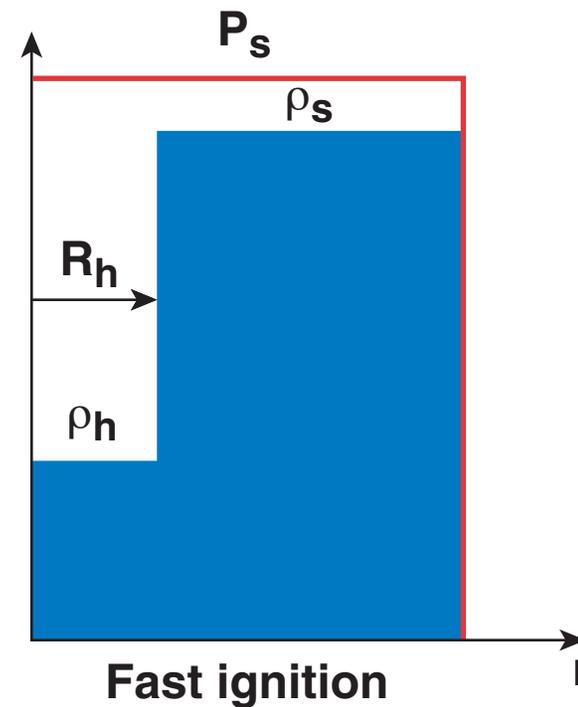
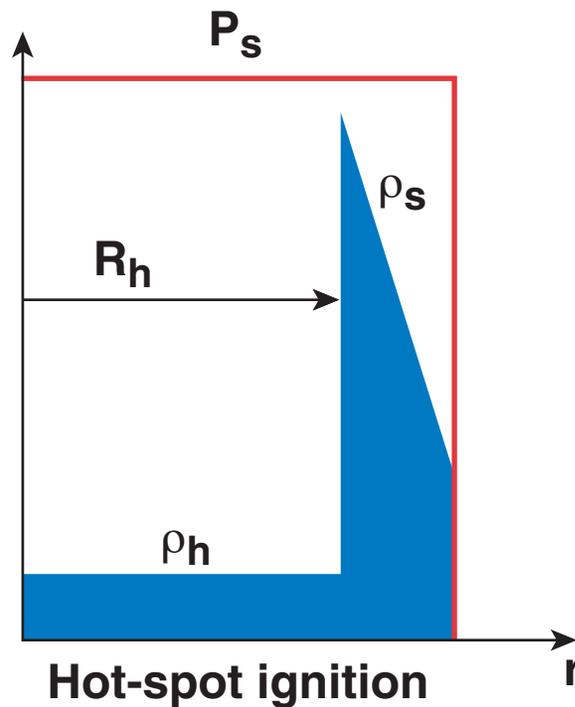
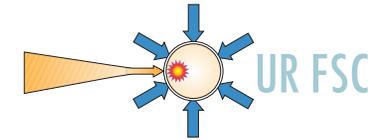


Fuel Assembly for Fast Ignition

UR
LLE



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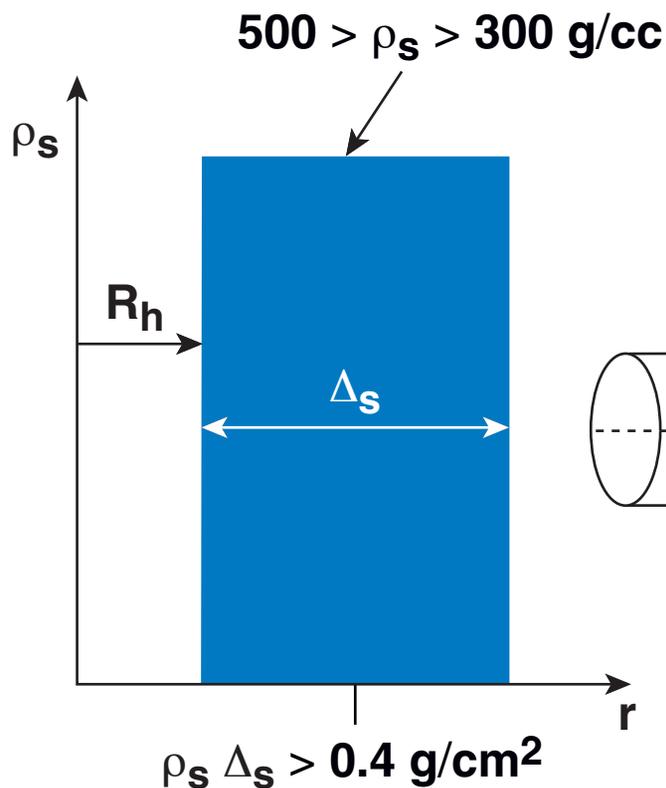
Summary

Significant progress has been made in the design of the fuel assembly for fast ignition using low-adiabat low-velocity implosions



- **A high-yield fuel assembly has been designed; it requires a 750-kJ driver to produce**
 - $\rho R \approx 3 \text{ g/cm}^2$
 - $300 < \rho < 500 \text{ g/cc}$
 - hot-spot volume/total volume $\sim 5\%$ to 7%
 - estimated yield $\sim 120 \text{ MJ}$ (if ignited)
- **A 25-kJ capsule can be assembled on the OMEGA laser to achieve**
 - $\rho R \approx 0.8 \text{ g/cm}^2$, $300 < \rho < 500 \text{ g/cc}$
 - hot-spot volume/total volume $< 15\%$
- **Planned plastic-shell implosion experiments on OMEGA are expected to achieve $\rho R \approx 0.3\text{--}0.5 \text{ g/cm}^2$ and $\rho \approx 200\text{--}300 \text{ g/cc}$.**

Ignition with fast ignition requires a fuel assembly with densities of $500 > \rho > 300 \text{ g/cc}$, $\rho R > 0.4 \text{ g/cm}^2$ and small hot-spot volume



$$E_{ig} \text{ (kJ)} = 11 \left[\frac{400}{\rho \text{ (g/cc)}} \right]^{1.85}$$

$$r_{beam} \text{ (\mu m)} = 15 \left[\frac{400}{\rho \text{ (g/cc)}} \right]^{0.95}$$

$$1 \text{ MeV e-stopping } \rho_s \Delta_s > 0.4 \text{ g/cm}^2$$

S. Atzeni, Phys. Plasmas 6, 3316 (1999).

C. K. Li and R. D. Petrasso, Phys. Rev. E. 70, 067401 (2004).

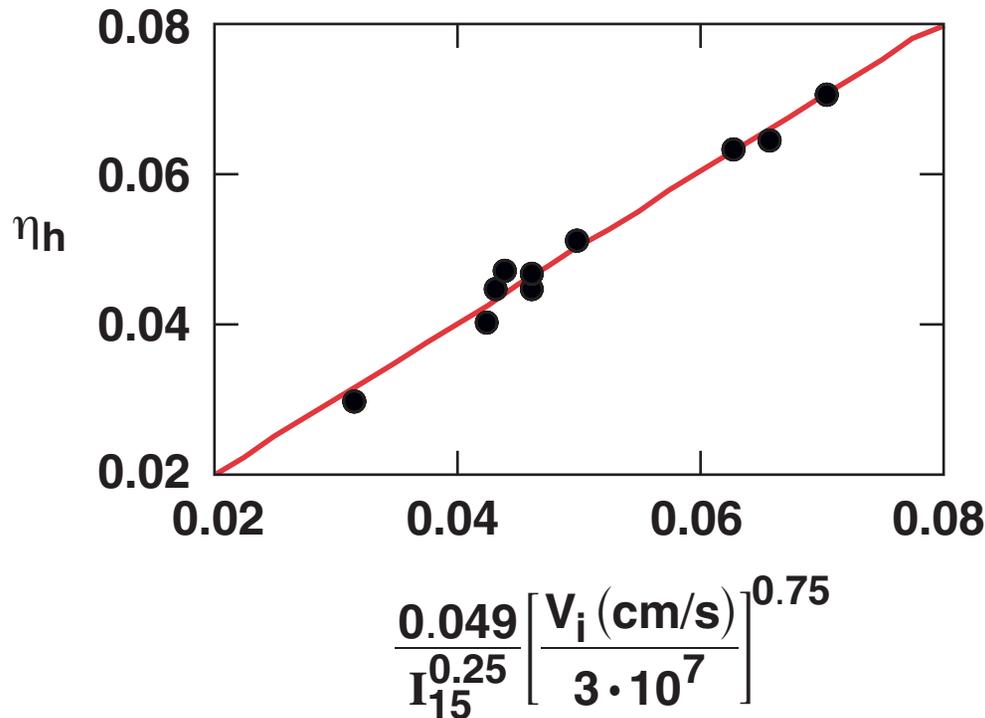
High yields with fast ignition require $\rho > 300 \text{ g/cc}$,
 $\rho R \sim 3 \text{ g/cm}^2$, small hot-spot volume, and gains > 100



$$\text{Gain} = \frac{\eta_h}{V_i^2} \frac{\theta E_f}{m_{\text{ion}}}$$

$$\text{Fraction burned} \rightarrow \theta \approx \frac{1}{1 + 7/\rho R}$$

$$\text{Hydro-efficiency} \rightarrow \eta_h = \frac{E_{\text{kinetic}}}{E_{\text{Laser}}}$$



$$\text{Gain} = \frac{73}{I_{15}^{0.25}} \left(\frac{3 \cdot 10^7}{V_i} \right)^{1.25} \left(\frac{\theta}{0.2} \right)$$

Scaling laws relating stagnation to in-flight hydro-variables are derived from conservation equations

(1) Mass $\rightarrow \quad \rho_s \Delta_s \sim \frac{M_{sh}}{R_h^2 \Sigma(A_s)} \sim \frac{E_K}{R_h^2 V_i^2 \Sigma(A_s)}$

(2) Energy $\rightarrow \quad E_K \sim P_s (R_h + \Delta_s)^3$

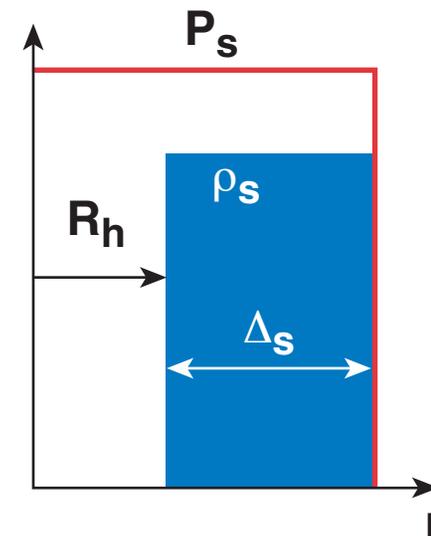
(3) Entropy $\rightarrow \quad \alpha_s \sim \alpha_{if} \text{Mach}_{if}^{2/3}$

$A_s \equiv R_h / \Delta_s \leftarrow$ Stagnation aspect ratio

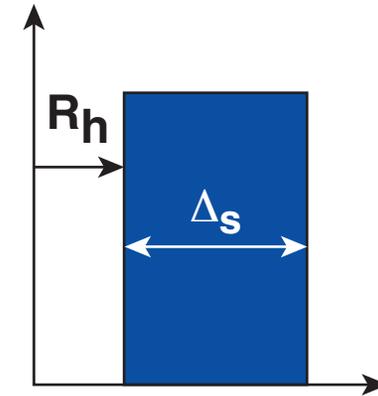
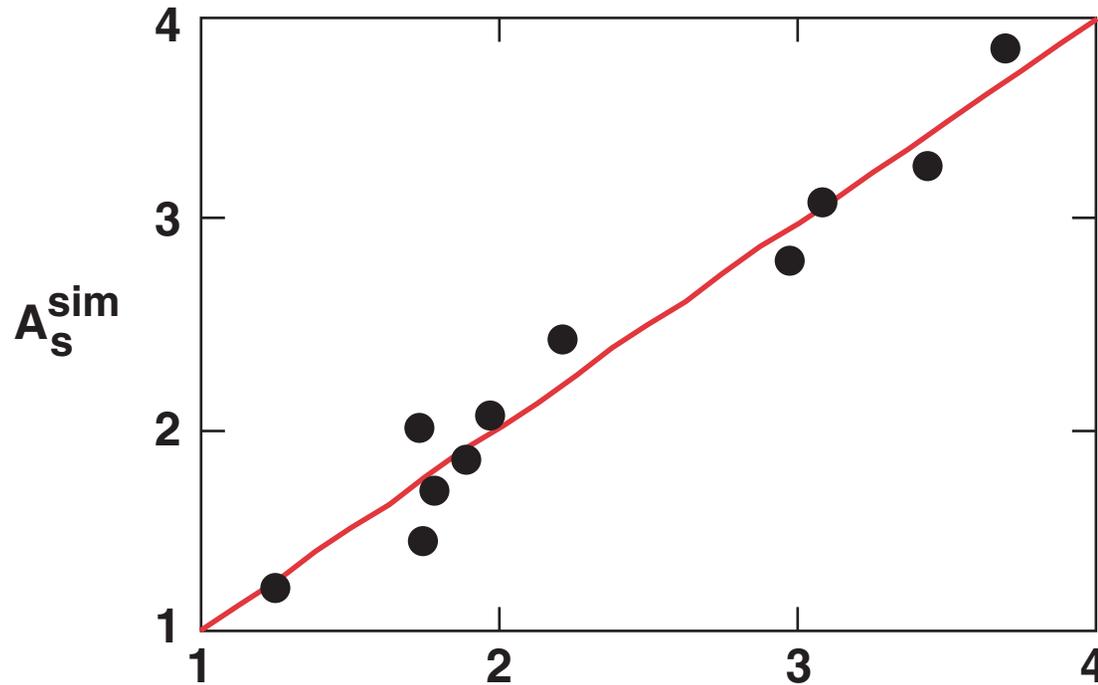
$\Sigma(x) \equiv 1 + 1/x + 1/(3x^2) \leftarrow$ Volume factor

$\alpha_{if} \equiv$ in-flight adiabat

Unknowns $\rightarrow P_s, \rho_s, A_s, \Delta_s$



Simulations of optimized implosions (max ρR and ρ) yield a scaling relation for the stagnation aspect ratio

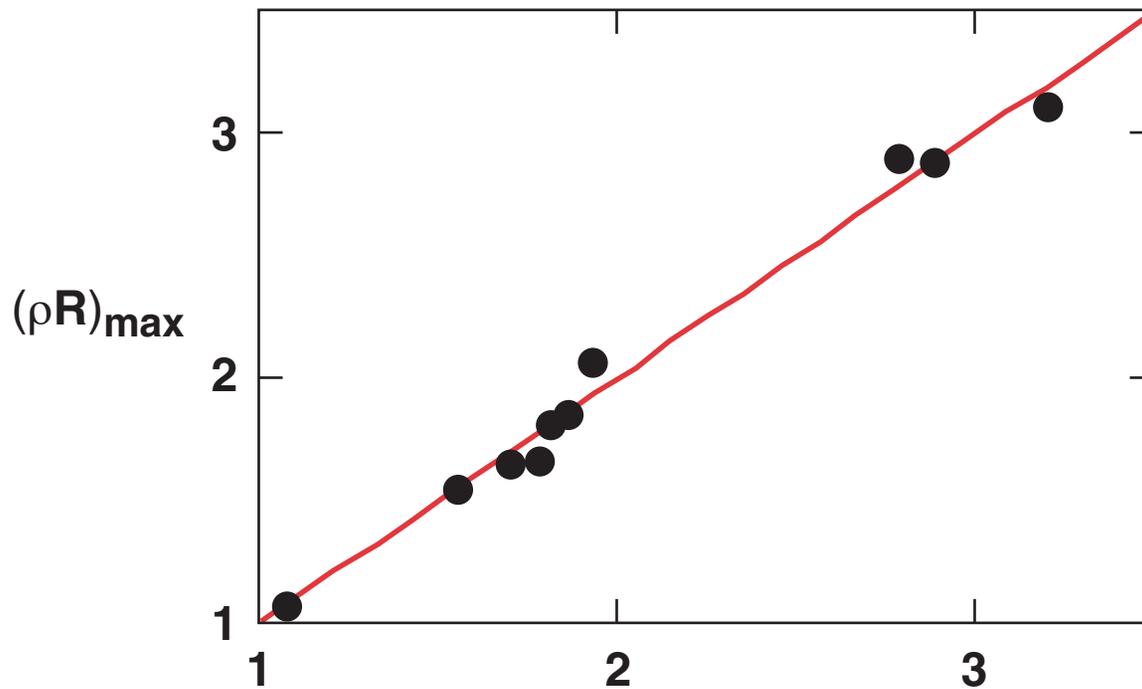


$$A_s \equiv \frac{R_h}{\Delta_s}$$

$$A_s^{fit} = 2.1 \left[\frac{V_i \text{ (cm/s)}}{3 \cdot 10^7} \right]^{0.96}$$

The areal density is weakly dependent on velocity;
it increases for lower adiabats and greater energies

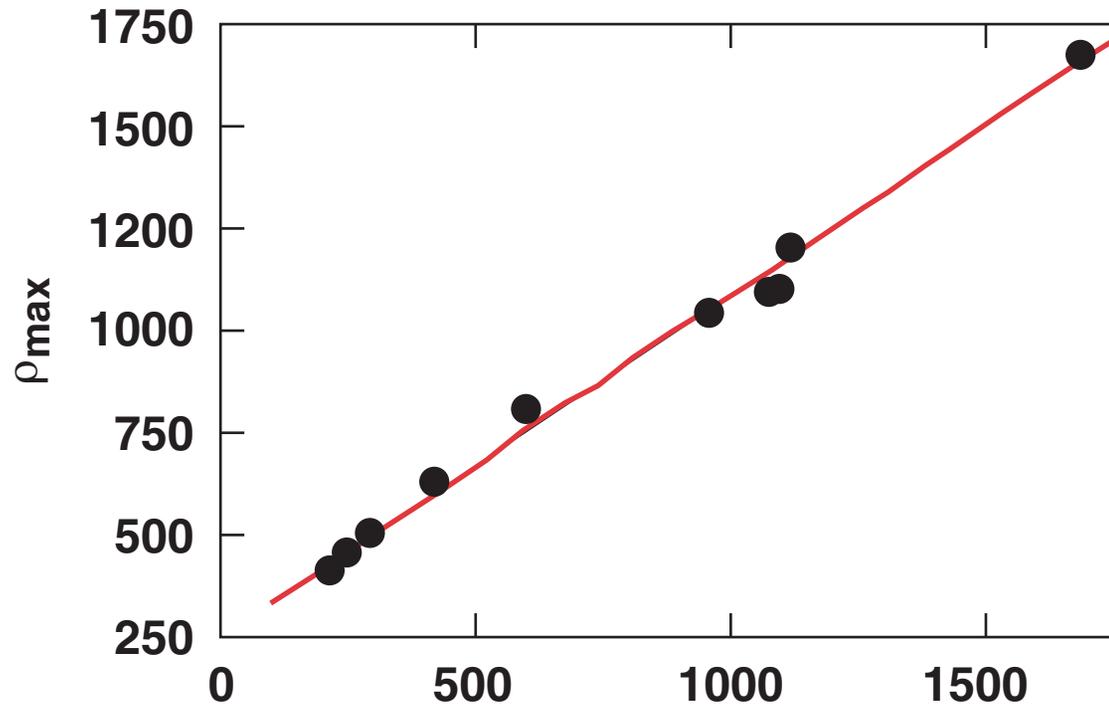
$$(\rho_s \Delta_s)^{\text{theory}} \sim E_L^{0.33} V_i^{0.03} \alpha_{if}^{-0.8}$$



$$(\rho R)_{\text{max}}^{\text{fit}} = \frac{1.3}{\alpha_{if}^{0.55}} \left[\frac{E_L \text{ (kJ)}}{100} \right]^{0.33} \left[\frac{V_i \text{ (cm/s)}}{3 \cdot 10^7} \right]^{0.06}$$

The density is independent of energy; it increases with the velocity and decreases with the adiabat

$$\rho_s^{\text{theory}} \sim V_i^{1.4} I_L^{0.13} \alpha_{if}^{-1.2}$$



$$\rho_{\text{max}}^{\text{fit}} = \frac{792}{\alpha_{if}} I_{15}^{0.13} \left[\frac{V_i \text{ (cm/s)}}{3 \cdot 10^7} \right]$$

High-gain fuel assemblies for fast ignition can be designed using the scaling formulas



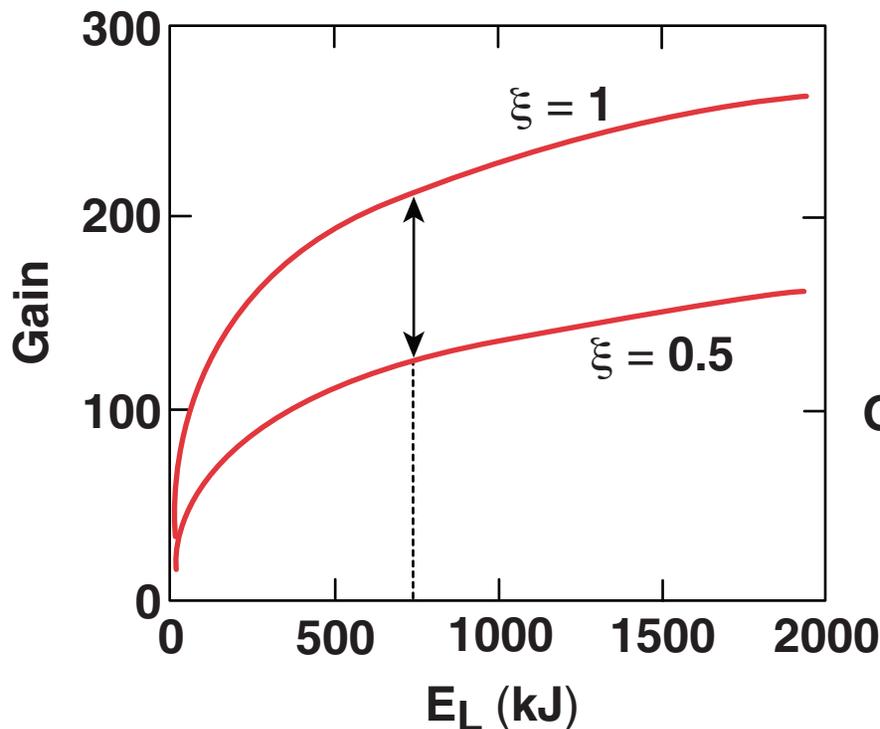
- Low adiabats enhance densities and areal densities: minimum practical adiabat $\alpha = 0.7$ to 0.8
- $\rho R(\alpha = 0.7) \approx 3 \Rightarrow E_{\text{Laser}} \approx 750 \text{ kJ}$
- $\rho_{\text{max}}(\alpha = 0.7) \approx 600 \text{ g/cc} \Rightarrow V_i(\text{cm/s}) \approx 1.7 \cdot 10^7 \text{ cm/s}$
- $V_i \approx 1.7 \cdot 10^7 \text{ cm/s} \Rightarrow R_h/\Delta_s \sim 1$

High-gain FI target: $E_L = 750 \text{ kJ}$, $\alpha = 0.7$, $V_i \approx 1.7 \cdot 10^7 \text{ cm/s}$

Estimated yield $\sim 120 \text{ MJ}$

In-flight aspect ratio (IFAR) = 18

For a fixed minimum adiabat and fixed peak density, the gain (without PW) depends only on the driver energy

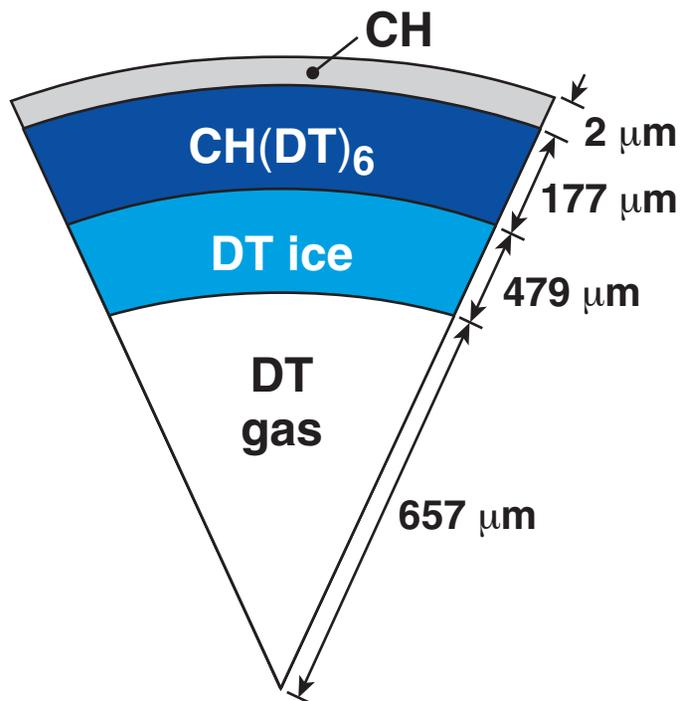


$$V_I \text{ (cm/s)} = \frac{1.7 \times 10^7}{I_{15}^{0.13}} \left(\frac{\rho_{\max}}{600} \right) \left(\frac{\alpha_{\text{if}}}{0.7} \right)$$

$$\text{Gain} \left[\begin{array}{l} \rho_{\max} = 600 \\ \alpha_{\text{if}} = 0.7 \\ I_{15} = 1 \end{array} \right] = \frac{742}{1 + 22/\xi E_L^{0.33} \text{ (kJ)}}$$

ξ = fraction of $(\rho R)_{\max}$ available for burn

A high-yield target has been designed for a 750-kJ laser driver

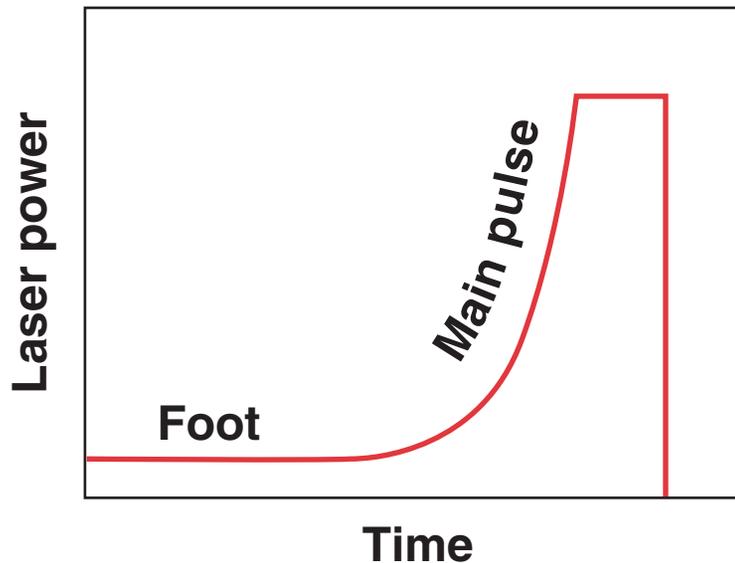


Energy	Implosion velocity	α	IFAR
750 kJ	$1.7 \cdot 10^7$ cm/s	0.7	18

Maximum averaged density	Peak Density	Maximum ρR
550 g/cc	670 g/cc	3 g/cm ²

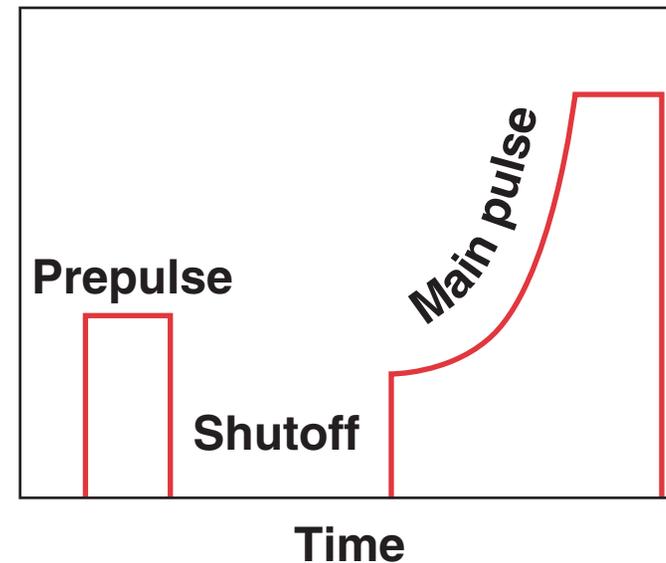
Using the relaxation-type laser pulse leads to improved hydrostability and a lower laser–power contrast ratio

Standard flat-adiabat pulse



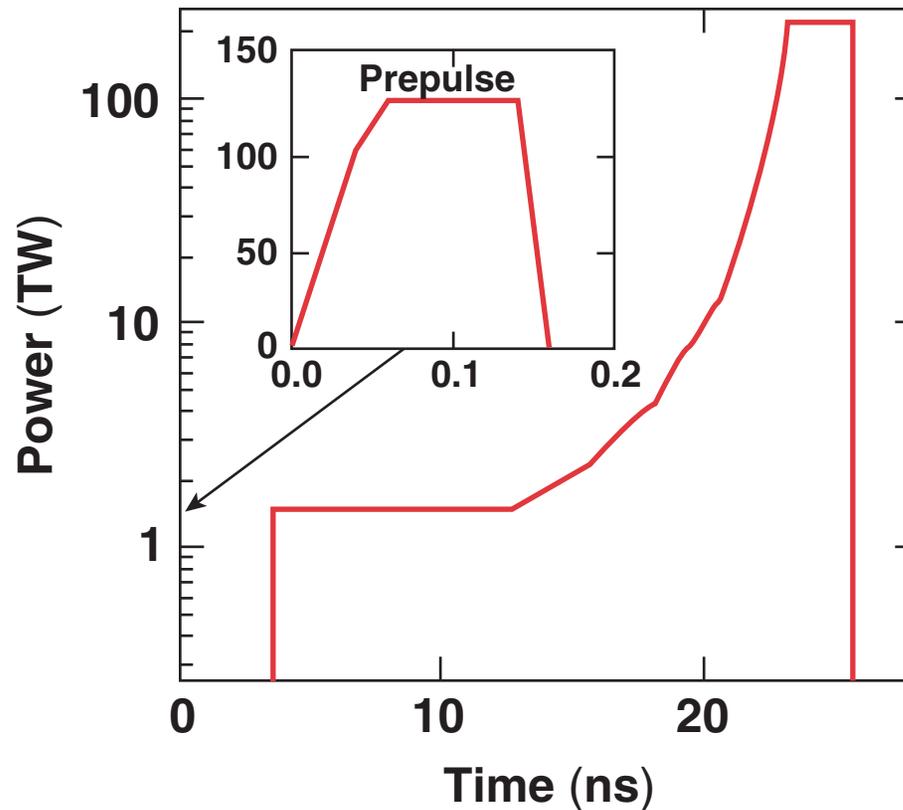
$$P_{\text{foot}}^{(\text{Mb})} \approx 2 \alpha_{\text{if}}$$

Relaxation shaped-adiabat pulse



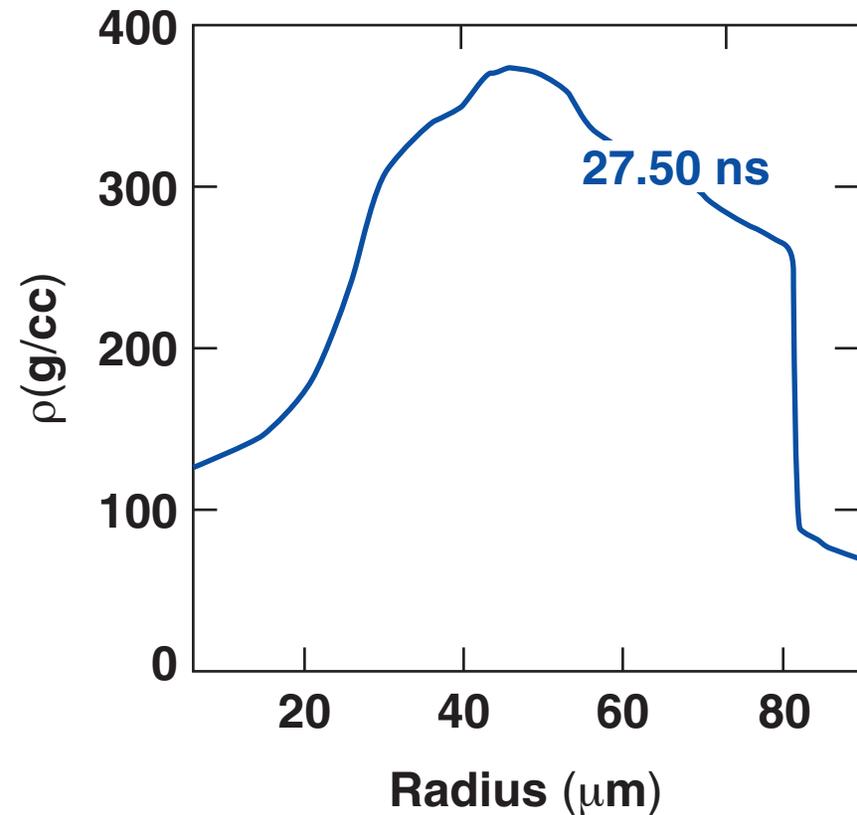
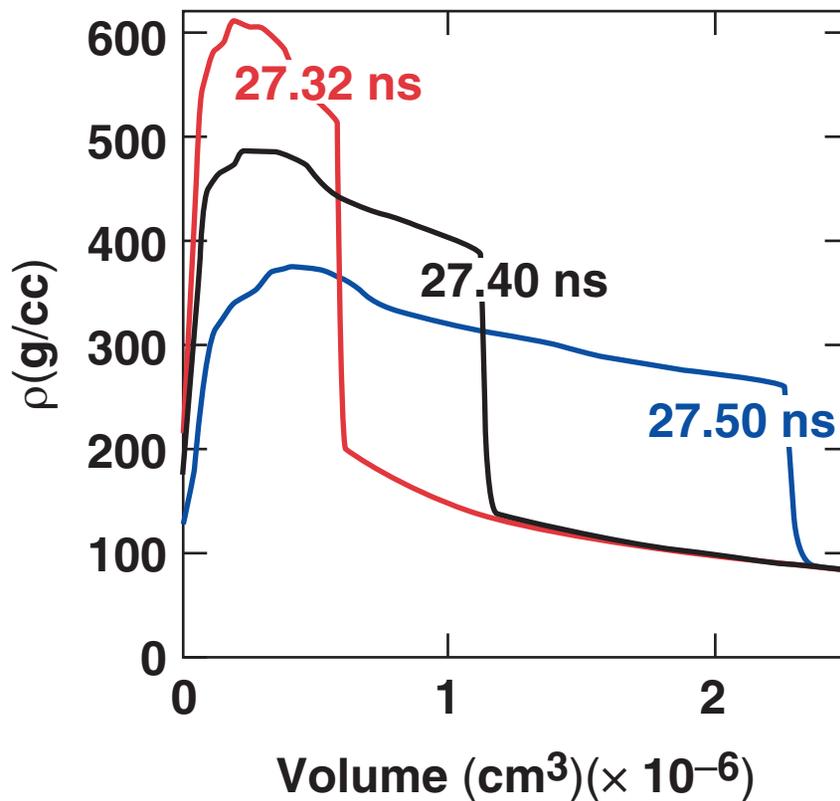
$$P_{\text{foot}}^{(\text{Mb})} \approx 6 \alpha_{\text{if}}$$

The 750-kJ capsule is driven by a relaxation laser pulse with a 22-ns main pulse and a contrast ratio of 150

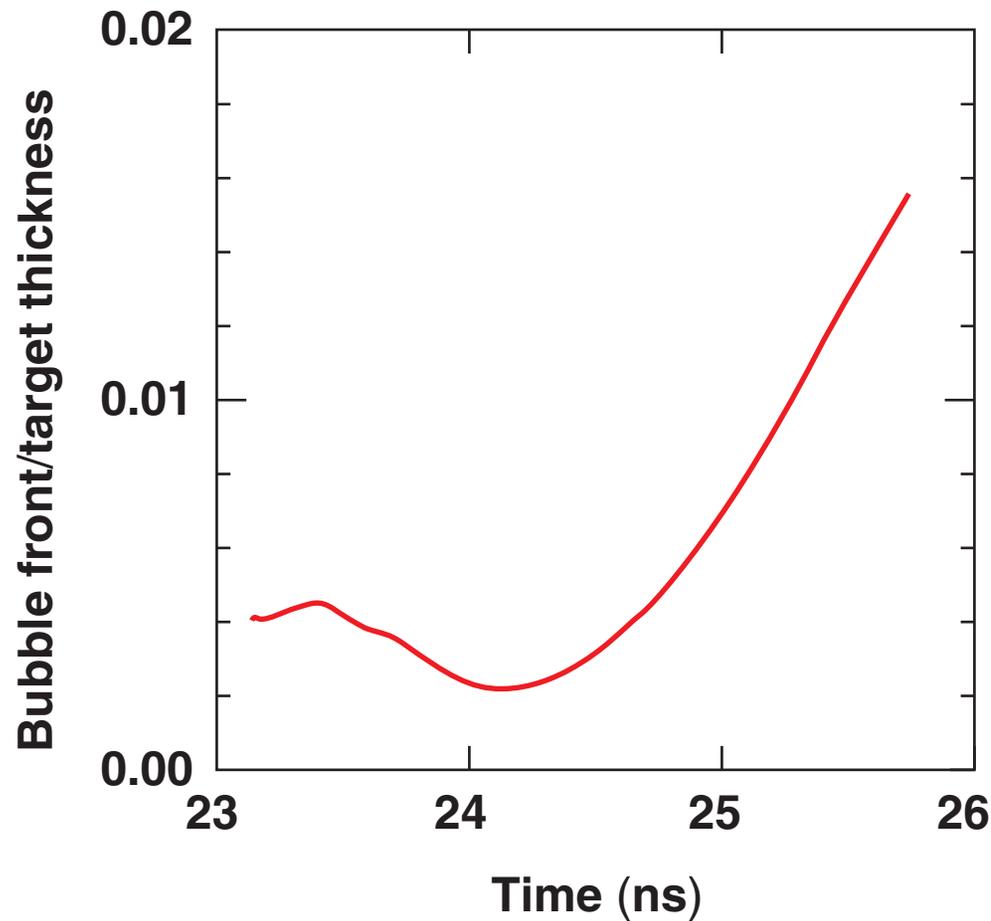


**Can NIF assemble high-gain FI targets?
Indirect-drive pulse is 18 ns with a contrast ratio of ~100**

The 750-kJ capsule yields a hot-spot volume < 8% of the compressed volume and a quasi-isochoric density profile



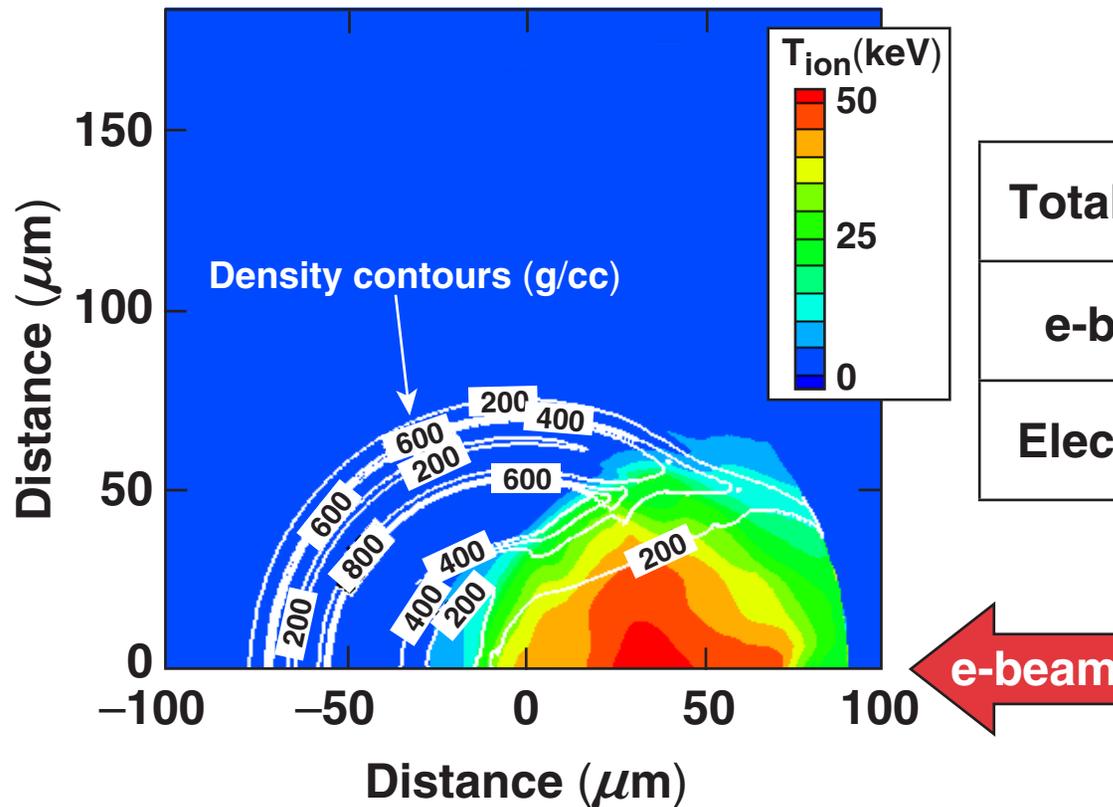
The slow implosion velocity leads to negligible Rayleigh–Taylor growth during the laser flat top



Results from RT postprocessor based on Haan–Goncharov models and NIF laser nonuniformities with 1-THz SSD.

V. N. Goncharov *et al.*, *Phys. Plasmas* **7**, 2062 (2000).
S. W. Haam, *Phys. Rev. A* **39**, 5812 (1989).

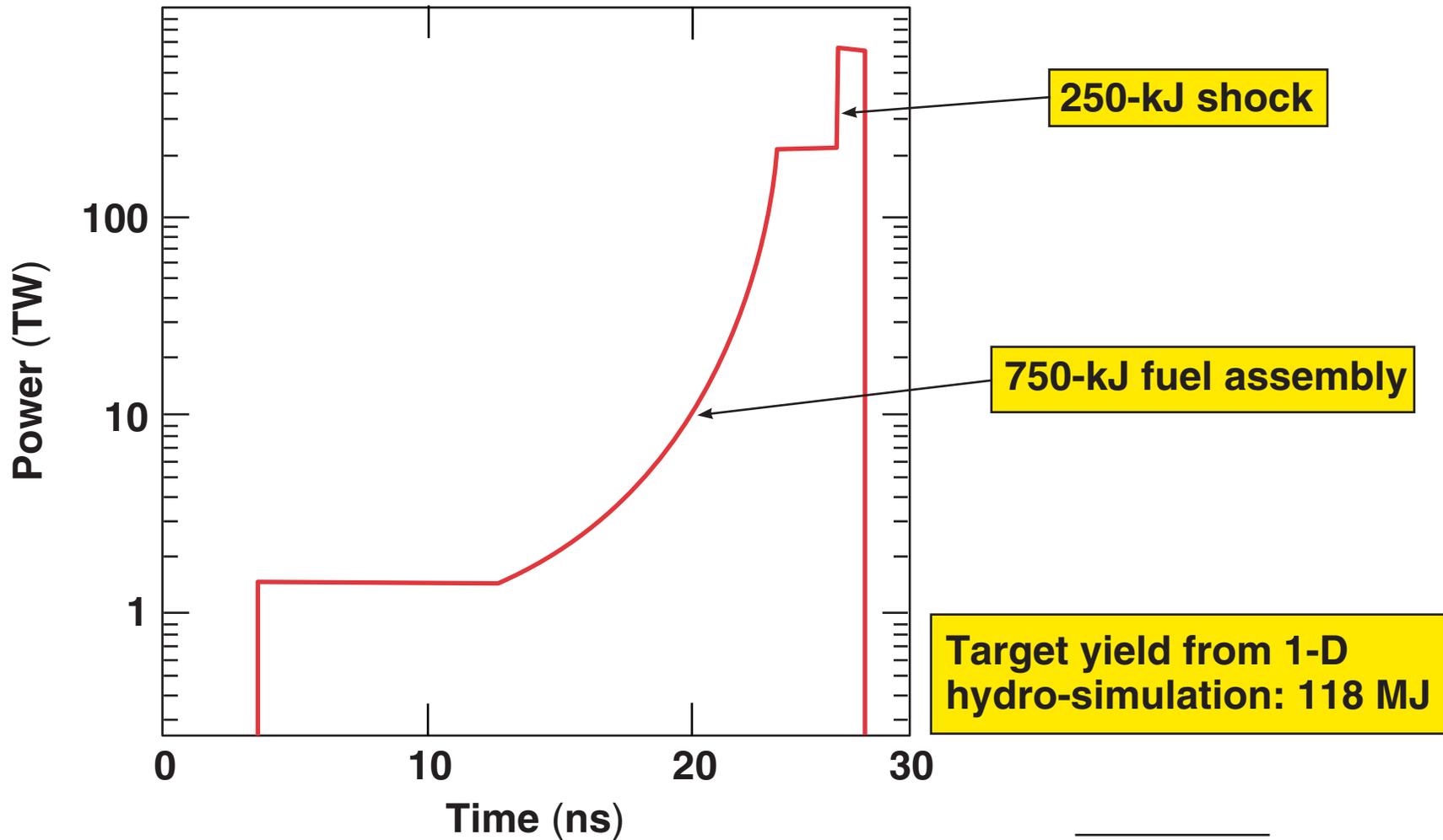
2-D hydro-simulations of ignition and burn of the 750-kJ target show energy yields >100 MJ



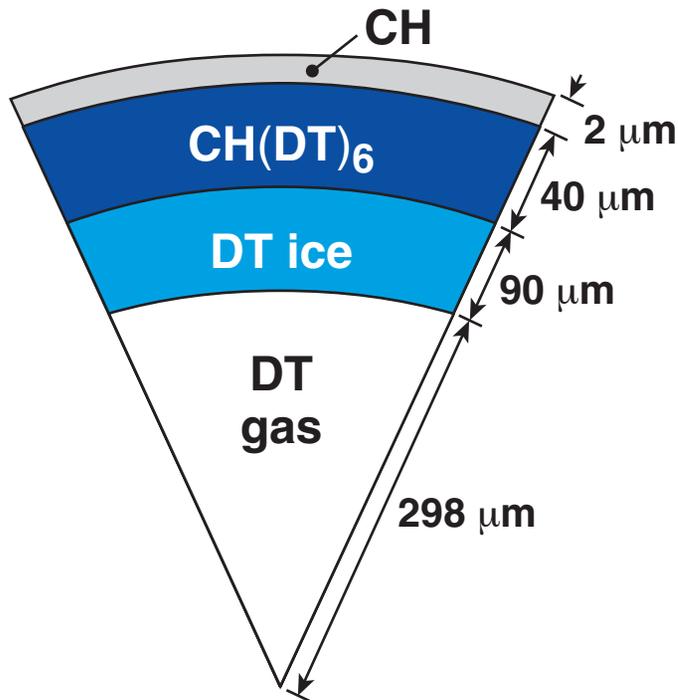
Total beam energy (kJ)	12–20
e-beam radius (μm)	20
Electron energy (MeV)	2–3

Energy yield \approx 116 MJ

The relatively cold hot spot of such fuel assemblies can also be ignited by a spherically convergent shock*



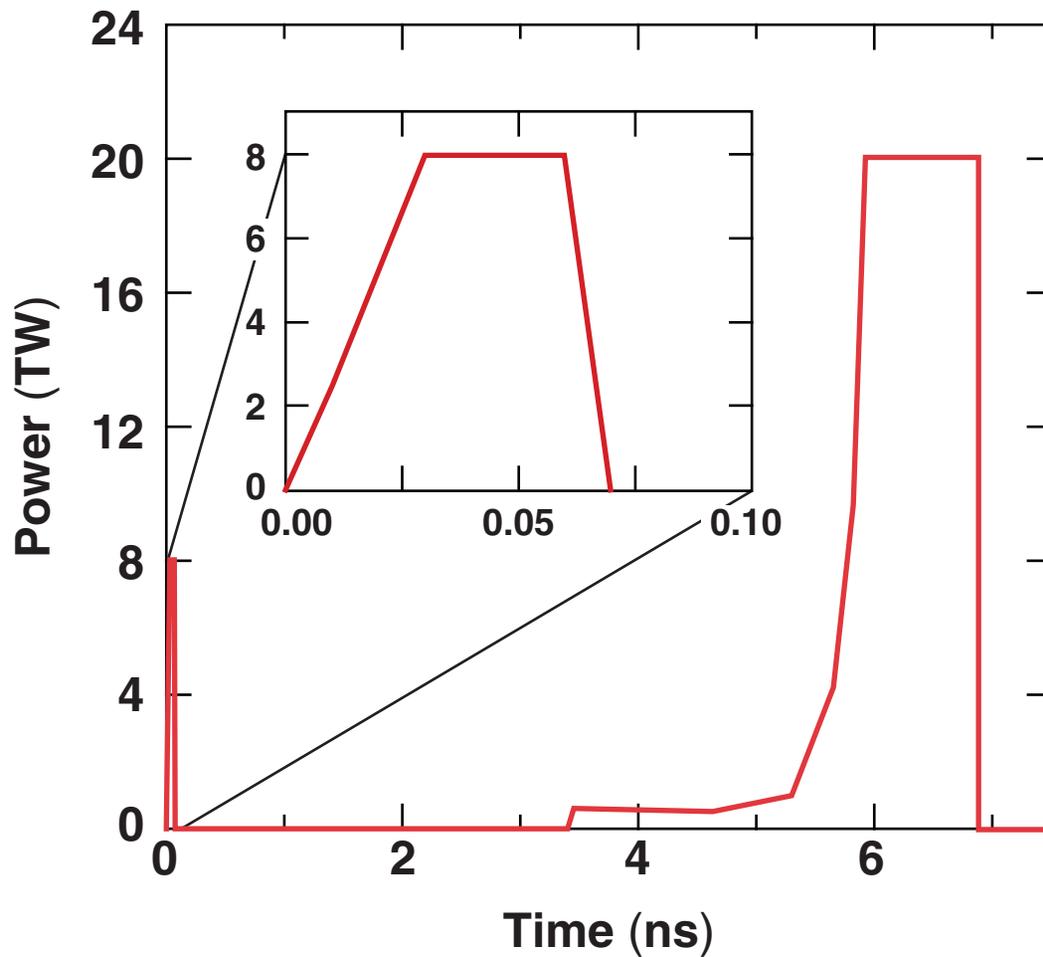
A 25-kJ driver can assemble fuel for fast ignition using low-adiabat implosions of thick shells with a pulse compatible with the OMEGA Laser System



Energy	IFAR	α
25 kJ	30	1

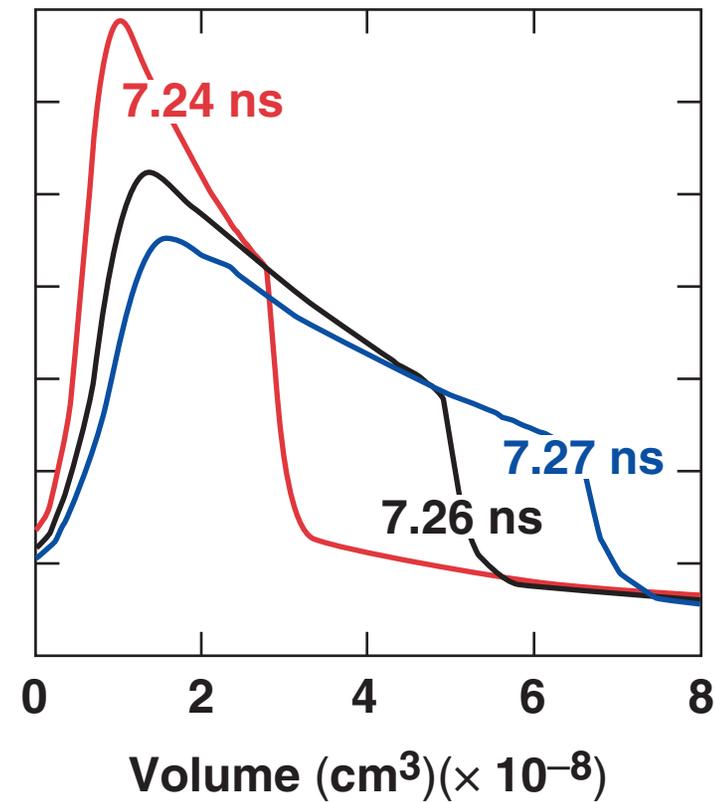
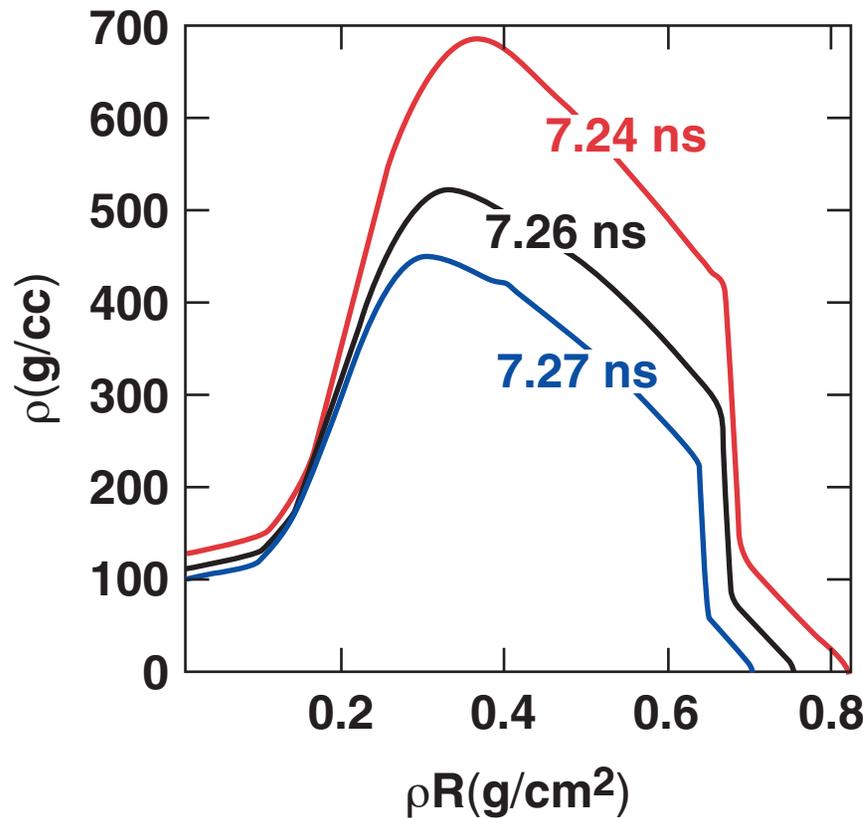
Implosion velocity	Maximum density	Maximum ρR
$2.6 \cdot 10^7$ cm/s	700 g/cc	0.8 g/cm ²

The 130- μm capsule is driven by a relaxation laser pulse within the capabilities of the OMEGA laser

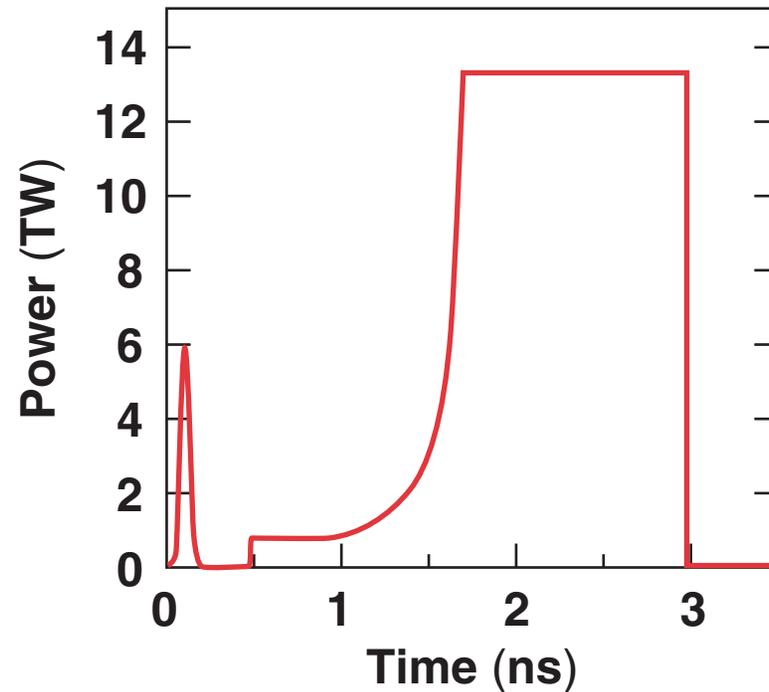
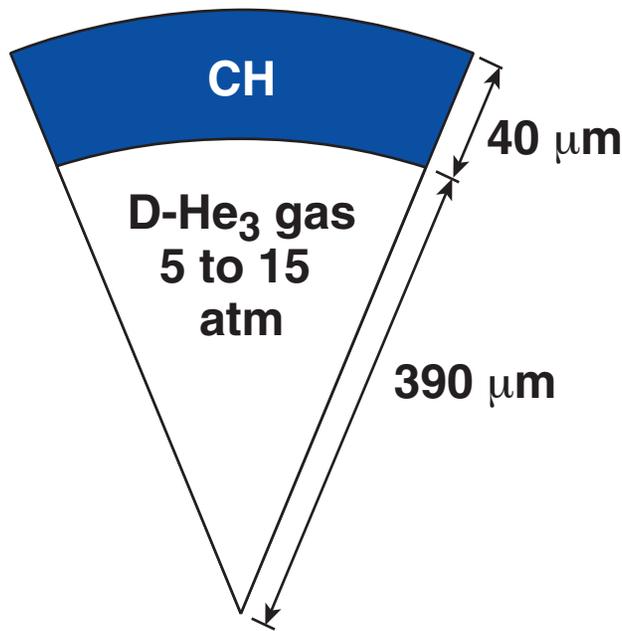


Main pulse length = 3.5 ns
Prepulse width = 70 ps
Prepulse/main = 0.4
Main contrast ratio = 32
Pulse energy = 25 kJ
Peak power = 20 TW

The 130- μm , $\alpha = 1$, OMEGA-compatible capsule yields a density > 300 g/cc over a $\rho R > 0.4$ g/cm² and a hot-spot volume $< 20\%$ of the total volume



This method for assembling FI fuel will be first tested on OMEGA with surrogate plastic-shell implosions



Energy	α	Implosion velocity	Maximum ρR (5–15 atm)	Maximum ρ (5–15 atm)	Proton yield (5–15 atm)
20 kJ	1.2	$2.1 \cdot 10^7$ cm/s	0.5–0.36 g/cm ²	276–190 g/cc	$1.2\text{--}2.3 \cdot 10^8$

Summary/Conclusions

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