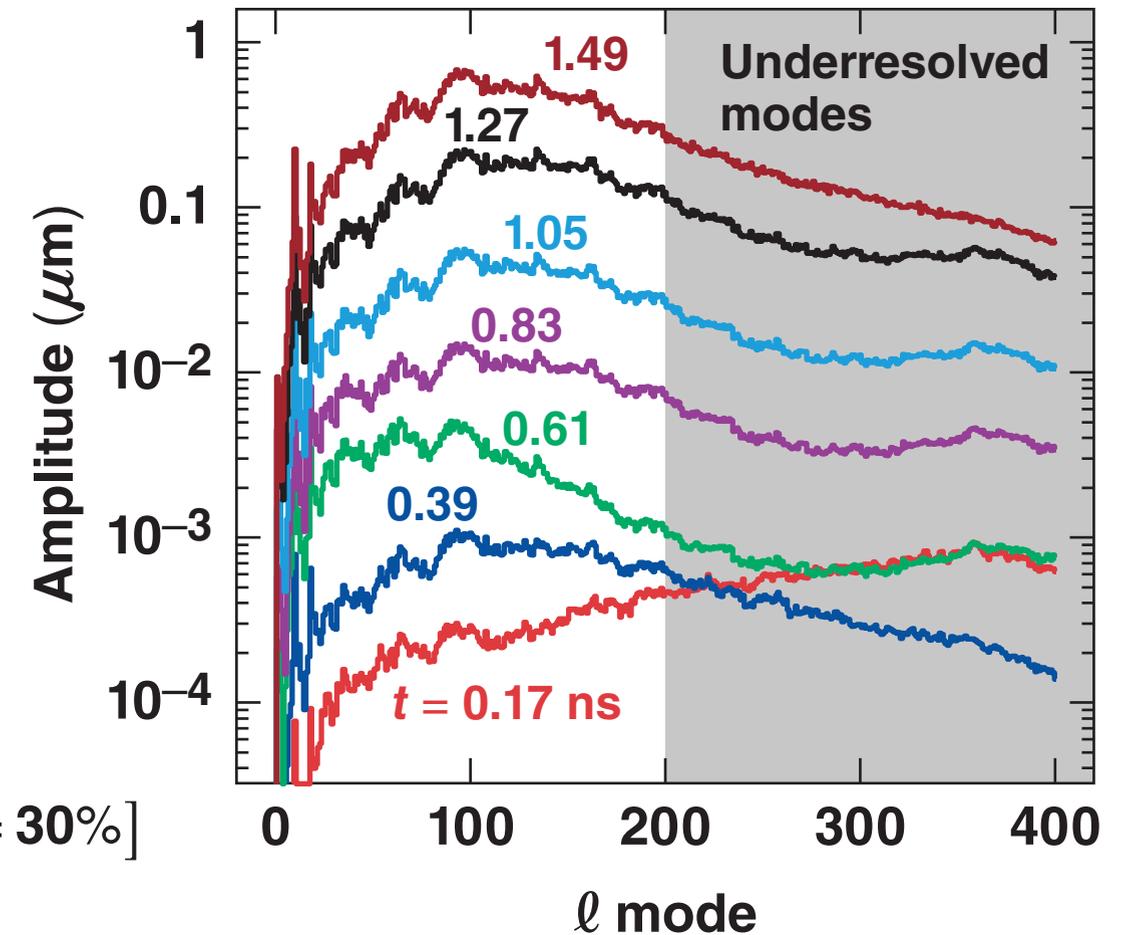
Target of shot 84629



Laser pulse and simulated shell trajectory



Evolution of  $\rho R$  modulation spectrum (up to end of linear RT\* growth, SSD on)

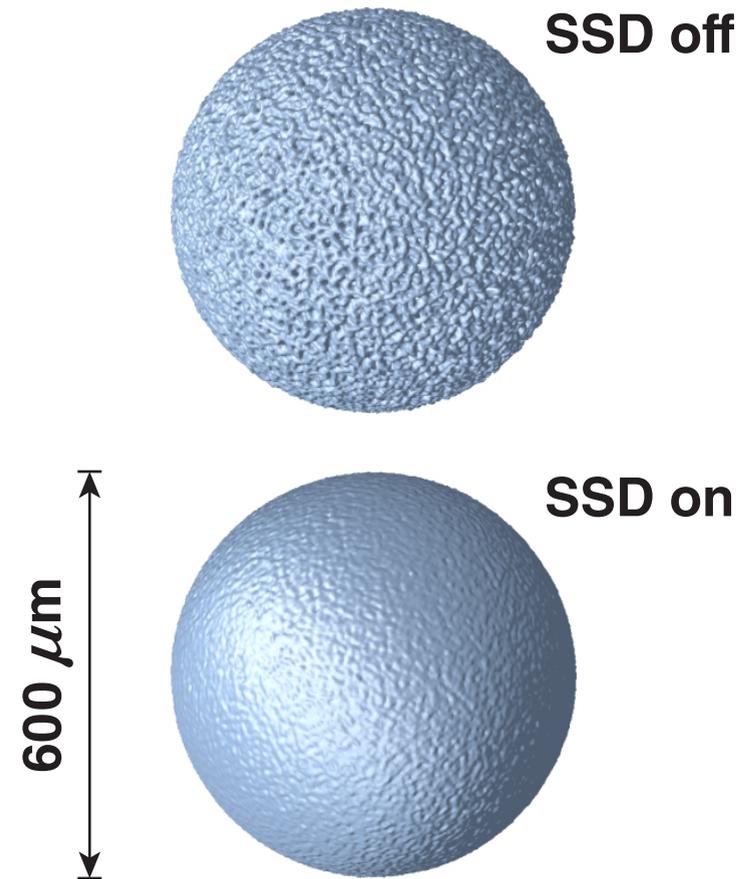


neutron yield<sub>exp</sub> =  $(1.05 \pm 0.03) \times 10^{11}$  [yield over clean (YOC) = 30%]

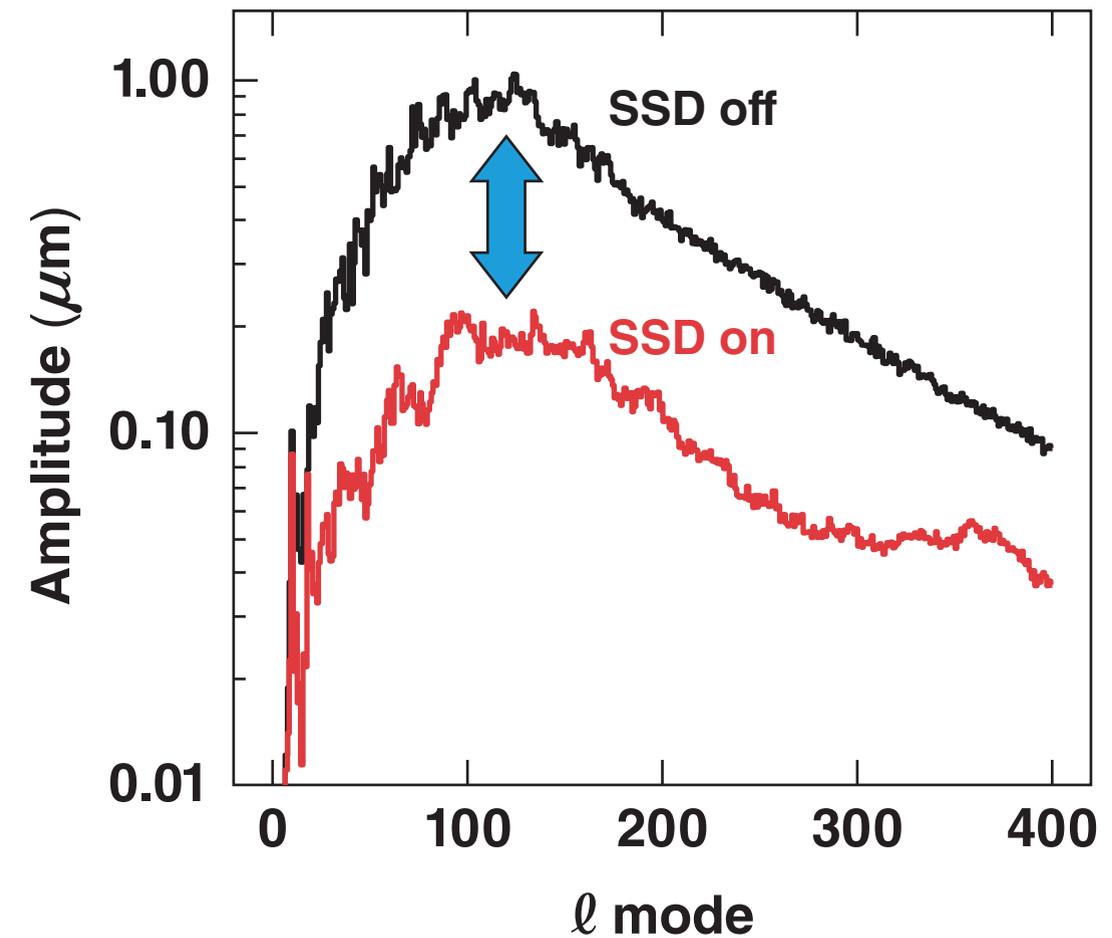
neutron yield<sub>3-D</sub> =  $1.63 \times 10^{11} / 0.30 \times 10^{11}$  (SSD on/off)

# OMEGA 2-D SSD reduces imprint perturbations by a factor of ~5

Density modulations at  
the ablation surface ( $\rho = 0.4 \text{ g/cm}^3$ )

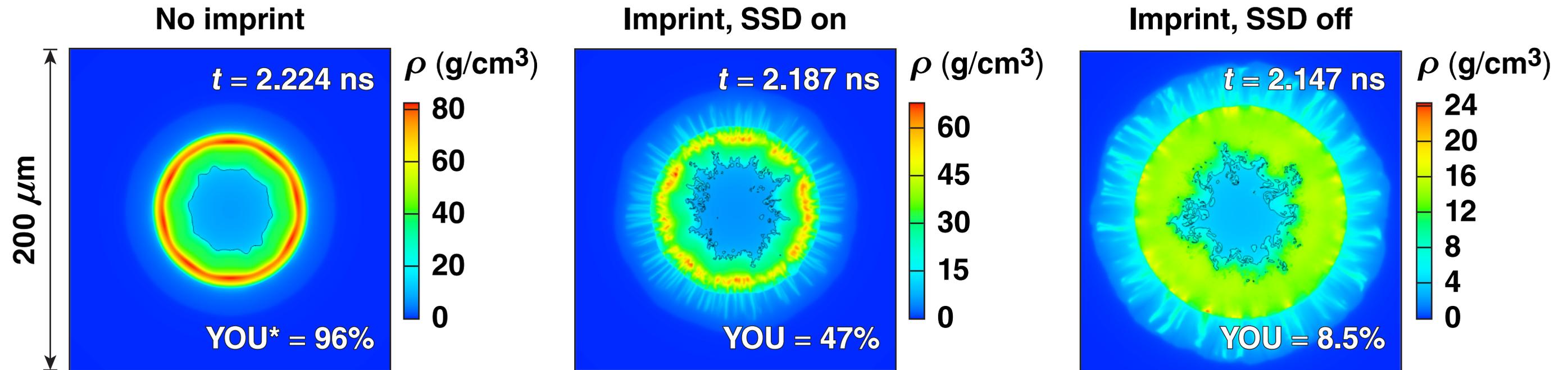


$\rho R$  modulation spectra ( $t = 1.268 \text{ ns}$ )

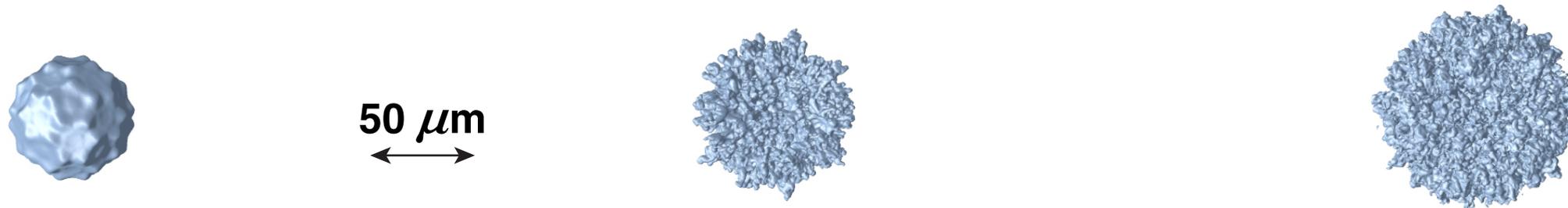


# Imprint introduces fuel–ablator material mix and increases the effective implosion adiabat by thickening the dense shell

Density map at neutron peak

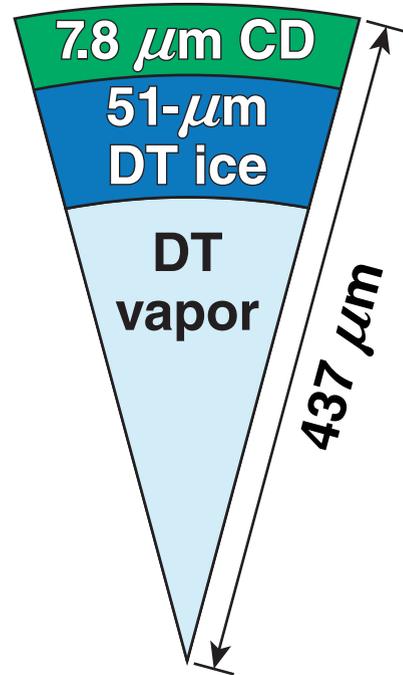


Fuel–ablator interface at neutron peak

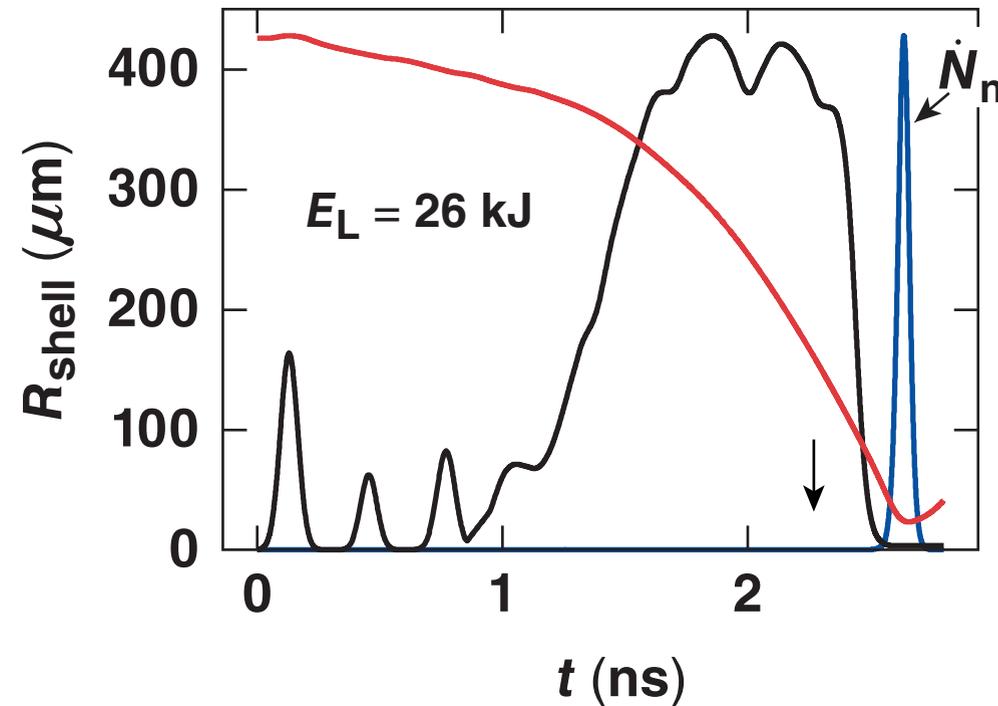


# Simulations indicate that cryogenic implosions are more affected by imprint at low modes ( $\ell < 50$ ) and less at high modes ( $\ell \sim 100$ )

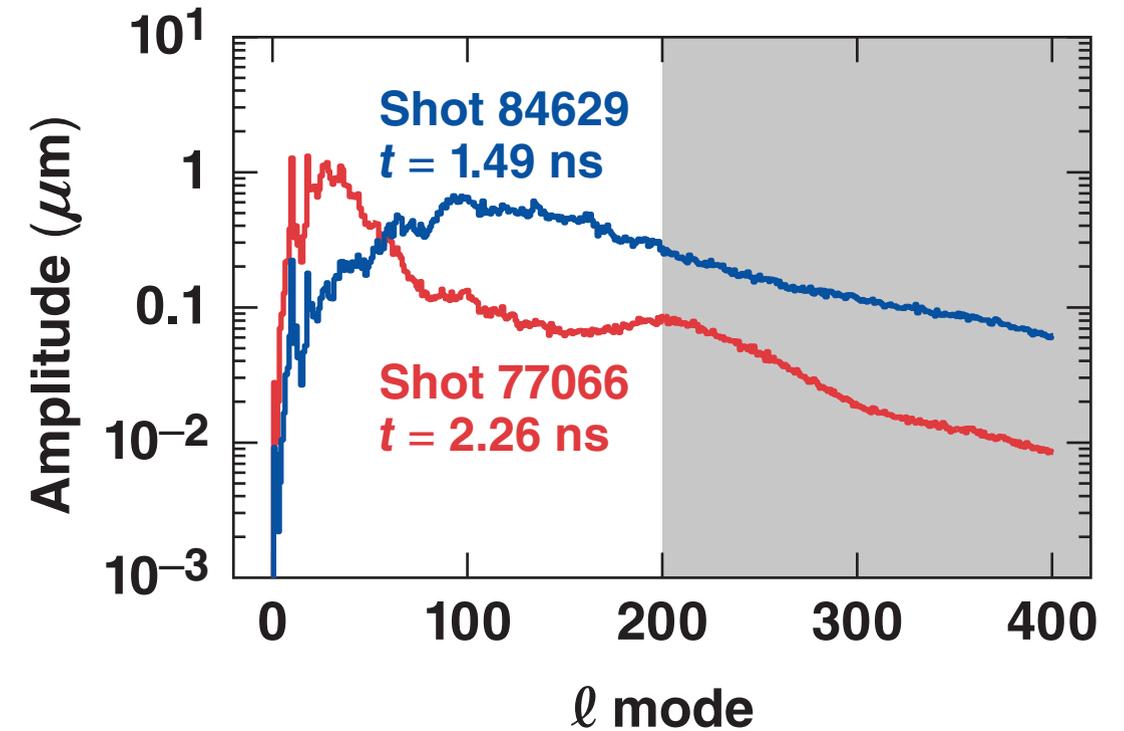
Target of shot 77066



Laser pulse and simulated shell trajectory



Spectrum of  $\rho R$  modulations at the end of linear RT growth (SSD on)



neutron yield<sub>exp</sub> =  $(3.93 \pm 0.2) \times 10^{13}$  (YOC = 30%)

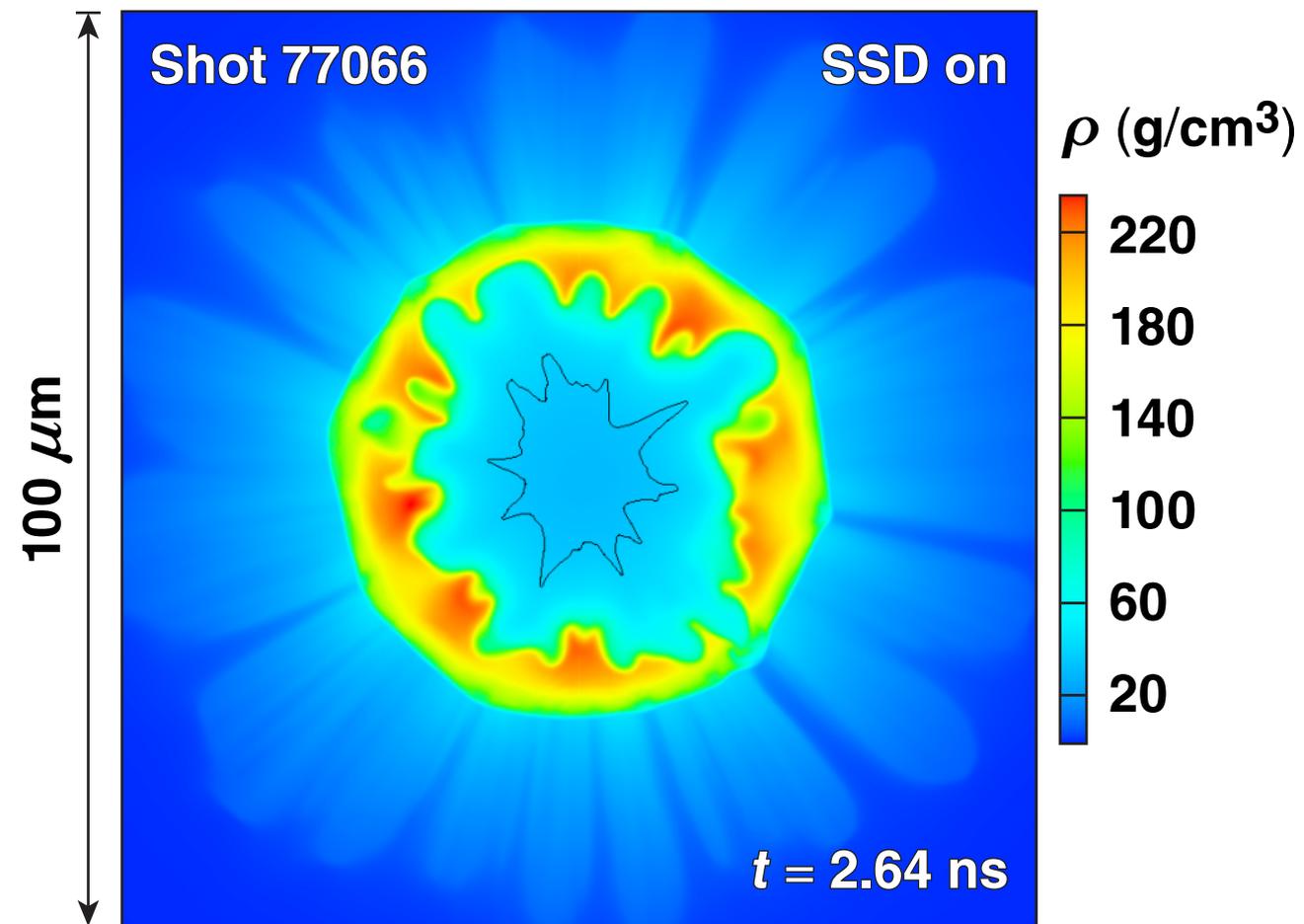
neutron yield<sub>3-D</sub> =  $9.785 \times 10^{13} / 4.09 \times 10^{13}$  (SSD on/off)

$P_{exp}^{hs} = 56 \pm 7$  Gbar

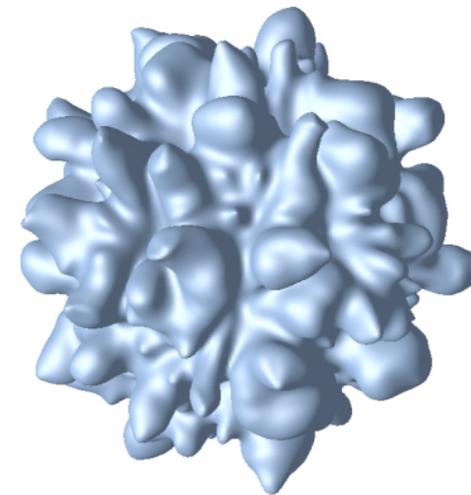
$P_{3-D}^{hs} = 107/56$  Gbar

# Imprint with applied 2-D SSD moderately affects mid-adiabat ( $\alpha \sim 4$ ) cryogenic OMEGA implosions

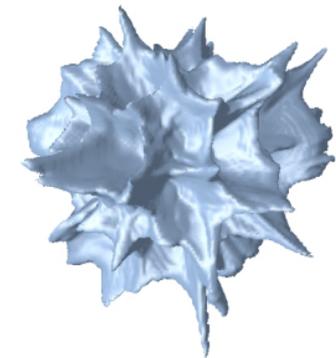
Density shell cross section  
at neutron peak



Hot-spot shape  
(surface  $T_i = 1$  keV)



Interface between  
the original DT ice  
and vapor materials



$\text{YOU} = 74\%$     $P_{3\text{-D}}^{\text{hs}}/P_{1\text{-D}}^{\text{hs}} = 0.94$

- Similar conclusions from *DRACO* simulations\*

\*P. B. Radha et al., CO8.00012, this conference.

# Effects other than imprint should be added into simulations to explain the underperformance of cryogenic OMEGA implosions



- **Low-mode nonuniformities\***
  - target offset
  - beam power imbalance
  - beam mispointing
  - beam mistiming
- **Surface defects (in progress)**
- **Target-engineering structures**
  - stalk mount\*\*
  - fill tube

} ← the effect of shadow could be important

\*I. V. Igumenshchev *et al.*, *Phys. Plasmas* **24**, 056307 (2017).

\*\*I. V. Igumenshchev *et al.*, *Phys. Plasmas* **16**, 082701 (2009).

# Simulations indicate that the effects of laser imprint alone are insufficient to explain the underperformance of cryogenic $\alpha \sim 4$ implosions on OMEGA



- The imprint model was developed and implemented in the 3-D hydrodynamics code *ASTER*\*
- Simulations reproduce observed improvement in implosion performance when polarization smoothing\*\* (PS) and smoothing by spectral dispersion† in two dimensions (2-D SSD) are applied
- Room-temperature targets suffer from imprint that introduces significant small-scale ( $\ell \sim 50$  to 150) modulations
- Imprint in cryogenic implosions develops broadband modulations with dominant  $\ell \sim 30$ ; these modulations have a moderate effect on the implosions

\*I. V. Igumenshchev *et al.*, Phys. Plasmas **23**, 052702 (2016).

\*\*T. R. Boehly *et al.*, J. Appl. Phys. **85**, 3444 (1999).

†S. Skupsky *et al.*, J. Appl. Phys. **66**, 3456 (1989).