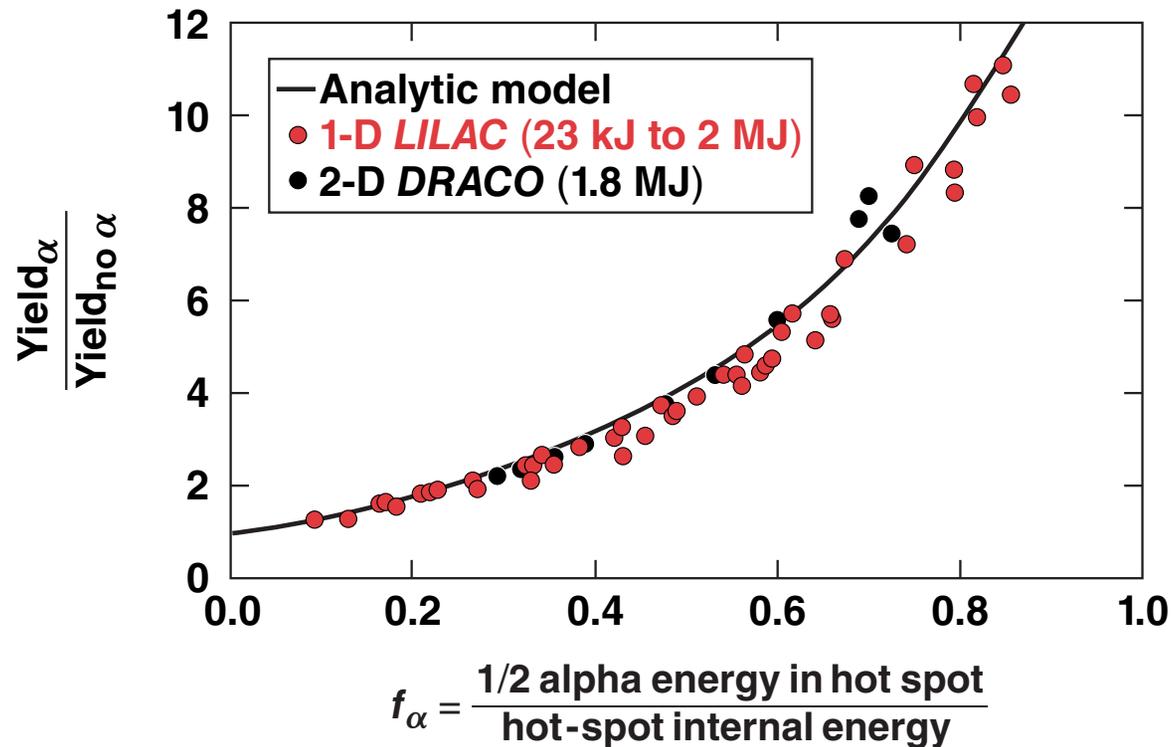


Measurements of Alpha Heating in Inertial Confinement Fusion



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Summary

Measurements of alpha heating are developed and used to determine the yield amplification caused by alphas and the requirements for ignition



- Alpha heating is estimated in two different ways using the Lawson parameter and the fractional alpha energy
- National Ignition Facility (NIF) shot N140120 exhibits an $\sim 2.5\times$ amplification of the neutron yield caused by alpha heating and a no-burn Lawson parameter of ~ 0.65
- Options for ignition: higher in-flight aspect ratios (IFAR's), improved yield-over-clean (YOC), and/or use of adiabat shaping

Collaborators



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Hot-spot evolution (including alpha heating) depends almost exclusively on the Lawson ignition parameter



- The model solves the mass, momentum, and energy conservation equations for the hot-spot pressure \hat{P} , temperature \hat{T} , and radius \hat{R}^*

- Energy conservation:
$$\frac{d}{d\tau}(\hat{P}\hat{R}^5) = \gamma\hat{P}^2\hat{R}^5\hat{T} - \beta\hat{P}^2\hat{R}^5\hat{T}^{-5/2}$$

↑
Alpha heating

↑
Radiation losses

- The solution is singular when $\chi_{no\ \alpha} = \left[\frac{P\tau}{(P\tau)_{ign}} \right]_{no\ \alpha} \equiv \underbrace{\frac{\gamma}{1.11 + 0.7\beta}}_{\text{Lawson parameter}}$

- Rewrite energy equation dependent approximately only on $\chi_{no\ \alpha}$

$$\frac{d}{d\tau}(\hat{P}\hat{R}^5) = 1.1\chi_{no\ \alpha}\hat{P}\hat{R}^5\hat{T} + \hat{P}^2\hat{R}^5(0.7\cancel{\chi_{no\ \alpha}}\beta\hat{T} - \cancel{\beta}\hat{T}^{-5/2})$$

↑
Quasi cancellation (~20% error)

*C. D. Zhou and R. Betti, Phys. Plasmas **15**, 102707 (2008);
P.-Y. Chang et al., Phys. Rev. Lett. **104**, 135002 (2010);
R. Betti et al., Phys. Plasmas **17**, 058102 (2010).

$\chi_{\text{no } \alpha}$ is the most useful ignition metric, but only χ_{α} can be measured



- The no- α ignition parameter is written in terms of the yield (Y) and ρR

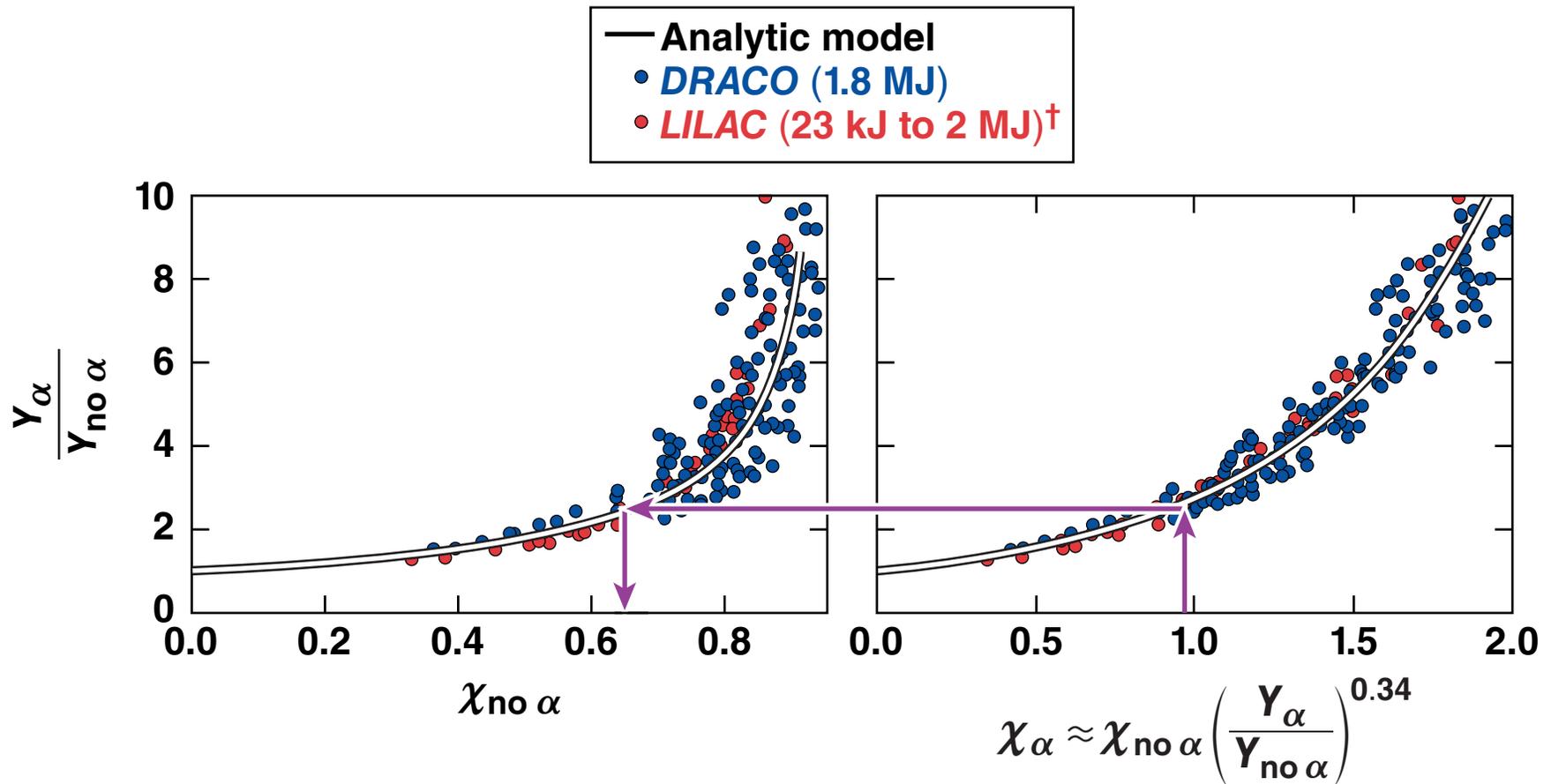
$$\chi_{\text{no } \alpha} = (\rho R_{\text{no } \alpha})^{0.61} \left(0.24 \frac{Y_{\text{no } \alpha}}{M^{\text{DT}}} \right)^{0.34}$$

- The measurable parameter must use quantities with alpha heating

$$\chi_{\alpha} = (\rho R_{\alpha})^{0.61} \left(0.24 \frac{Y_{\alpha}}{M^{\text{DT}}} \right)^{0.34} \approx \chi_{\text{no } \alpha} \left(\frac{Y_{\alpha}}{Y_{\text{no } \alpha}} \right)^{0.34}$$

- ρR = total areal density in g/cm²
- Y = neutron yield in units of 10¹⁶ neutrons
- M^{DT} = unablated DT mass in mg

The yield-enhancement curves are used to measure both the yield enhancement caused by alphas and the no-alpha Lawson parameter



[†]In general agreement with B. Spears' (LLNL) simulation results for NIF-ID point design target:

B. Spears and J. Lindl, "Ignition Metrics and Their Role in Setting Specifications and Evaluating Progress Toward Ignition on the NIF," LLNL, Livermore, CA, UCRL report under review (2014).; P. Patel *et al.*, *Bull. Am. Phys. Soc.* **58**, 193 (2013).

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Another way of inferring alpha heating is through the fractional alpha energy deposited in the hot spot



$$f_{\alpha} \equiv \frac{0.5 E_{\alpha}^{\text{abs}}}{\frac{3}{2} \langle P \rangle V_{\text{hs}}}$$

- $E_{\alpha}^{\text{abs}} = \theta_{\alpha} \times \text{Yield} \times 3.5 \text{ MeV} = \text{absorbed alpha energy in the hot spot}$
- $\theta_{\alpha} \approx 0.9 \left(1 - \frac{1}{3.4 \tau} + \frac{1}{160 \tau^2} \right) = \text{fraction of alphas absorbed in the hot spot}^*$

$$\text{where } \tau = \frac{\text{hot-spot radius}}{\text{alpha-particle range}} = \left(\frac{P_{\text{Gbar}}}{100} \right) \left(\frac{R_{\mu\text{m}}}{50} \right) \left(\frac{5}{T_{\text{keV}}} \right)^{5/2}$$

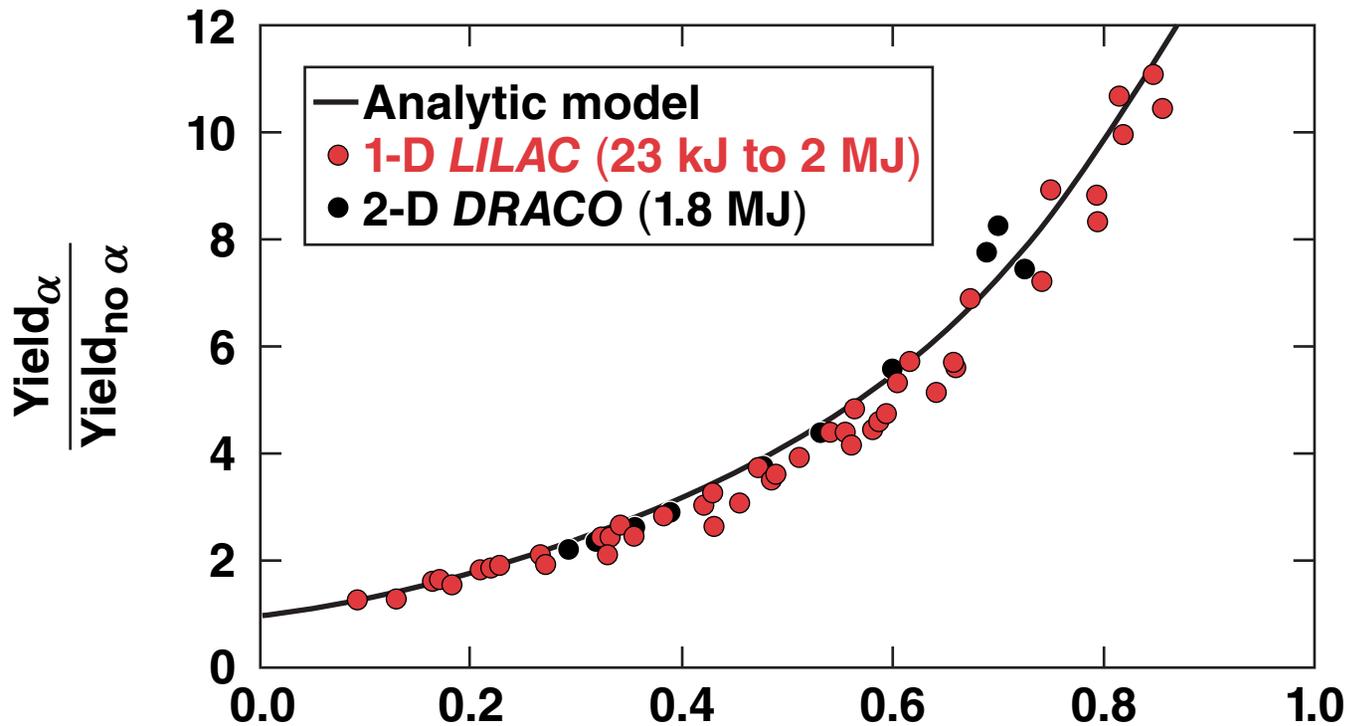
- Pressure and volume are inferred from observables according to C. Cerjan *et al.***

f_{α} can be directly inferred from experimental observables.

* O. N. Krokhin and V. B. Rozanov, Sov. J. Quantum Electron. **2**, 393 (1973).

C. Cerjan, P. T. Springer, and S. M. Sepke, Phys. Plasmas **20, 056319 (2013).

The yield enhancement is an almost unique function of the fractional alpha energy f_α



$$f_\alpha = \frac{1/2 \alpha\text{-energy deposited in hot spot}}{\text{hot-spot internal energy}}$$

Alpha-heating analysis of NIF high-foot (HF) shot N140120



- Yield = 9.2×10^{15} , $\rho R = 0.8$ g/cm², $M_{DT} = 0.18$ mg, burnwidth = 161 ps, $T = 4.9$ keV, and $R_{HS} = 35.2$ μ m*
- The χ_{α} analysis gives $\chi_{\alpha} \approx 1$, a yield amplification of 2.5, and $\chi_{no \alpha} = 0.65$
- The f_{α} analysis gives $f_{\alpha} \approx 0.38$, a yield amplification of 2.7, and $\chi_{no \alpha} = 0.67$

Both the χ_{α} and the f_{α} analyses give similar results.

*P. K. Patel, Lawrence Livermore National Laboratory, private communication (2014); O. A. Hurricane and H. S. Park, presented at the IDI Web Meeting, February 2014.

Options for ignition: higher IFAR, improved YOC, and/or use of adiabat shaping

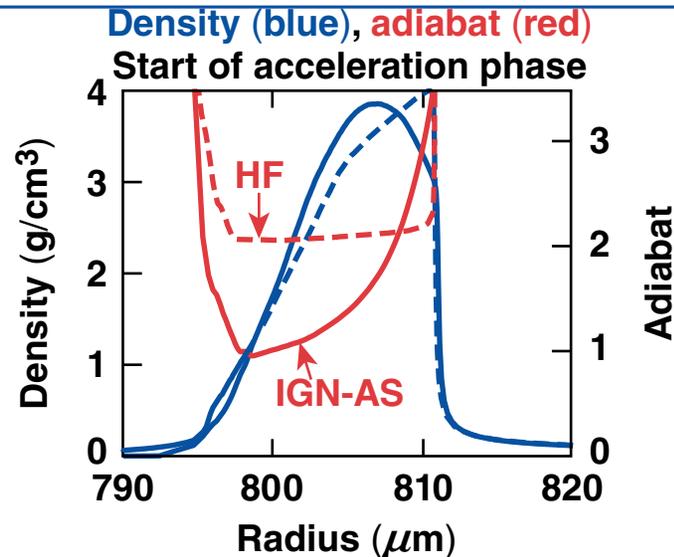
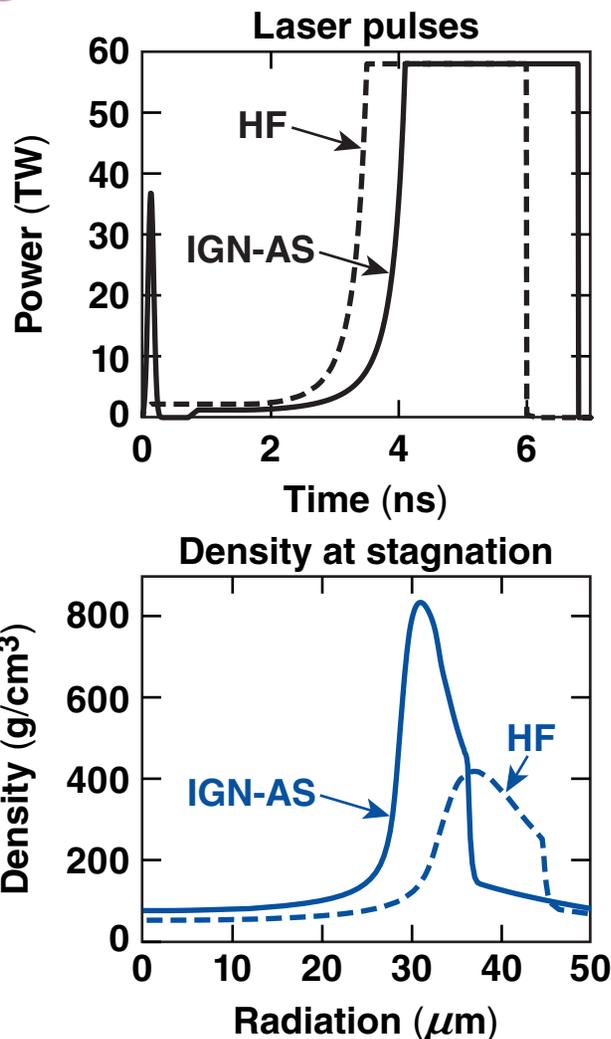


$$\chi_{\text{no } \alpha} \sim \theta_{\alpha} E_{\text{kin}}^{0.37} \text{YOC}^{0.4} P_{\text{abl}}^{0.4 \text{ to } 0.6} \text{IFAR} S_{\text{adiabat}}^{0.6}$$

- θ_{α} fraction of absorbed alphas ≈ 0.7
- S_{adiabat} = adiabat shaping factor = $\langle \alpha \rangle / \alpha_{\text{inner}}$
- Best shot to date $\rightarrow \chi_{\text{HF}} \approx 0.65$ ($\chi = 1$ for ignition)
- Options for achieving ignition
 - improve YOC (shape, IFAR, S_{adiabat}) but YOC is already high
 - increase θ_{α} (compressed B field at 0.5 Gauss?)*
 - increase IFAR (but YOC may go down)
 - use adiabat shaping \rightarrow increase S_{adiabat}

*L. J. Perkins *et al.*, Phys. Plasmas **20**, 072708 (2013);
G. Logan, Lawrence Berkeley National Laboratory, private communication (2014).

Direct-drive simulations of the high foot (HF) show a possible ignition path with adiabat shaping (AS) and modest IFAR, CR, and pulse-length increase



- Ignition pulse with AS is 13% longer
- IFAR at 2/3 radius is 20% higher in ignition pulse with AS
- Expected similar ablation-front growth factors
- Convergence ratio (CR) is 15% higher in ignition pulse with AS

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