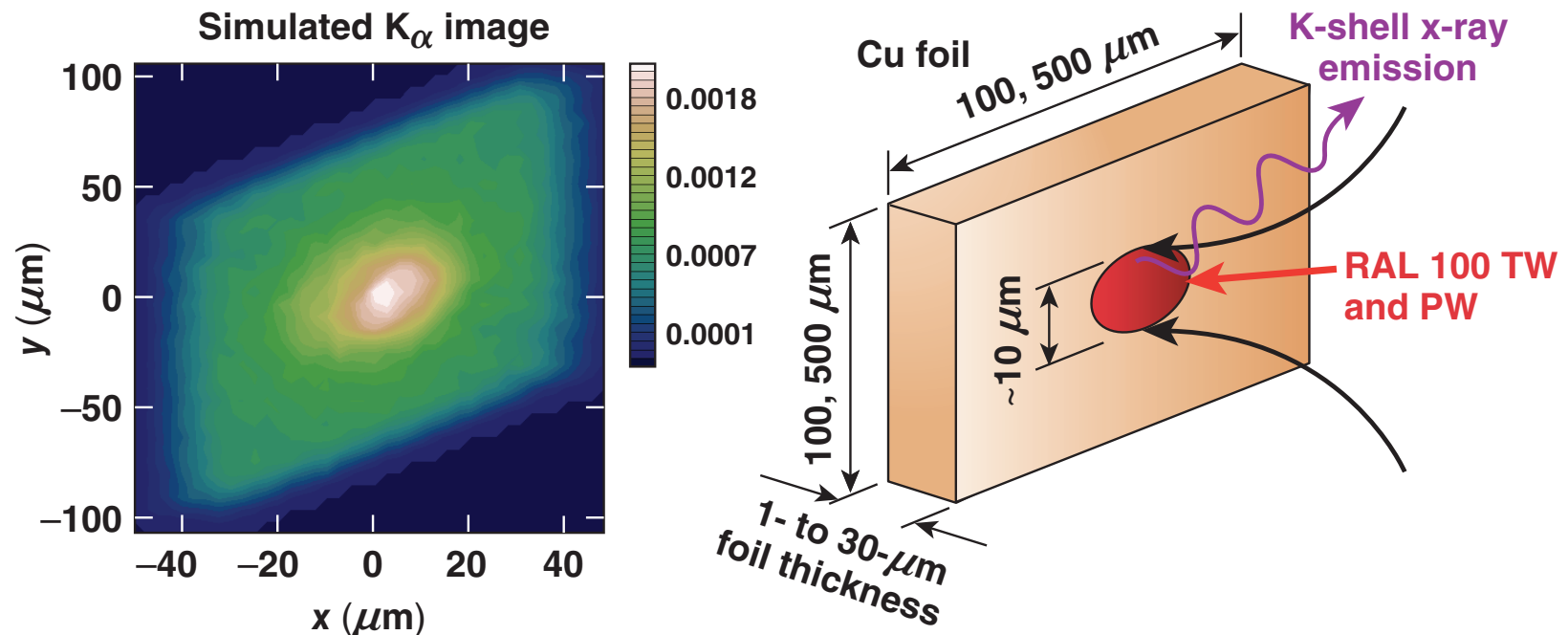


High-Intensity Laser Interactions with Solid Targets and Implications for Fast-Ignition Experiments on OMEGA EP



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Summary

K_α and K_β emission from PW laser interaction with mass-limited foils at RAL is well predicted



- A semi-analytic model describes the K_α data well and is supported by implicit-hybrid PIC calculations (*LSP**).
- Mass-limited targets simplify the modeling (good “test bed” for FI).
- The experimental K_α yields are consistent with classical stopping and a constant hot-electron conversion efficiency of $\sim 10\%$.
- The ratio of K_β to K_α signal is sensitive to target heating and provides a consistency check on conversion efficiency.

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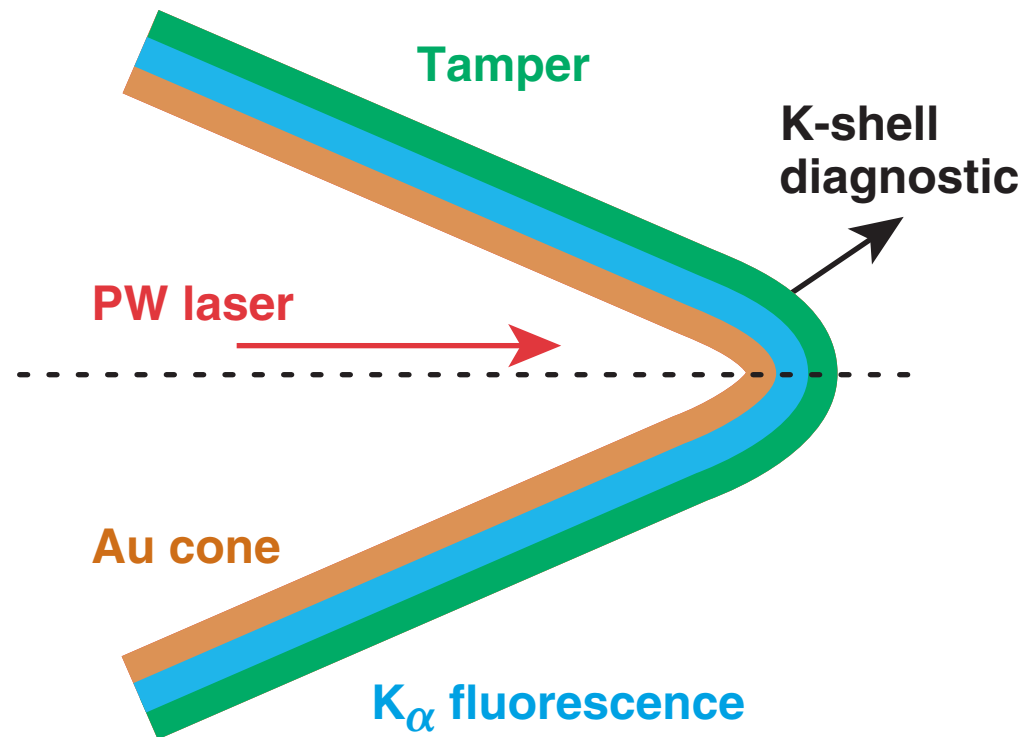
Outline



- **Petawatt-foil experiments are relevant to FI scheme, particularly cone-in-shell**
- **Experiments**
- **Mass-limited targets**
- **Conversion efficiency**
- **Target heating**
- **Future experiments**

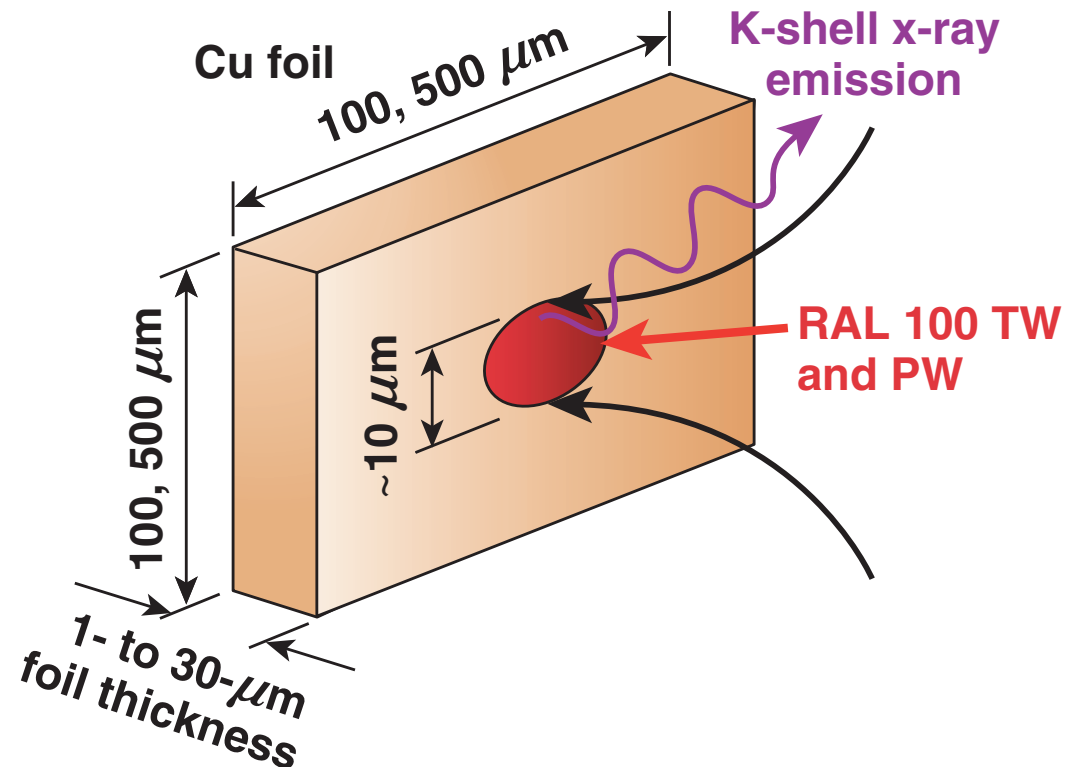
PW-foil experiments are relevant to cone-in-shell fast ignition and high-brightness backlighting

- In the cone-in-shell scheme, the PW laser interacts with a solid-density interface
- Foils are a simple test bed for:
 - hot-electron spectrum
 - conversion efficiency
 - transport
- Interaction conditions are similar



K-fluorescence from high-intensity laser-irradiated mid-Z foils provides an important diagnostic of hot electrons*

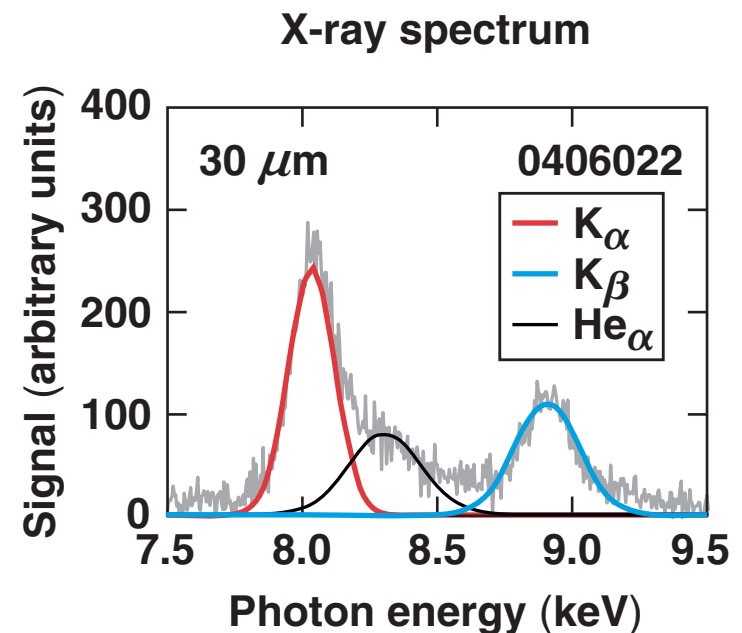
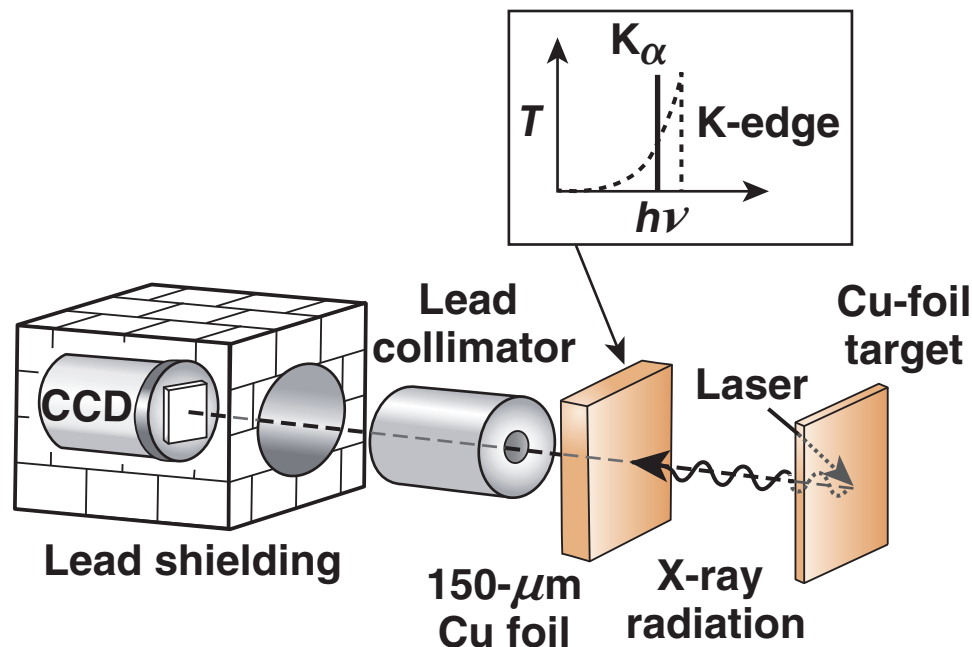
- Hot electrons create K-shell vacancies
- $L \rightarrow K$ ($M \rightarrow K$) transition gives K_α (K_β) emission
- Emission occurs during hot-electron lifetime (few ps)



* R. B. Stephens *et al.*, Phys. Rev. E **69**, 066414 (2004).
J. D. Hares *et al.*, Phys. Rev. Lett. **42**, 1216 (1979).

Measurements of K-shell emission from mass-limited targets have been performed at the RAL 100 TW and PW facilities*

- Laser intensities of $10^{18} < I < 10^{20}$ W/cm²
- A range of target volumes: $10^{-5} < V < 10^{-1}$ mm³
- Solid copper targets



Mass-limited targets have some remarkable properties that simplify the modeling



- The majority of hot electrons stop in the target due to space-charge (refluxing)
- Secondary radiation production is simple to compute (as in infinite medium)
- Efficient K_{α} , K_{β} radiators
- Transparent to K-shell x rays
- Access high temperatures at solid density
- Can be used to benchmark codes—*LSP* can model the whole target in 3-D (implicit-hybrid mode)
- Good test bed for hot-electron conversion and volumetric heating

A combination of semi-analytic modeling and implicit-hybrid PIC calculations have been used to investigate RAL 100 TW and PW experiments

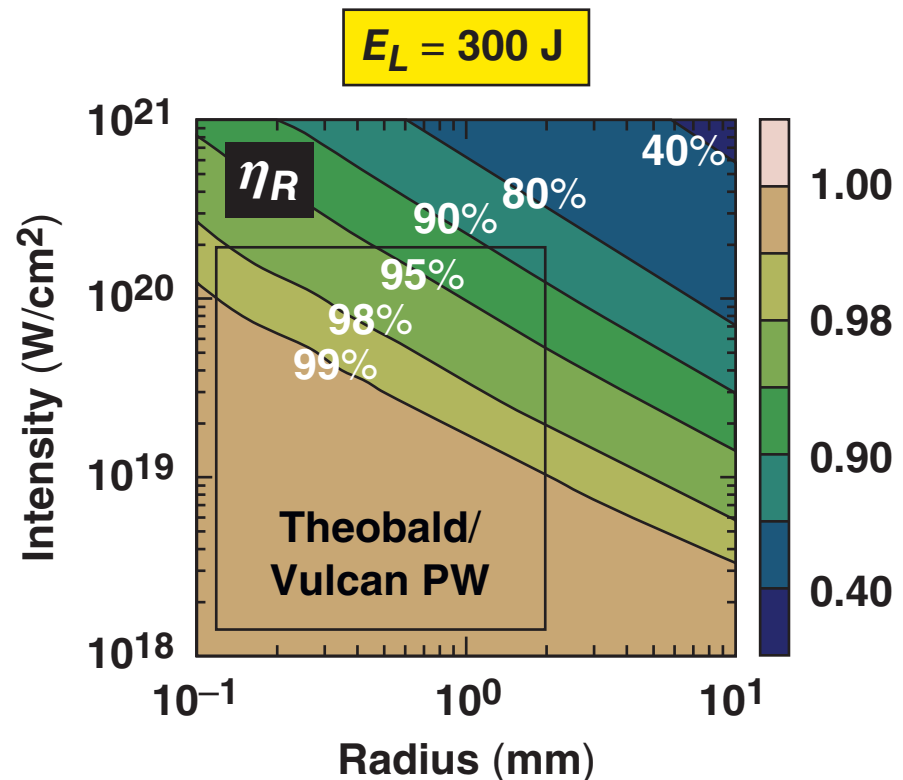


- The semi-analytic model is a modification of calculation in Green and Cosslett*
 - classical CSDA stopping, relativistic K-ionization cross sections
 - refluxing effect is included
 - re-absorption included
- *LSP* contains additional physics
 - collision model requires tuning
 - self-generated fields
 - resistive inhibition and instability of hot current
 - target heating

The refluxing efficiency is a parameter in the model, but in *LSP* the trajectories are computed self-consistently



- For the parameters of the RAL experiments, the hot electron refluxing is nearly perfect.
- A capacitor/Boltzmann electron model is used to estimate the importance of refluxing.



Both the semi-analytic model and *LSP* require the hot-electron spectrum and its dependence on laser intensity to be specified



- We chose an exponential distribution for hot-electron energy characterized by a temperature T_{hot} .
- We assume that for different laser intensities the distribution is changed through only the temperature moment.
- Intensity is allowed to be a function of space.
- The temperature is given by the “ponderomotive scaling” (Wilks*).

$$T_{\text{hot}} (\text{MeV}) = 0.511 \left[\left(1 + I_{18} \lambda_{\mu\text{m}}^2 / 1.37 \right)^{1/2} - 1 \right]$$

To complete the model K-shell ionization cross section and hot-electron stopping power are required

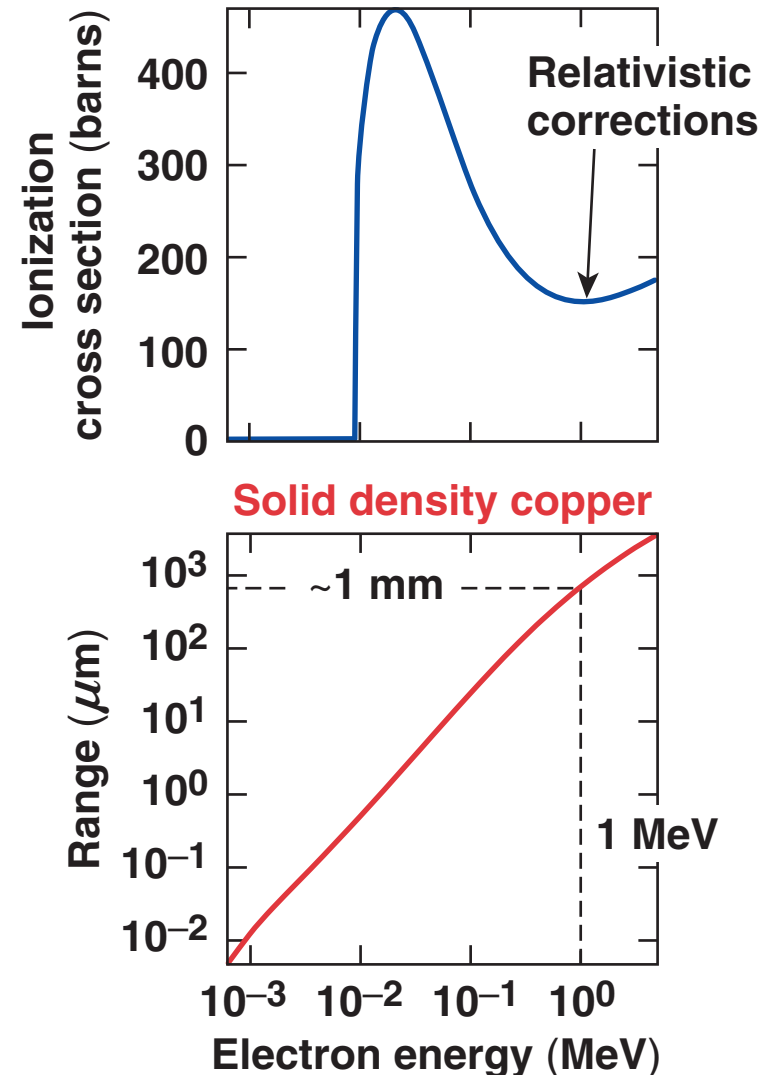


- $\sigma_K(E)$ [taken from Kolbenstvedt¹], relativistic corrections are essential
- The fluorescence probability ω_K is taken for cold matter at solid density
- CDSA range²

Need to specify: $E_e (= \eta_{L \rightarrow e} E_L)$

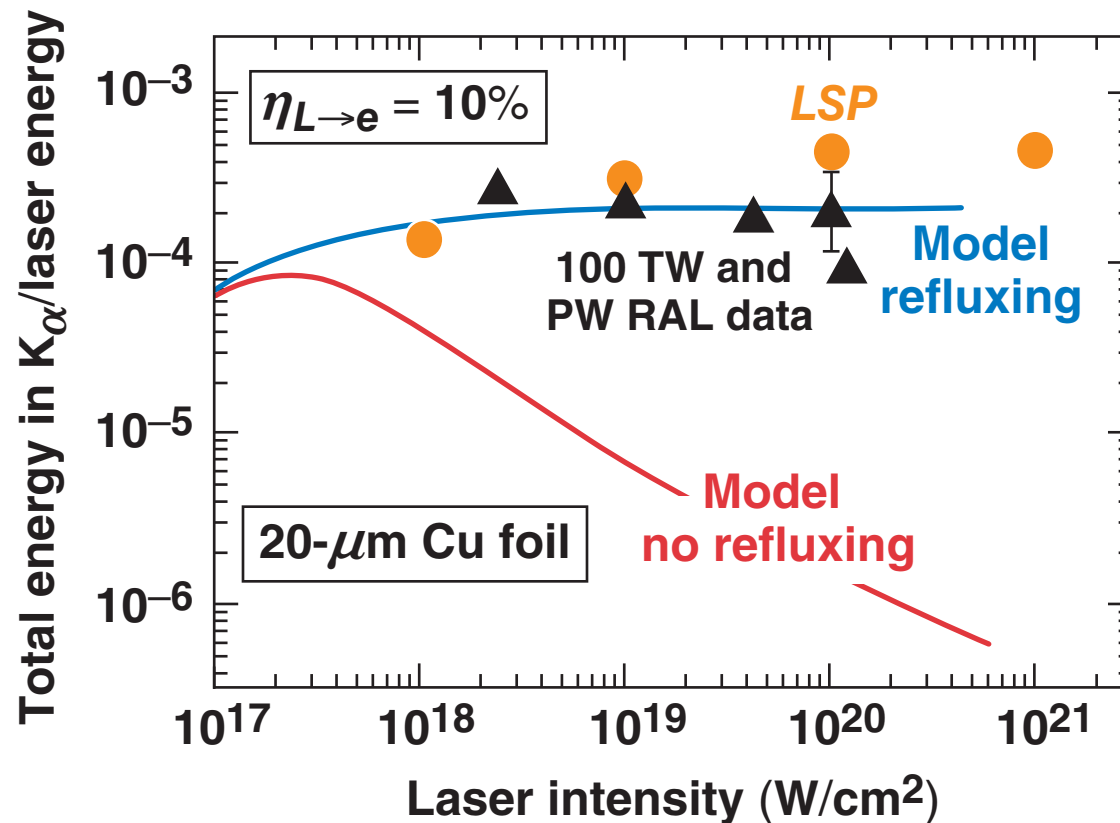
¹ H. Kolbenstvedt, J. Appl. Phys. **38**, 4785 (1967).

² H. O. Wyckoff, *ICRU Report 37*, Intern. Comm. on Radiation Units and Measurements, Inc., Bethesda, MD (1984).



The semi-analytical model and the *LSP* calculations agree on the K_{α} yield

- Hot electron range in *LSP* is the same as CSDA range
- Insensitivity to the laser energy is a direct result of refluxing



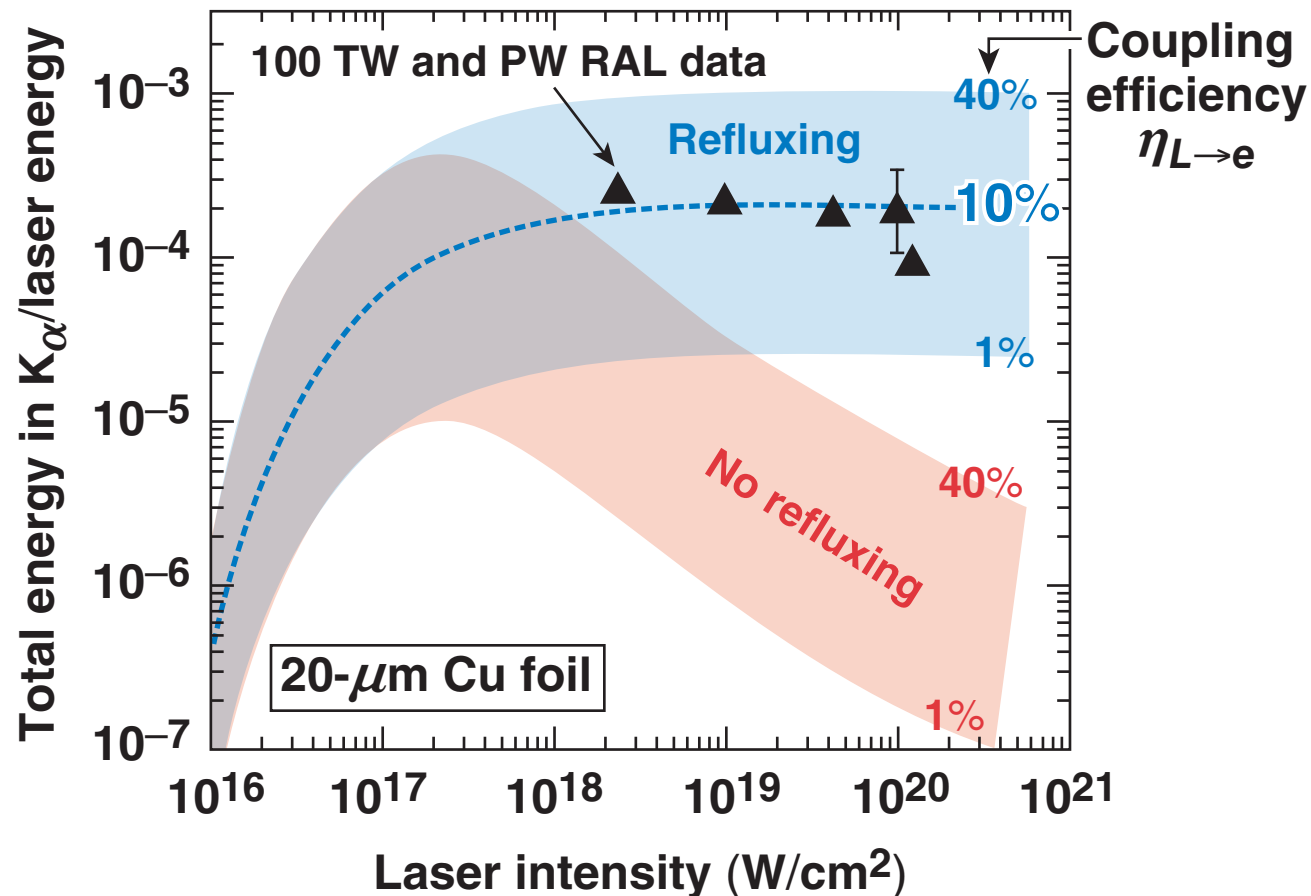
Resistive inhibition would reduce the hot-electron range and lead to a lower K-fluorescence production efficiency



- *LSP* calculations, using Lee and Moore electrical conductivity, show that resistive inhibition is unimportant for the parameters of the RAL experiments
- Compare classical range with “resistive range”
- The electric field $E (= j/\sigma_{Cu}) \sim 2 \times 10^5$ kV/cm and would stop a 1-MeV electron in $50 \mu\text{m}$.
- This is much less than the collisional range $\sim 700 \mu\text{m}$ at 1 MeV
- Greater than target thickness
 - refluxing hot electrons provide the return current (much less collisional)
- Transit time is only 0.1 ps

The hot-electron conversion efficiency can be determined by a fit to the RAL data

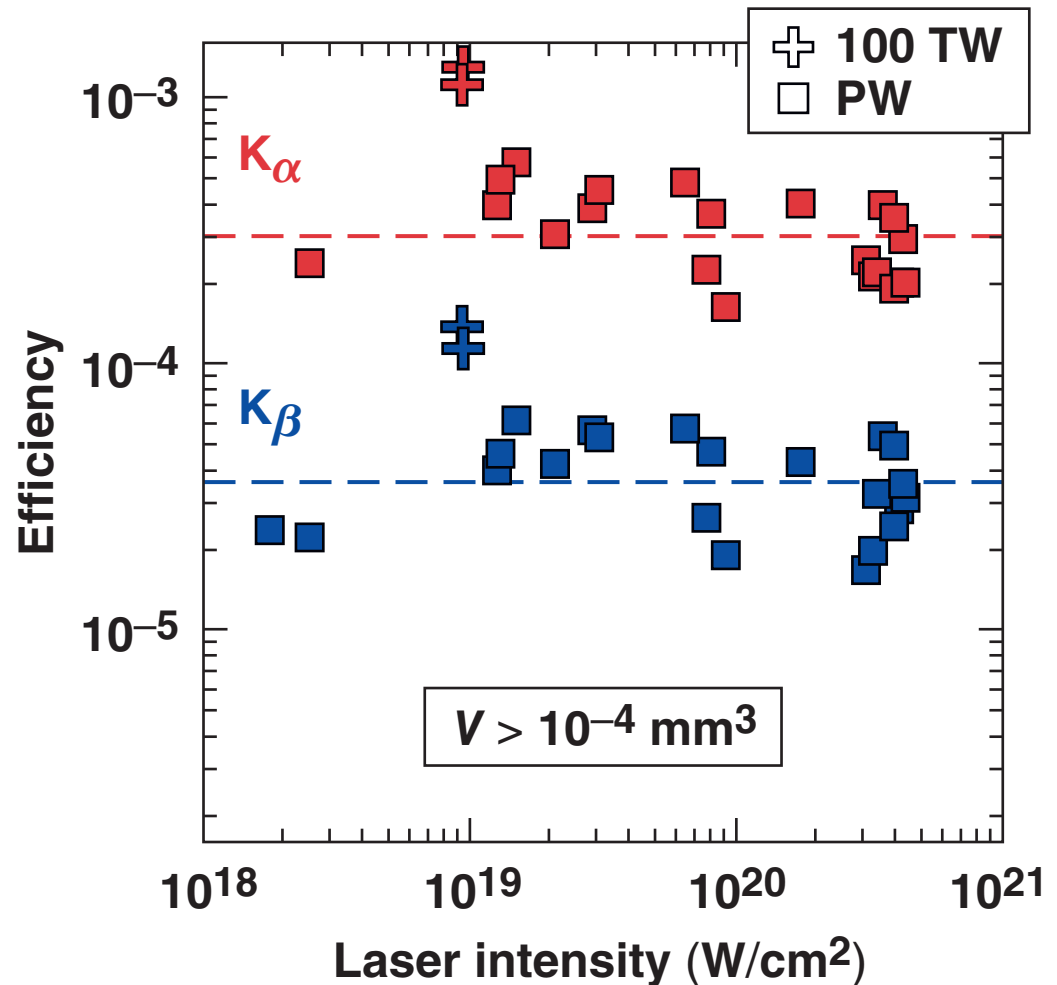
- The experimental data (here for a particular geometry) is consistent with a constant conversion efficiency of $\sim 10\%$.



The experimental K_α and K_β yields are essentially independent of the target geometry and laser intensity

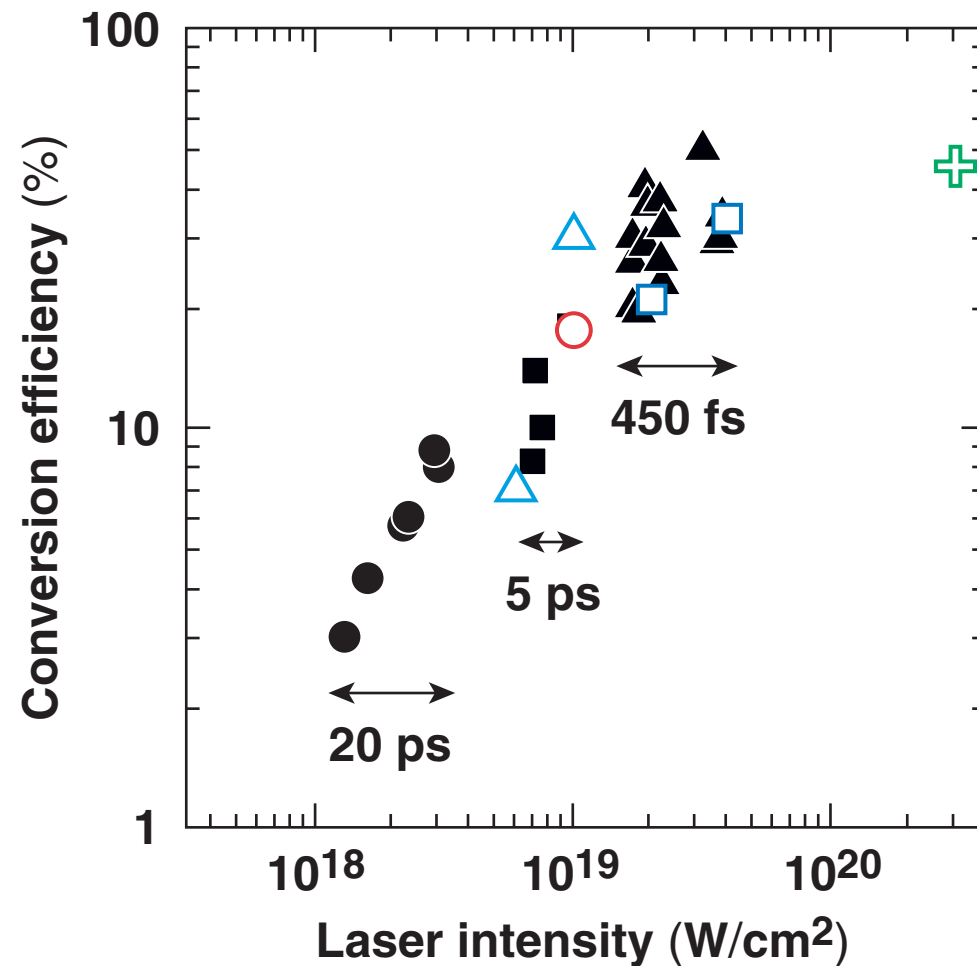


- This is a direct prediction of refluxing plus cross-section and stopping power
- E_L varies by a factor of 5
- Foil geometry varies widely $10^{-4} < V < 10^{-1} \text{ mm}^3$
- Very small mass targets ($V < 10^{-4} \text{ mm}^3$) are excluded

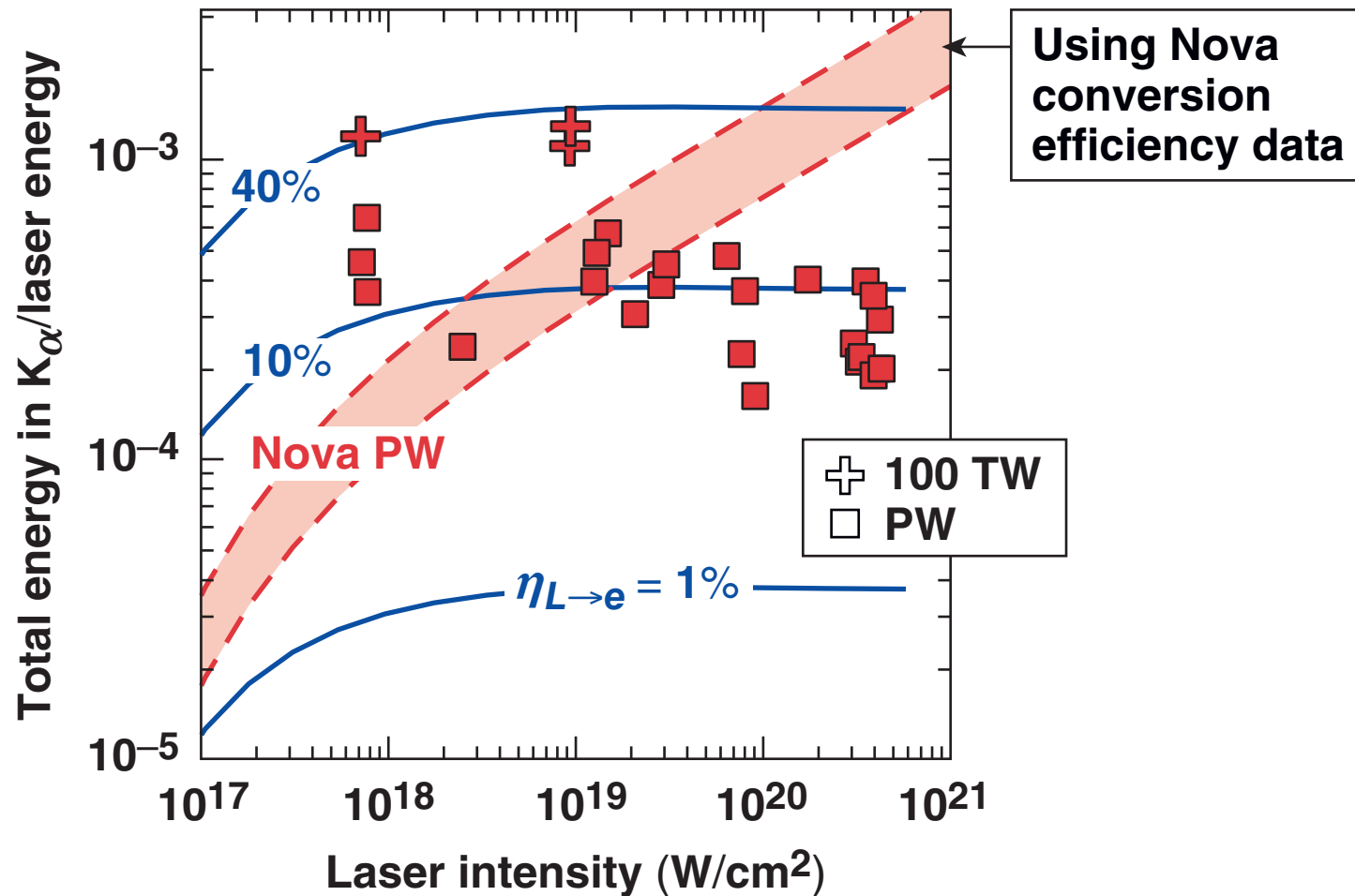


Previous experiments* indicated that the hot-electron conversion efficiency increases dramatically with laser intensity

- These experiments are compromised by poorly characterized interaction conditions.
- These are indirect measurements and rely on model calculations.

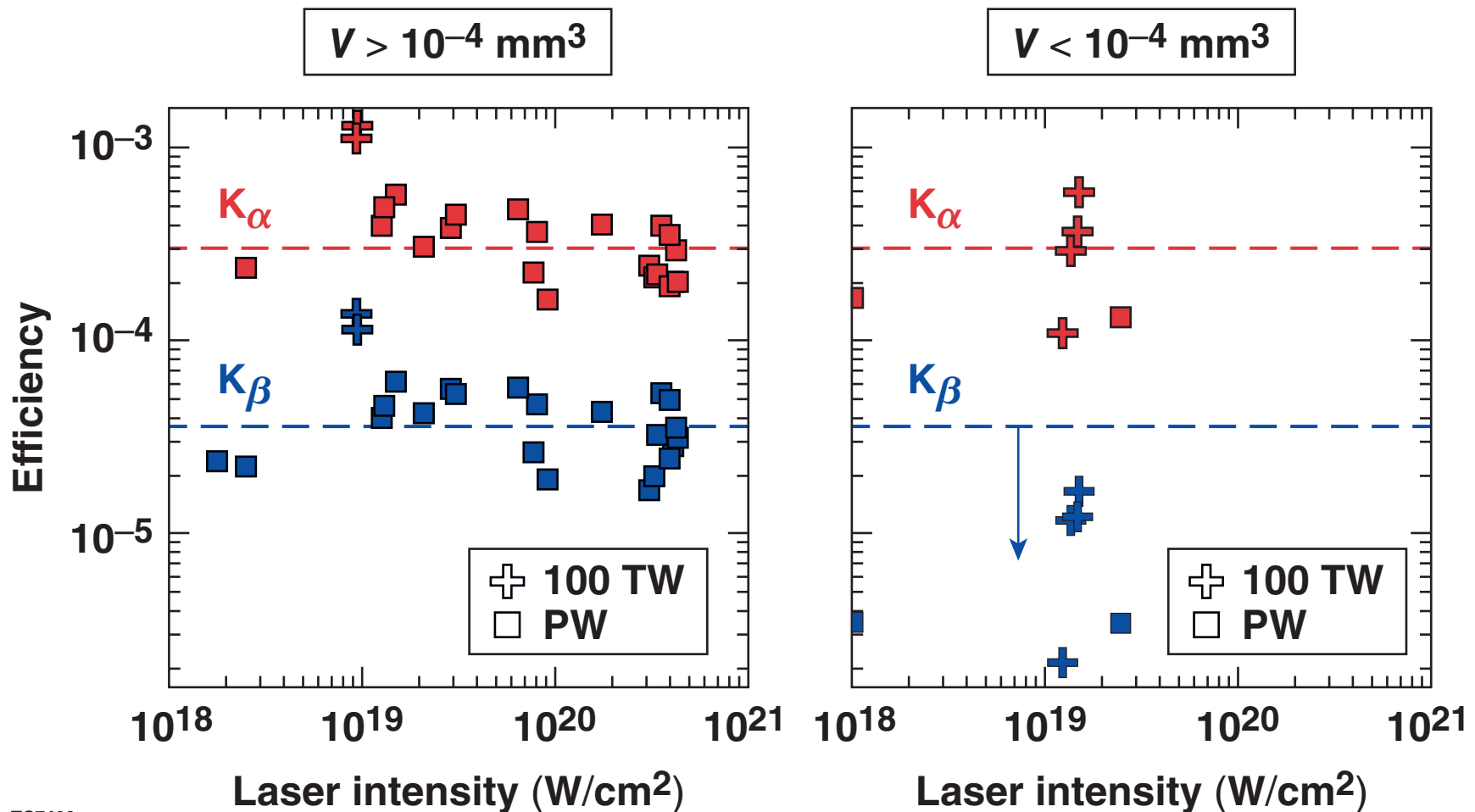


The model is completely specified if the hot-electron conversion efficiency is known



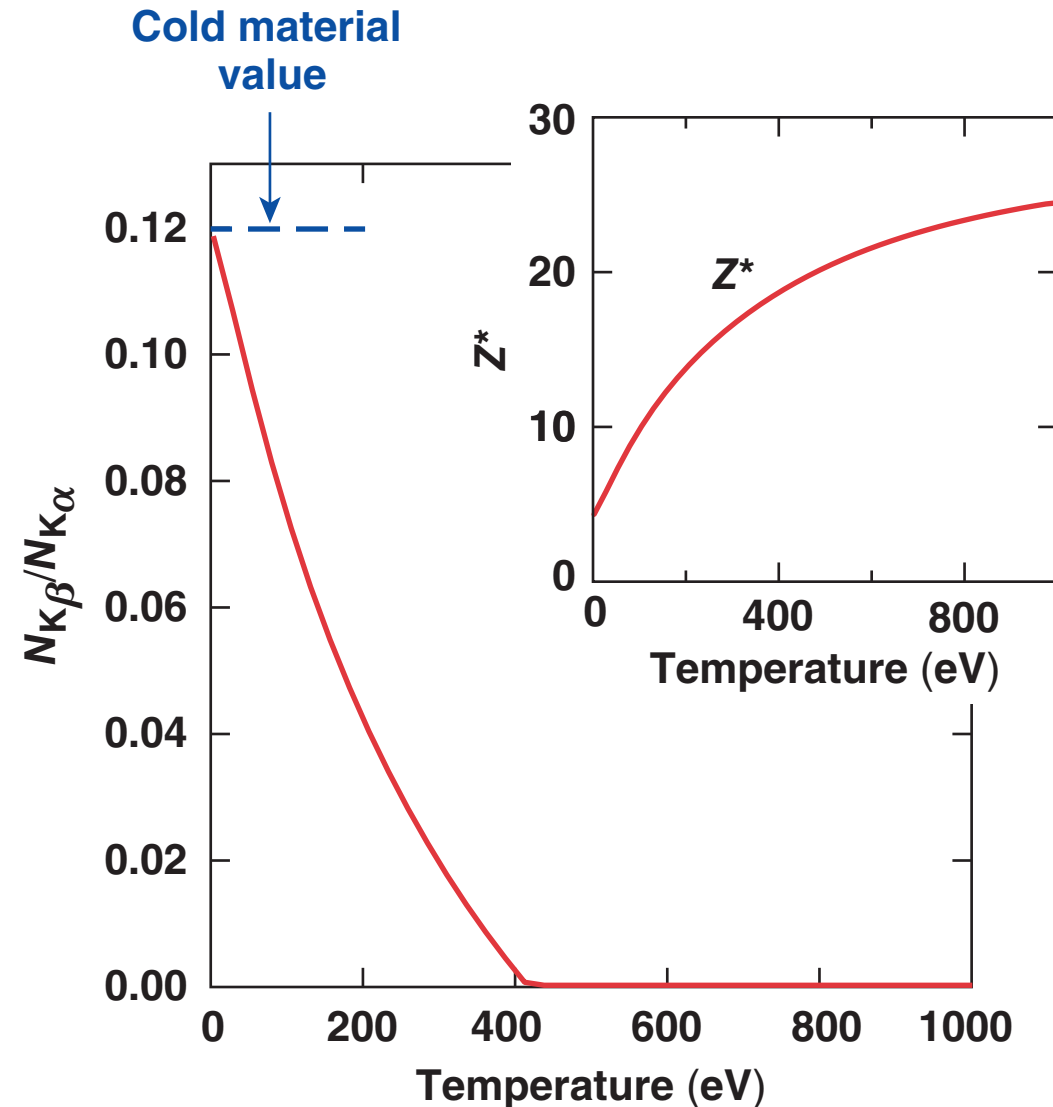
The K_β/K_α ratio can be used as a consistency check on the hot-electron conversion efficiency

- For small volume targets, $V < 10^{-4} \text{ mm}^3$, the K_β efficiency is dramatically reduced



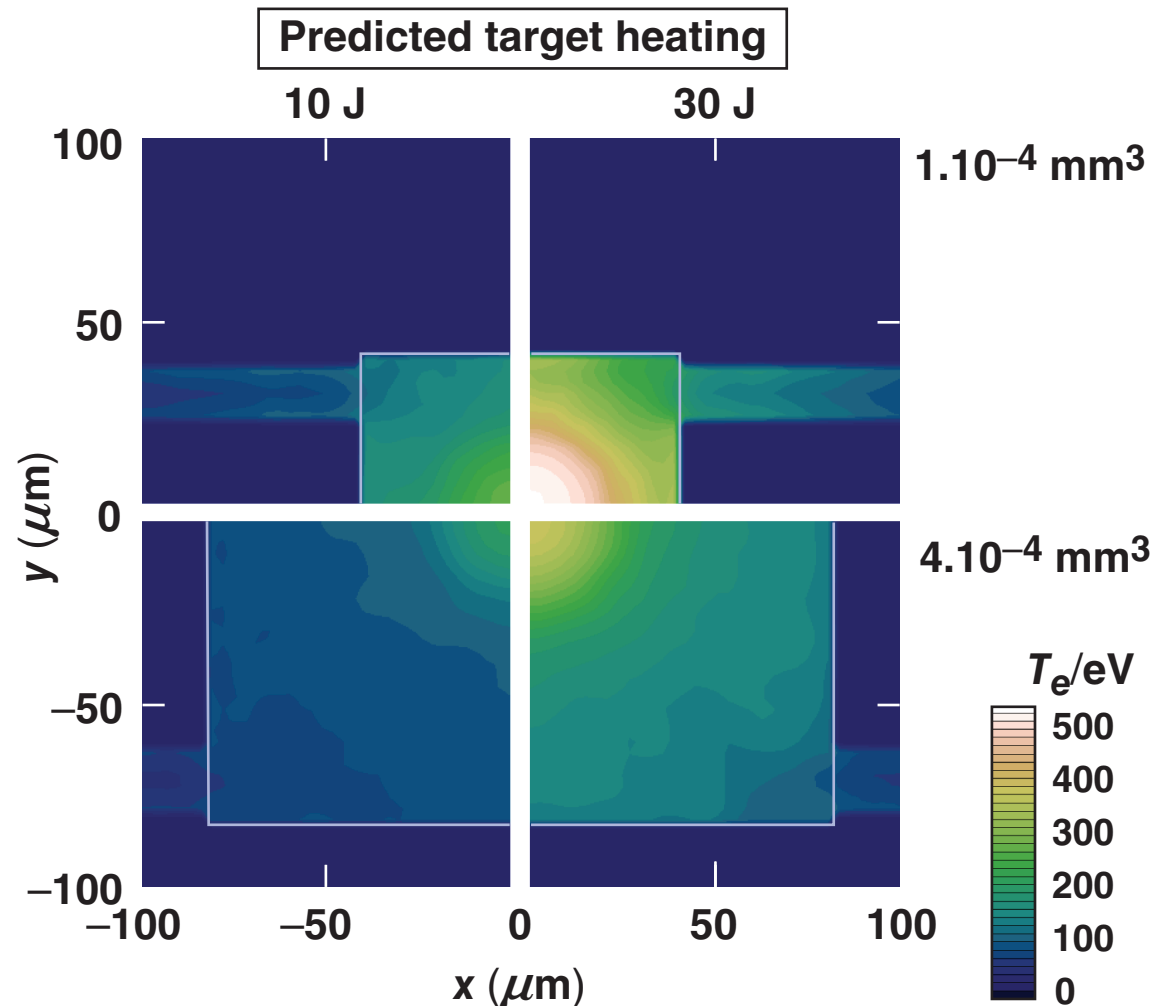
The reduction in K_{β}/K_{α} ratio for small target volumes is expected due to target heating

- Small volume targets can reach temperatures of several hundred eV
- Depletion of the copper M-shell reduces the K_{β} emission relative to K_{α}
- We use T-F to get Z^* from local $[n(x,t), T(x,t)]$
- Modify $(p_{K_{\alpha}}, p_{K_{\beta}})$ appropriately



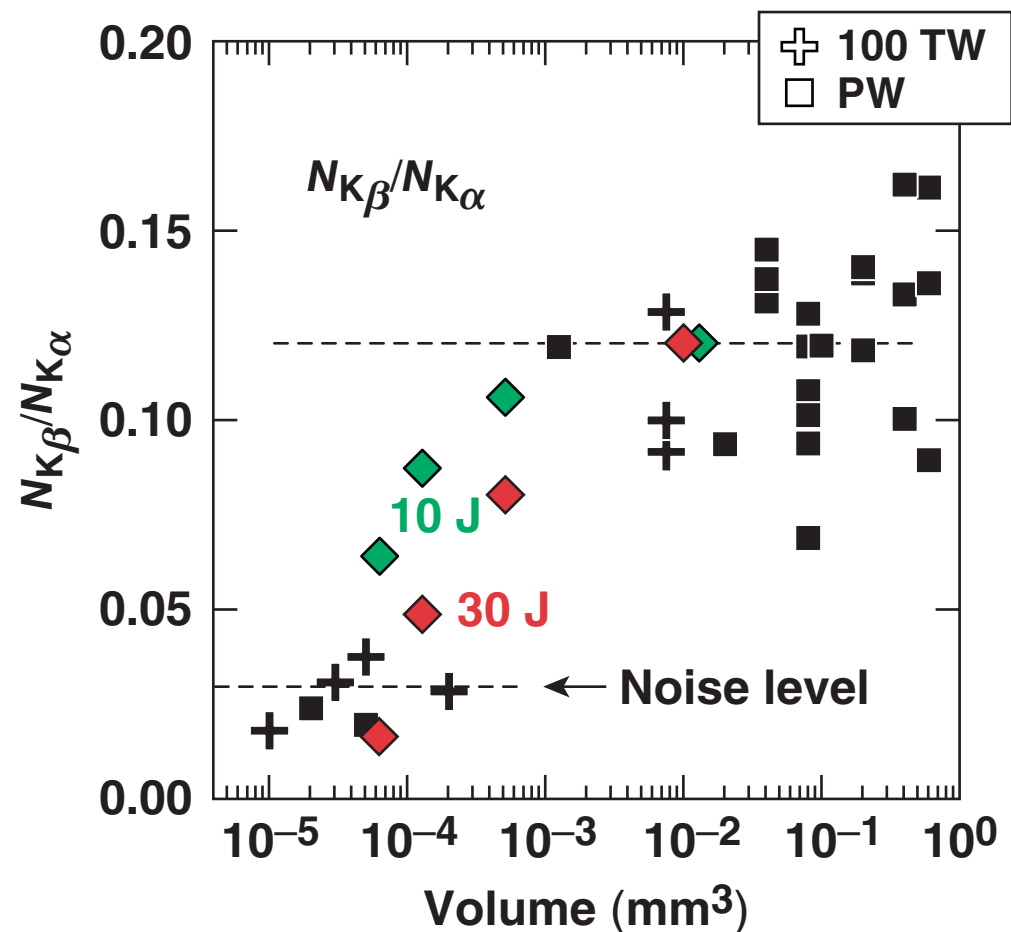
3-D *LSP* calculations predict target heating for a given conversion efficiency

- Essentially all of the hot-electron energy is deposited in the target
- Temperature rise depends upon the energy deposited
- Assuming Thomas–Fermi EOS model
- Smallest volume $V \sim 10^{-4} \text{ mm}^3$ targets are predicted to reach several hundred eV



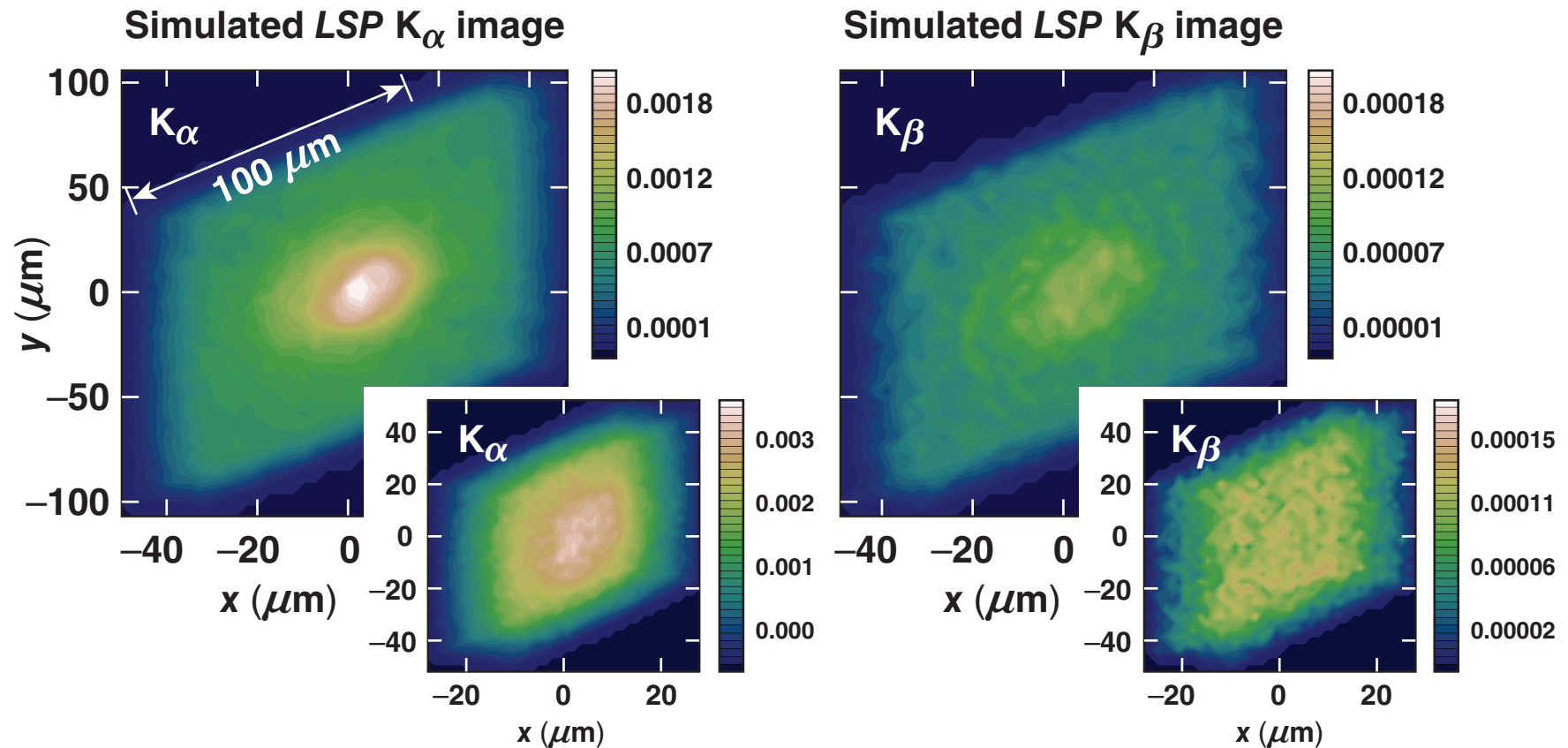
The K_{β}/K_{α} ratio is computed in *LSP* by modifying the K-photon emission probabilities using the local temperature

- The data is not sufficiently precise to either confirm or reject the conversion efficiency obtained by fitting the absolute K_{α} yield
- Spatial/temporal profile has to be taken into account (*LSP*)



LSP calculations predict the spatial distribution of K-emission which will be compared with experiment

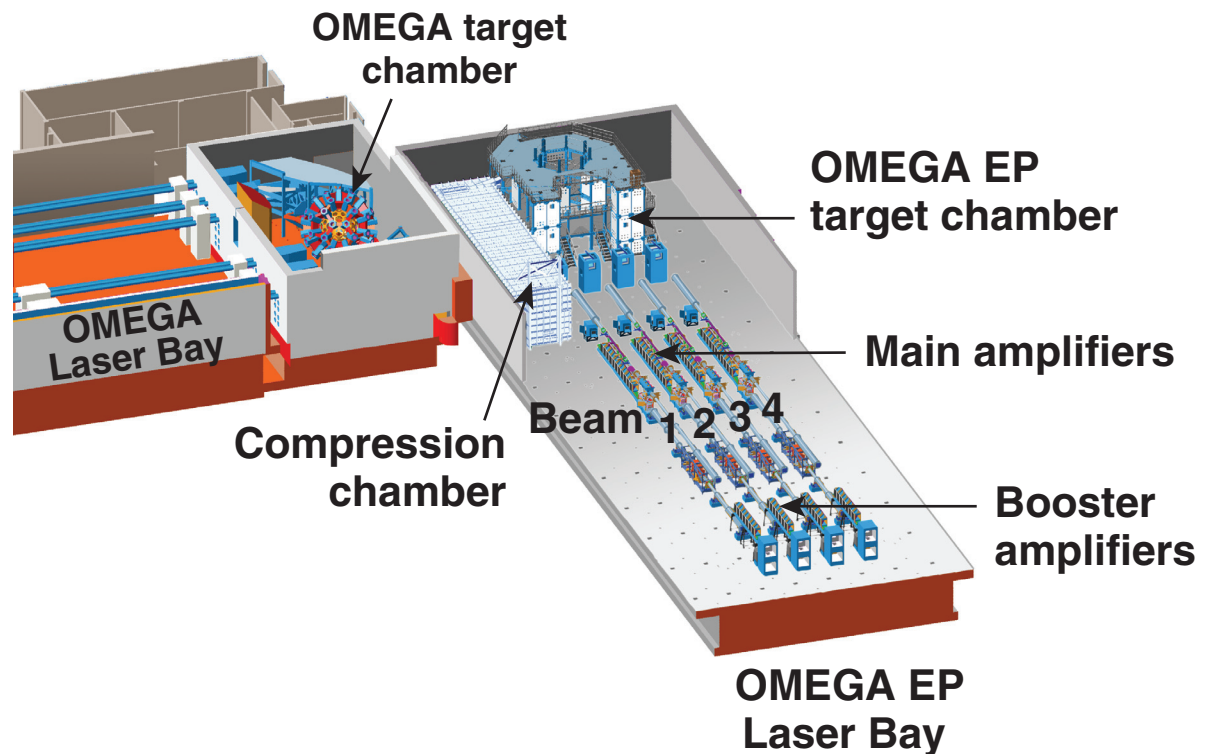
Three-dimensional *LSP* calculations



The results suggest future mass-limited foil experiments on OMEGA EP



- Dedicated experiments with well-characterized interaction conditions
- K_{β}/K_{α} ratio probes bulk temperature
 - relevant for FI
 - imaging
- Consistency between absolute yield and line ratio
- XUV probes surface temperature
- Benchmark codes



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