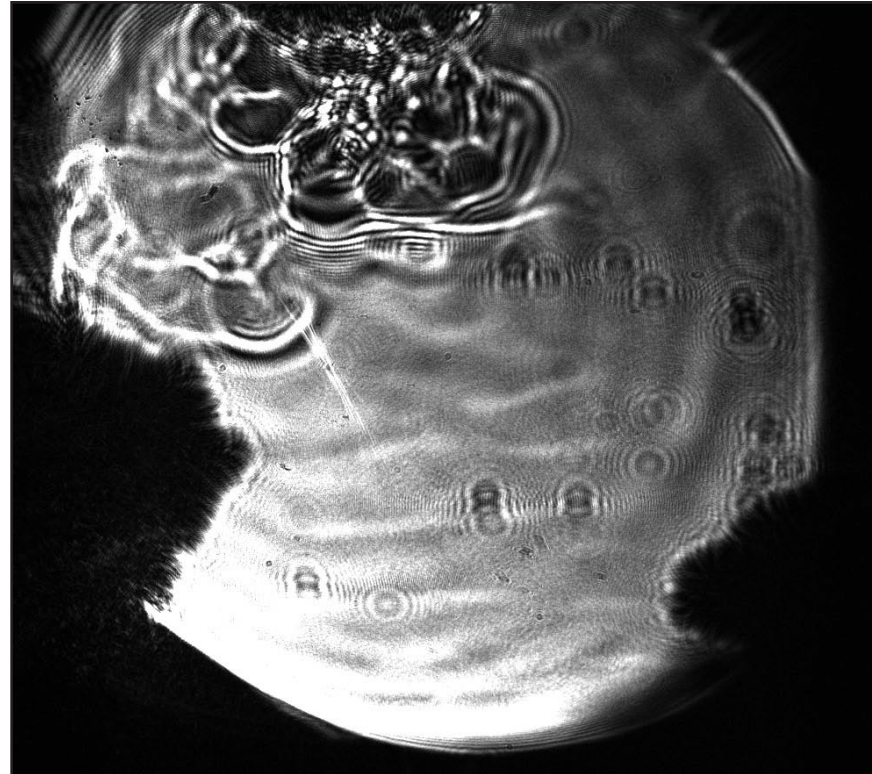


# First Tests on OMEGA of a Bubble Chamber for Neutron Detection



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## Summary

# The 14-MeV-neutron sensitivity of a freon-based bubble detector has been tested on OMEGA

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- The chamber has detected 14-MeV DT neutrons at yields of  $\sim 10^{13}$ .
- The measured sensitivity agrees with that calculated for neutron–freon bubble formation/growth.
- The sensitivity is too low for neutron imaging on OMEGA, but more than adequate for the higher neutron yields expected at the NIF.

# Collaborators

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# Neutron imaging can provide data that show why an ICF capsule fails to ignite<sup>1,2</sup>

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- Neutron images of ICF capsules provide a direct measurement of the fusion burn region within a compressed target.
- The radiation symmetry can be inferred from a neutron image of the hot-spot fusion region (where the fusion processes occur).
- Bubble chambers are detectors with a high potential in achieving high resolution neutron images<sup>3</sup>.

<sup>1</sup>L. Disdier *et al.*, Nucl. Instrum. Methods Phys. Res. A 489, 496 (2002).

<sup>2</sup>R. A. Lerche *et al.*, Rev. Sci. Instrum. 74, 1709 (2003).

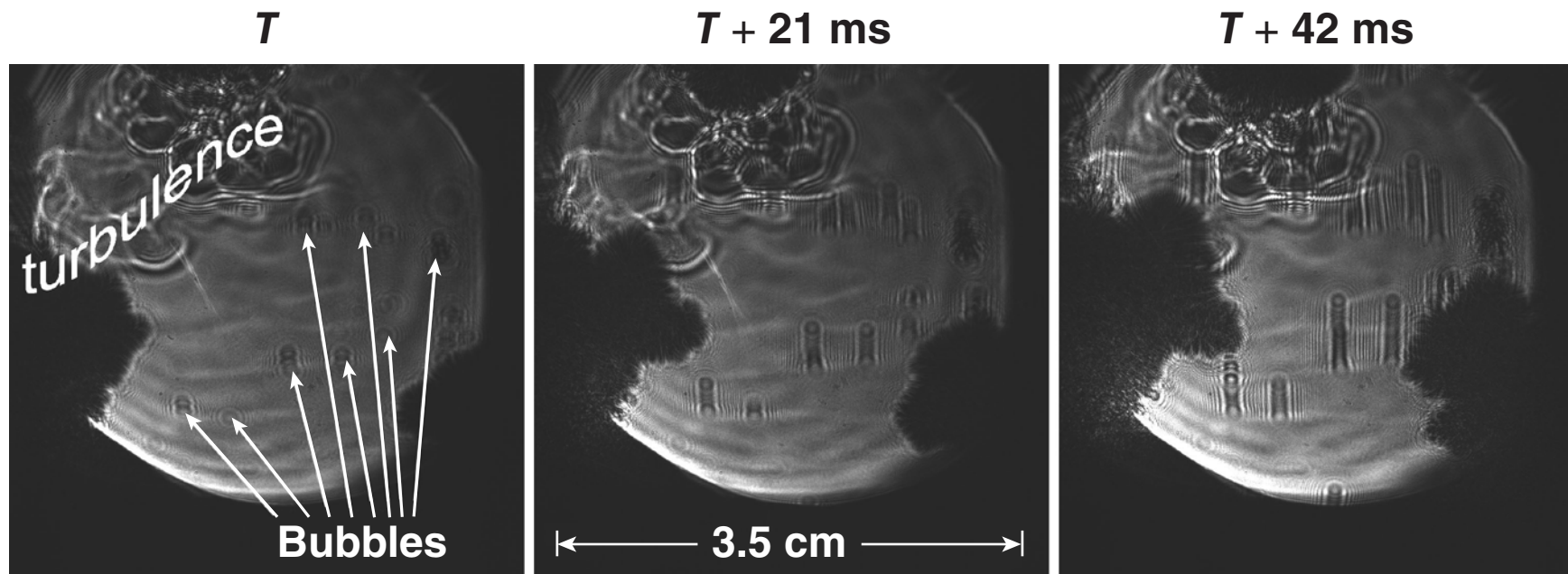
<sup>3</sup>R. K. Fisher *et al.*, Phys. Plasmas 9, 2182 (2002).

# The bubble chamber is a fully self-contained platform located in the OMEGA Target Bay



- Neutron interactions in the superheated freon create bubbles that are counted/imaged.
- The bubbles are detected in parallel, monochromatic light.
- For imaging, distribution of bubbles ~ neutron spatial distribution.

# The number of observed bubbles inside freon confirmed the theoretical calculations



- Successive images (21-ms difference) of neutron-induced cavitation inside the BUBDET: about 14 bubbles can be counted in the area not affected by turbulence.

# Nucleation occurs inside a bubble chamber when the deposited energy reaches a threshold value

- Thermodynamics of the superheated liquid gives the threshold energy to create a bubble:\*

$$W_{\text{bubble}} = \frac{16 \pi \gamma^3}{3(\rho_v - \rho_0)^2},$$

$\gamma$ : surface tension of the active medium

$\rho_v$ : co-existence phase pressure

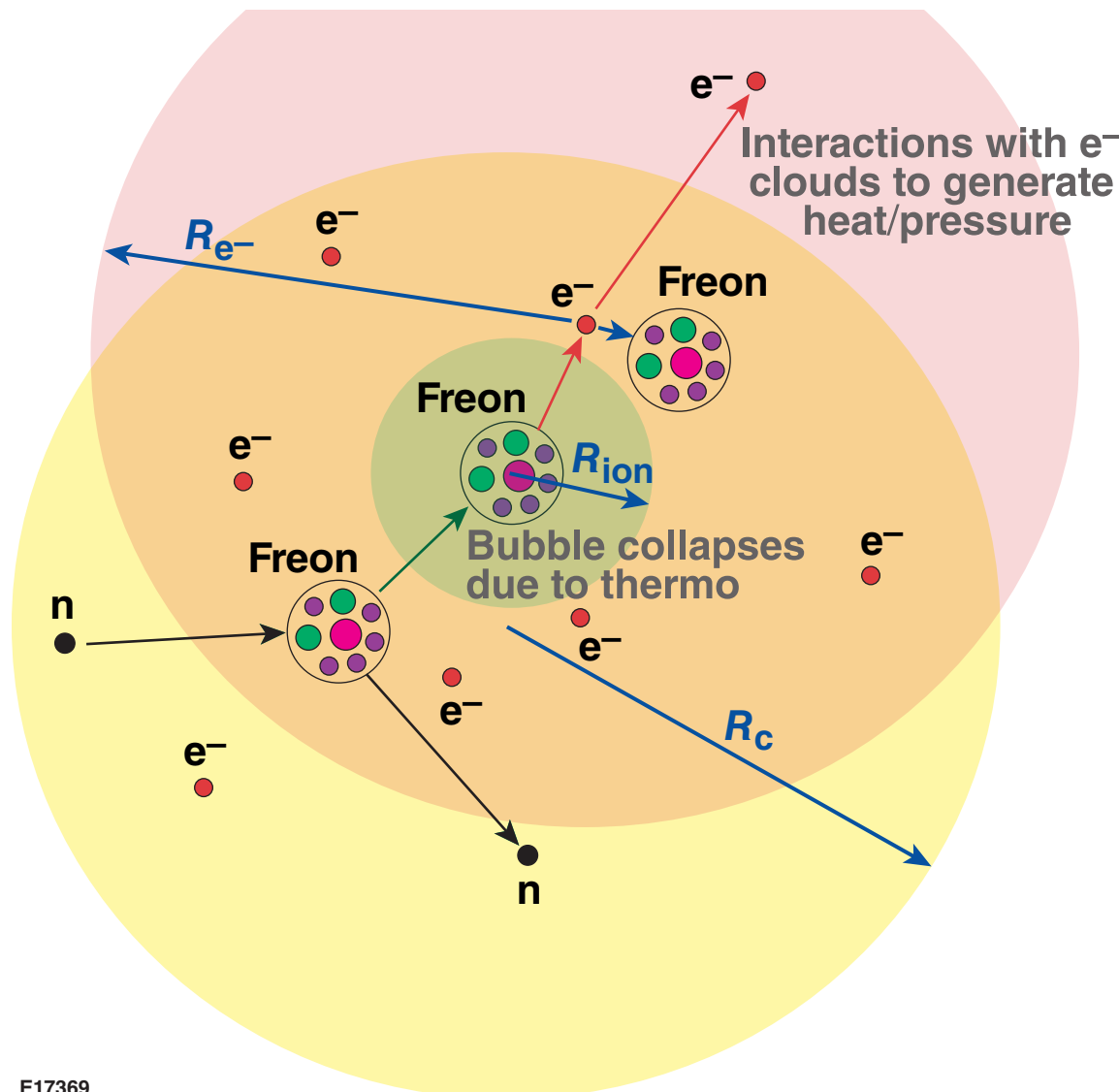
$\rho_0$ : superheated state pressure

- Maximizing the free energy necessary to form a bubble, the critical radius from which a bubble does not collapse but continues to grow is

$$R_c = \frac{2 \gamma}{\rho_v - \rho_0}.$$

- Therefore, for a thermodynamically viable bubble, the energy  $W_{\text{bubble}}$  must be deposited over a volume  $\sim R_c^3$ .
- However, the ion-recoil range for  $(n, \text{freon}) \rightarrow (n', \text{freon}')$  is  $\ll R_c$ .

# The sensitivity of bubble formation can be understood by examining the details of ion recoil



- Since the recoil-ion range is short, energy must be deposited in a volume  $\sim R_c^3$  by energetic (several hundred eV) electrons.
- Furthermore, only a small fraction of the recoil ions have energies  $> W_{bubble}$  ( $\sim 2$  keV for freon 115).



# The number of bubbles generated per source neutron can be expressed by a simple equation

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- For a given solid angle of the bubble detector  $d\Omega$ , the number of bubbles per neutron source can be expressed as

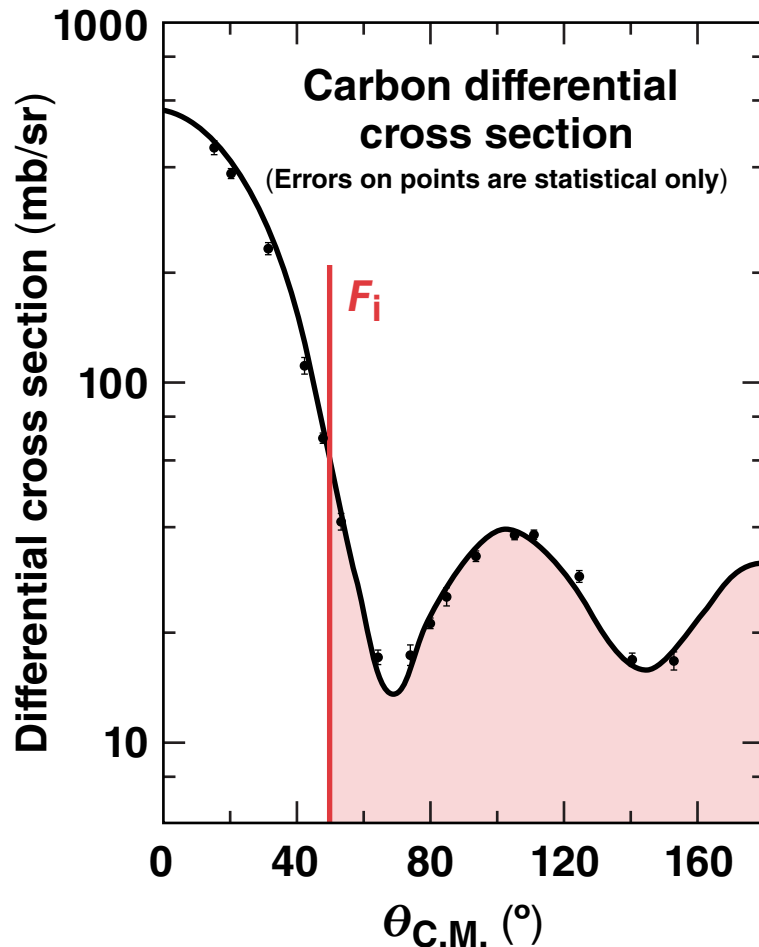
$$\frac{\text{number of bubbles}}{n_{\text{source}}} = F_n \cdot F_i \cdot F_e \cdot d\Omega.$$

$F_n$  is the fraction of the incident neutrons interacting with the superheated liquid.

$F_i$  is the fraction of the ejected ions with energy  $\geq W_{\text{bubble}}$ .

$F_e$  is the fraction of the ejected ions that induce, on scattered electrons, an energy  $\geq W_{\text{bubble}}$  and range  $\sim R_C$ , and produce bubbles.

# The interaction coefficients from the previous slide can be easily calculated



- $F_n$  (the fraction of incident neutrons) can be calculated based on the total scattering cross section.
- $F_i$  (the fraction of the ejected ions with energy  $\geq W_{\text{bubble}}$ ) can be calculated based on the differential cross section ( $>50^{\circ}$  for the case of freon).
- $F_e$  (the fraction of the ejected ions that produce bubbles) can be calculated from the interaction cross section between a recoil nucleus and an electron\*\*

$$\sigma_{i,e} = 18.74 \times 10^{-21} \frac{Z_e R_h}{W_{\text{bubble}}} (\text{cm}^2),$$

where  $Z_e$  is the number of electrons/molecule and  $R_h$  is the Rydberg energy.

\*A. J. Frasca *et al.*, Phys. Rev. **144**, 854 (1966).  
\*\*F. Seitz, Phys. Fluids **1**, 2 (1958).

# The number of bubbles/neutron sources can be estimated for the LLE freon bubble detector

- For the case of freon 115 ( $\text{CCl}_2\text{F}_5$ ) at  $50^\circ\text{C}$ 
  - $F_n = 0.243$
  - $F_i = 0.08$
  - $F_e = 5.719 \times 10^{-5}$
- Therefore, for the detector's solid angle

$$\text{Neutron sensitivity} = \frac{\text{number of bubbles}}{n_{\text{source}}} = F_n \cdot F_i \cdot F_e \cdot d\Omega = 1.33 \times 10^{-12}$$

- For the set of images shown earlier (subtracting the turbulence area from  $d\Omega$ ), the neutron yield is  $y_n = 10^{13}$ .
- After subtracting the turbulence area, expected number of bubbles  $\approx 12$ .

# Neutron yield at the NIF will be sufficient to obtain a high-resolution neutron image



Penumbral/pinhole imaging with bubble chambers requires at least  $10^3$  to  $10^4$  bubbles inside the detector for a  $4\text{-}\mu\text{m}$  to  $1\text{-}\mu\text{m}$  resolution of the neutron source image (for a magnification  $M = 30$ ).

OMEGA	NIF
Source-detector distance – 8 m	Source-detector distance – 16 m
FOV – $200\ \mu\text{m}$	FOV – $200\ \mu\text{m}$
Neutron yield $\sim 10^{13}$	Neutron yield $\sim 10^{19}$
No. of bubbles observed $\sim 14$	No. of bubbles expected $\sim 3 \times 10^6$

**Neutron yield at the NIF will reach  $y_n = 10^{19} \rightarrow 10^6$  bubbles can be produced, more than enough for a high-resolution neutron image.**

## The 14-MeV-neutron sensitivity of a freon-based bubble detector has been tested on OMEGA



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