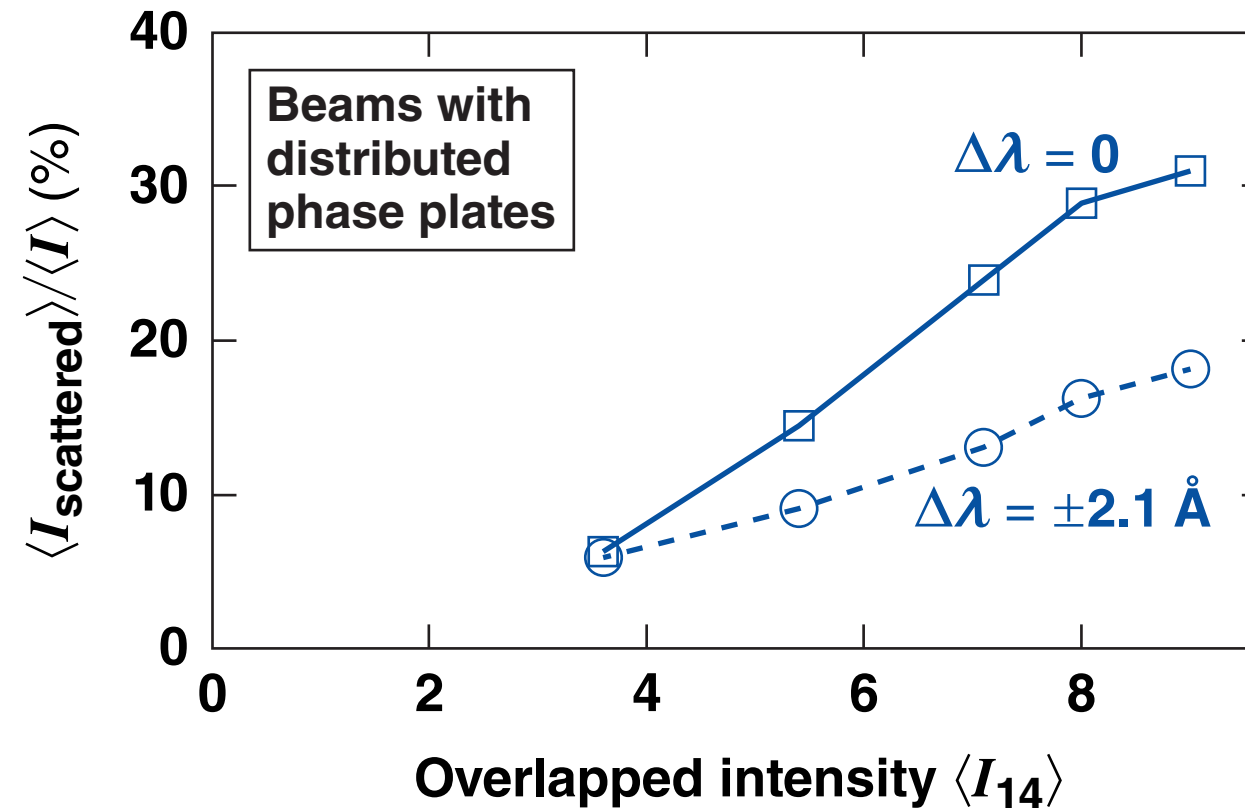


# Beam-Energy Exchange Driven by Incoherent Laser Beams with Frequency Detuning



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## Summary

**In direct-drive inertial confinement fusion (ICF) plasmas, the frequency detuning in the driving incoherent laser beams mitigates the beam-energy exchange**



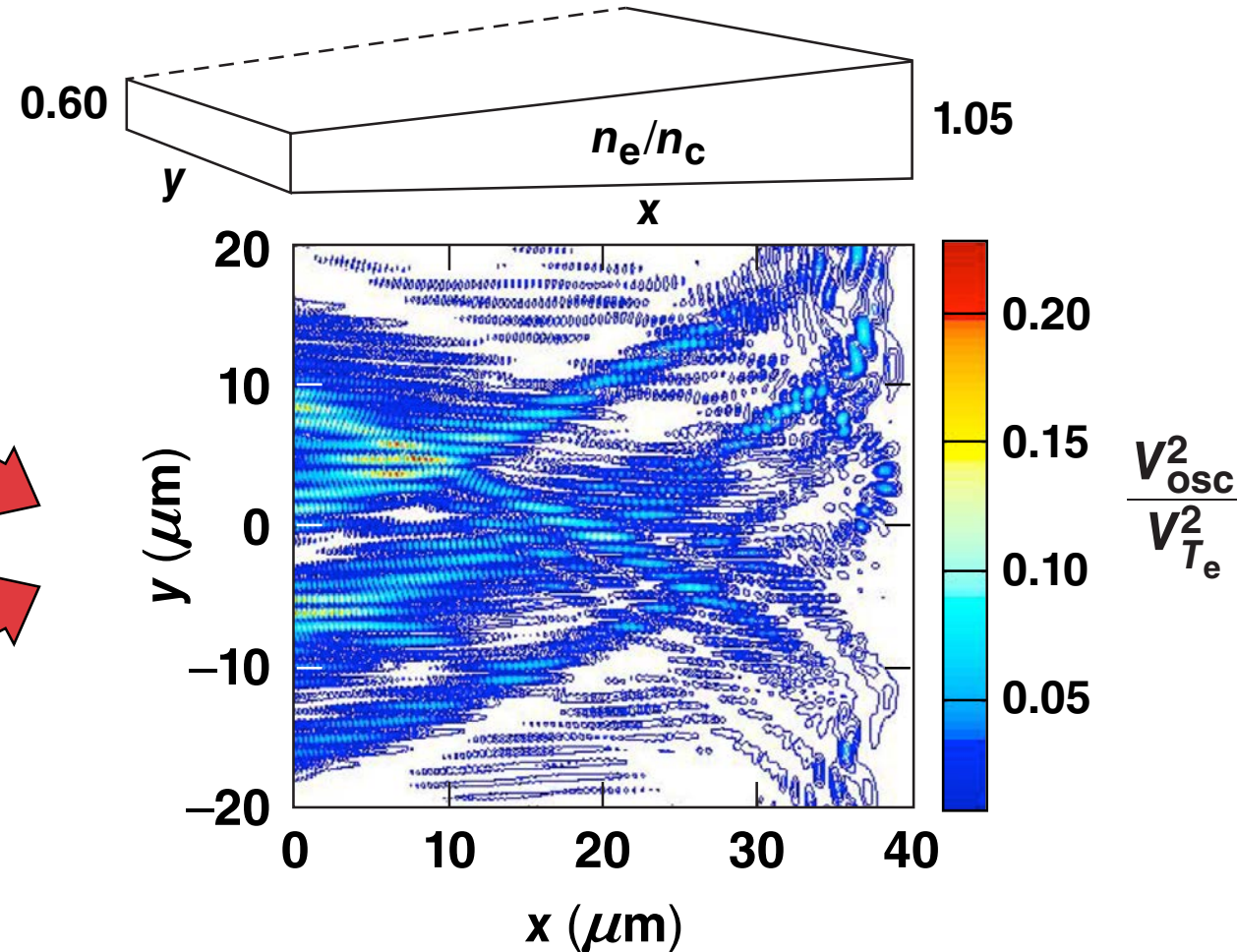
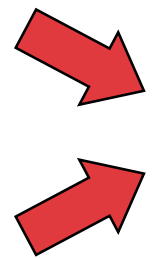
- **In dense plasmas (including beam-turning points) nonlinear laser–plasma interactions (LPI’s) can lead to angular broadening of scattered light**
- **The use of frequency detuning in laser beams can limit the beam-energy exchange in plasmas because it**
  - **reduces the role of common ion-acoustic waves (IAW’s)**
  - **increases the frequency broadening of scattered light**
- **The scaling of the beam-energy exchange with intensity and frequency detuning in an LPI model for incoherent beams is different from the plane-wave model**

# Nonlinear propagation of laser beams with frequency detuning is modeled in dense plasmas

- Beams with wavelength detuning are used to limit beam-to-beam coupling
  - two-dimensional nonparaxial model near turning points
  - related to parameters from simulations of OMEGA experiments (flow velocity  $\ll c_s$ )

$T_e = 2 \text{ keV}$   
 $\Delta\lambda_0 = 0, \pm 2.1 \text{ \AA}$

$\theta = \pm 15^\circ, \pm 30^\circ$   
 $f/6$   
 $\langle I_{14} \rangle = 8$



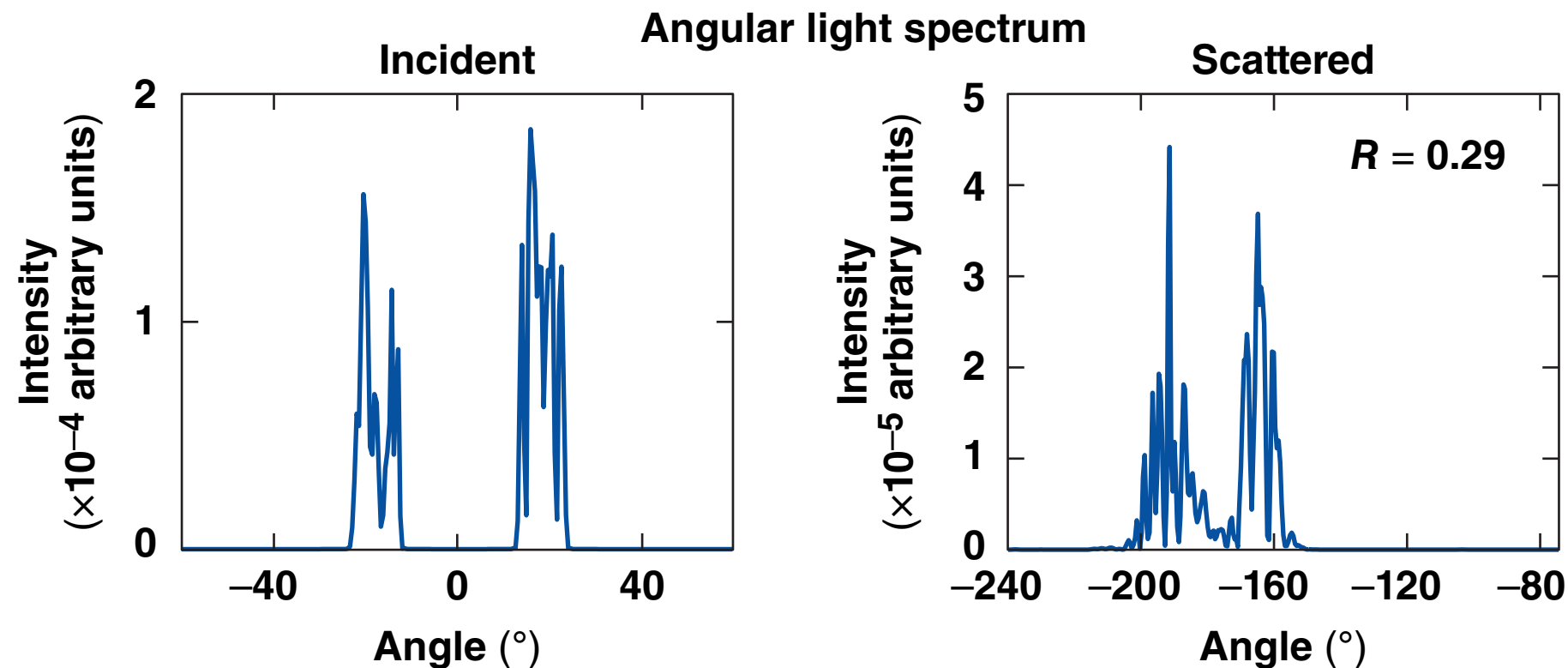
# The LPI model includes backward- and forward-stimulated Brillouin scattering (SBS), beam self-focusing, field swelling, and absorption

- Beam-to-beam coupling can be described by backward SBS gains

$$\frac{dG_{\text{SBS}}}{d\ell} = \frac{\omega_0^2}{2c^2 n_c} \text{Re} \left\{ \frac{n_e k_s^2 c_s^2 \times I_0}{2\nu_i \omega_s + i[(\omega_s + k_s v_0)^2 - k_s^2 c_s^2]} \times \frac{1}{2k_{0x}} \right\},$$

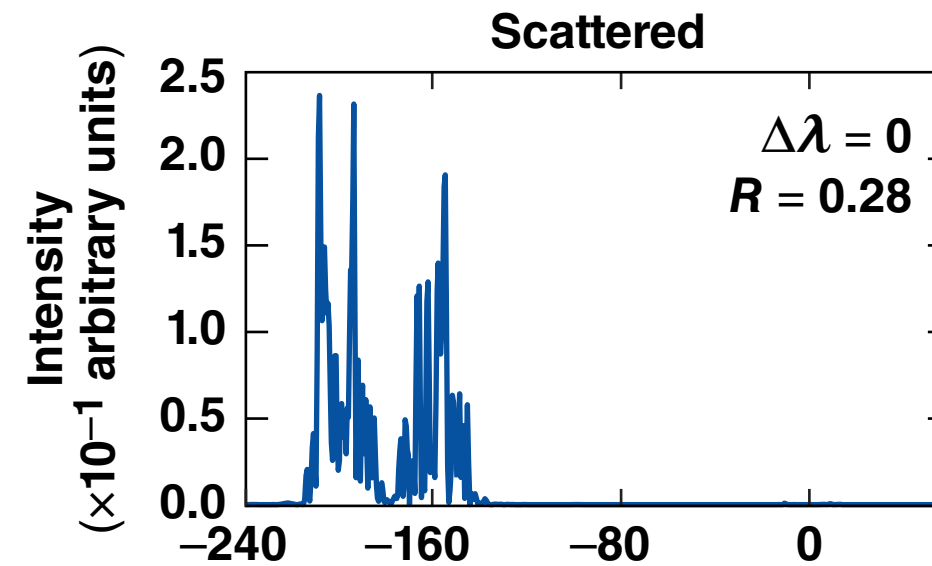
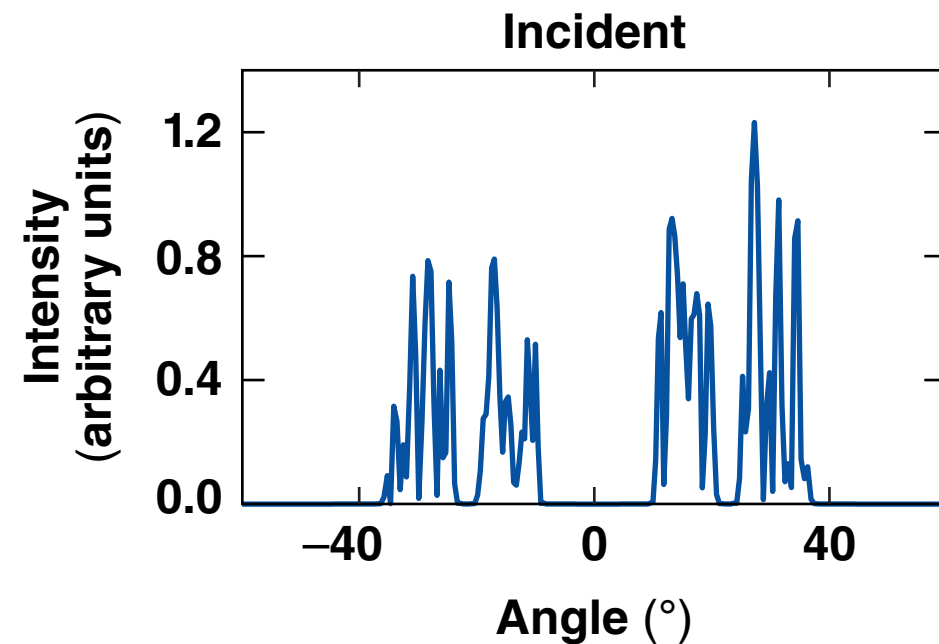
$$I_0 = |E|^2 / 4\pi n_c T_e$$

$$\langle I_{14} \rangle = 8$$



The angular width of scattered light is increased

# The beam-energy exchange is significantly reduced when frequency detuning (color) is applied to crossing laser beams

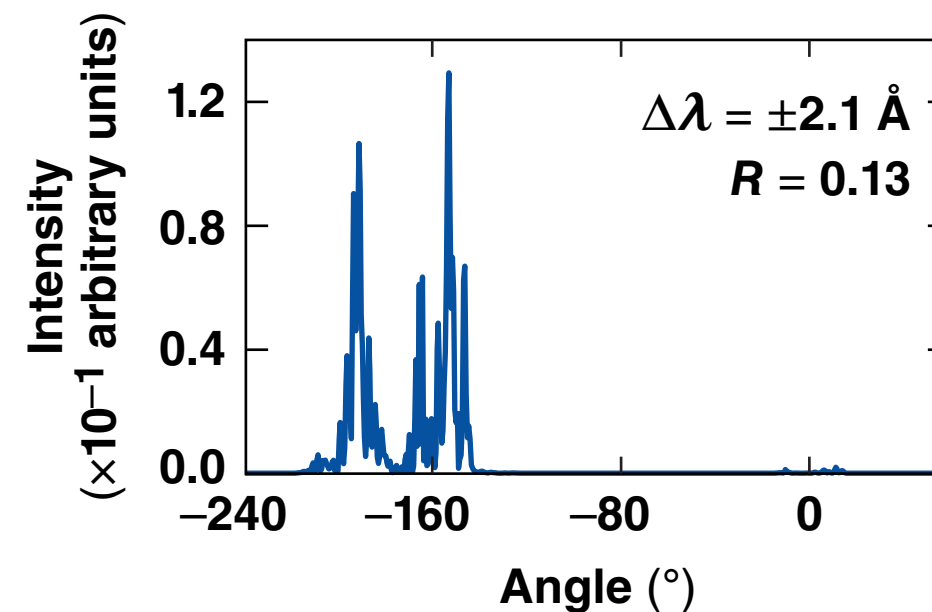


Four laser beams

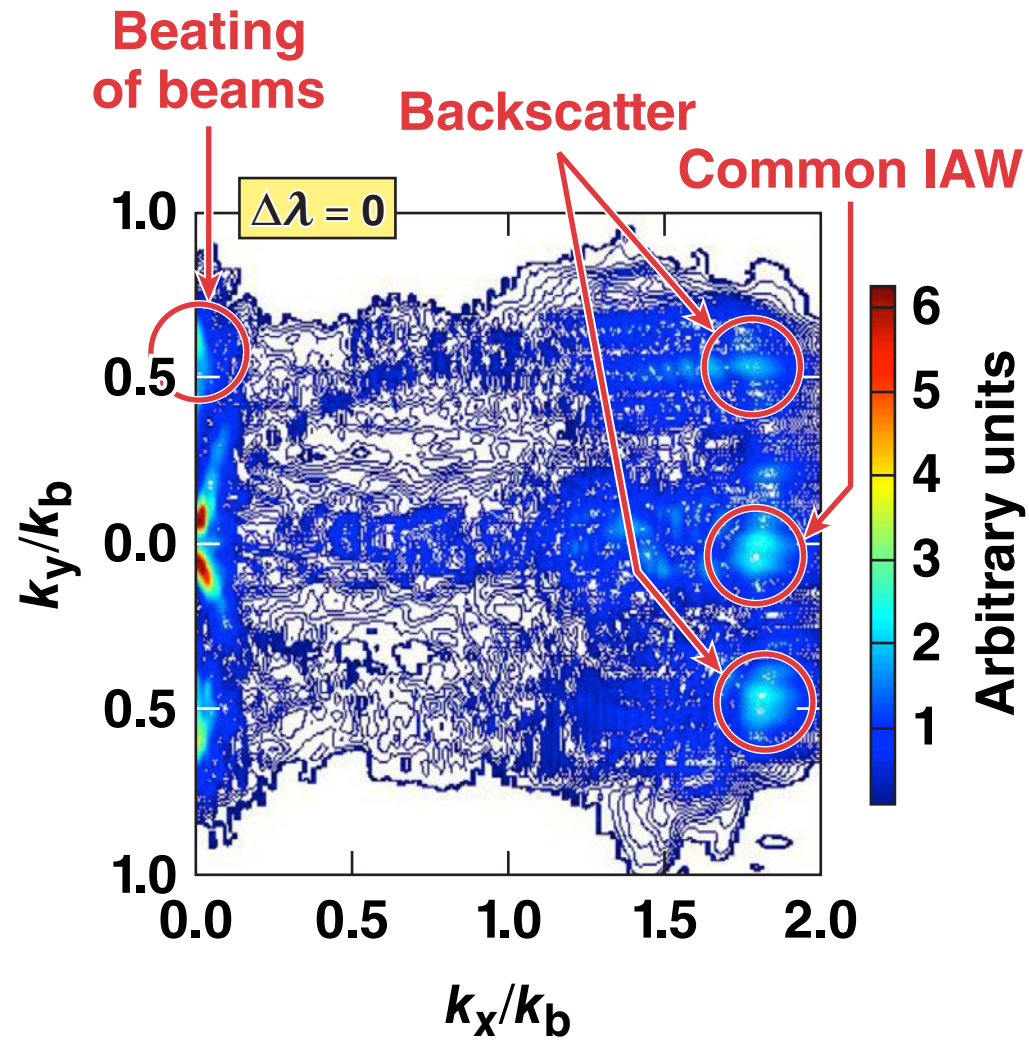
$$\theta = \pm 15^\circ, \pm 30^\circ$$

$$\langle I_{14} \rangle = 8$$

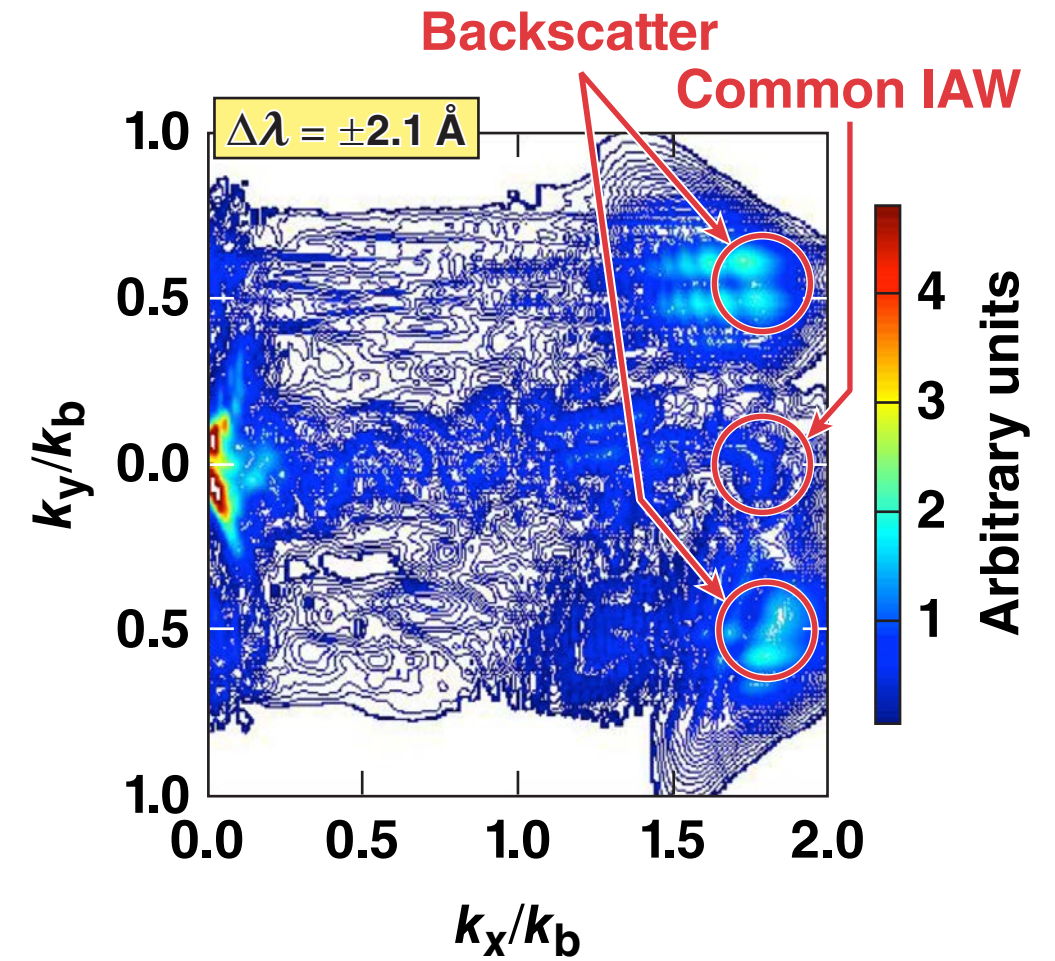
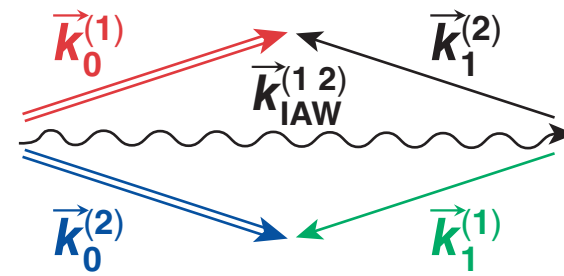
$$I_{\text{outer}} / I_{\text{inner}} = 1$$



# The spectra of density perturbations show different ion-acoustic waves driven in dense plasmas

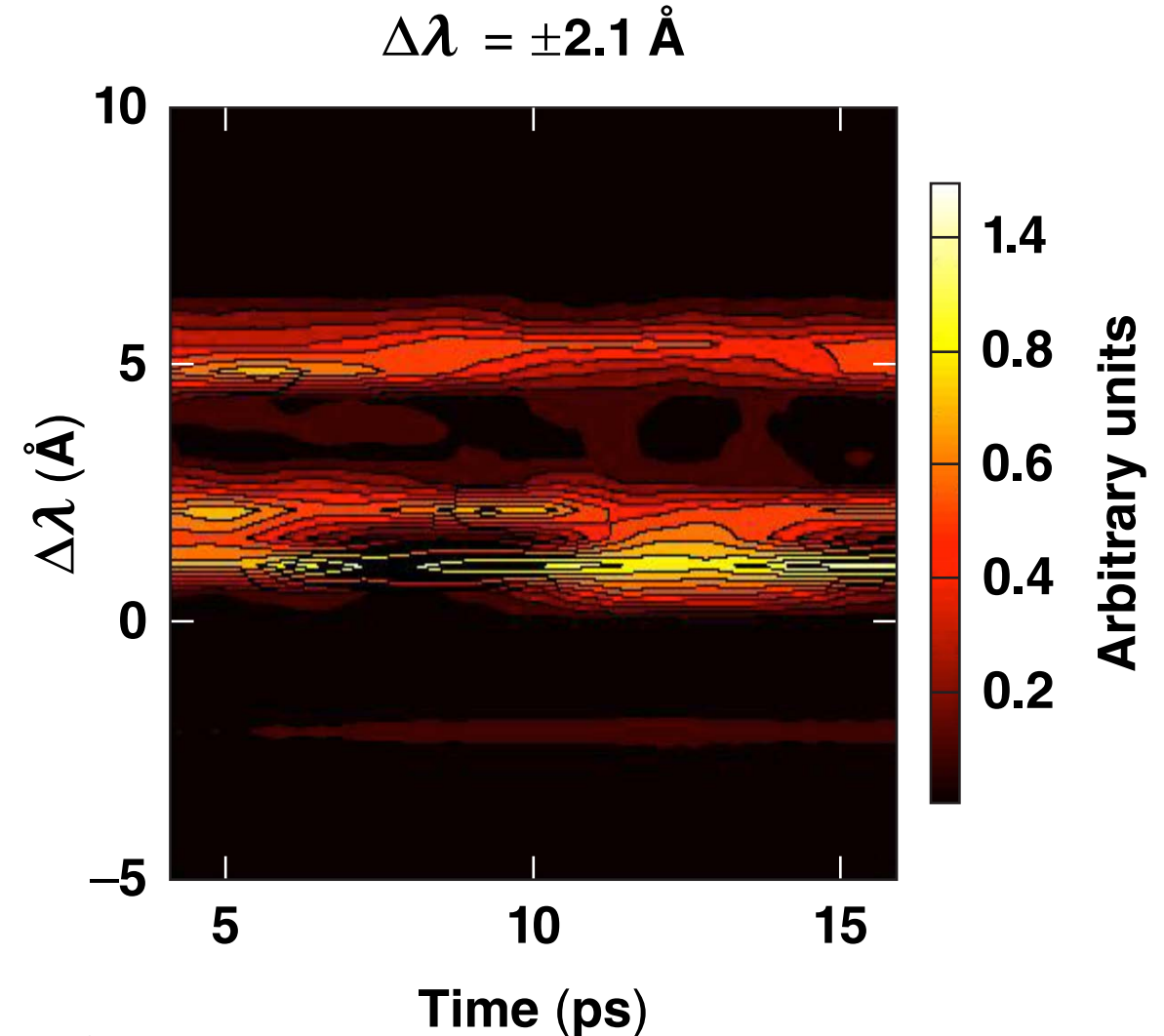
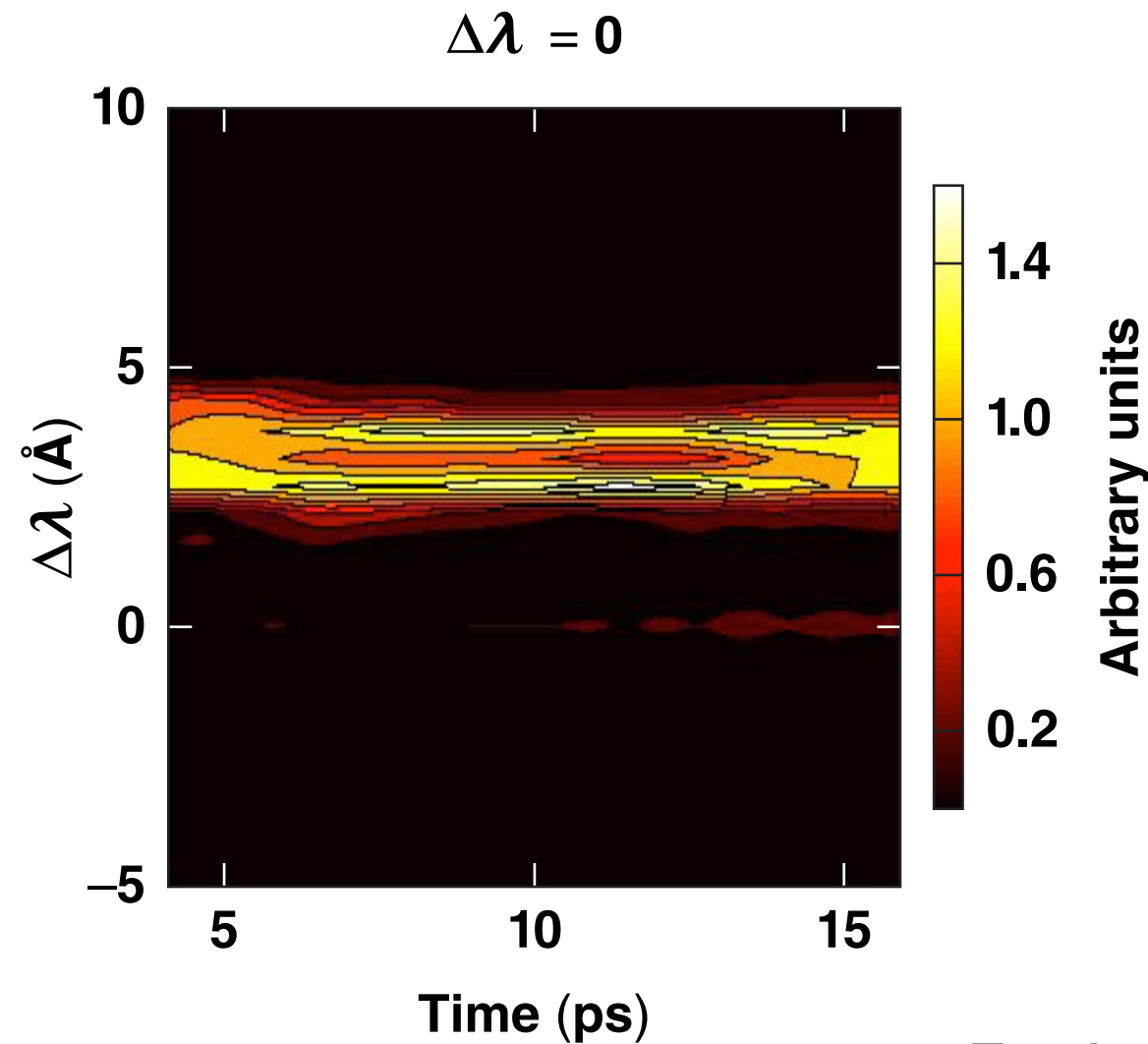


Two beams  $\pm 15^\circ$ ,  $\langle I_{14} \rangle = 8$



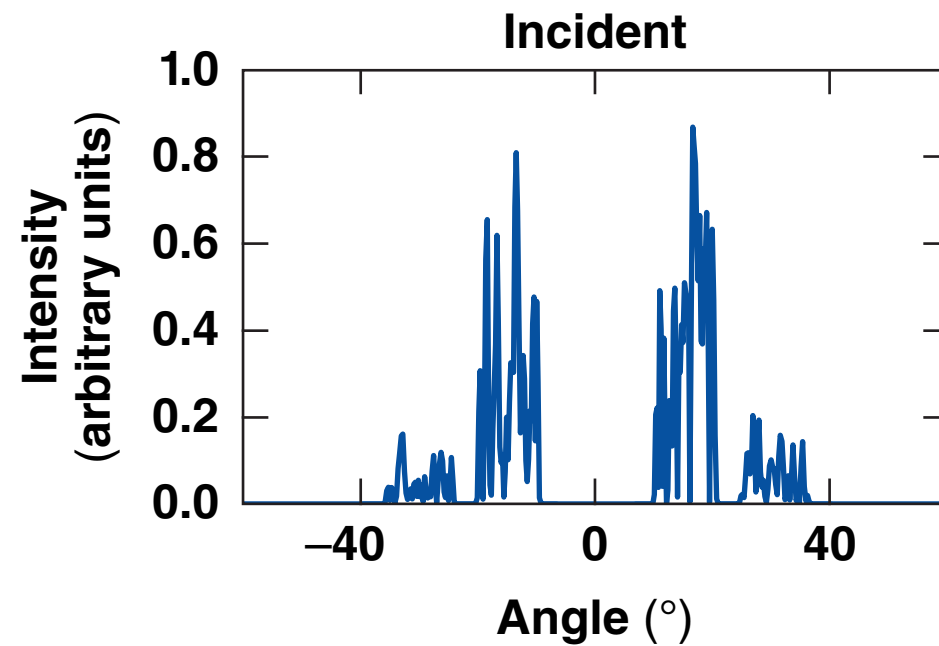
$$k_b = \frac{\omega_0}{c} \sqrt{1 - \frac{n_b}{n_c}}$$

# The frequency detuning in driving laser beams strongly modifies the frequency spectra of scattered light



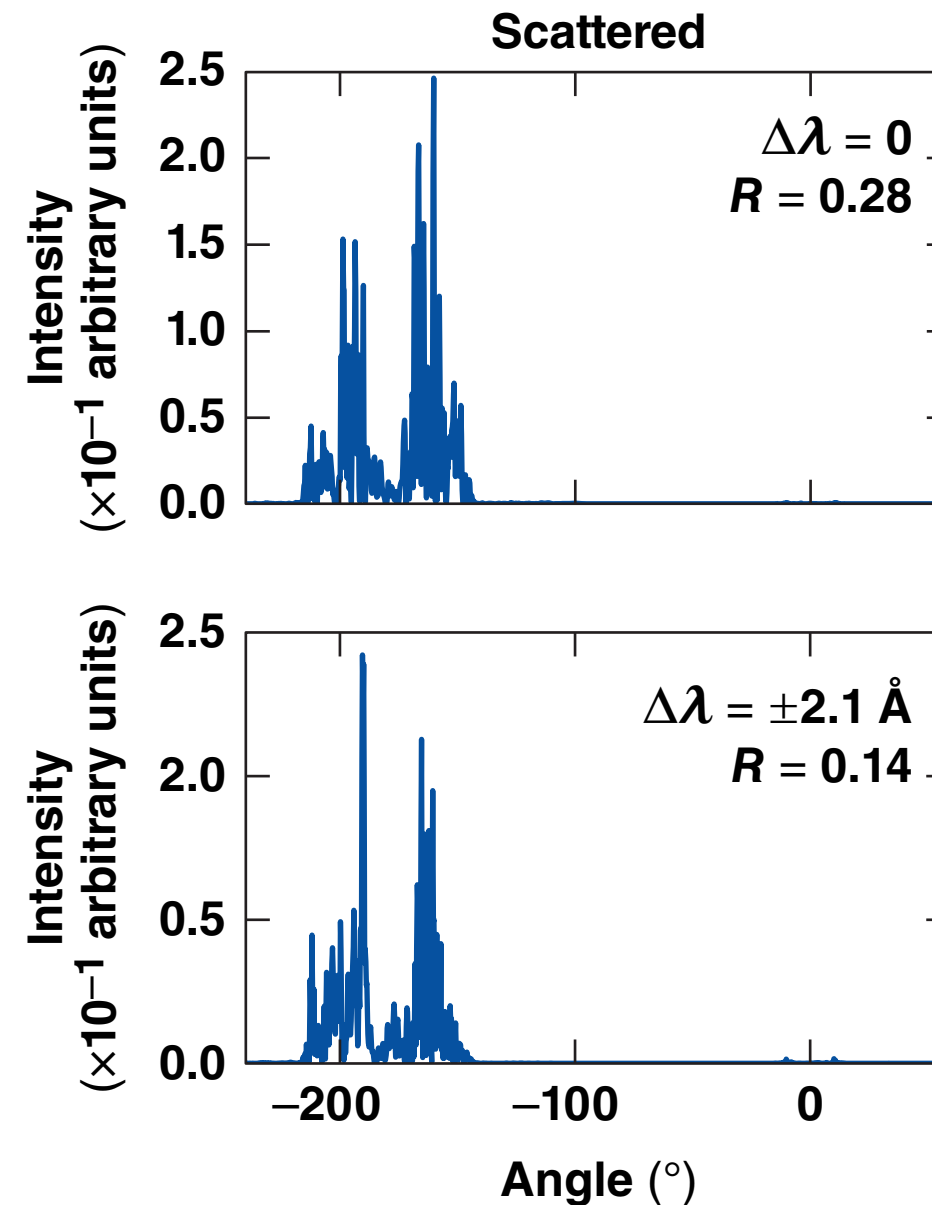
Two beams  $\pm 15^\circ$   
 $\langle I_{14} \rangle = 8$

# The beam-energy exchange driven by incoherent beams with an intensity contrast leads to the angular broadening of scattered light



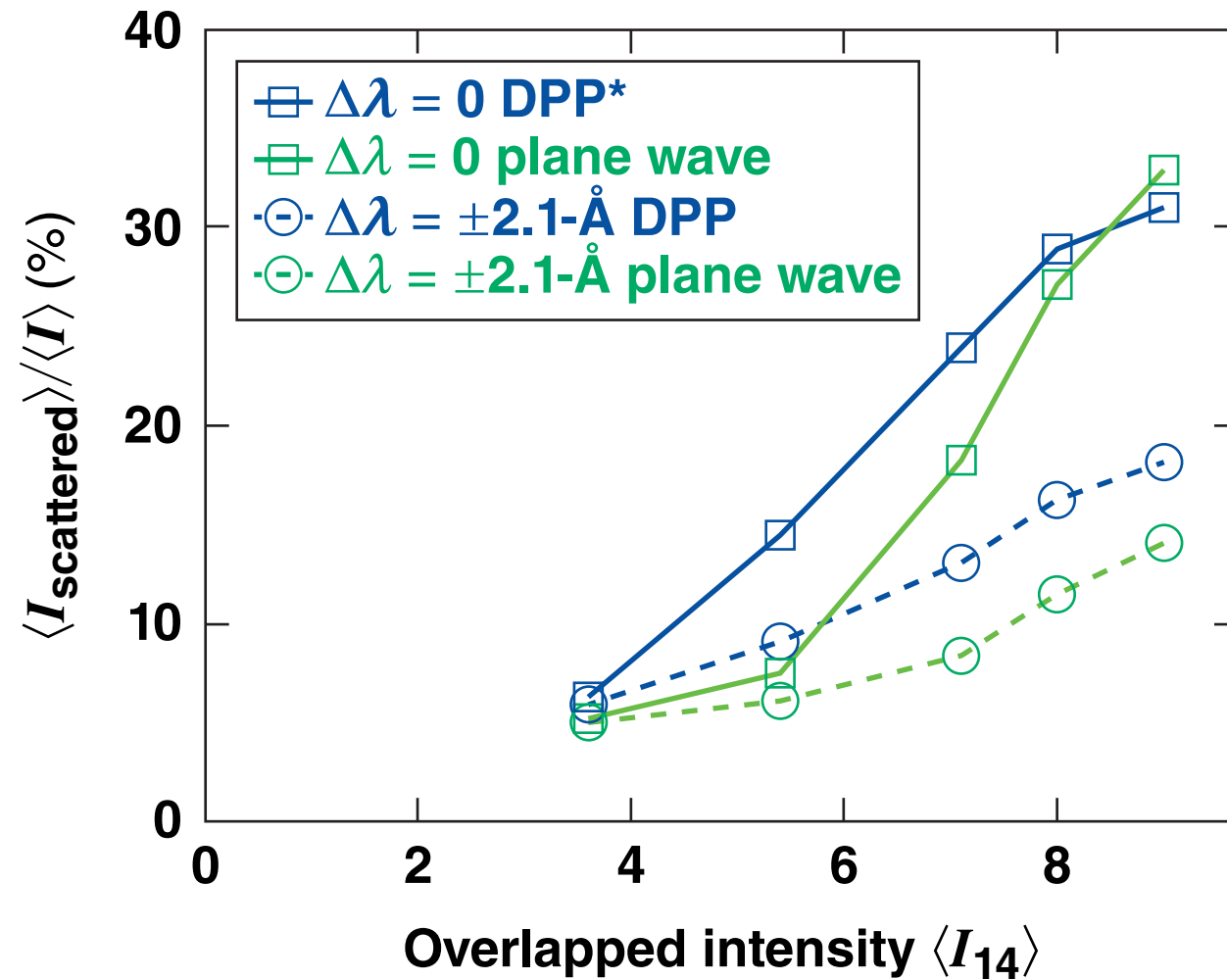
Driving beams with large intensity contrast

$$\langle I_{14} \rangle = 8$$
$$I_{\text{outer}} / I_{\text{inner}} = 0.2$$





# The frequency detuning in incoherent laser beams leads to a significant reduction in the scattered-light intensity



The ray-based beam-energy exchange models do not account for the speckled distribution of the laser intensity in plasmas.

\* DPP: distributed phase plate

## Summary/Conclusions

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