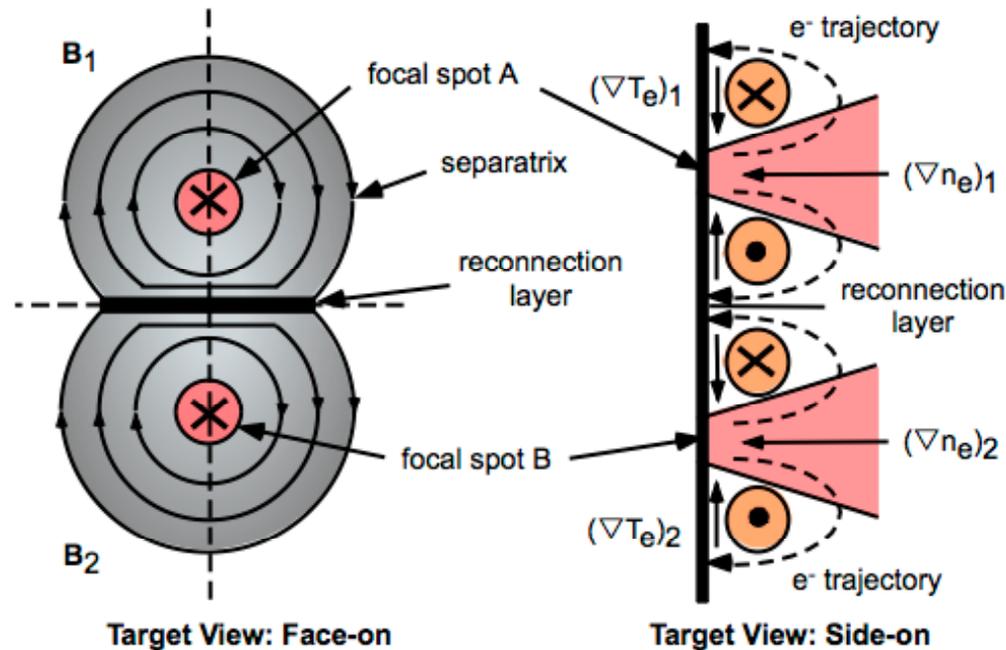


# Magnetic Reconnection and Plasma Dynamics in Two-Beam Laser–Solid Interactions

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48th Annual Meeting of the  
American Physical Society  
Division of Plasma Physics  
Philadelphia, PA  
30 October–3 November 2006

## Summary

**A magnetic-reconnection geometry has been studied using two closely focused laser beams ( $10^{15}$  W/cm<sup>2</sup>, 1 ns)**

- **The plasma dynamics created by two heater beams at the surface of a planar solid target (Al or Au) has been studied.**
- **Observations consistent with a magnetic reconnection have been made**
  - **the formation of a driven magnetic-reconnection field distribution**
  - **the interaction of 0.7- to 1.3-MG magnetic fields in a reconnection layer**
  - **collimated, high-velocity jet formation**
  - **high electron temperature ( $T_e = 1.7$  keV) in the reconnection layer**
- **Reconnection rates predicted by the Sweet–Parker model of reconnection (assuming Spitzer resistivity) are too slow to explain these observations.**

# Collaborators

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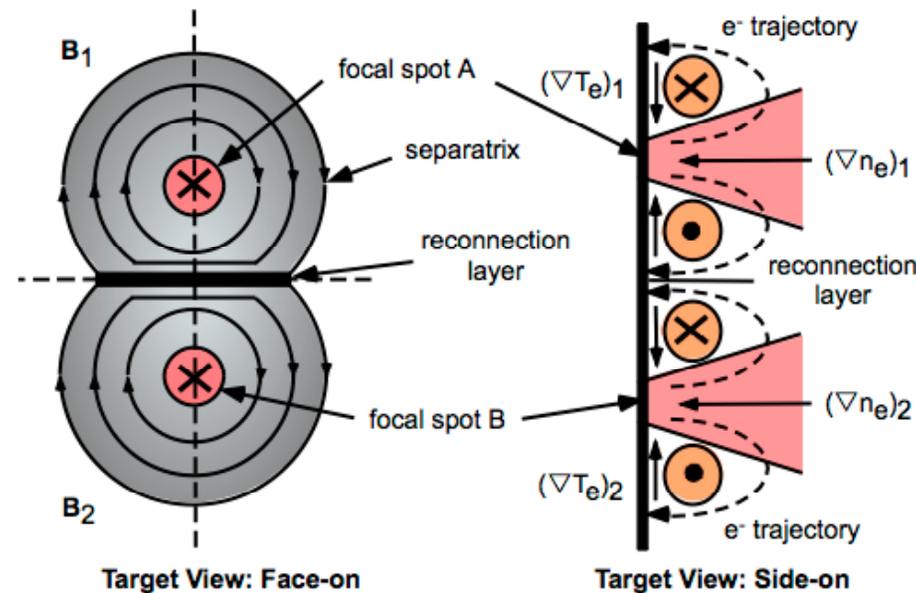
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**M. Notley and M. Sherlock**

**CCLRC, Central Laser Facility, Rutherford Appleton Laboratory, Oxon, UK**

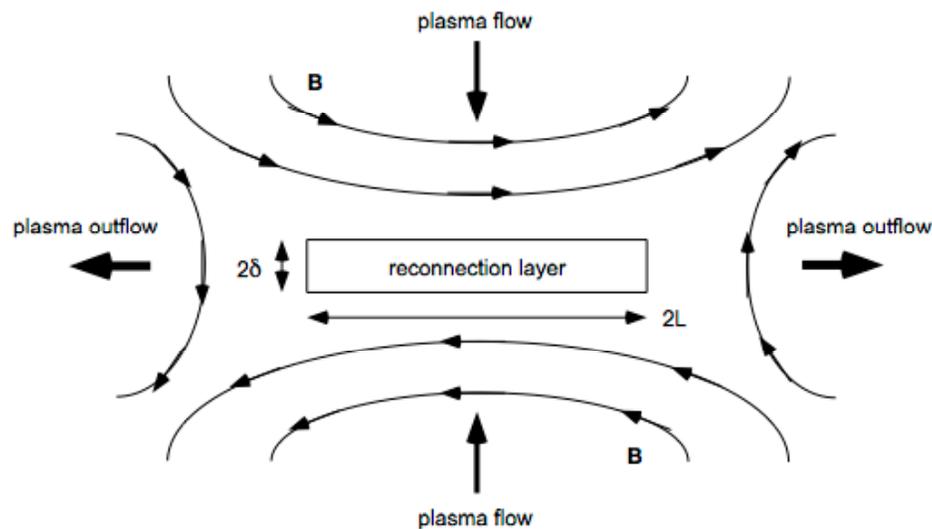
# Magnetic fields are generated by nonparallel electron-density and temperature gradients

- Laser-pulse characteristics: 1 ns,  $10^{15}$  W/cm<sup>2</sup>



$$\partial_t \mathbf{B} = \nabla \times (\mathbf{v} + \mathbf{B}) + (k_B / n_e e) \nabla T_e \times \nabla n_e + (1 / \mu_0) \nabla \times (\eta \nabla \times \mathbf{B})$$

# The self-generated magnetic-field distribution is similar to the Sweet–Parker model of magnetic reconnection



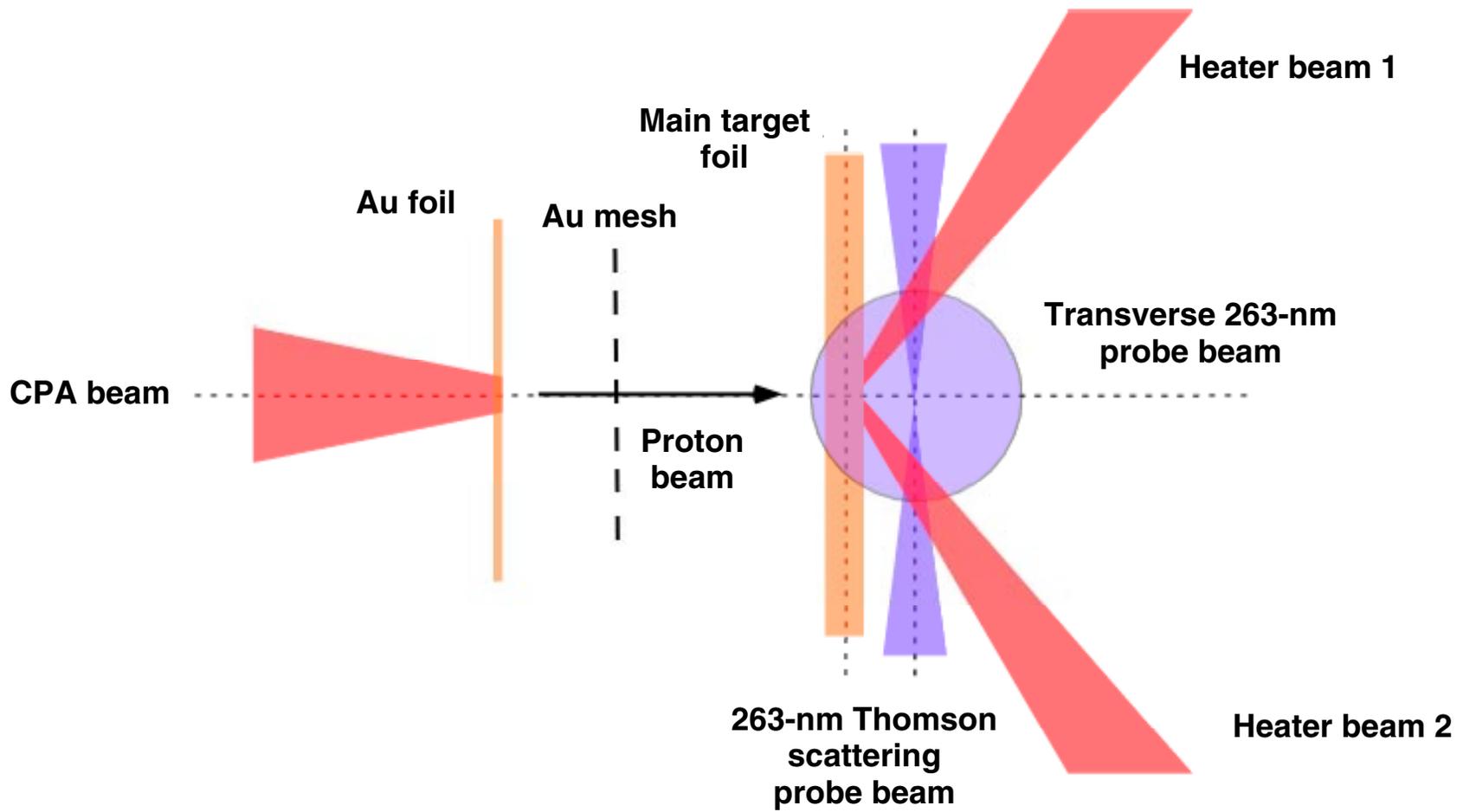
- Magnetic Reynolds number:

$$\frac{1}{S^{1/2}} = \left( \frac{\eta}{\mu_0 L \nu_A} \right)^{1/2}$$

- Reconnection time:

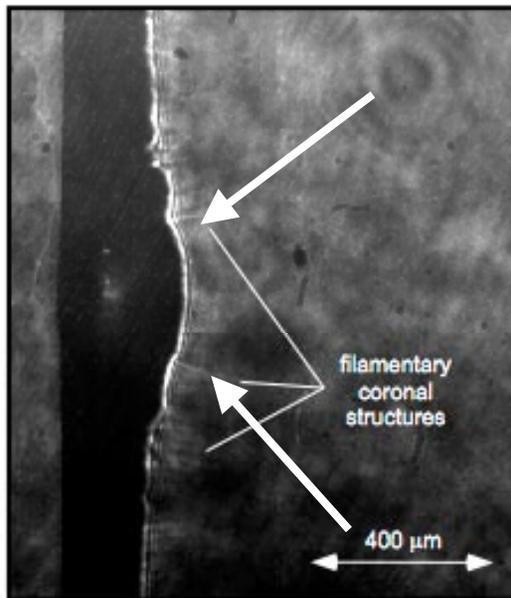
$$\tau_{sp} = \frac{L}{\nu_r} = \frac{L}{\nu_A S^{1/2}} = (\tau_A \tau_R)^{1/2} \approx 4 \text{ ns}$$

# Experimental setup: The Vulcan Laser Facility

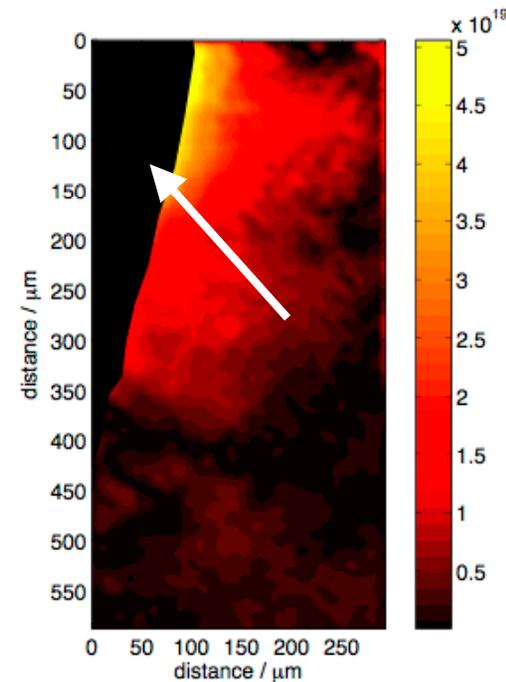


# A single region of expanding plasma develops when the two heater beams are separated by approximately $200\ \mu\text{m}$

- Shadowgram
- $t = t_0 + 1.5\ \text{ns}$



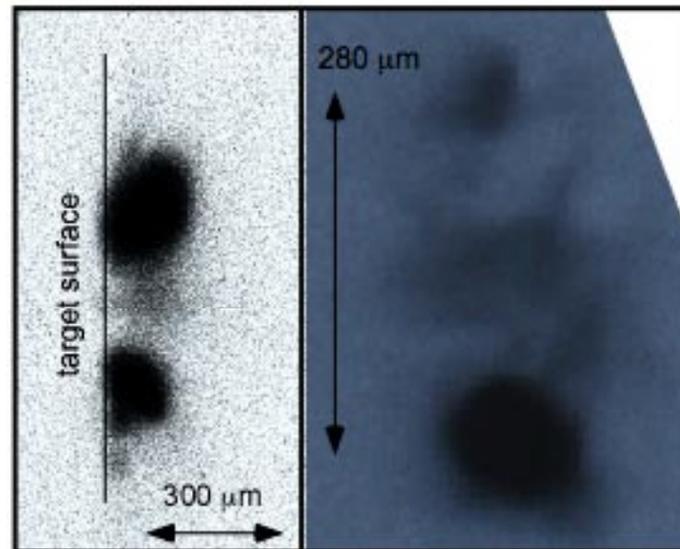
- Interferogram
- $t = t_0 + 1.5\ \text{ns}$



- Aluminum target foil

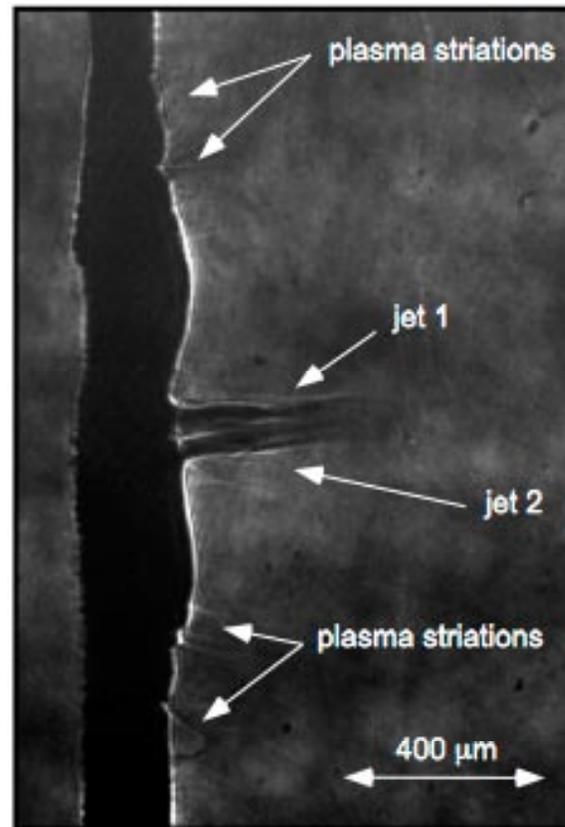
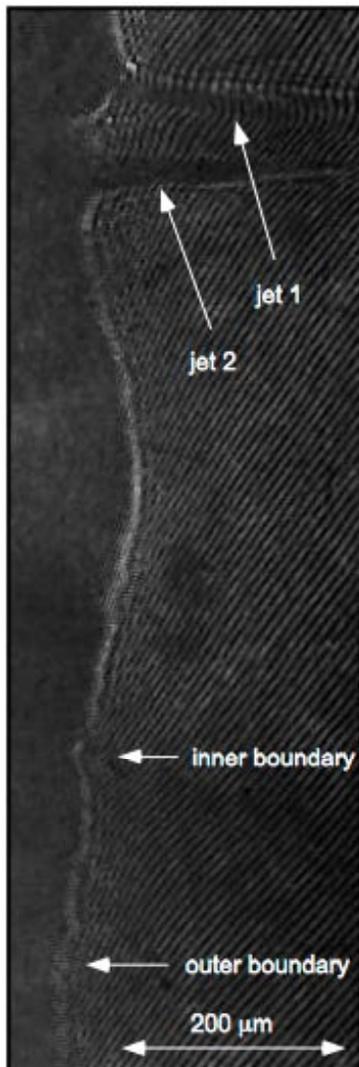
# Time-integrated x-ray pinhole imaging shows heating due to a plasma collision

- X-ray pinhole images (time-integrated)
- Al target (1-keV emission)



- Conversion of streaming ion kinetic energy into ion thermal energy
- Electrons gain energy through electron-ion equilibration

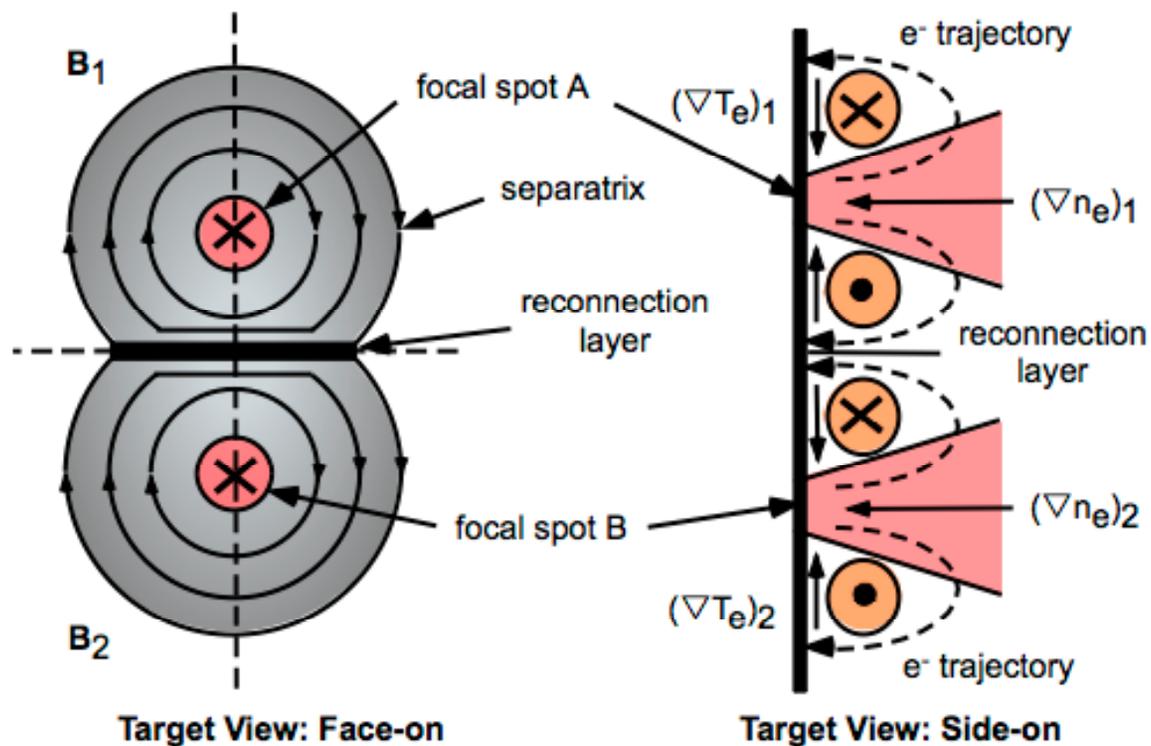
# Sudden jet formation occurs for larger laser-spot separations of around $400\ \mu\text{m}$



- Aluminum target foil
- $t = t_0 + 1.5\ \text{ns}$

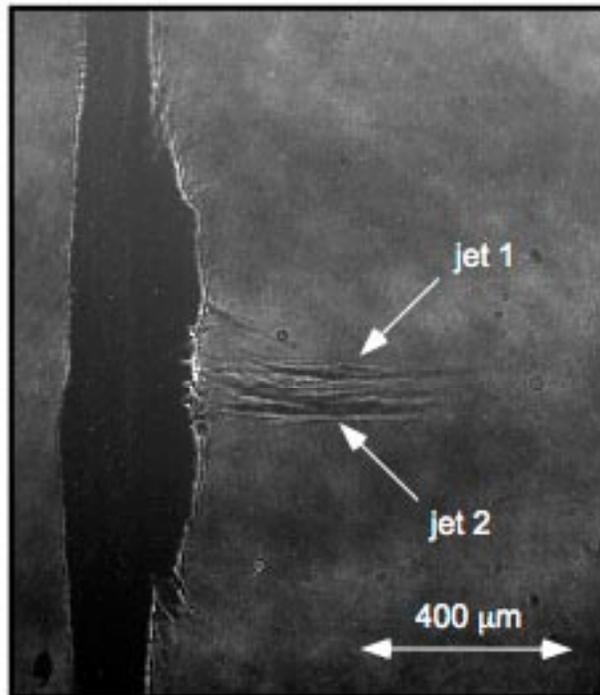
- Jet formation occurs at an angle to the target surface

# Jet formation occurs at an angle to the target surface

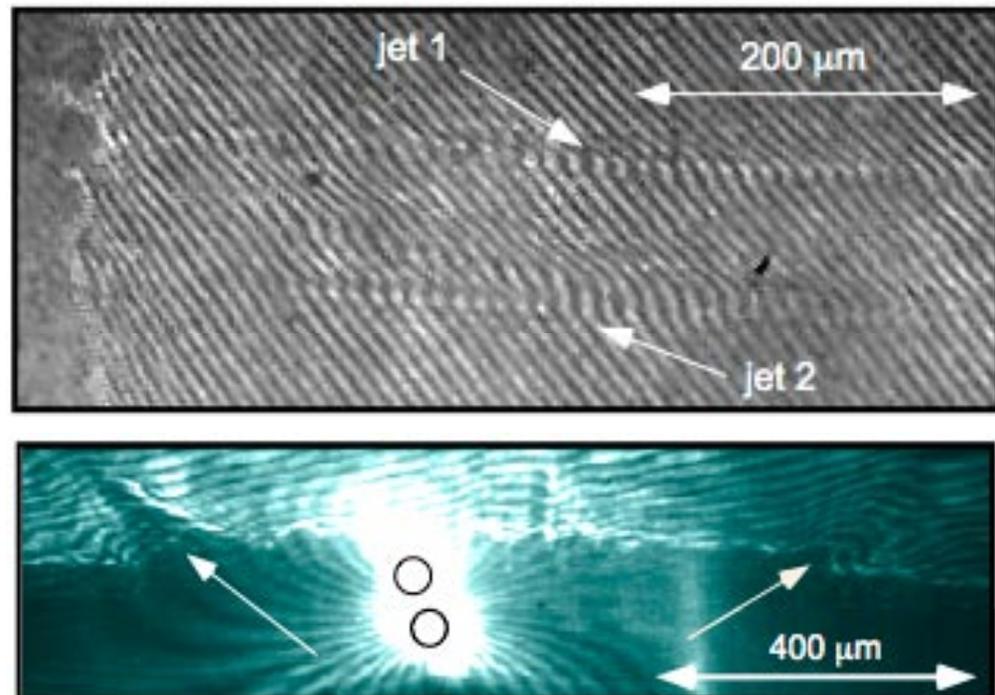


# Greater jet collimation is observed in gold target interactions consistent with radiative-cooling effects

$t = t_0 + 2.5 \text{ ns}$

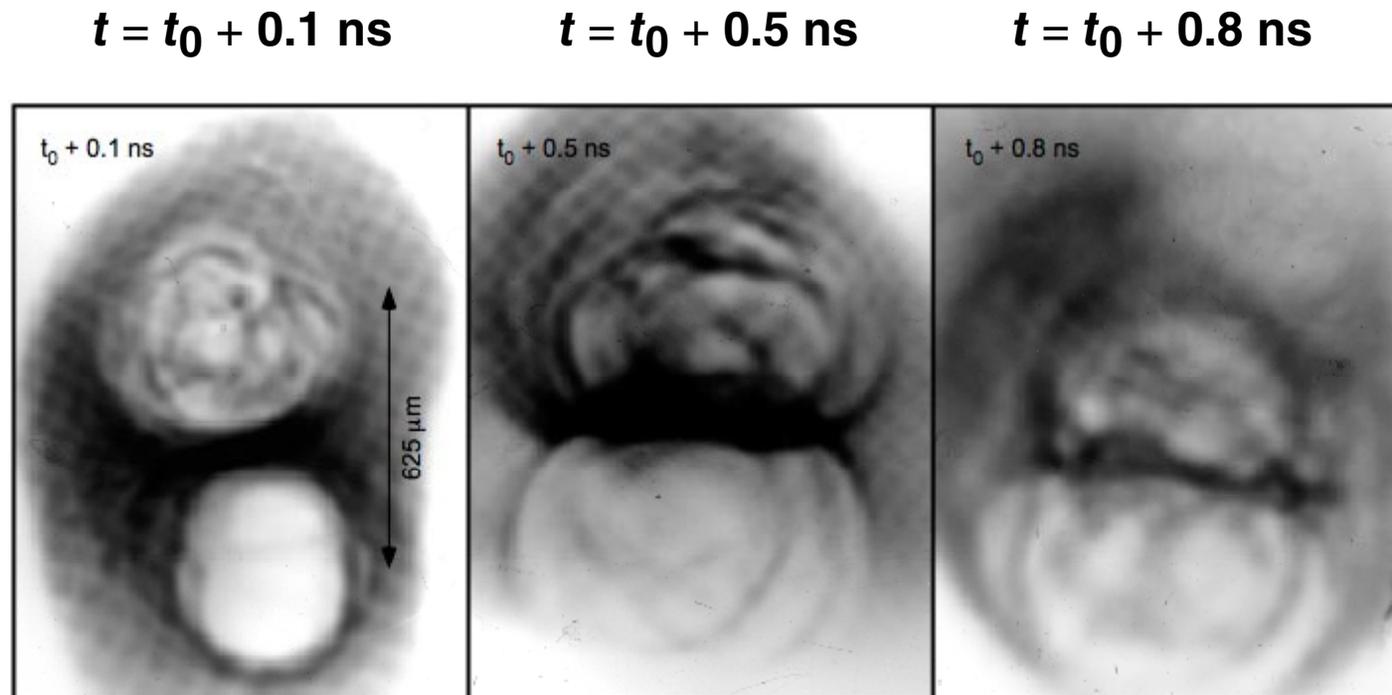


$t = t_0 + 0.7 \text{ ns}$



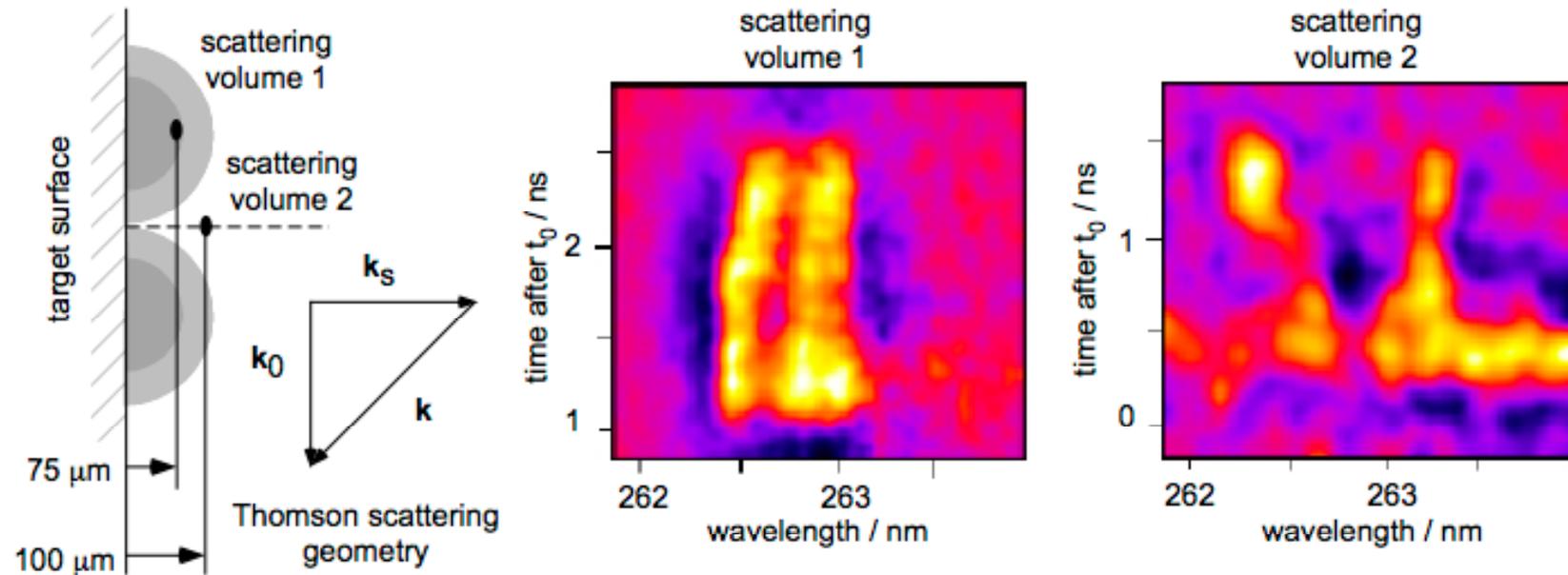
Gold target foil

# Interacting MG-level azimuthal magnetic fields are measured with proton deflectometry

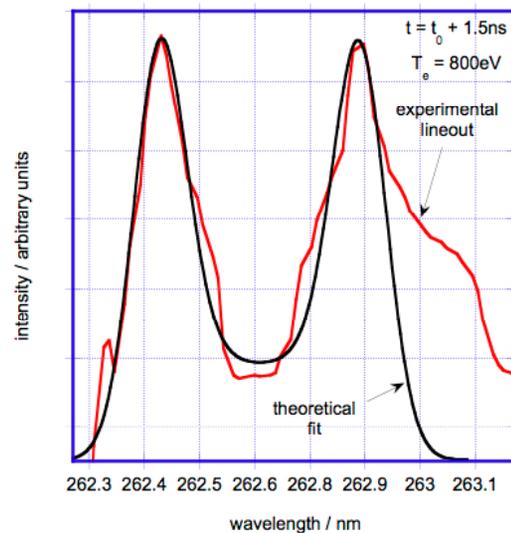


- 0.7- to 1.3-MG magnetic fields
- Stressed magnetic-field lines

# Thomson scattering (TS) measurements show kilo-electron-volt electron temperatures in the aluminum interaction layer

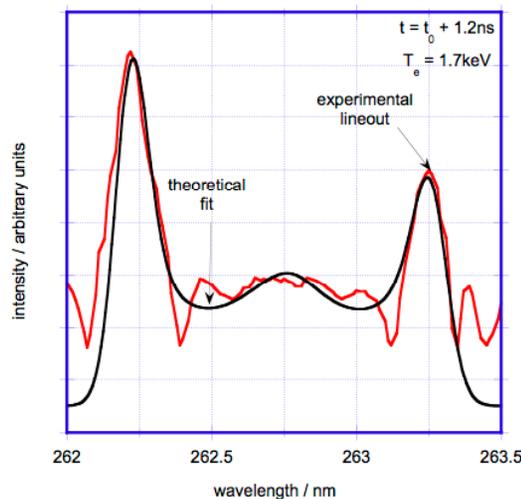


# Thomson scattering (TS) measurements show kilo-electron-volt electron temperatures in the aluminum interaction layer



- Scattering volume 1: plume
- $T_e = 800 \text{ eV}$  at  $t = t_0 + 1.5 \text{ ns}$

• Cooling due to hydrodynamic expansion



- Scattering volume 2: interaction region
- $T_e = 1700 \text{ eV}$  at  $t = t_0 + 1.2 \text{ ns}$

• Assumes an ion-distribution function represented by the sum of two Maxwellians shifted by the beam flow velocity

## Summary/Conclusions

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