Simulations of Self-Generated Magnetic Fields in Implosion Experiments on OMEGA



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Summary

Nonuniformities in directly driven ICF targets generate magnetic fields that affect target performance

- Self-generated magnetic fields are developed at perturbed ablation surfaces and material interfaces. These fields are not dynamically important (P_m << P_g), but suppress heat fluxes.
- Magnetic fields of several MG originate in plasma corona from various sources of nonuniformities
 - laser imprint: the fields developed at material interfaces can increase perturbations and reduce the laser drive
 - dust particles: the field effects significantly modify perturbations induced by the particles
- Magnetic fields up to 30 MG are generated at peak compression in hot spots of cryogenic targets and up to 100 MG in hot spots of warm targets

Collaborators



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Magnetic fields are generated mainly at perturbed ablation surfaces and material interfaces



The field strength is limited by magnetic diffusivity.

Self-generated magnetic fields and magnetic-dependent heat fluxes are calculated using the Braginskii formulation*

• The induction equation:

$$\frac{\partial \vec{\mathbf{B}}}{\partial t} = \nabla \times (\mathbf{V} \times \vec{\mathbf{B}}) + \frac{\mathbf{c}}{\mathbf{e}} \left[\nabla \times \frac{\nabla \mathbf{P}_{\mathbf{e}}}{n_{\mathbf{e}}} - \nabla \times \frac{(\nabla \times \vec{\mathbf{B}}) \times \vec{\mathbf{B}}}{4\pi n_{\mathbf{e}}} - \nabla \times \frac{\vec{\mathbf{R}}_{\mathbf{T}} + \vec{\mathbf{R}}_{\mathbf{L}}}{n_{\mathbf{e}}} \right]$$

• The electron heat flux:
$$\vec{q}_e = \vec{q}_T^e + \vec{q}_u^e$$

The field affects the heat flux when $\Omega_{Le} \tau_e \ge 1$.

 The MHD equations with thermal conduction terms are solved using the modified 2-D hydrocode DRACO**

^{*}S. I. Braginskii, in Reviews of Plasma Physics (Consultants Bureau, NY, 1965). **P. B. Radha et al., Phys. Plasmas 12, 032702 (2005).

Perturbations from laser imprint are the source of self-generated magnetic fields

Cryogenic target in the middle of CH is all ablated: magnetic the main pulse: magnetic fields at fields at CH–D interface CH ablation surface* t = 3.2 ns $t = 3.5 \, \text{ns}$ B (MG) B (MG) 1.5 4 0.0 1 -1.5 **Density** Density (g/cm^3) (g/cm^3) 4 2 2 0 0 100 200 300 0 Average Z CH D $R(\mu m)$ These fields at ablation surface have a small effect on target performance. 0 *J. R. Rygg et al., Science 319, 1223 (2008).





Radially extended magnetic fields at perturbed CH–D interface efficiently suppress the inward heat flux

The temperature change in the corona due to magnetic fields



 The laser drive is reduced due to suppression of the heat flux

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• The fields enhance perturbations with *L*-modes ~20, which are the most dangerous modes at the deceleration phase

For best performance, plastic shells should all be ablated only after the end of the laser pulse.

Dust particles on the target surface are the source of localized MG field structures

 The surfaces of warm and cryogenic targets are typically contaminated by several tens of ~10-μm-sized dust particles

Cryogenic DT target with a 10- μ m dust particle





Perturbations induced by a 10- μ m dust particle evolve differently with and without magnetic fields



- *R-T* instability develops a hole and some mass is injected inside the shell
- Magnetic fields thermally isolate the perturbed area resulting in a locally reduced drive

Simulations predict that magnetic fields will reduce the damage to hot spots from dust particles



• Up to 30-MG fields are generated in hot spots of cryogenic targets

• Larger fields up to 100 MG are generated in hot spots of warm targets due to the presence of plastic-fuel interfaces

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