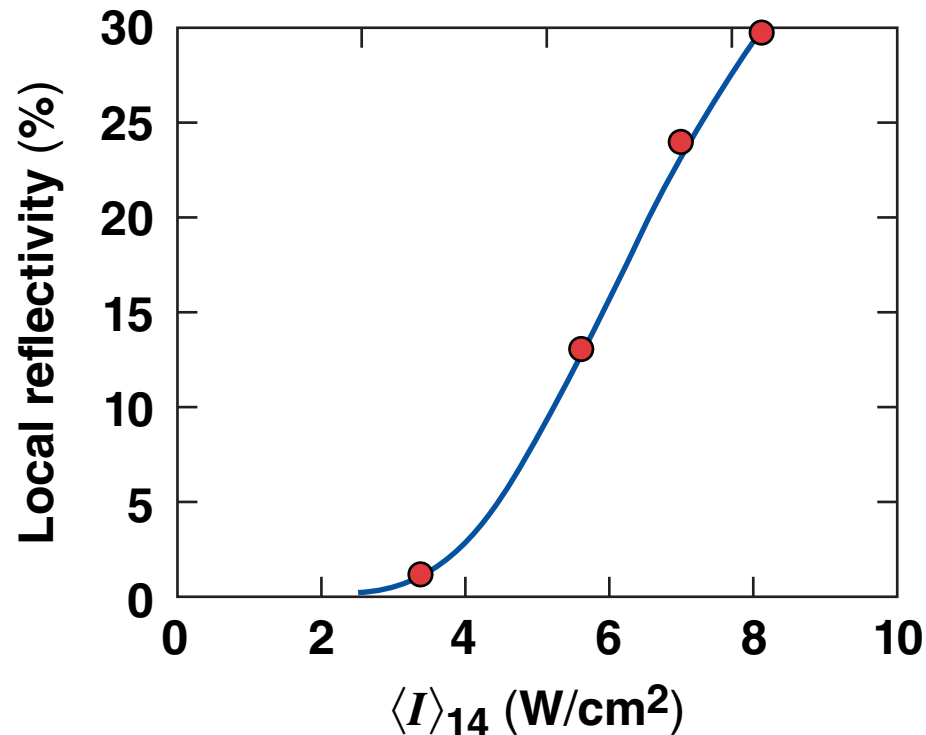


Energy Transfer Between Crossing Laser Beams in the Plasmas of Direct-Drive ICF



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Summary

In direct-drive ICF plasmas, nonlinear interaction between crossing laser beams leads to significant scattering

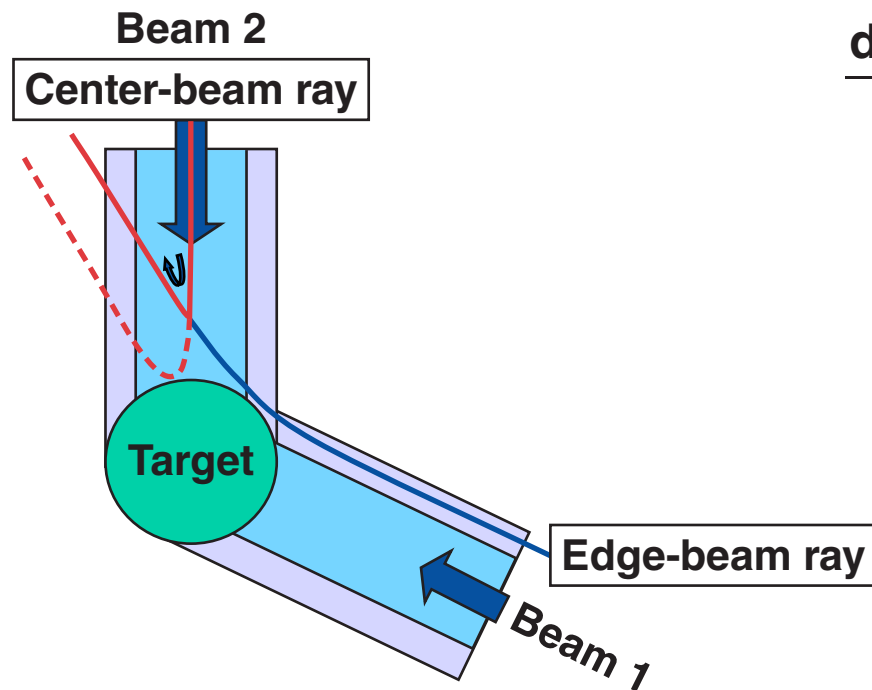


- **At moderate plasma densities (0.3 to 0.6 critical density), interaction between incident laser beams and a counter-propagating seed can lead to a local reflectivity exceeding 20%**
- **The direction of scattered light is determined by the structure of laser speckles**
- **Interaction in intense speckles leads to a power scaling of reflectivity with intensity**
- **Nonlinear propagation of laser beams near their turning points provides a seed for scattering in the lower-density region**

In large-scale hydrodynamic simulations, crossed-beam energy transfer is shown* to significantly influence the laser absorption

- For direct-drive ICF plasmas, the interaction between rays is

$$\frac{dI_1}{d\ell} = I_1 I_0 \frac{\omega_0^2}{2c^2 n_c} \operatorname{Re} \left\{ \frac{n_e k_s^2 c_s^2}{2\nu_i \omega_s + i[(\omega_s + k_s v_0)^2 - k_s^2 c_s^2]} \times \frac{1}{2k_{0x}} \right\}$$



- Crossed-beam energy transfer reduces the energy of incoming center-beam light and increases the energy of outgoing edge-beam light

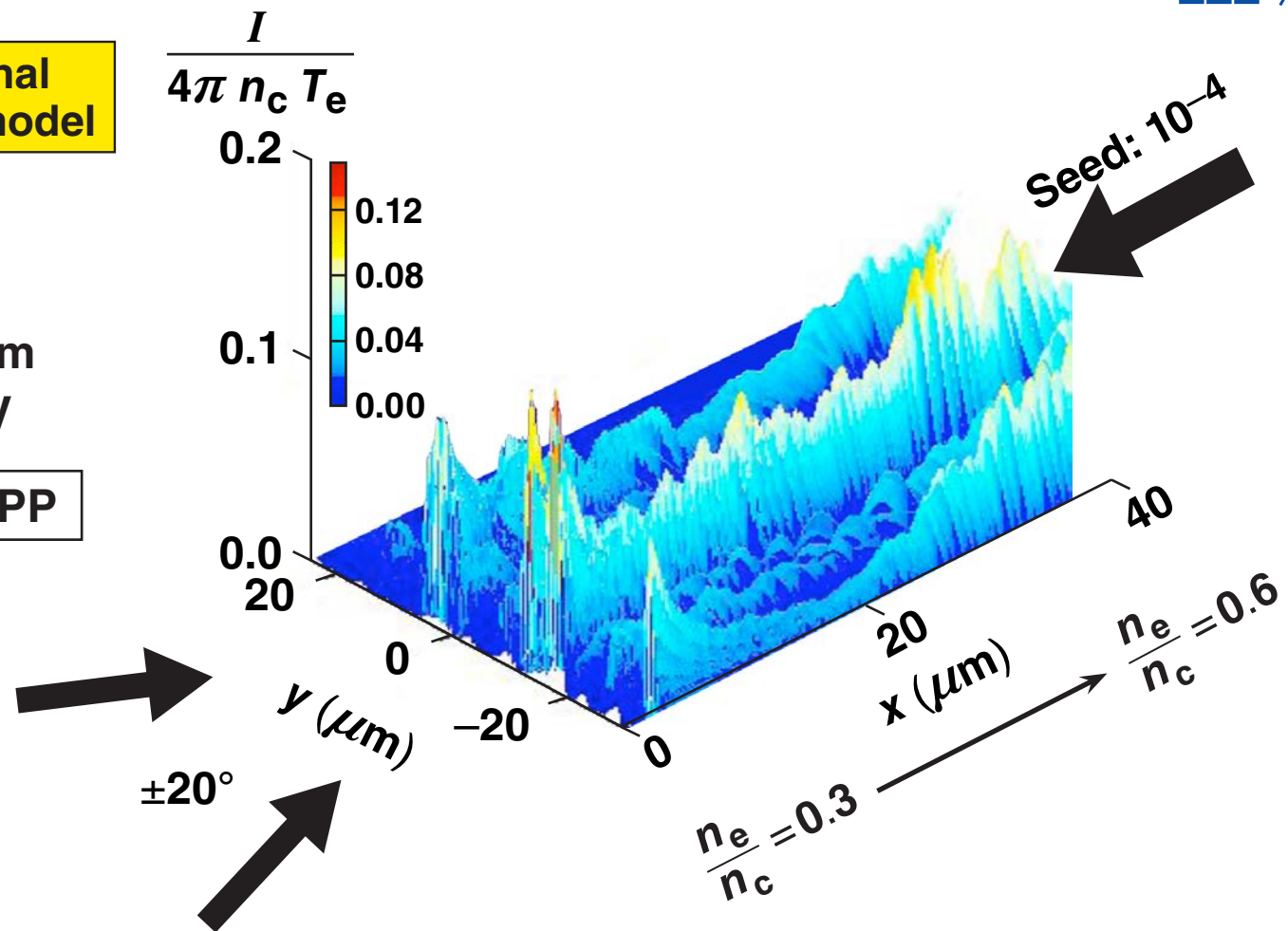
*I. V. Igumenshchev *et al.*, Phys. Plasmas **17**, 122708 (2010).
I. V. Igumenshchev *et al.*, YI3.00001

The nonlinear propagation of crossing laser beams has been modeled in the region of moderate plasma density, about $0.3 n_c$ to $0.6 n_c$

Two-dimensional
non-paraxial model

$L_n = 140 \mu\text{m}$
 $T_e = 2 \text{ keV}$

$f/6$ or $f/15$ DPP



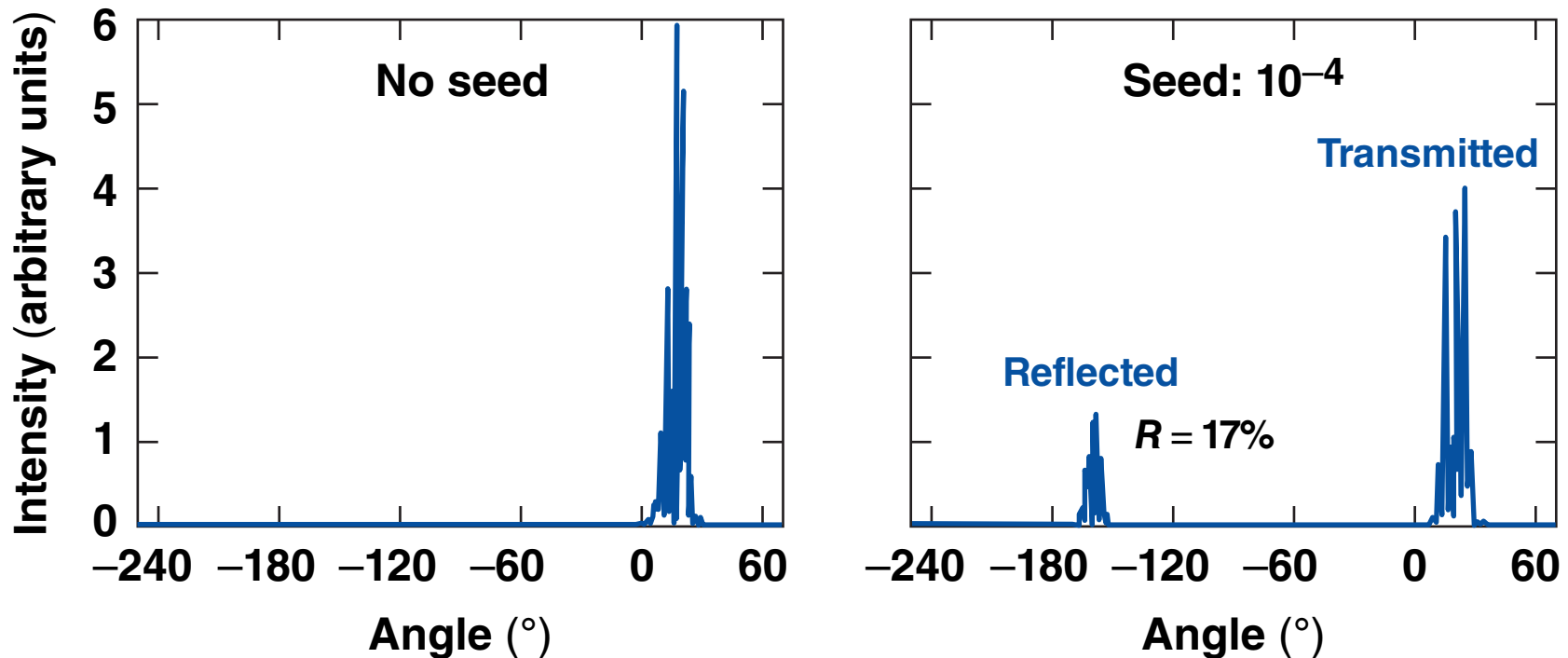
The backscatter depends on the electromagnetic seed, which is caused by opposing beams or turning beams



Scattering from a single DPP beam*

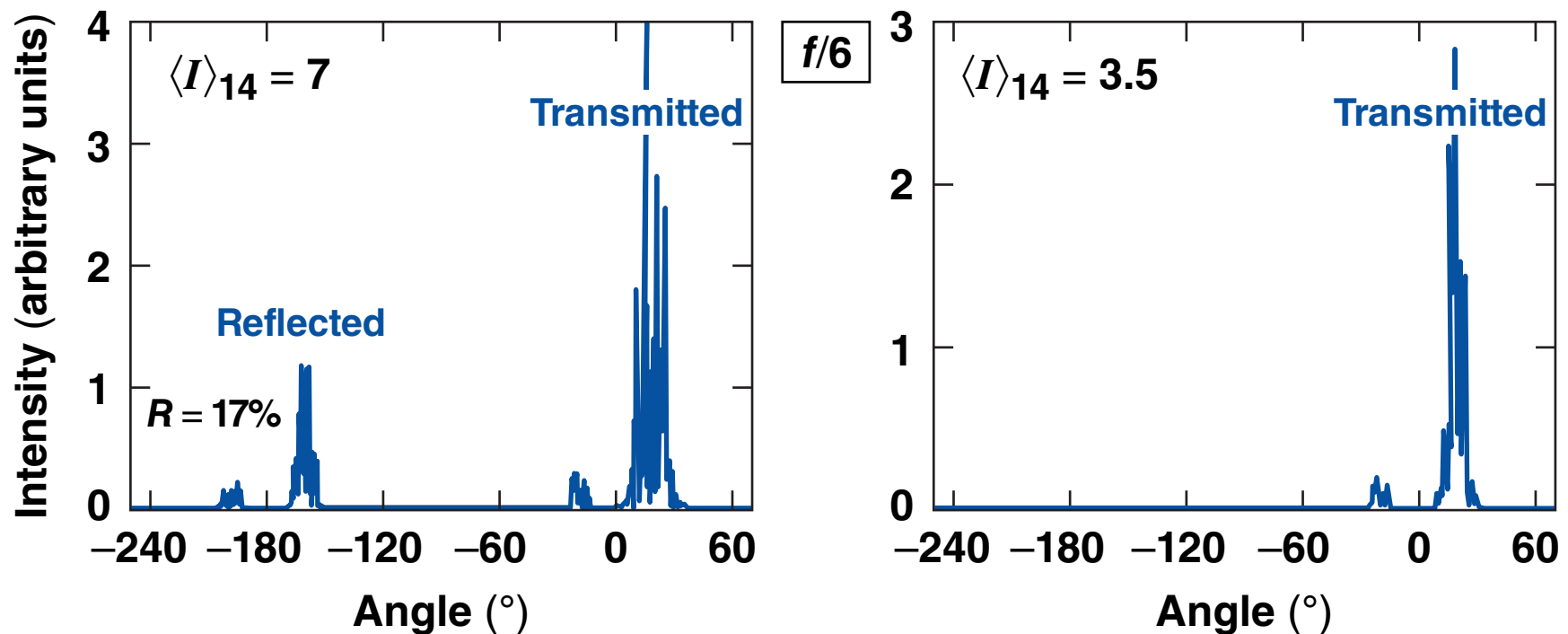
$f/6$

$\langle I \rangle_{14} = 7$



The threshold for the backscattering driven by crossing laser beams has been found at moderate laser intensities

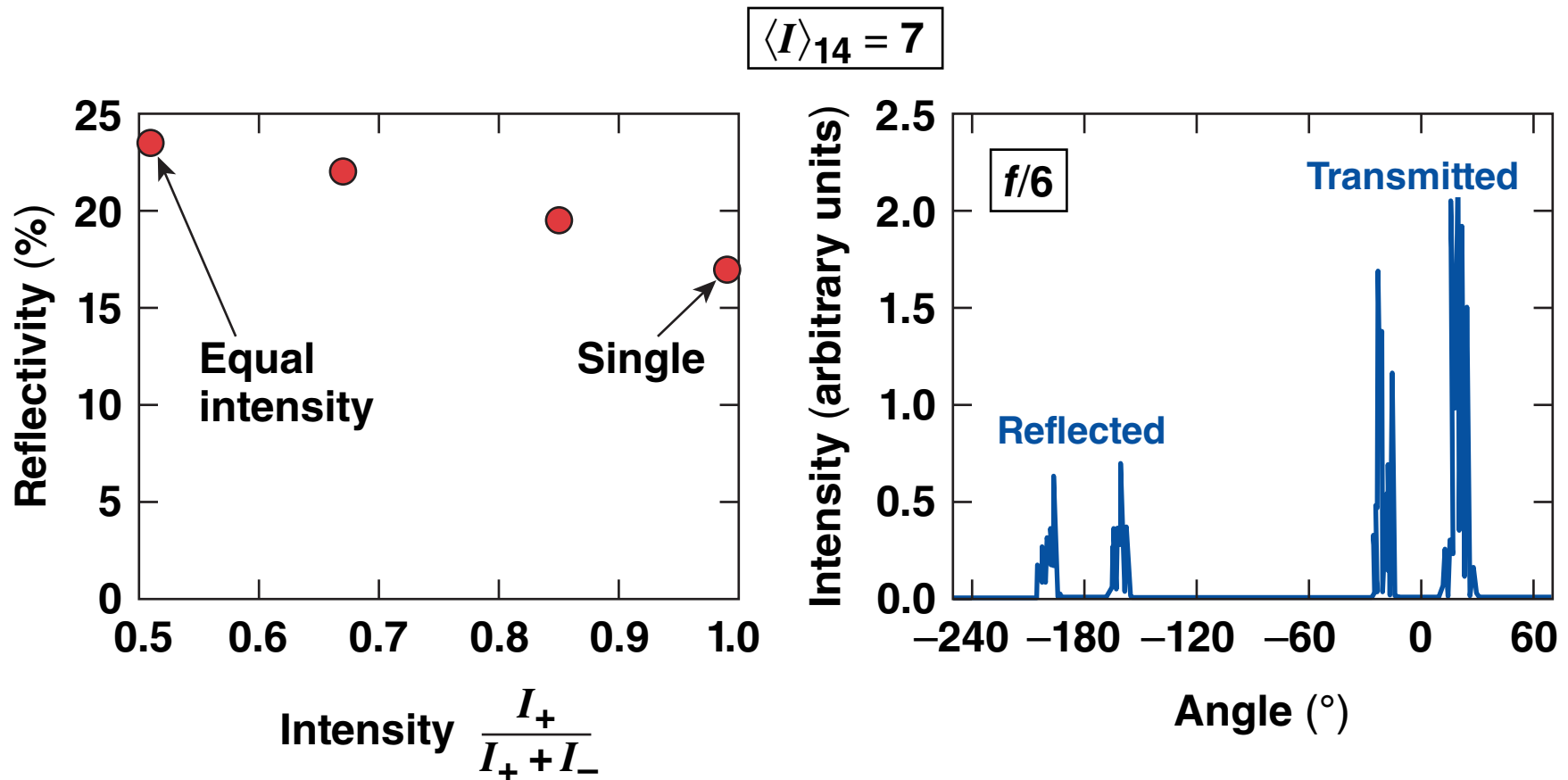
- The intensities of the two driving beams are different by a factor of 10.



$$G_{\text{SBS}} = 0.24 \langle I \rangle_{14} \left(\frac{I_{\text{max}}}{\langle I \rangle} \right)$$

Interaction in intense hot spots

The reflectivity has a moderate dependence on the distribution of intensity between the driving laser beams

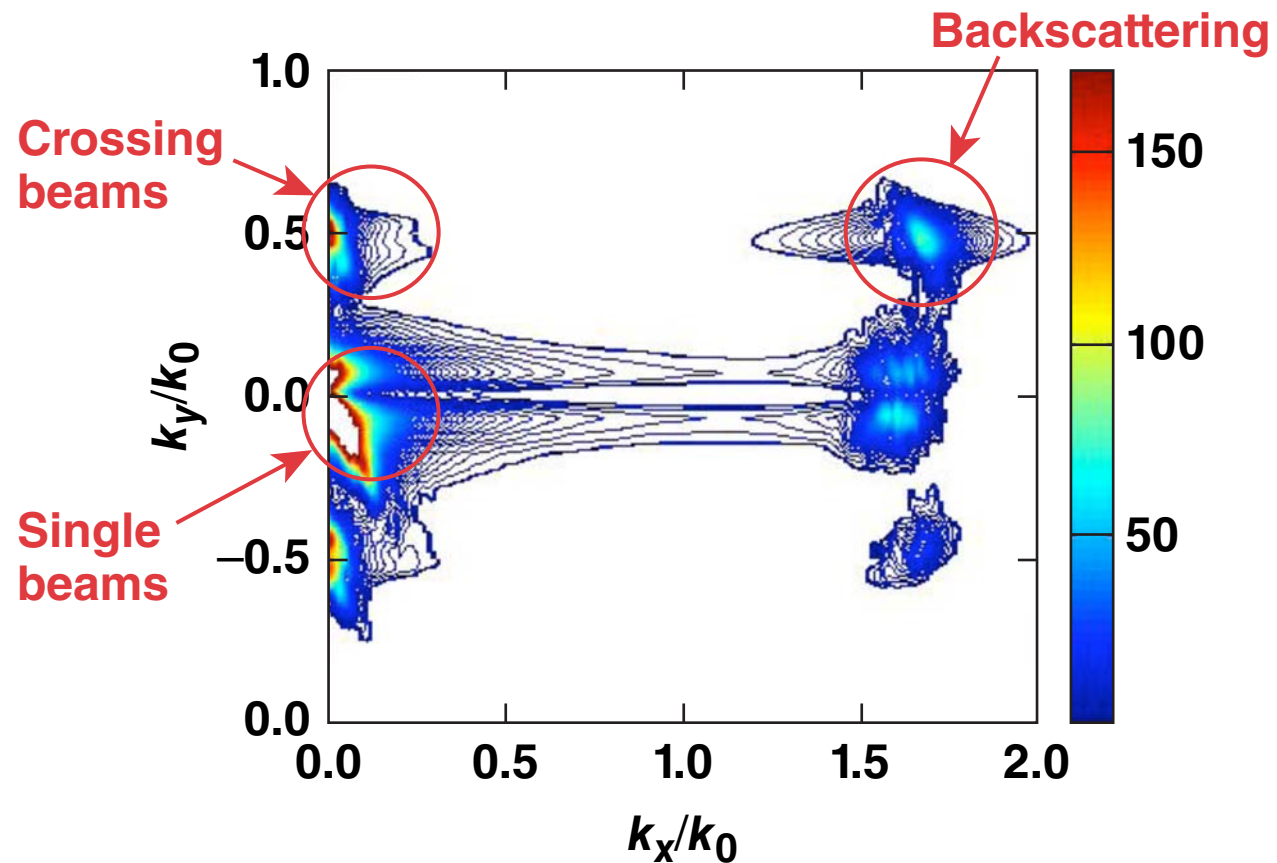


The hot-spot structure determines the direction of scattered light.

The interaction of incoherent crossing laser beams with plasmas produces a broad spectrum of low-frequency density perturbations

$L_n = 140 \mu\text{m}$
 $T_e = 2 \text{ keV}$

$f/6$



Laser beams can share density perturbations.

Nonlinear interaction in intense laser speckles determines the scaling of reflectivity with intensity

Reflectivity

$$\frac{d\langle R \rangle}{dx} \sim U_m^3 e^{-U_m},$$

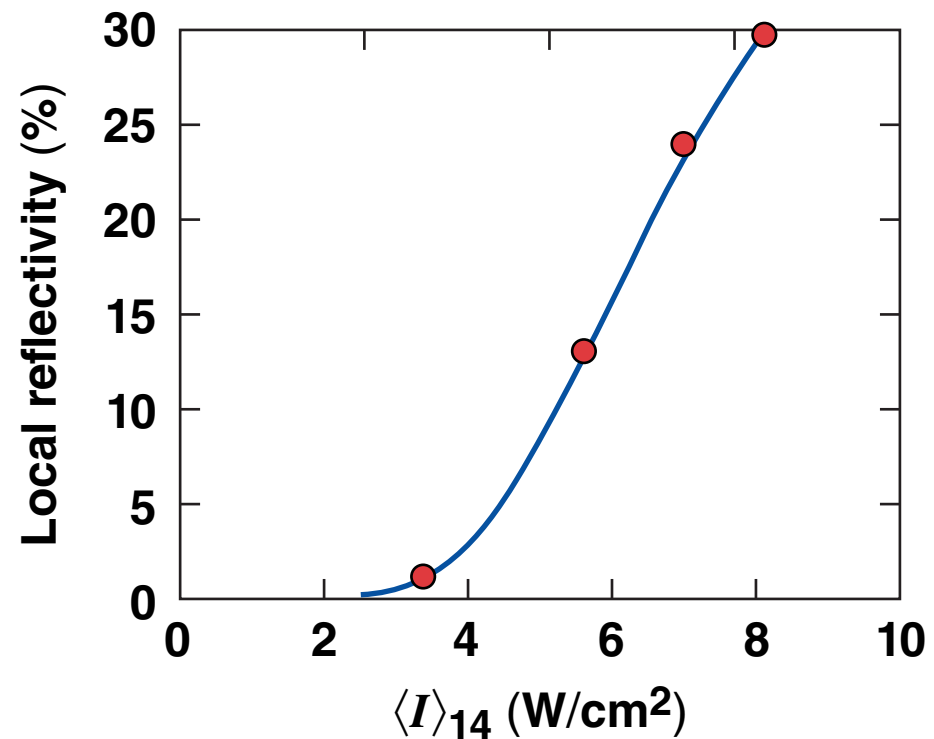
where

$$U_m \equiv \frac{I_m}{\langle I \rangle}$$

for the saturation R_{sat}

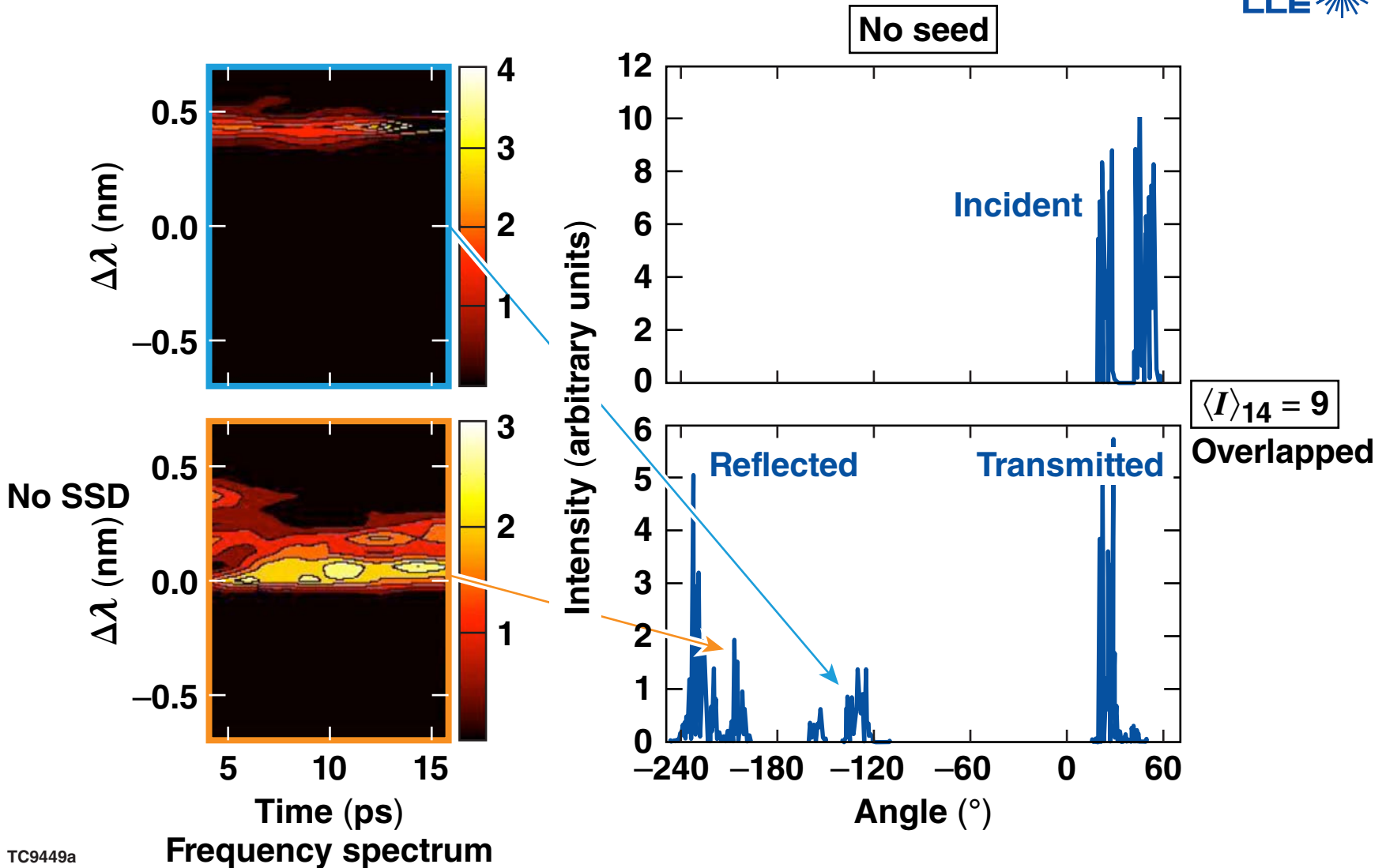
$$R_{\text{sat}} = \varepsilon e^{\langle G_{\text{SBS}} \rangle U_m}$$

ε – seed



The reflectivity scaling is influenced by the hot-spot structure.

At moderate plasma densities, the interaction between beams incident at different angles leads to a broad spectrum of backscattered light



Summary/Conclusions

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