Intense Electron-Beam Transport in Dense, Cryogenic, DT, Fast-Ignition Fusion Targets



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Summary

The PIC/hybrid approach is a promising technique for modeling electron transport in very overdense plasmas

- Two- and three-dimensional simulations of fast (1-MeV) electron beams have been made in imploded cryogenic fusion targets.
- The PIC/hybrid approach (LSP^{*}) has several advantages over traditional PIC that allow large volumes of plasma to be simulated.
- A self-generated azimuthal magnetic field collimates the electron beam in 2-D, but not in 3-D.
- We observe filamentation of the beam current in both 2-D and 3-D.
 - Filamentation is shown to depend on the background plasma density and beam temperature.
- The target is primarily heated by the return current for the chosen conditions.

We treat the fast-electron transport but not the generation mechanism



Large contrast in plasma density

Electron-beam parameters are relevant to future fast-ignition studies on OMEGA EP

- An electron beam is generated by promotion from background over a 20- μ m spot with a pulse duration of 10 ps.
- FI-relevant parameters are chosen for the beam source

$$n_b = 2 \times 10^{20} \eta_{eff} \frac{I}{10^{19} \text{W cm}^{-2}} \frac{1 \text{ MeV}}{\epsilon_b} \text{ cm}^{-3}$$

- Unlike simulations in near-critical plasmas, the beam is "weak" in the sense that $n_b/n_e << 1$

$$\begin{split} \mathbf{I}_{b} &= \mathbf{30} \ \eta_{eff} \ \frac{\mathbf{I}}{\mathbf{10^{19} W cm^{-2}}} \ \frac{\mathbf{A}_{spot}}{\mathbf{300 \ \mu m^{2}}} \left(\frac{\mathbf{1 \ MeV}}{\epsilon_{b}} \right) \mathbf{MA} \\ \mathbf{I}_{b} &>> \mathbf{I}_{Alfvén} = \mathbf{17} \gamma \beta \ \mathbf{kA} \end{split}$$

• Self-generated fields are therefore important for transport.

In 2-D, r-z geometry, the electron beam breaks into filaments and contracts radially



In three dimensions the electron beam breaks up into filaments

200 _ n_b (cm^{_3}) 10²⁰ $n_{e} (cm^{-3})$ 20 10²⁵ 1022 **Υ** (μm) 0 150 -20 (**m**ḿ) **Z** 10¹⁸ 100 1020 10²³ n_e (cm⁻³) 20 **Υ (μm)** 50 0 Background plasma -20 $n_b (cm^{-3})$ **10**¹⁸ 1022 0 -40 40 0 20 0 -20 **Χ** (μ**m**) **Χ** (μ**m**)

Homogeneous simulations show that beam filamentation depends on the plasma density



The filaments are smoothed-out in the high density shell

200 (\mathbf{A}) n_b (cm⁻³) **10**²⁰ $n_e (cm^{-3})$ 20 10²⁵ 1022 **Υ** (μm) 0 150 Α -20 **Z** (μm) 1018 100 1020 **(B**) В _ n_b (cm^{_3} 10²⁰ 20 **Υ (μm) 50** 0 Background plasma $n_b (cm^{-3})$ -20 1018 **10**¹⁸ 0 -40 0 40 -20 0 20 \boldsymbol{X} ($\mu \boldsymbol{m}$) \boldsymbol{X} ($\mu \boldsymbol{m}$)

Filamentation is suppressed by a large beam "temperature"

High temperature Low temperature $T_b = 10 \text{ keV}$ T_b = 200 keV 200 150 1020 1020 (**m**ḿ) **Z** 100 **10**¹⁸ 1018 50 $n_b (cm^{-3})$ 0 40 -40 40 -40 0 0 \boldsymbol{X} ($\mu \boldsymbol{m}$) **Χ** (μ**m**)

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