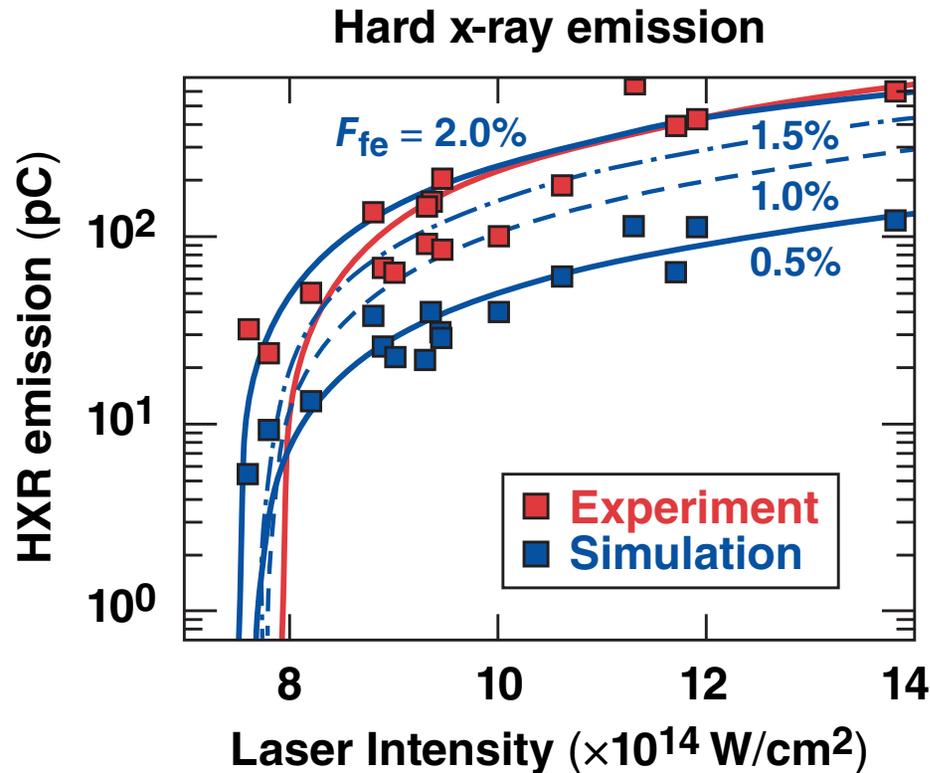


# Transport of Energetic Electrons from Two-Plasmon Decay in the 1-D Hydrodynamic Code *LILAC*



J. A. Deletrez  
University of Rochester  
Laboratory for Laser Energetics

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## Summary

# The fast electrons from the two-plasmon-decay instability have little effect on two-picket cryogenic implosions

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- Relativistic fast-electron transport is modeled in *LILAC* with a radial straight-line model.
- The characteristics of the electron source were determined from warm CH shell implosions.
- A fractional energy of  $\sim 2\%$  absorbed into fast electrons was determined to match the warm CH shells' areal densities.
- Simulations of cryogenic targets resulted in a different fractional energy between continuous and two-picket pulses to match the measured areal densities.

# Collaborators

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**V. N. Goncharov**

**A. V. Maximov**

**J. F. Myatt**

**P. B. Radha**

**T. C. Sangster**

**W. Seka**

**V. A. Smalyuk**

**C. Stoeckl**

**B. Yaakobi**

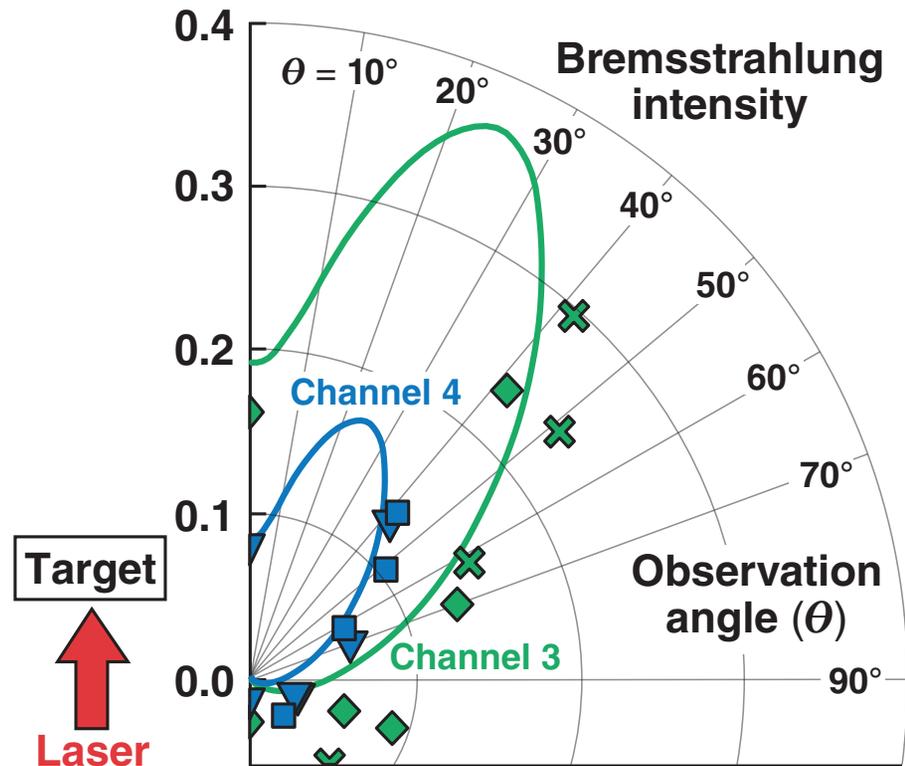
**University of Rochester  
Laboratory for Laser Energetics**

**J. A. Frenje**

**Massachusetts Institute of Technology**

# The fast electrons are transported with a straight-line model

- Fast-electron transport modeling depends on the source angular distribution and energy.
- Experimental results with planar targets suggest that electrons are produced nearly normal to the target surface.<sup>1</sup>
- The electrons are created at the  $N_c/4$  surface and travel in the radial direction.



# The fast-electron source parameters are normalized to the hard x-ray (HXR) emission from warm CH targets

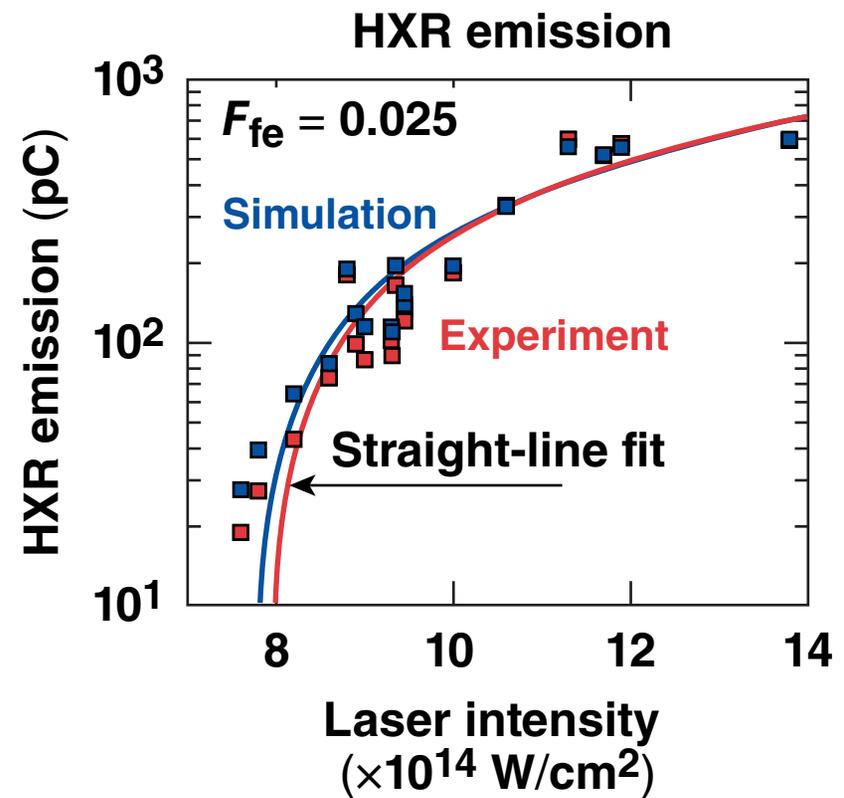
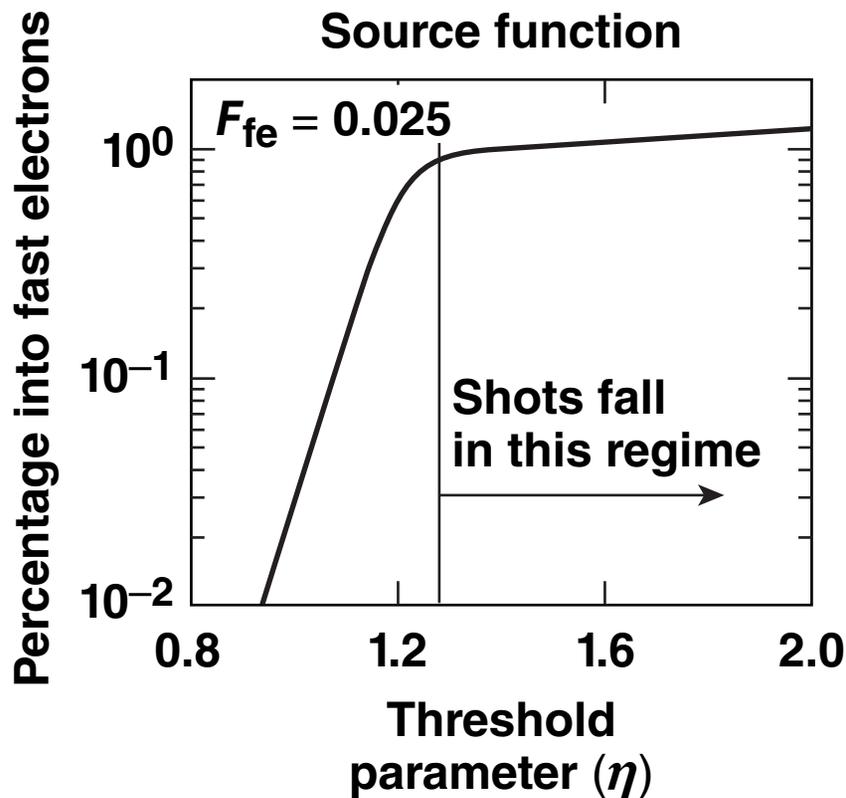


- The energy source is taken to scale as

$$\frac{E_{\text{fast}}}{E_{1/4 N_C}} = F_{\text{fe}} S(\eta).$$

- $F_{\text{fe}}$ , the energy fraction taken from each ray in the laser ray trace, is a free parameter.
- $\eta = I_{14} L_{\mu\text{m}} / 233 T_C$  (keV) is the threshold parameter<sup>1</sup> evaluated at the  $N_C/4$  surface.
- $S(\eta)$  is a source function determined from experiment results.

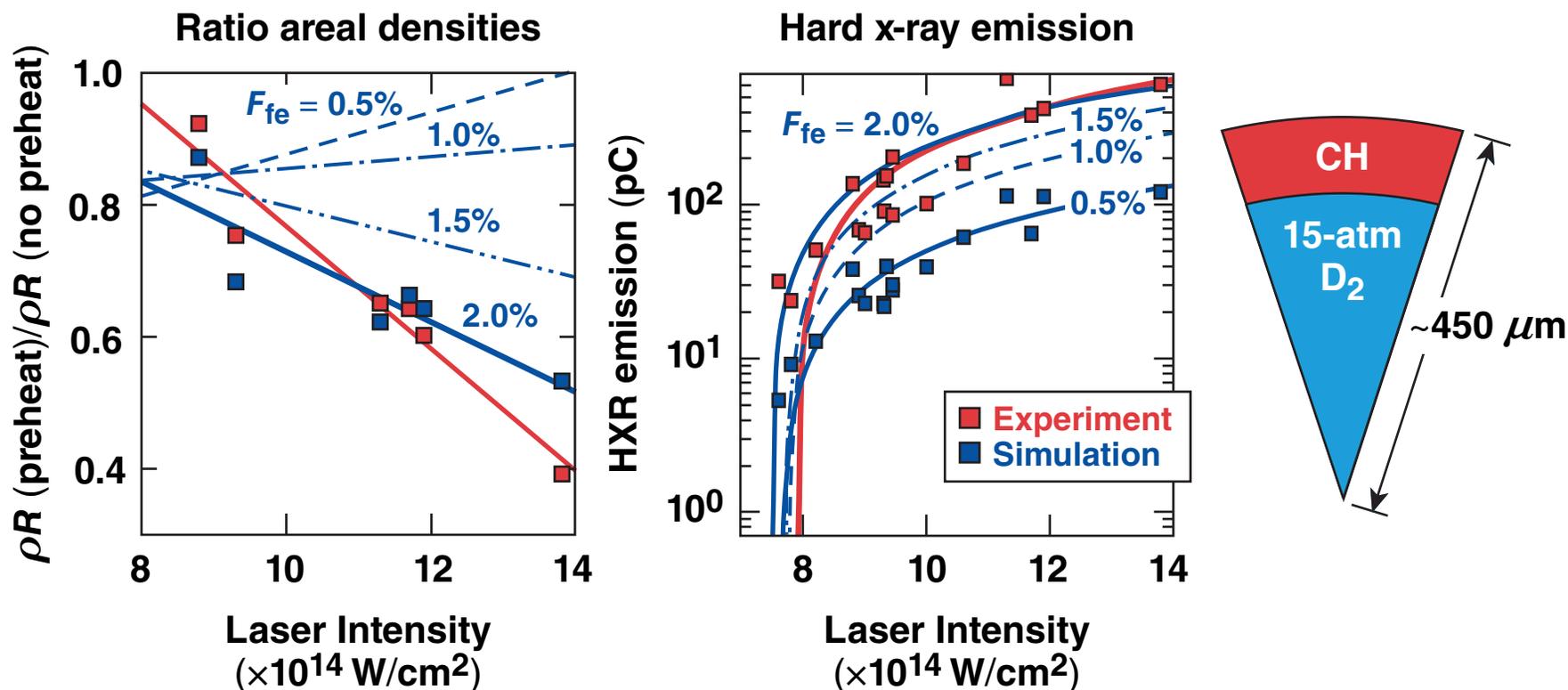
# The source function was chosen to match the integrated HXR emission from warm CH targets over all intensities



32 pC/mJ used to convert simulation emission\*

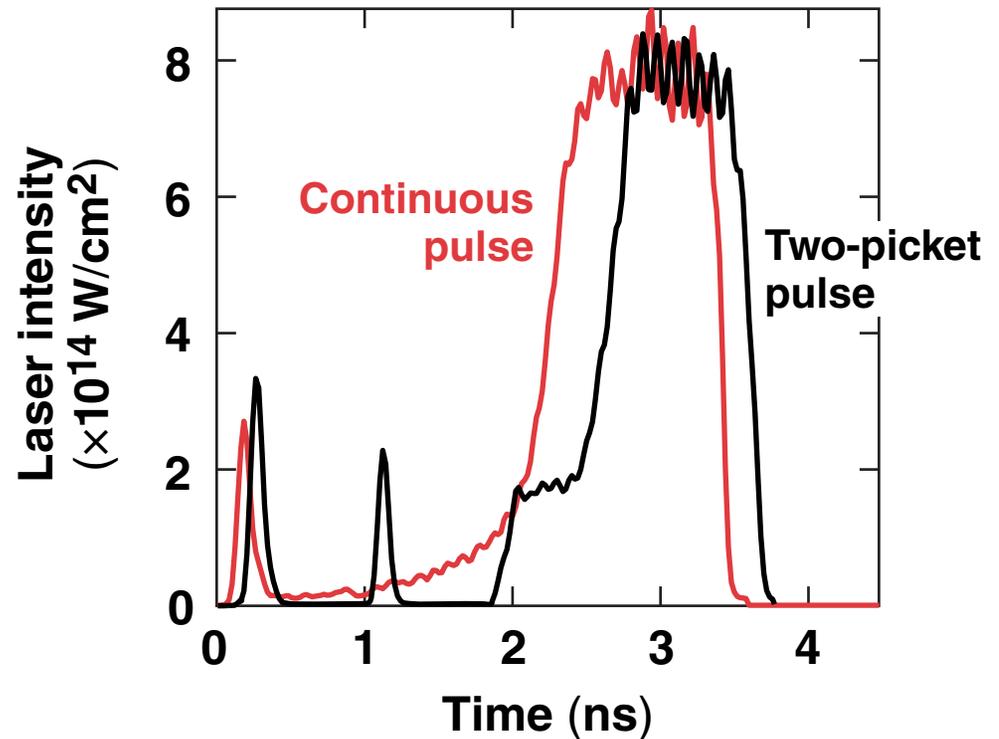
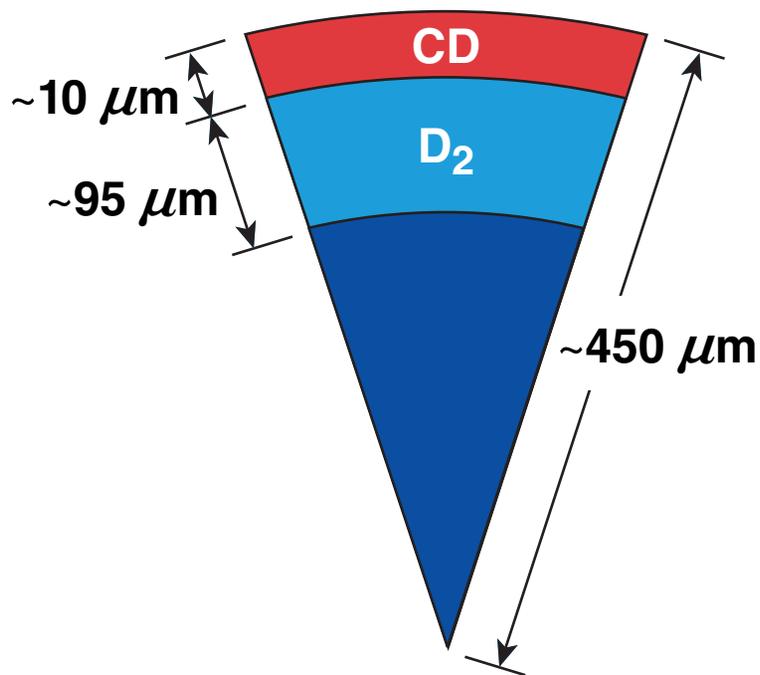
\*B. Yaakobi *et al.*, Phys. Plasmas 12, 062703 (2005) and private communication.

# A free parameter $F_{fe}$ value $\sim 2\%$ fits the measured areal densities and the HXR emission for warm CH implosions

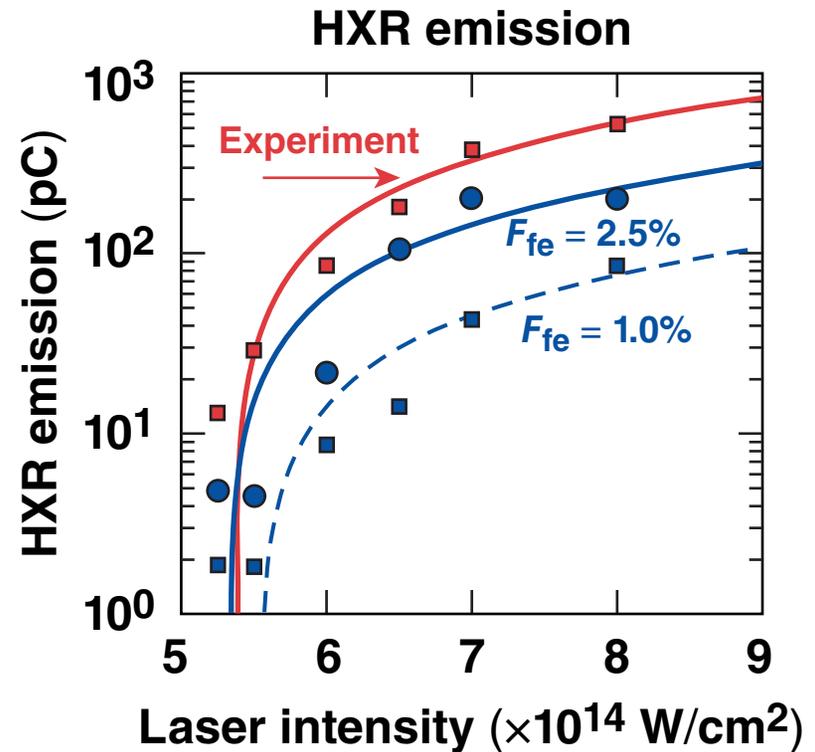
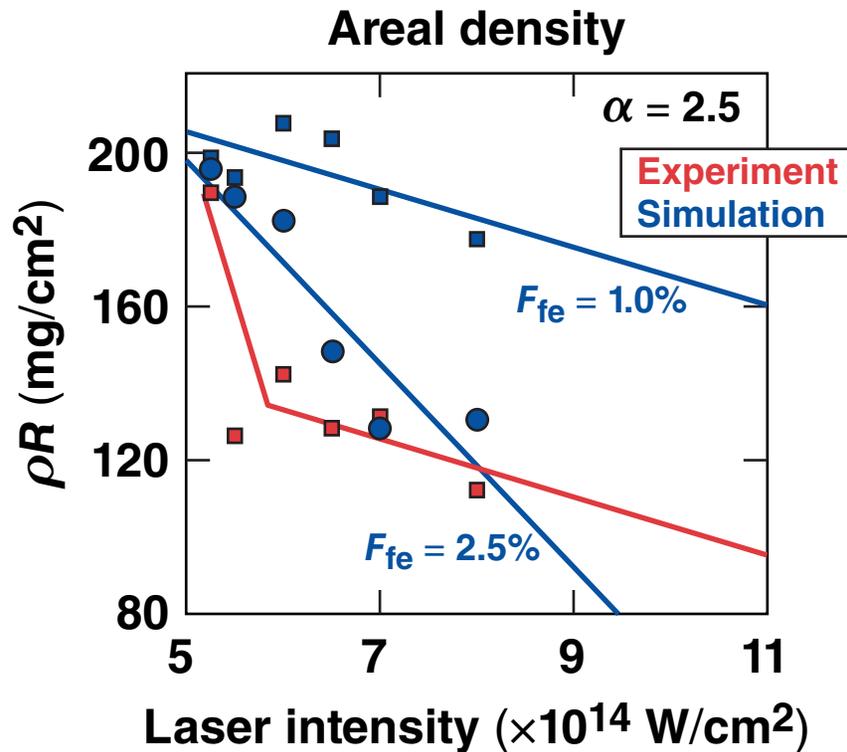


Data points from shots with various shaped pulses and target thicknesses (15 to 28  $\mu\text{m}$ ).

# Thick CD shell cryo targets implosions were carried out with continuous and multipicket pulses



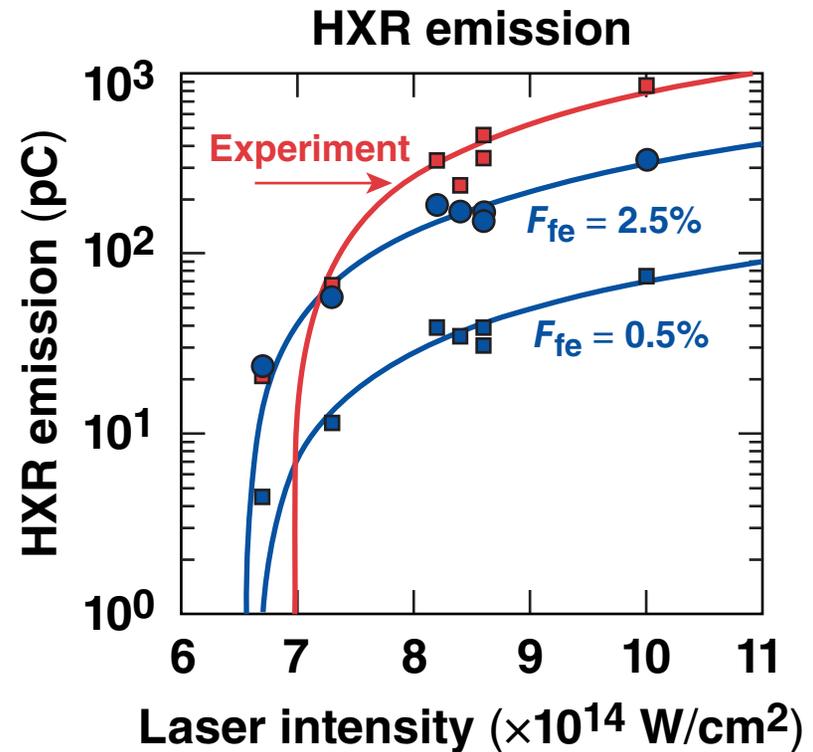
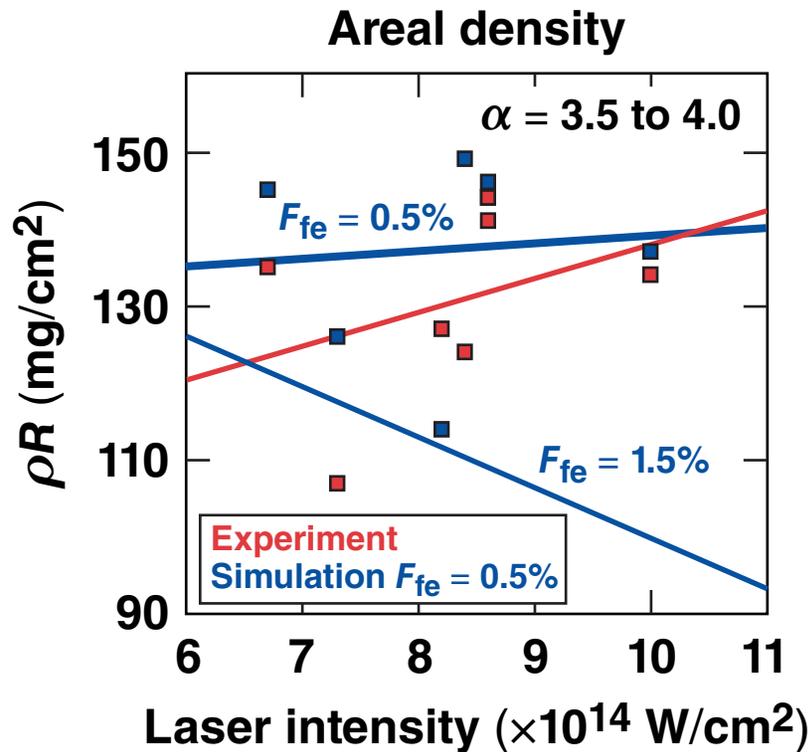
For the continuous pulse a value of 2.5% for  $F_{fe}$  gives good agreement for both  $\rho R$  and HXR emission in cryo implosions



- The  $\rho R$  is affected by shock timing<sup>1</sup> and sampling<sup>2</sup>.
- The measured HXR emission is a factor of two to three too high due to x rays produced outside the target.

<sup>1</sup>V. N. Goncharov (TO5.0006).  
<sup>2</sup>P. B. Radha (NO5.0003).

For the two-picket pulse,  $F_{fe} = 0.5\%$  reproduces the insensitivity of the  $\rho R$  to the two-plasmon decay instability



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