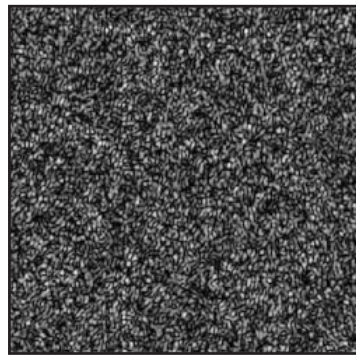


Rayleigh–Taylor Growth Measurements of 3-D Modulations in Nonlinear Regimes



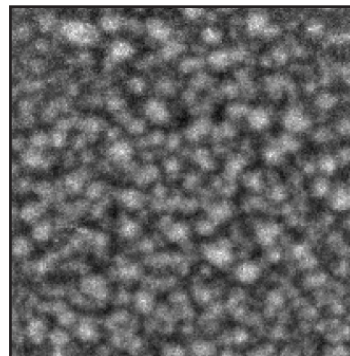
**SG8 DPP
laser modulation
(input)**



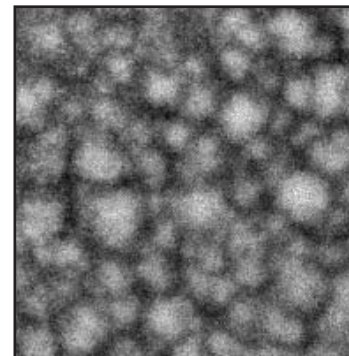
← 333 μm →

X-ray radiographs at

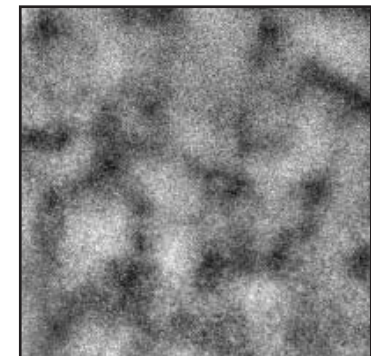
4 ns



6 ns



10 ns



**V. A. Smalyuk
University of Rochester
Laboratory for Laser Energetics**

**47th Annual Meeting of the
American Physical Society
Division of Plasma Physics
Denver, CO
24–28 October 2005**

Summary

The measured nonlinear Rayleigh–Taylor growth is in good agreement with theoretical models



- Direct-drive, planar through-foil radiography is analyzed in real and Fourier spaces.
- Nonlinear velocities and spectral shapes are consistent with Haan-model* predictions.
- Late nonlinear growth is relatively insensitive to initial conditions.
- The self-similar regime** of bubble size distributions was observed with the modulation σ_{rms} growing as $\alpha_{\sigma}gt^2$, with $\alpha_{\sigma} = 0.027$.

⇒ Bubble front $\alpha_{\text{B}} \simeq \sqrt{2} \cdot \alpha_{\sigma} = 0.04$ is consistent with Haan's saturation parameter $\left(S_{\text{k}} = \frac{2}{Lk^2}\right)$ and previous classical RT results.

*S. Haan, Phys. Rev. A 39, 5812 (1989).

**U. Alon *et al.*, Phys. Rev. Lett. 74, 534 (1995).

Contributors



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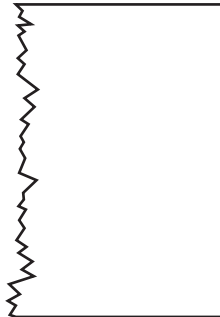
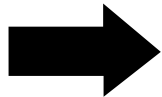
**University of Rochester
Laboratory for Laser Energetics**

**O. Sadot, D. Shvarts
Nuclear Research Center
Negev, Israel**

The Rayleigh–Taylor instability has linear, nonlinear, and turbulent stages

Linear stage

Acceleration



Small-amplitude modulations grow exponentially in the linear regime.

Nonlinear stage

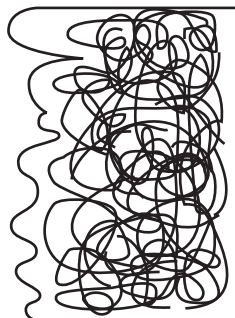
Acceleration



Bubbles compete and merge in the nonlinear regime.

Turbulent-mixing stage

Acceleration



Chaotic mixing is characteristic in the turbulent regime of classical Rayleigh–Taylor instability.

In the Haan saturation model, 3-D broadband modulations saturate at $S_k = \frac{2}{Lk^2}^*$

- For a single-mode modulation Z_k the saturation level is 0.1λ

$$S_k = 0.1 \lambda$$

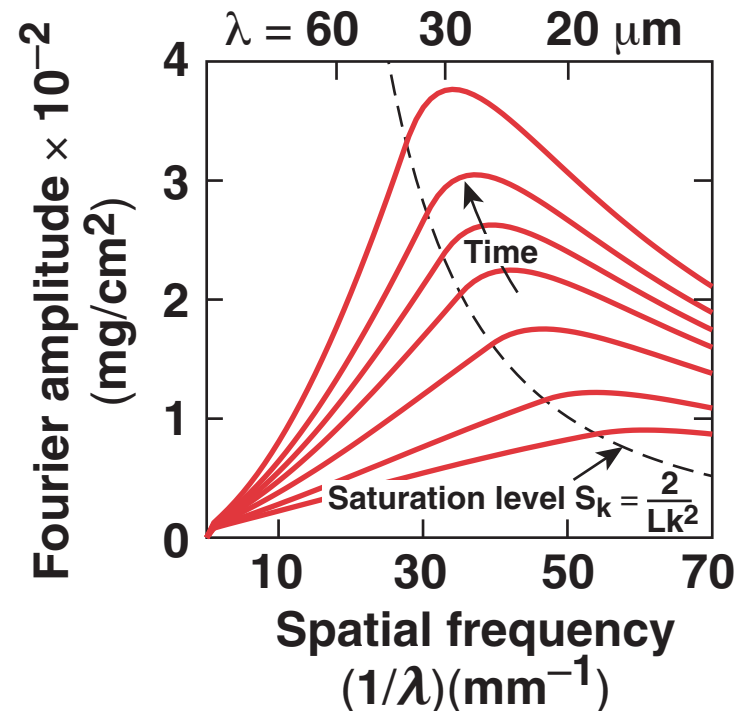
- For broadband modulations saturation, the σ_{rms} of modes near Z_k is 0.1λ

$$\sqrt{\sum_{k-\delta k}^{k+\delta k} S_k^2} = 0.1 \lambda$$

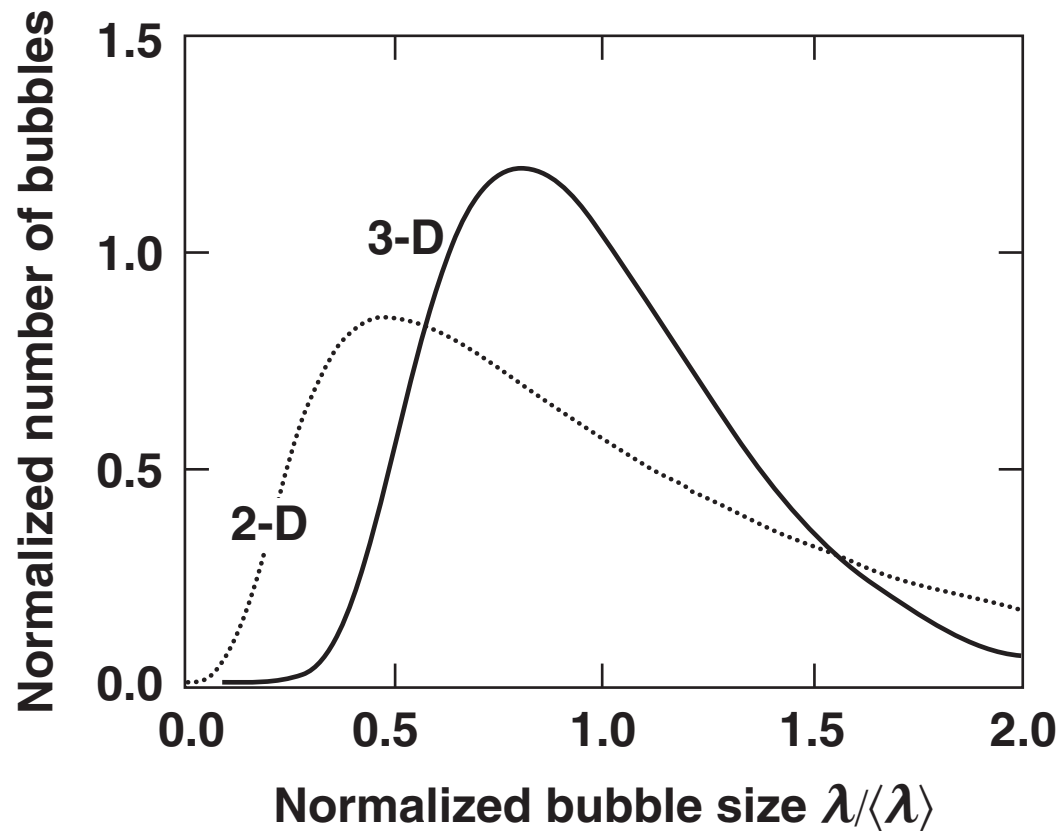
- $S_k = \frac{\epsilon}{Lk^2}$, where $\epsilon = 2$

In the Haan saturation model,* the Fourier amplitudes of broadband modulations saturate at $S_k = \frac{2}{Lk^2}$

- Fourier amplitudes grow exponentially up to the saturation levels $S_k = \frac{2}{Lk^2}$
- Subsequent growth is linear in time with velocities $V_s = S_k \times \gamma(k)$ [$\gamma(k)$ is the linear growth rate].
- The shape of the late modulation spectrum is relatively insensitive to the initial modulation spectrum and to growth history (growth rates).



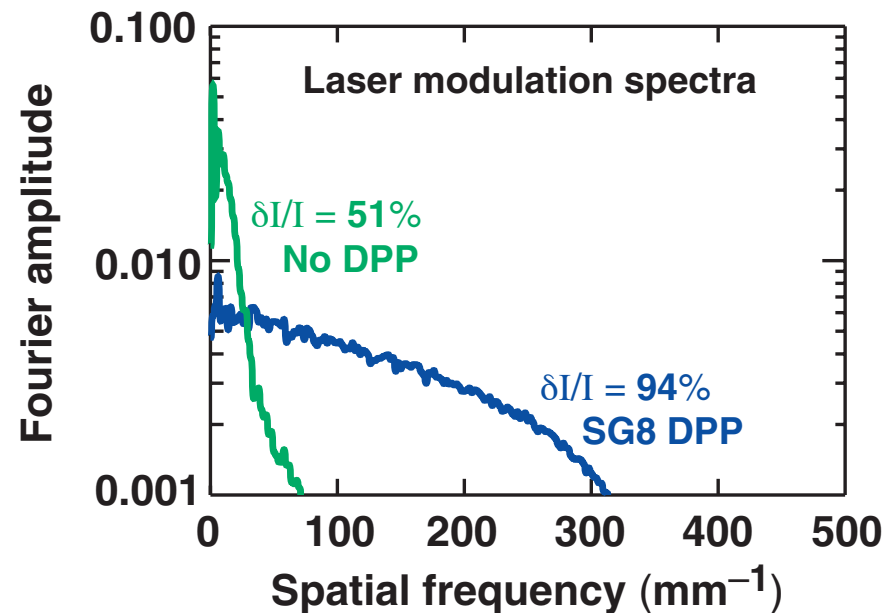
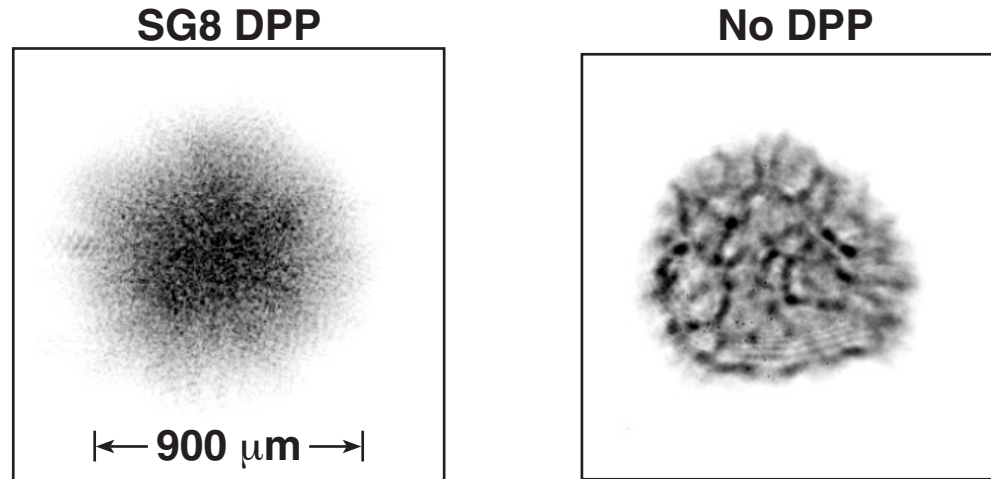
In real-space analysis, bubble competition models predict that bubble size distributions evolve in self-similar regimes*



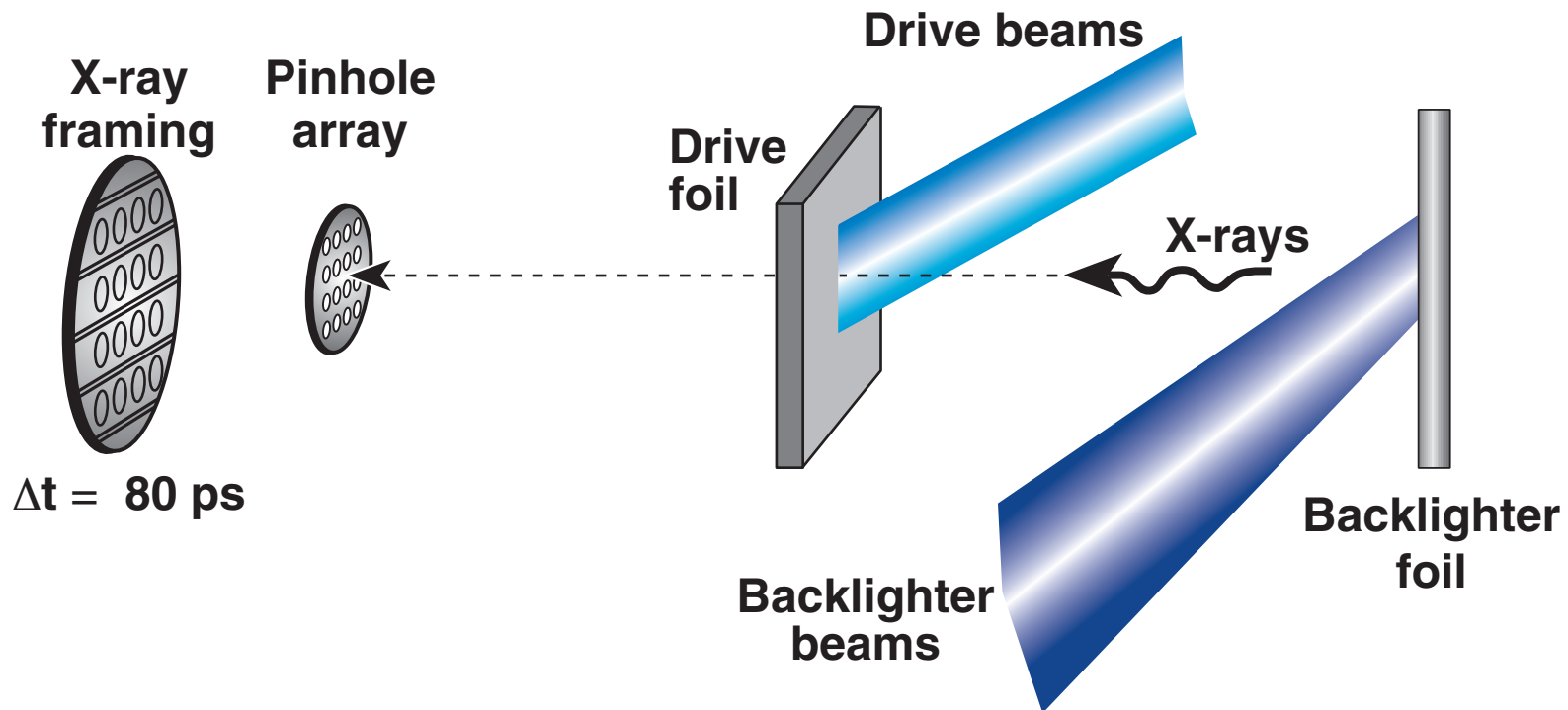
*U. Alon *et al.*, Phys. Rev. Lett. 74, 534 (1995).

*D. Oron *et al.*, Phys. Plasmas 8, 2883 (2001).

Initial seeds for the Rayleigh–Taylor instability were produced by laser imprinting



X-ray framing cameras are the primary diagnostics of instability growth



- 20- μm thick foils driven by 3-ns pulses at $2 \times 10^{14} \text{ W/cm}^2$
- 50- μm thick foil driven by 12-ns pulses at $5 \times 10^{13} \text{ W/cm}^2$

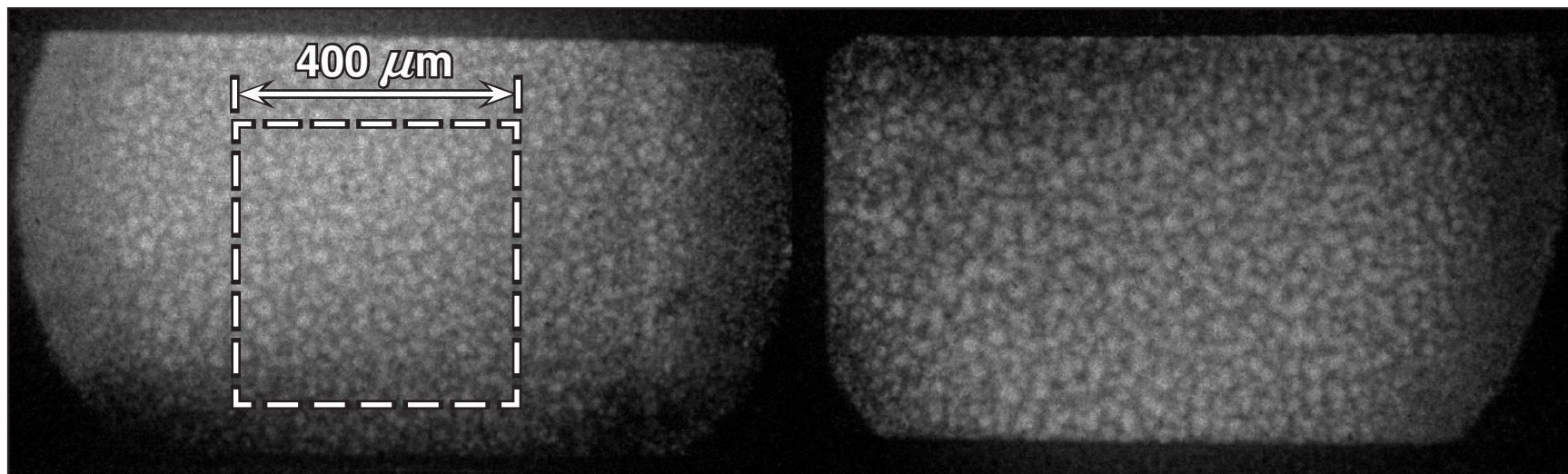
Thin Target, Early-Time Experiment

The growth of imprinted broadband modulations has been measured with a spatial resolution of $10\ \mu\text{m}$



$t = 1.8\ \text{ns}$

$t = 2.3\ \text{ns}$

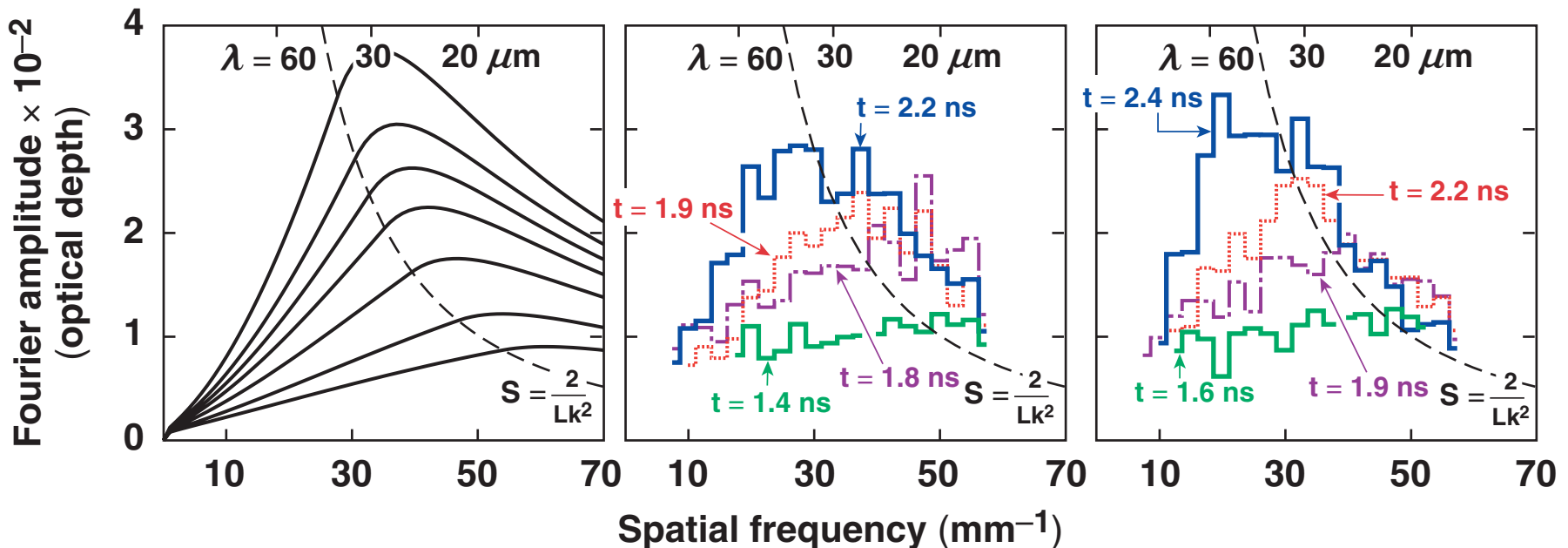


← 1 mm →

- $20\text{-}\mu\text{m}$ thick CH foils were driven with 3-ns square laser pulses at $2 \times 10^{14}\ \text{W}/\text{cm}^2$.
- Fourier amplitudes were obtained by taking azimuthal averages of the 2-D Fourier images.

Thin Target, Early-Time Experiment

The measured spectral evolution is in agreement with Haan's saturation level, evolving toward longer wavelengths*

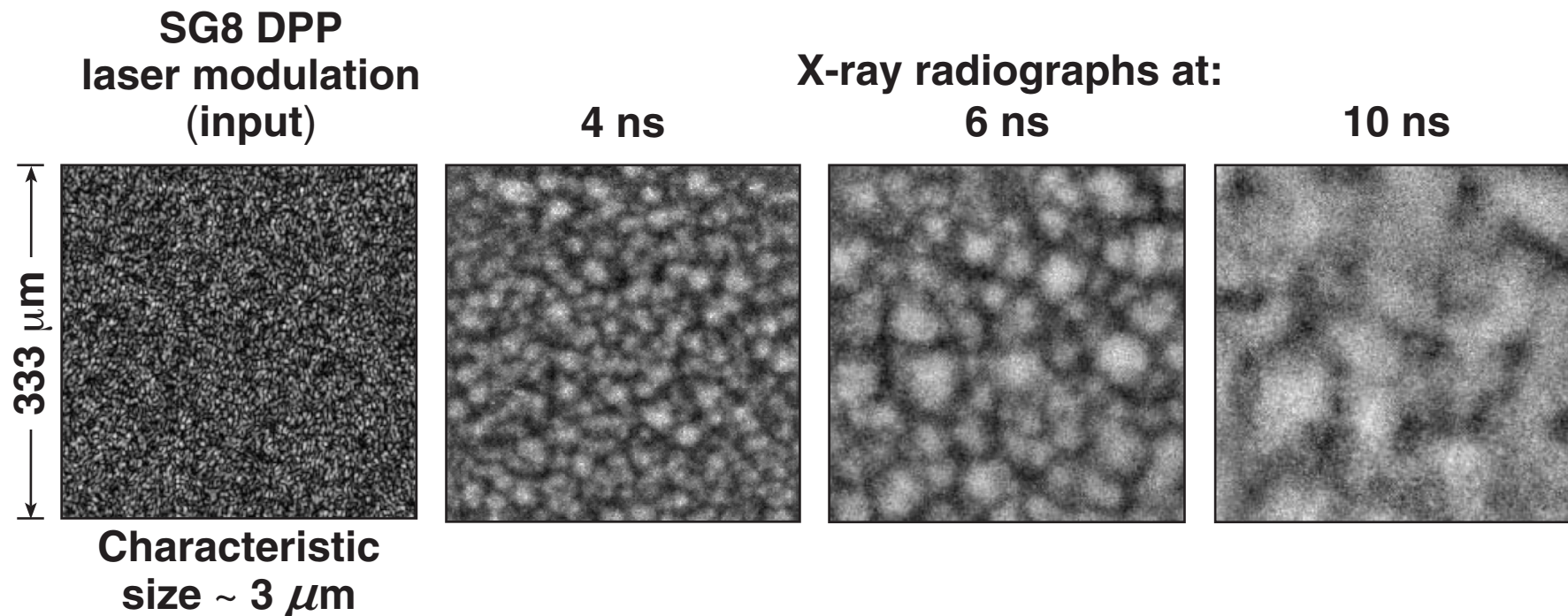


- Questions left:

1. What are the post-saturation velocities?
2. When do bubble competition and merger happen?

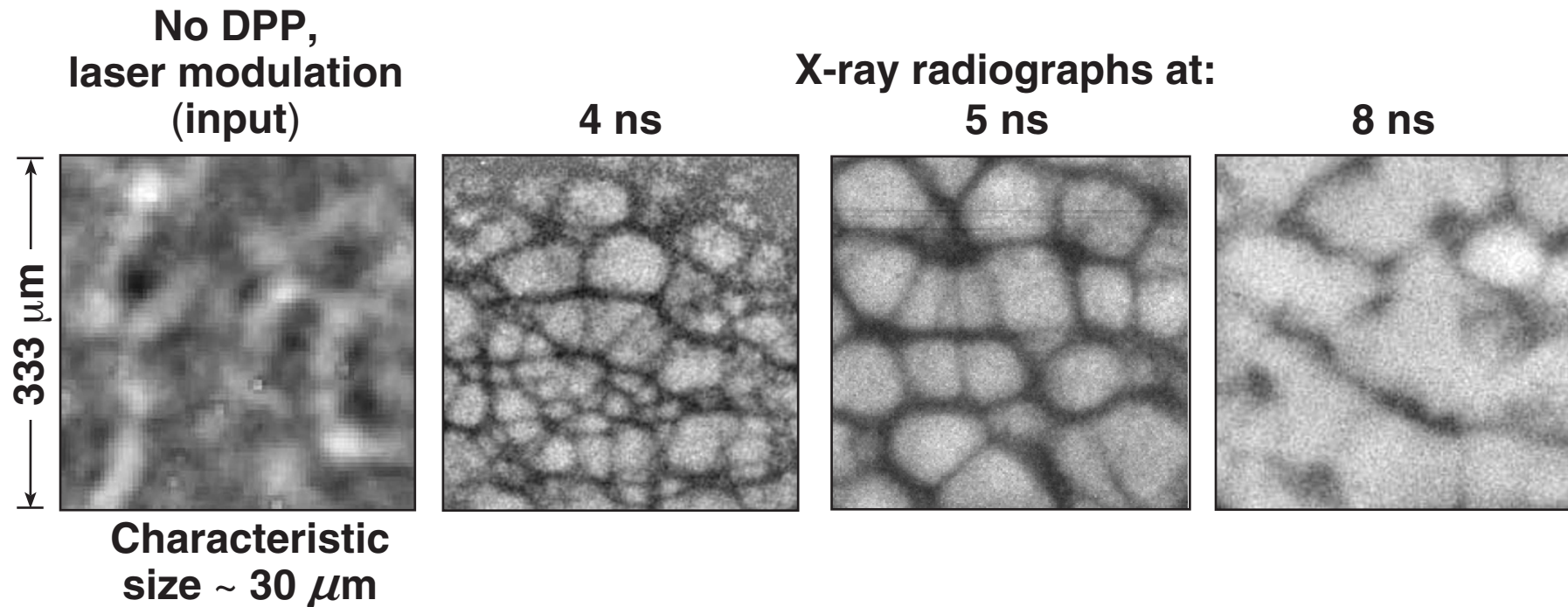
Thick Target, Late-Time Experiment

Broadband modulations become larger as they grow nonlinearly



- 50- μm -thick CH foils were driven with 12-ns-square laser pulses at $5 \times 10^{13} \text{ W/cm}^2$

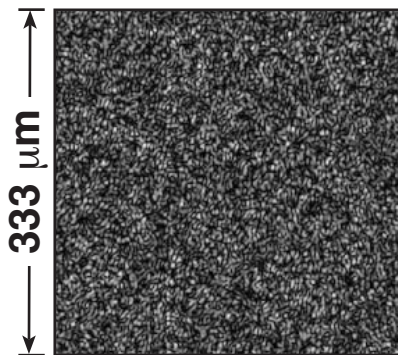
The late nonlinear evolution is relatively insensitive to initial conditions



- 50- μm thick CH foils were driven with 12-ns-square laser pulses at $5 \times 10^{13} \text{ W/cm}^2$

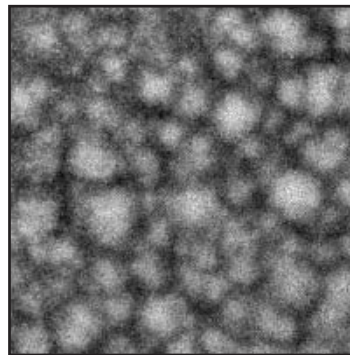
The nonlinear evolution is relatively insensitive to initial conditions

SG8 DPP
laser modulation

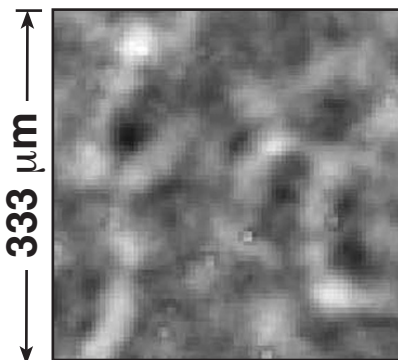


Characteristic size $\sim 3 \mu\text{m}$

X-ray radiograph
at $\sim 6 \text{ ns}$

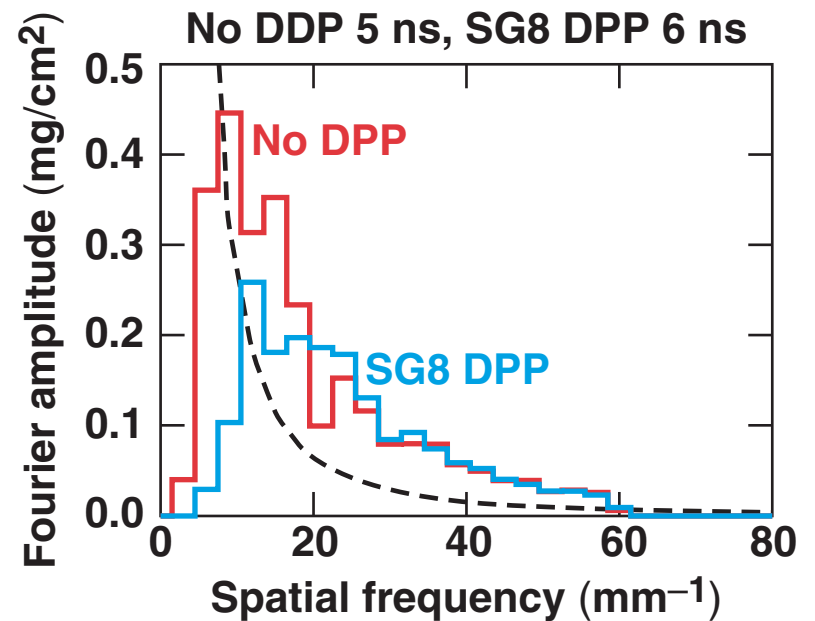
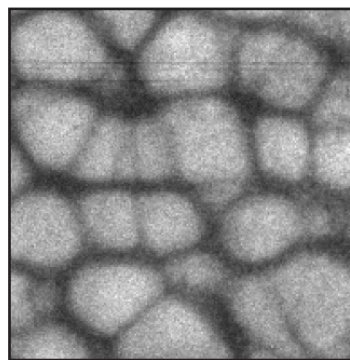


No DPP
laser modulation

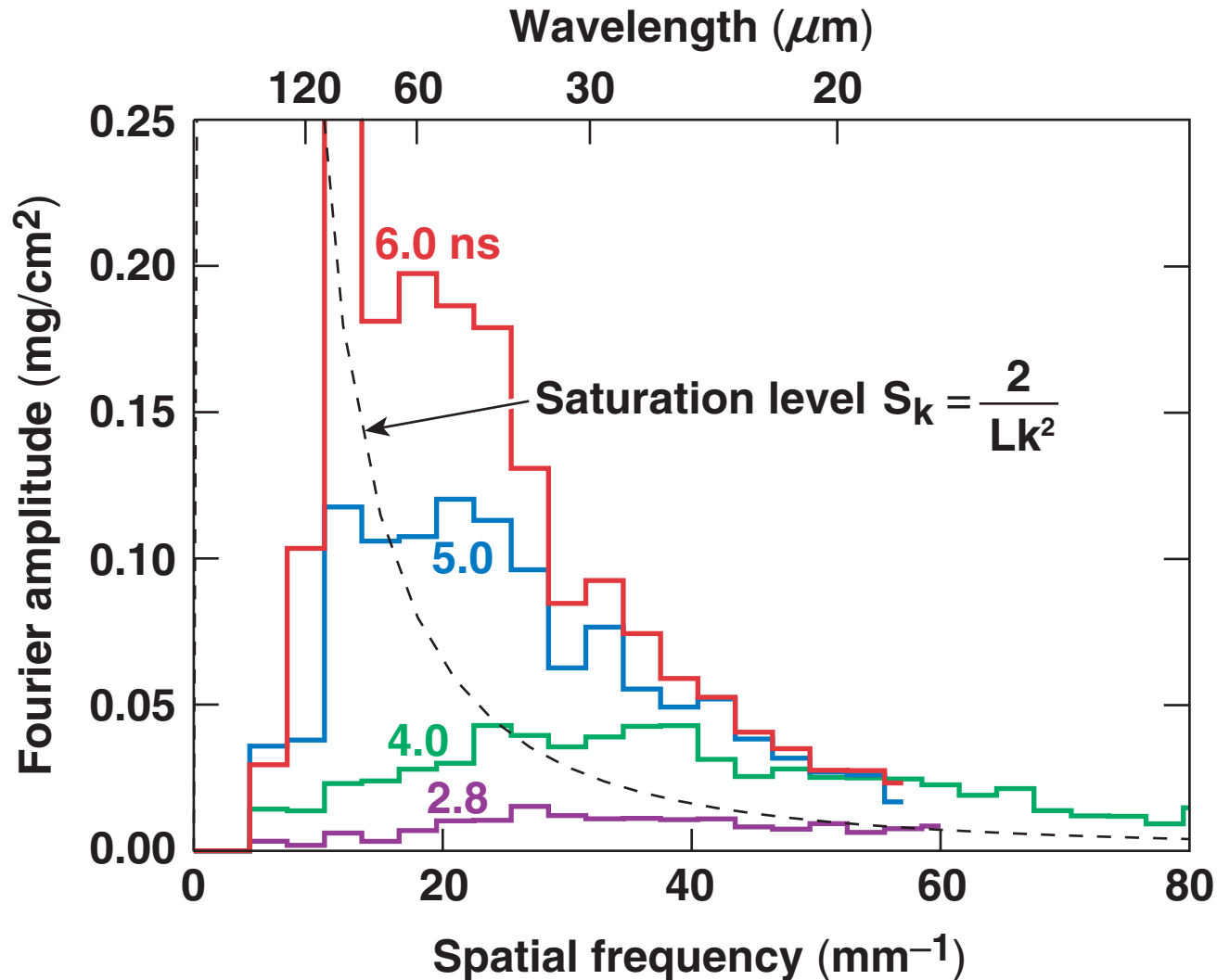


Characteristic size $\sim 30 \mu\text{m}$

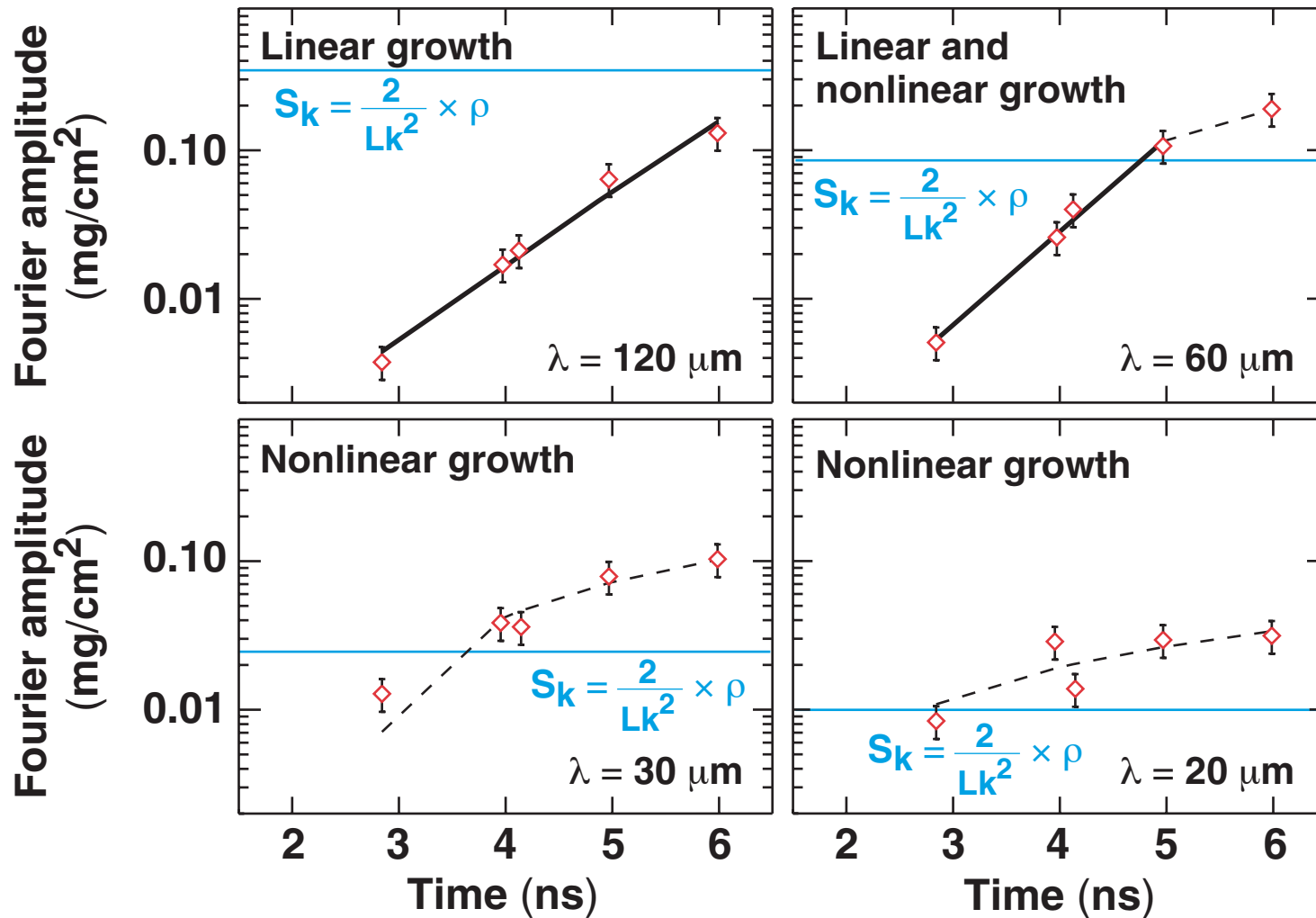
X-ray radiograph
at $\sim 5 \text{ ns}$



The modulation spectra shift to longer wavelengths as they grow, similar to Haan model predictions

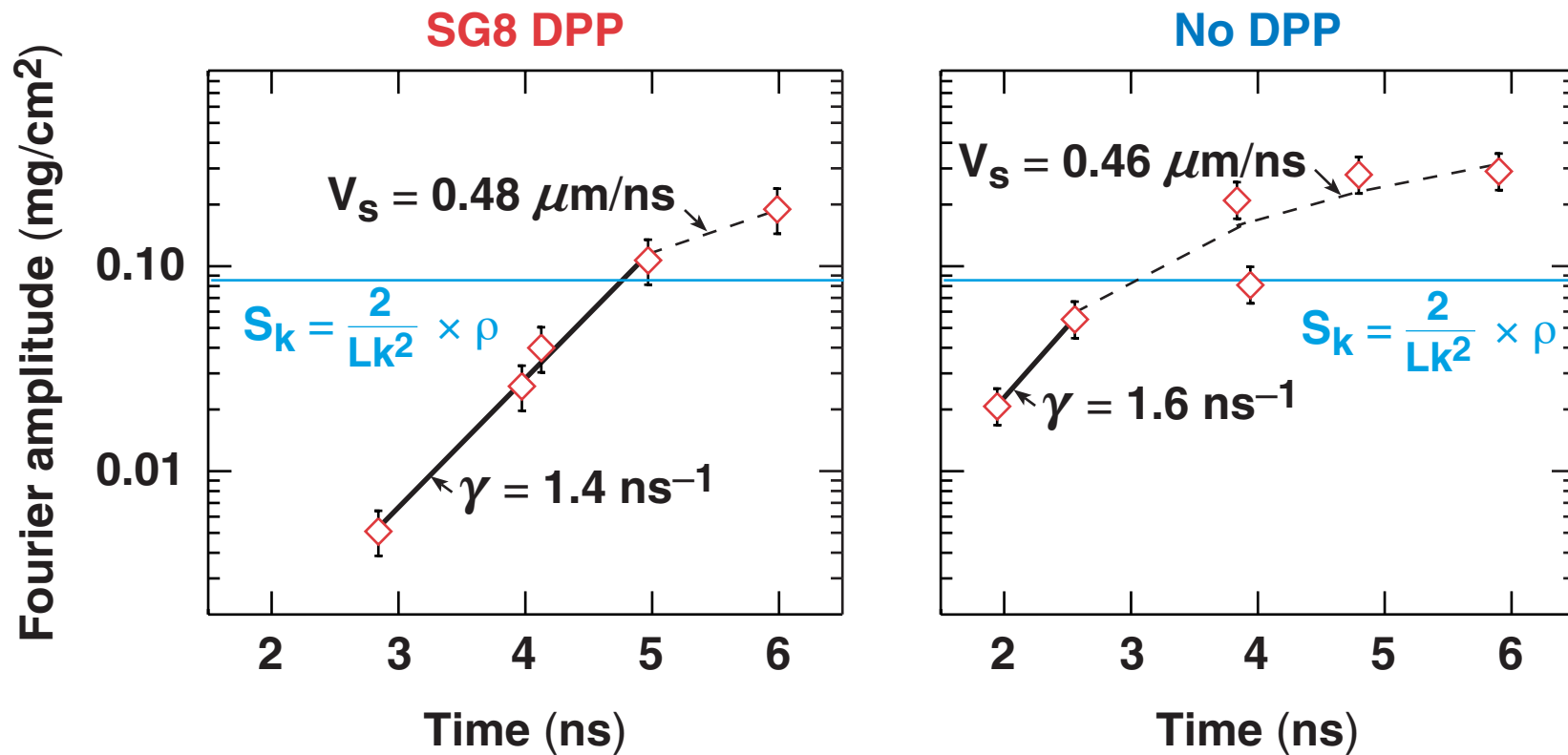


Nonlinear velocities and growth rates were determined by fitting the experimental data at various wavelengths

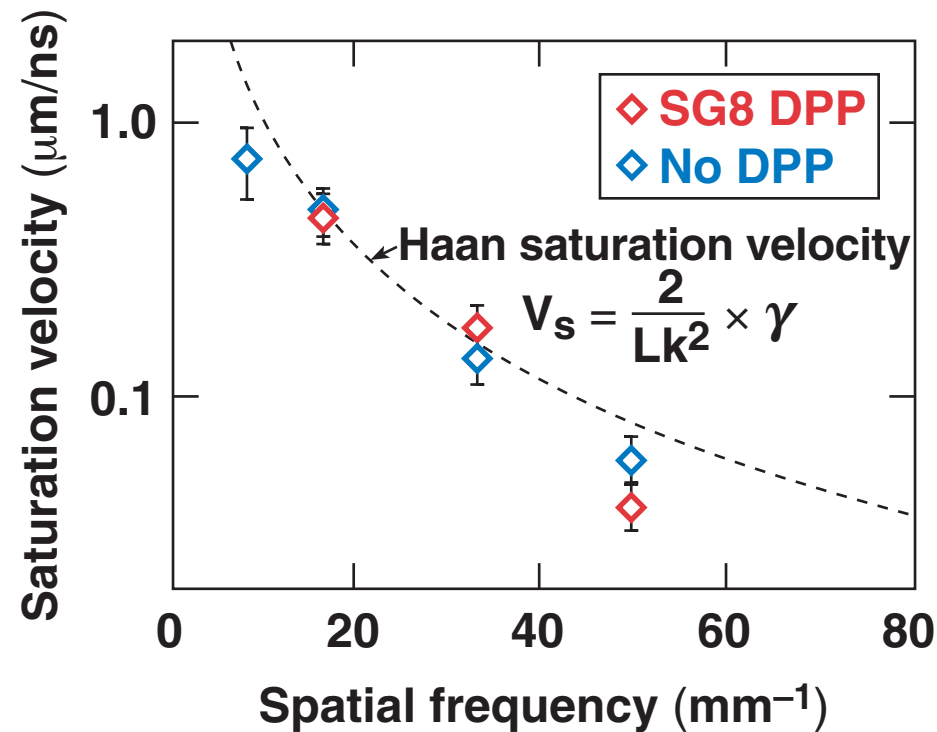


- Initial modulations were imprinted using a beam with SG8 DPP.

Saturation levels, pre-saturation growth, and post-saturation growth are insensitive to the initial spectrum

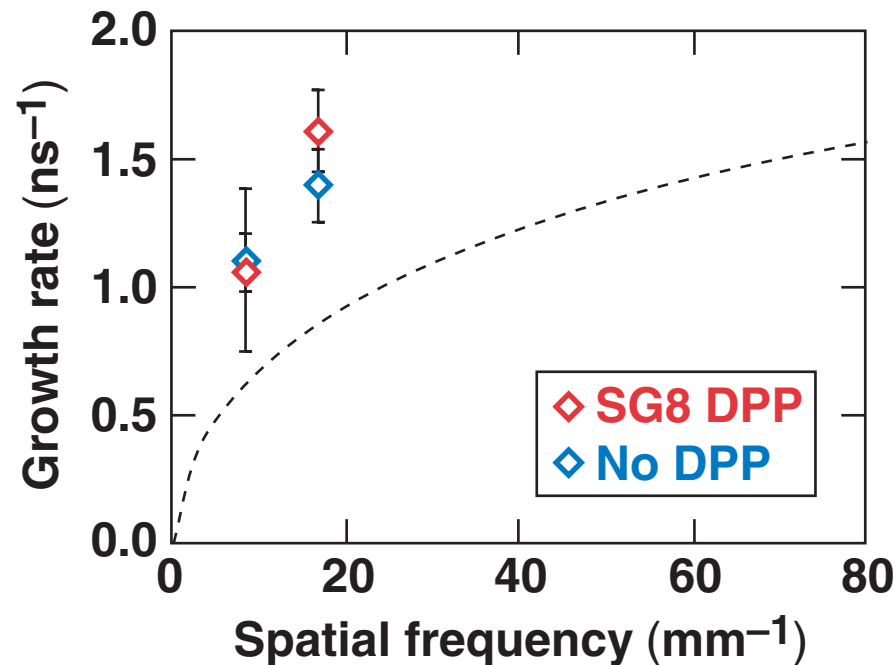


Measured nonlinear velocities are in good agreement with Haan model predictions*



- Betti–Goncharov growth rate $\gamma = 0.94 \sqrt{\frac{kg}{1 + kL_m}} - 1.5 V_a k$

The measured growth rates of longer-wavelength modes are higher than single-mode predictions



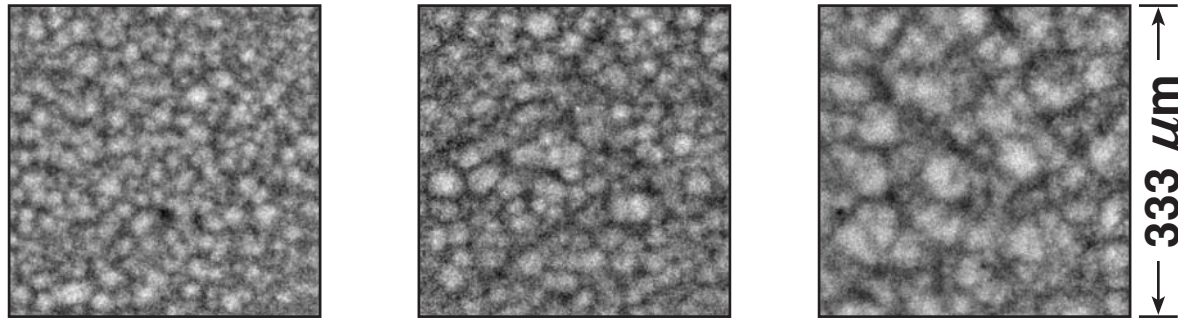
- Betti–Goncharov growth rate $\gamma = 0.94 \sqrt{\frac{kg}{1 + kL_m}} - 1.5 V_a k$
- S. W. Haan* and D. Ofer *et al.*** predict no significant mode coupling in the ablative case.
- J. Sanz *et al.*† predicted enhanced mode coupling in the ablative case compared to the classical case.

*S. W. Haan, Phys. Fluids B 3, 2349 (1991), **D. Ofer *et al.*, Phys. Plasmas 3, 3073 (1996),

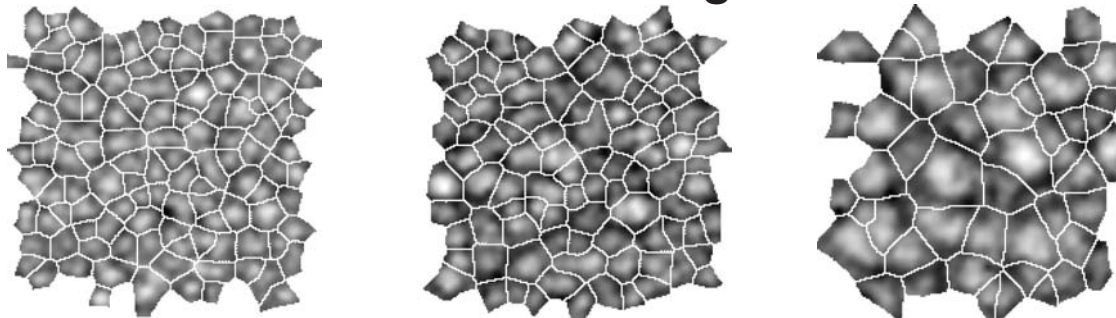
†J. Sanz *et al.*, Phys. Rev. Lett. 89, 195002 (2002).

Real-space bubble competition models describe Rayleigh–Taylor evolution more naturally

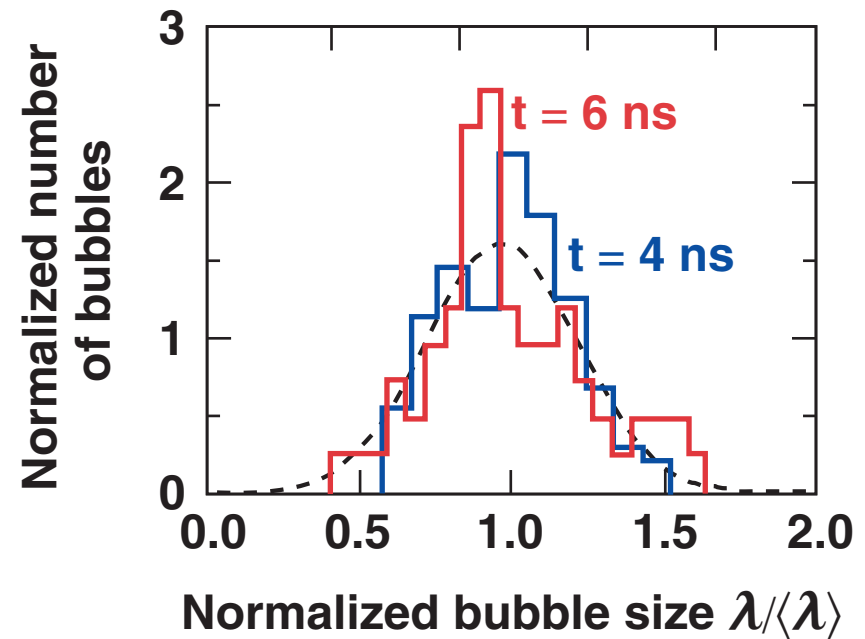
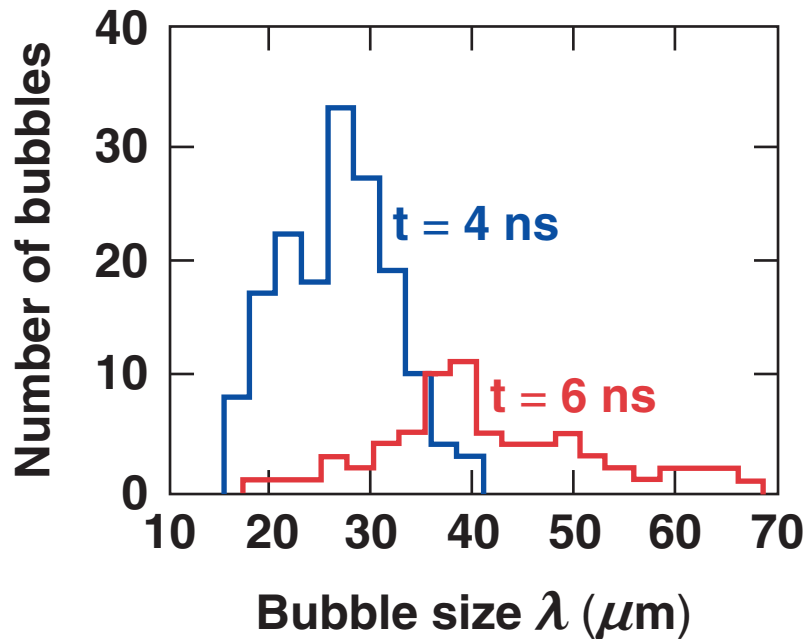
Raw images



Processed images



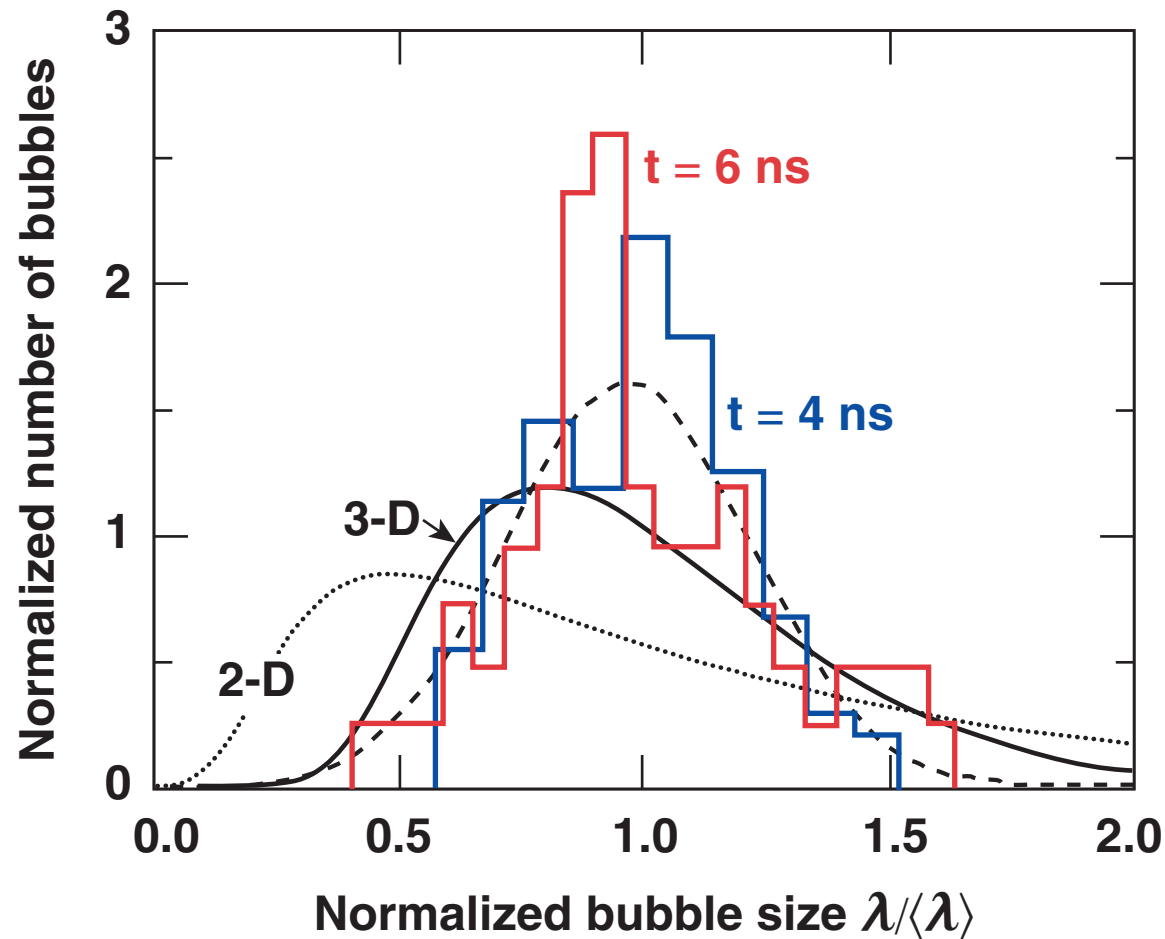
The nonlinear bubble evolution is self-similar in the nonlinear regime



Measured distributions were fit with a normal distribution function.

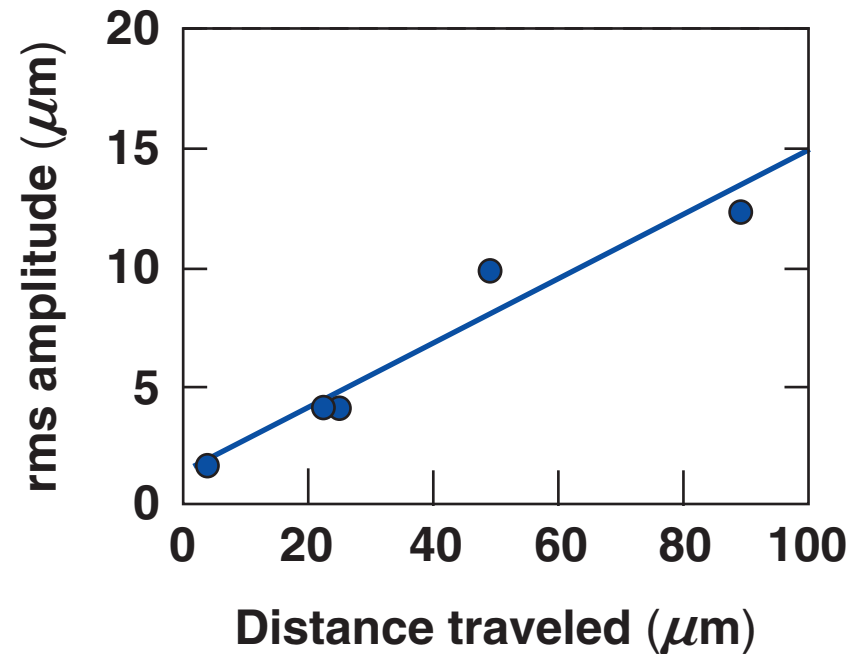
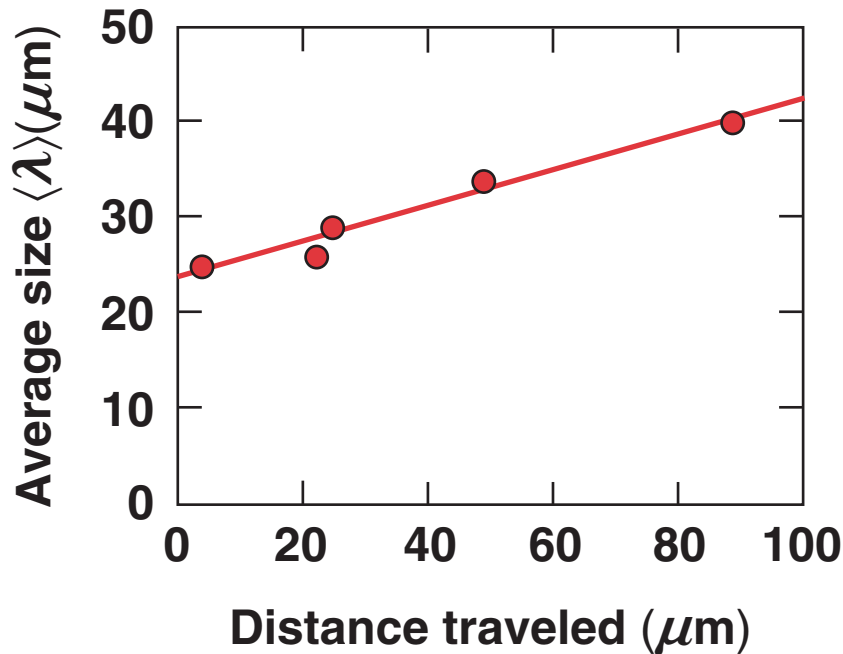
$$f\left(\frac{\lambda}{\langle\lambda\rangle}\right) = \frac{1}{\sqrt{2\pi} C_\lambda} \exp\left[-\frac{\left(\frac{\lambda}{\langle\lambda\rangle} - 1\right)^2}{\sqrt{2} C_\lambda^2}\right], \quad C_\lambda = 0.24 \pm 0.01.$$

Measured bubble size distributions are in better agreement with 3-D models than with 2-D models*



In the self-similar regimes, both the average bubble size and rms amplitude grow linearly with the distance traveled

Real-Space Analysis



- The modulation σ_{rms} grows as $\alpha_{\sigma} g t^2$, with $\alpha_{\sigma} = 0.027 \pm 0.003$.

Measured bubble-front evolution is consistent with $\alpha_B \sim 0.04$



- $h_{\text{bubble}} \simeq \sqrt{2} \alpha_{\sigma} \cdot gt^2 = 0.04 \text{ gt}^2$
- Haan model
 $S_k = \frac{2}{Lk^2}$ is consistent with $\alpha_B \sim 0.035$
if $S_k = \frac{4}{Lk^2}$ is consistent with $\alpha_B \sim 0.07$
- Measured ablative RT α_B is similar to classical RT α_B .*

**The bubble merger is evident from the evolution
of the same area of the target***



The measured nonlinear Rayleigh–Taylor growth is in good agreement with theoretical models



- Direct-drive, planar through-foil radiography is analyzed in real and Fourier spaces.
- Nonlinear velocities and spectral shapes are consistent with Haan-model* predictions.
- Late nonlinear growth is relatively insensitive to initial conditions.
- The self-similar regime** of bubble size distributions was observed with the modulation σ_{rms} growing as $\alpha_{\sigma}gt^2$, with $\alpha_{\sigma} = 0.027$.

⇒ Bubble front $\alpha_{\text{B}} \simeq \sqrt{2} \cdot \alpha_{\sigma} = 0.04$ is consistent with Haan's saturation parameter $\left(S_{\text{k}} = \frac{2}{Lk^2}\right)$ and previous classical RT results.

*S. Haan, Phys. Rev. A 39, 5812 (1989).

**U. Alon *et al.*, Phys. Rev. Lett. 74, 534 (1995).