Theoretical Investigation of Strong Coupling and Degeneracy Effects in ICF Implosions



Summary

An equation-of-state table of deuterium from firstprinciples calculations has been established and has indicated significance in ICF implosions

• We have established a first-principles equation-of-state (FP-EOS) table for deuterium, using the restricted path-integral Monte Carlo (PIMC) method, which covers typical ICF fuel conditions of $\rho \approx 0.002 \sim 673$ g/cm³ and $T \approx 1.35$ eV ~ 5.5 keV

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- The FP-EOS table shows discrepancies in pressures and energies in the moderate coupling (Γ ≥ 1) and degenerate (θ ≤ 1) regimes when compared to the SESAME table
- Hydrodynamics simulations of cryogenic ICF implosions using the FP-EOS table have indicated significant differences in peak density, ρR , and neutron yield relative to SESAME-EOS simulations



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Low-adiabat ICF implosion designs routinely access to strongly coupled and degenerate plasma conditions!



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The SESAME-EOS table is based upon the "chemical model" of matter*, which does not fully account for strong-coupling and degeneracy effects

The essential physical models in the SESAME-EOS table:

- The molecular solid phase: is valid only in low-T and not-too-high density ($\rho \le 4 \text{ g/cm}^3$) regimes where molecular identity still exists.
- <u>The atomic solid phase:</u> The zero-Kelvin cold curve used in the atomic-solid model is still uncertain*.
- <u>The molecular/atomic fluid phase</u>: The first-order expansion was used in the liquid perturbation theory to only account for the nearest neighbor interactions.
- Strong-coupling and many-body degeneracy effects are not fully accounted.

The first-principles method of path-integral Monte Carlo (PIMC) was used to calculate the FP-EOS of deuterium

For a many-body system in thermal equilibrium at temperature *T*, the density matrix *ρ*(*R*,*R*';*T*) uniquely contains all thermodynamic information of the system; *ρ*(*R*,*R*';*T*) is defined as (in configuration representation)

$$\rho(\mathbf{R},\mathbf{R}';\mathbf{T}) = \left\langle \mathbf{R} \left| \mathbf{e}^{-\vec{\mathbf{H}}/k\mathbf{T}} \right| \mathbf{R}' \right\rangle = \sum_{n} \varphi_{n}(\mathbf{R}) \varphi_{n}(\mathbf{R}') \mathbf{e}^{-\mathbf{E}_{n}/k\mathbf{T}} \qquad \left(\vec{\mathbf{H}} \varphi_{n} = \mathbf{E}_{n} \varphi_{n}; \vec{\mathbf{H}} = \vec{\mathbf{K}} + \vec{\mathbf{V}} \right)$$

• Using the convolution property of $\rho(R,R';\beta)$, one can compute the density matrix from high-T ρ_0

$$\rho(\mathbf{R}, \mathbf{R}'; \boldsymbol{\beta}) = \int d\mathbf{R}_1 d\mathbf{R}_2 \dots d\mathbf{R}_{M-1} \rho_0(\mathbf{R}, \mathbf{R}_1; \Delta \boldsymbol{\beta})$$
$$\times \rho_0(\mathbf{R}_1, \mathbf{R}_2; \Delta \boldsymbol{\beta}) \dots \rho_0(\mathbf{R}_{M-1}, \mathbf{R}'; \Delta \boldsymbol{\beta})$$

with
$$\rho_0(\mathbf{R}_i, \mathbf{R}_{i+1}; \Delta \beta) = \exp\left[\frac{-\pi}{\lambda_\Delta^2}(\mathbf{R}_i - \mathbf{R}_{i+1})^2 + \Delta \beta \times \mathbf{V}(\mathbf{R}_i)\right], \beta = 1/kT, \Delta \beta = \beta/M,$$

and $\lambda_D = \sqrt{2\pi \hbar^2 \Delta \beta / m}$ is the DeBroglie wavelength

• Once we know $\rho(R,R';\beta)$, we can calculate the thermodynamic properties of the system with the corresponding operators \hat{O} :

$$\langle \hat{\mathbf{O}} \rangle = \frac{\int d\mathbf{R} d\mathbf{R}' \left\langle \mathbf{R} | \hat{\mathbf{O}} | \mathbf{R}' \right\rangle \left\langle \mathbf{R}' | \boldsymbol{\rho}(\boldsymbol{\beta}) | \mathbf{R} \right\rangle}{\int d\mathbf{R} \left\langle \mathbf{R} | \boldsymbol{\rho}(\boldsymbol{\beta}) | \mathbf{R} \right\rangle}$$

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The principal Hugoniot comparison shows that deuterium is slightly softer under ~2 Mbar, but stiffer in FP-EOS than SESAME for ~2 Mbar < *P* < ~100 Mbar

*B. Militzer and D. M. Ceperley, Phys. Rev. Lett. <u>85</u>, 1890 (2000).

In comparison to the SESAME table, the FP-EOS table has shown large differences, especially in internal energy

Hydro-simulations of low-adiabat implosions using the FP-EOS has indicated significant differences in ρ_{peak} , ρR , and neutron-yield compared to SESAME simulations

Simulations for NIF direct-drive designs also show significant differences in peak density, yield, and gain

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