#### Modeling of Two-Plasmon-Decay Instability in OMEGA Plasmas



A. V. Maximov, J. Myatt, R. W. Short, W. Seka, C. Stoeckl, and J. A. Delettrez University of Rochester Laboratory for Laser Energetics 50th Annual Meeting of the American Physical Society Division of Plasma Physics Dallas, TX 17–21 November 2008

Summary

### The multispecies composition of plasmas has a strong effect on the two-plasmon-decay instability

- The multispecies composition modifies the ion-ion collision properties and, therefore, the ion heat conductivity and ion viscosity.
- The presence of a high-Z dopant in the multispecies plasma changes the dispersion relation for the ion-acoustic modes, and decreases the ion-wave damping.
- The increased level of ion-acoustic perturbations limits the growth of the two-plasmon-decay instability.

## In plasmas with multispecies ions, the ion–ion collision frequency includes collisions with all ion species

• Collision integral  $\delta J_{i\Sigma} = \sum_{A} \delta J_{iA}, \left( \delta J_{iA} \sim \frac{\mathbf{e}_{i}^{2} \mathbf{e}_{A}^{2}}{T_{i}^{3/2} m_{eff}^{1/2}} \right) \text{ and } \nu_{i\Sigma} = \frac{4 \ln \Lambda}{3 \sqrt{\pi}} \frac{\mathbf{e}_{i}^{2}}{T_{i}^{3/2} \sum_{A} \left( \frac{n_{A} \mathbf{e}_{A}^{2}}{m_{eff}^{1/2}} \right)$ The ion heat flux  $\vec{q}^{\Sigma} = \sum_{i} \vec{q}^{i} = -\kappa_{\Sigma} \frac{\partial T_{i}}{\partial \vec{r}}$ Ion heat conductivity\*  $\kappa_{\Sigma} = \sum_{i} 3.9 \frac{n_{i} T_{i}}{m_{i} \nu_{i\Sigma}} \sim \frac{\sum_{i} \left( n_{i} / \mathbf{e}_{i}^{2} m_{i} \right)}{\sum_{i} \left( n_{i} \mathbf{e}_{i}^{2} / m_{eff}^{1/2} \right)} T_{i}^{5/2}$ 

• Ion viscosity\* 
$$\eta^{\Sigma} = \sum_{i} 0.96 \ \frac{n_{i}T_{i}}{\nu_{i\Sigma}} \sim \frac{\sum_{i} \left(n_{i} / e_{i}^{2}\right)}{\sum_{i} \left(n_{i} e_{i}^{2} / m_{eff}^{1/2}\right)} T_{i}^{5/2}$$

\*S. I. Braginskii, in *Reviews of Plasma Physics*, edited by Acad. M. A. Leontovich (Consultants Bureau, New York, 1965), Vol. 1, p. 205.

# The coefficients for ion heat conductivity, and ion viscosity are significantly modified when the plasma composition changes

	$\frac{\nu_{\rm e\Sigma}}{\sqrt{2}\nu_{\rm ee}} = Z_{\rm eff}$	$\frac{\nu_T^{\Sigma}}{\sqrt{2}\nu_{ee}}$	$\frac{\kappa_{\Sigma}}{\kappa_{D_2}}$	$\frac{\eta^{\Sigma}}{\eta_{D_2}}$
D <sub>2</sub>	1.0	0.50 ( <i>m/m<sub>p</sub></i> )	1.000	1.000
CD	5.3	0.50 ( <i>m/m<sub>p</sub></i> )	0.032	0.027
CHSi(0.06)	7.1	0.56 ( <i>m/m<sub>p</sub></i> )	0.032	0.017

• The rate of energy transfer between electrons and ions

$$v_T^{\Sigma} = \sum_i 2 \frac{m}{m_i} v_{ei}$$

(~10% to 20% of ohmic heating for  $D_2$ )

### The influence of multispecies effects on the plasma parameters has been studied with *LILAC*





#### The dispersion equation for ion-acoustic waves in multispecies plasmas has several solutions

 $\frac{k^2 c_{s0}^2}{n_e} \sum_{i} \frac{Z_i^2 n_i / M_i}{(\omega - \vec{k} \vec{V})^2 + 2i\gamma_i (\omega - \vec{k} \vec{V}) - (5/3) k^2 V_{Ti}^2} = 1$ Fluid model

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- $\vec{V}$ -flow velocity;  $Z_i$  and  $M_i$ -ion charge and mass;  $c_{s0}^2 = (T_e/m_p)$
- In the collisional (i–i) regime, the ion-acoustic damping is determined by the ion viscosity and ion heat conductivity.

$$\gamma_i = k^2 \left( 0.64 + 0.87 \frac{V_{Ti}^2}{c_s^2} \right) \frac{V_{Ti}^2}{v_{i\Sigma}}$$



In OMEGA experiments, the hard x-ray production depends on the overlapped intensity of multiple laser beams and on the ion composition of plasmas



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#### Low-frequency perturbations in electron density are produced by the interaction of incoherent laser beams with plasmas



 $\langle I \rangle = 9 \times 10^{14} \text{ W/cm}^2, \ T_e = 2 \text{ keV}, \ n_0 \approx \frac{n_c}{4}, \text{CH} \qquad \left(\frac{n_e}{n_0} - 1\right) \sim \frac{I}{\langle I \rangle}$ 

### With the decrease of the ion-acoustic damping, the level of low-frequency density perturbations increases



CHSi

# The low-frequency perturbations in the electron density can detune the TPD resonance and reduce the TPD growth



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### The multispecies composition of plasmas influences the hydrodynamic equations

 $\partial_t \rho + \vec{\nabla}(\rho \vec{V}) = 0$ ion density  $n_{\Sigma} = \sum_{i} n_{i}$  | mass density  $\rho = \sum_{i} n_{i} m_{i}$  $\rho \Big[ \partial_t \mathbf{V}_j + (\vec{\mathbf{V}} \cdot \vec{\nabla}) \mathbf{V}_j \Big] = -\nabla_j \Big( n_e \mathbf{T}_e + n_{\Sigma} \mathbf{T}_j \Big) - \nabla_k \sigma_{jk}^{\Sigma} - \frac{e^2}{4m\omega_0^2} \Big\{ n_e \nabla_j |\mathbf{E}|^2 - \nabla_k \Big[ n_e \Big( \mathbf{E}_k \mathbf{E}_j^* + \mathbf{E}_j \mathbf{E}_k^* \Big) \Big] \Big\}$ lon Electron  $\partial_t T_e + (\vec{V} \cdot \vec{\nabla}) T_e + \frac{2T_e}{3} \vec{\nabla} \cdot \vec{V} + \frac{2}{3n_e} \vec{\nabla} \cdot \vec{q}^e = \frac{e^2 |E|^2 v_{e\Sigma}}{3m \omega_0^2} - \sum_i \frac{2m v_{ei}}{m_i} (T_e - T_i)$ Electron-lon  $\partial_t T_i + (\vec{\nabla} \cdot \vec{\nabla}) T_i + \frac{2T_i}{3} \vec{\nabla} \cdot \vec{\nabla} + \frac{2}{3n_{\Sigma}} \vec{\nabla} \cdot \vec{q}^{\Sigma} = \sum_i \frac{2m \nu_{ei}}{m_i} (T_e - T_i)$ 

• The ion temperatures of different species are assumed to be the same.

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