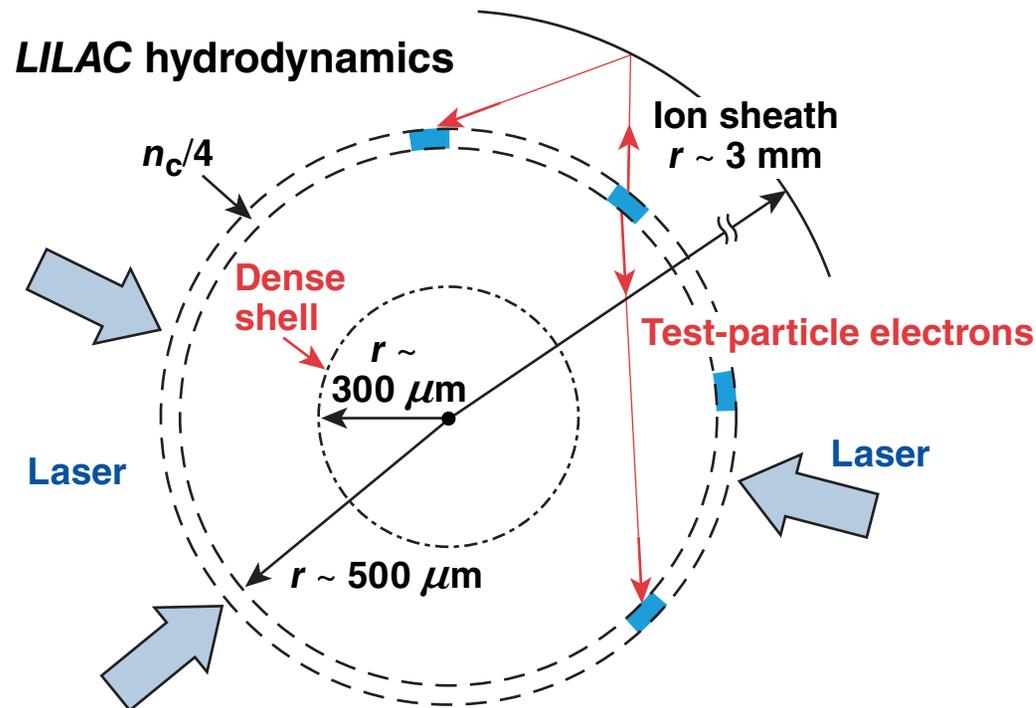


Two-Plasmon-Decay Calculations for OMEGA and Ignition-Scale Direct-Drive Inertial Confinement Fusion



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Summary

A physical model of TPD in direct-drive targets that gives scalings for hot-electron temperature has been constructed



- The scheme combines *LILAC* hydrodynamics, an extended Zakharov model of TPD growth and saturation, and hot-electron generation and transport in the test-particle approximation
- The scaling of hot-electron temperature is obtained for electron-density scale lengths in the range $L_n = 100$ to $350 \mu\text{m}$ and laser intensities between $I_{15} = 0.1$ to 2 that corresponds to either OMEGA- or ignition-scale targets
- The hot-electron temperature is well determined by the value of the threshold parameter η^* alone and is found to vary in the range $T_e = 50$ to 150 keV

Collaborators



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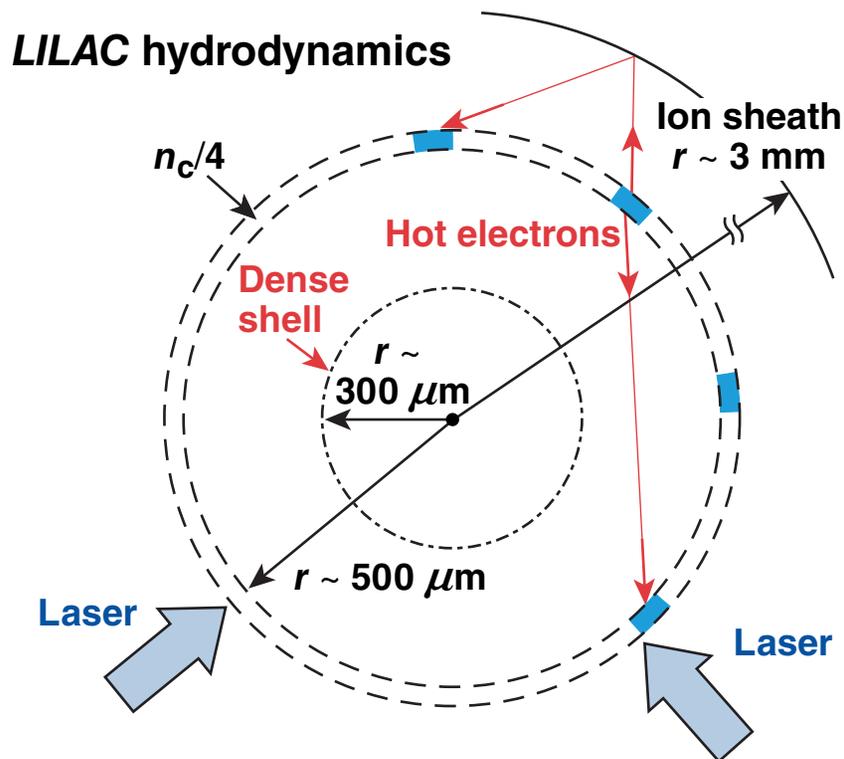
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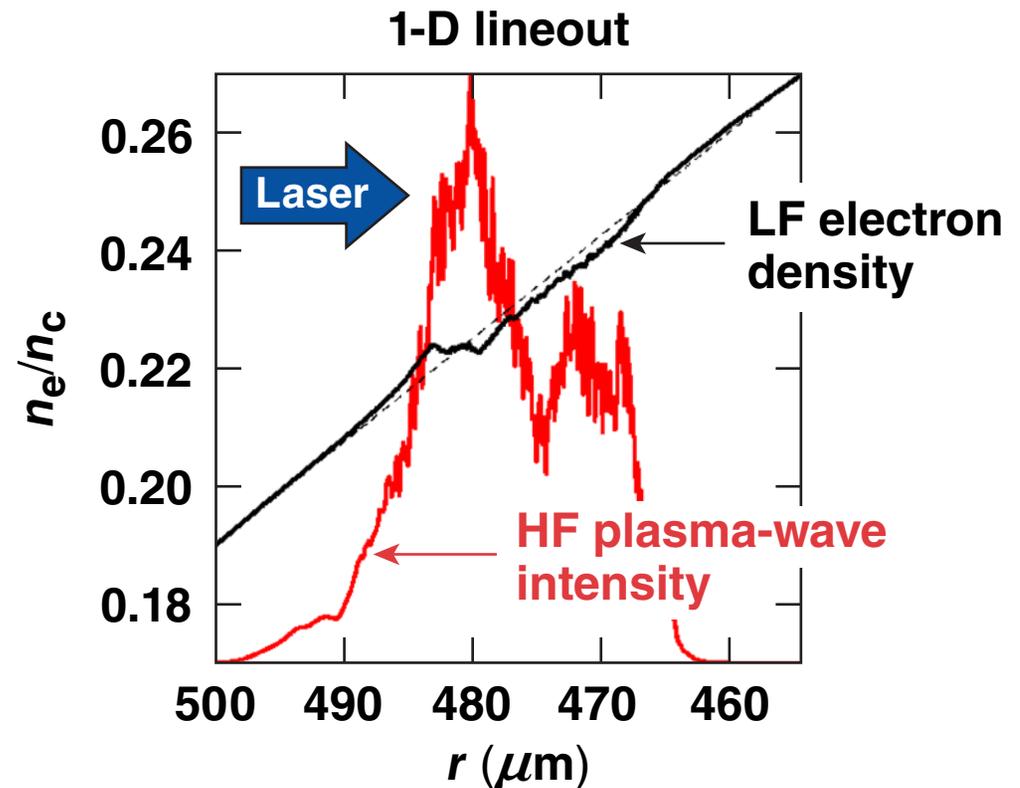
The model combines the effects of target hydrodynamics, TPD saturation, hot-electron production, and transport



- For a given implosion, *LILAC* calculations define the plasma parameters
- The extended Zakharov equations are solved to saturation
- Electron test particles are integrated to give heated distribution function
- Test particle transport is included

The nonlinear evolution of TPD is determined using an extended Zakharov model

- For a given implosion, *LILAC* calculations are used to determine the laser intensity and plasma parameters at quarter critical
- The Zakharov equations adopt these parameters and are run to full saturation (tens of picoseconds)



The “Zakharov” equations are extended when applied to the two-plasmon-decay problem

- “Extended” Zakharov equations used in Zak*

$$\nabla \cdot \left[D_{\text{LW}} - \omega_0^2 (\delta n + \delta N) / n_0 \right] \mathbf{E} = \left(e / 4 m_e \right) \nabla \cdot \left[\nabla (\mathbf{E}_0 \cdot \bar{\mathbf{E}}) - \mathbf{E}_0 \nabla \cdot \bar{\mathbf{E}} \right] + \mathbf{S}_E$$

$$D_{\text{IAW}} \delta n = \nabla^2 |E|^2 / (16 \pi m_i) + \mathbf{S} \delta n$$

TPD source term

Dispersion relations
for LW and IAW

Wave envelopes

$$D_{\text{LW}} = \left[2i\omega_{p0} (\partial_t + \nu_e^*) + 3\nu_e^2 \nabla^2 \right] \quad \tilde{E} = 1/2 E(x, y, t) \exp \left[-i(\omega_{p0} t) \right] + \text{c.c.}$$

$$D_{\text{IAW}} = \left(\partial_t^2 + 2\nu_i^* \partial_t - c_s^2 \nabla^2 \right) \quad \tilde{E}_0 = \mathbf{e}_y \sum_i |E_{0i}| \exp \left[i \vec{k}_{0i} \cdot \vec{x} - i(\omega_0 - 2\omega_{p0}) t \right]$$

* D. F. DuBois *et al.*, Phys. Rev. Lett. **74**, 3983 (1995);
D. A. Russell and D. F. DuBois, Phys. Rev. Lett. **86**, 428 (2001).

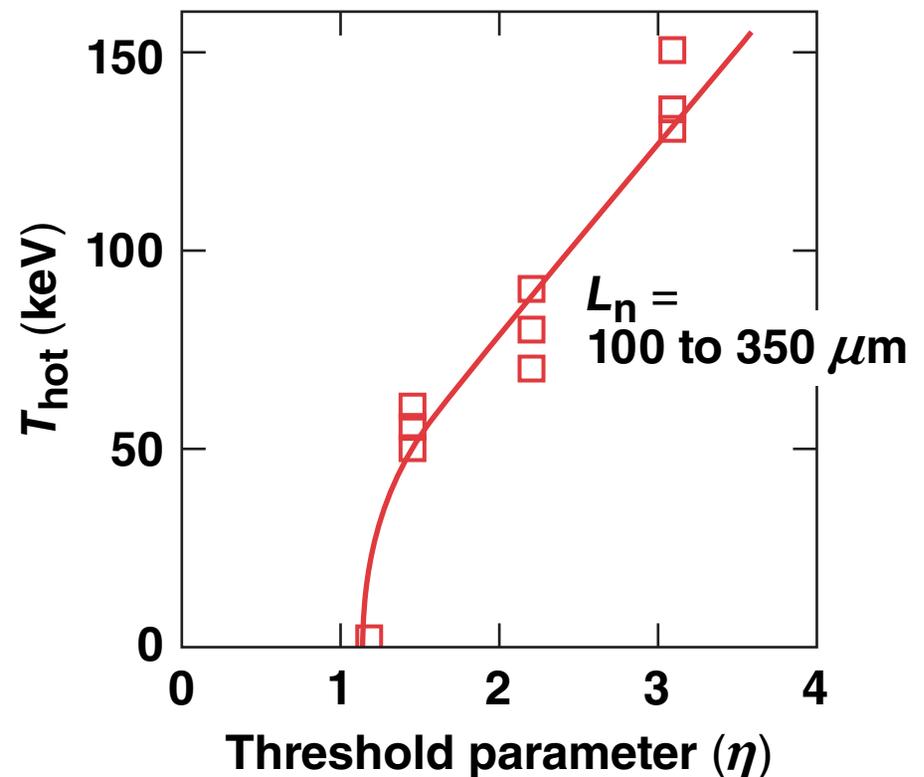
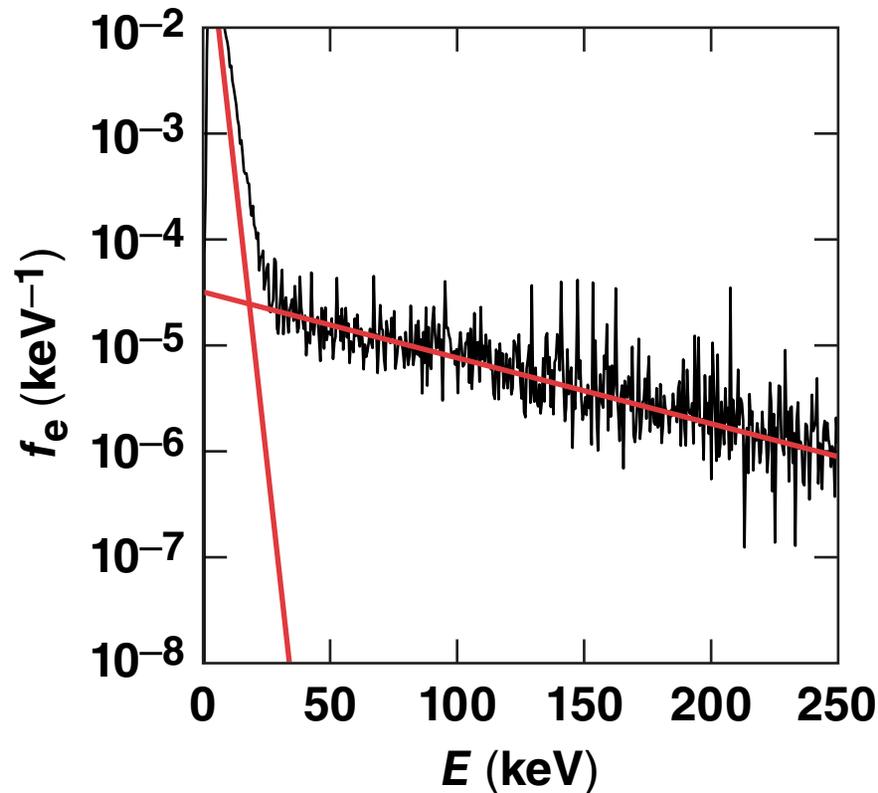
Hot-electron generation is computed in the test-particle approximation



- A large ensemble of electron test-particle trajectories are run to determine the hot-electron distribution function by randomizing over the initial position while drawing the initial energy from a thermal distribution
- Inside the Zakharov simulation volume, the electron trajectories are integrated in the microscopic electrostatic fields
- Outside of the simulation volume, the trajectories are subject to specular reflection at the ion sheath and energy loss caused by collisional stopping in the target
- These effects are implemented as boundary conditions¹

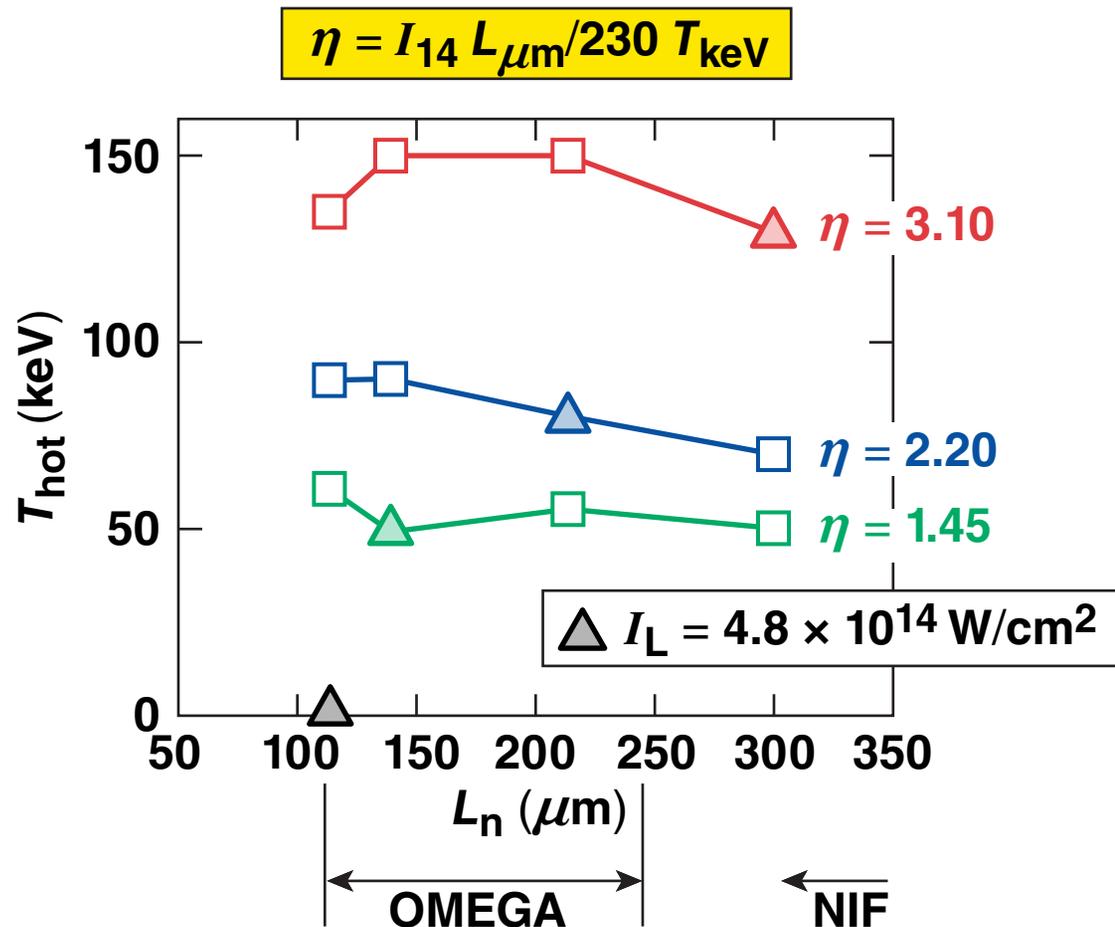
Heated electrons are observed once the TPD threshold is exceeded and the temperature scales with the “ η ” parameter

The threshold parameter* is given by $\eta = I_L/I_{thr} = I_{14}L_{\mu m}/(230 T_{keV})$.



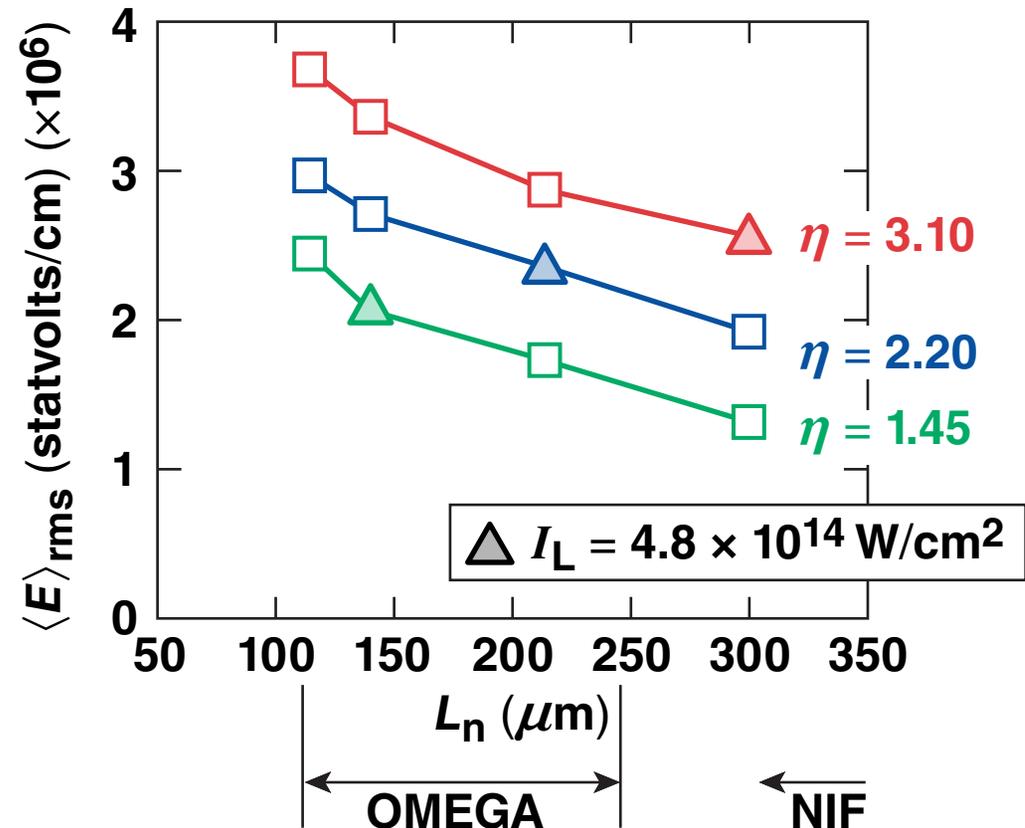
For a fixed value of the threshold parameter η , the scaling of hot-electron temperature with electron-density scale length is weak

- For all cases, the initial electron temperature $T_e = 2$ keV
- As L_n is increased, I_L is decreased to keep the product constant



The weak dependence of the hot-electron temperature on density scale-length arises due to the reduction in rms electric field

- The hot-electron temperature depends upon the electric field excitation level and the acceleration length



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