Effects of Resonant Absorption in Direct-Drive Target Designs on OMEGA



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Summary

Resonant absorption is important during the rapid increase of laser power in direct-drive target designs on OMEGA

- Resonant absorption on OMEGA is determined by linear effects.
- Resonant absorption is important during laser pickets or at the beginning of long laser pulses, when the density length scale near the critical surface is relatively small, $L < 2 \ \mu$ m.

- In spherical implosions, resonant absorption can enhance the earliertime laser absorption up to 20%.
- Planar OMEGA experiments will validate theoretical predictions with the use of inclined *s* and *p*-polarized laser beams.



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A simplified model of resonant absorption predicts a very large electric field at the resonance peak

An oblique ray incident onto a cold inhomogeneous plasma slab z > 0



• Ponderomotive force $\left[\sim (\nabla \vec{E})^2 \right]$ can be important in the resonance region.

- Mechanisms limiting the field:
 - electron-ion collisions
 - thermal convection
 - nonlinear wave breaking

The effect of Langmuir waves has been included in the calculation of laser absorption in the 1-D code LILAC

• The linearized electron-momentum equation

$$\frac{\partial \vec{u}_{e}}{\partial t} = -\frac{e}{m_{e}}\vec{E} - \nu \vec{u}_{e} - \frac{1}{m_{e}n_{e}}\nabla P_{e}$$

combined with Maxwell's equations,¹

$$\Delta \vec{E} - \nabla (\nabla \cdot \vec{E}) + \frac{\omega^2}{c^2} (\vec{E} + i\frac{4\pi}{\omega}\vec{j}) = 0,$$

$$\vec{j} = i\frac{\omega_{pe}^2 \vec{E}}{4\pi(\omega + i\nu_{em})} - \frac{i\frac{3v_{Te}^2 \nabla (\nabla \cdot \vec{E})}{4\pi(\omega + i\nu_{w})}}{i\frac{3v_{Te}^2 \nabla (\nabla \cdot \vec{E})}{4\pi(\omega + i\nu_{w})}}.$$



 $v_{em} = (collisional damping)$

 $v_{W} = ($ collisional damping) + (Landau damping)

Generation of Langmuir waves is the dominant mechanism that limits the amplitude of resonant fields



 Under typical conditions on OMEGA,

 $\lambda_{\text{las}}^2 I \approx (1 \text{ to } 10) \times 10^{13} \cdot \text{W/cm}^2,$ $L \approx 1 \text{ to } 2 \,\mu\text{m}, T_e \approx 0.5 \text{ to } 1 \text{ keV},$ $P_e \gg E_{\text{fs}}^2 / 8\pi, \text{CH} - \text{targets};$

the convection of the Langmuir waves reduces the amplitude of the resonance field below the wave-breaking limit.

- Landau damping of Langmuir waves produces hot electrons with $T_h \approx 5$ to 10 keV.
- The ponderomotive force does not exceed the pressuregradient force and has a small dynamic effect.

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Reflected laser-light measurements in the planar OMEGA experiments will validate theoretical predictions



 θ = 23°, 48°, and 62° I = 10¹⁴ to 10¹⁵ W/cm² Picket or square laser pulses

Several laser beams can be used simultaneously on OMEGA

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- Five laser beams at 23.2°.
- Background beams have a mixed *s* and *p* polarization.
- Diagnostics will also include shock breakout and x-ray measurements.

Simulation of 23°, 1-ns laser beams predict different absorption of *s*- and *p*-polarized light



 The effect of resonant absorption is smaller for beams with larger incident angles.



Planar simulations with angle-dependent laser intensity estimate absorption in spherical implosions

Illumination of a Simulated absorption for a 25-kJ implosion spherical target $I(\mathbf{r}) = I_0 \mathbf{e}^{-(\mathbf{r}/\mathbf{r}_0)^n}$ Intensity (10¹⁵ W/cm²) 1.0 Laser pulse 0.5 $n = 412 \ \mu m$ **H** $r_0 = 352 \ \mu m$ 0.0 0.20 $(Q^{SP}-Q^{S})/Q^{S}=0.045$ $Q_{W}^{SP} - Q^{SP} = 0.013$ QSF 0.15 0.10 QSP_ QSP_ $\hat{R} = 430 \ \mu m$ 0.05 $I(\mathbf{r}) \Rightarrow I(\theta), \theta = 0...\pi/2$ 0.00 0.5 0.0 1.0 Angle-dependent intensity Time (ns) in planar geometry

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