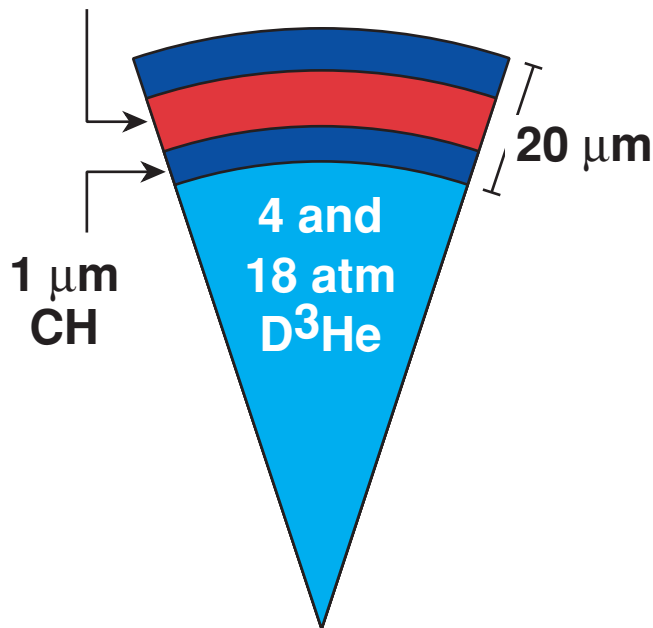
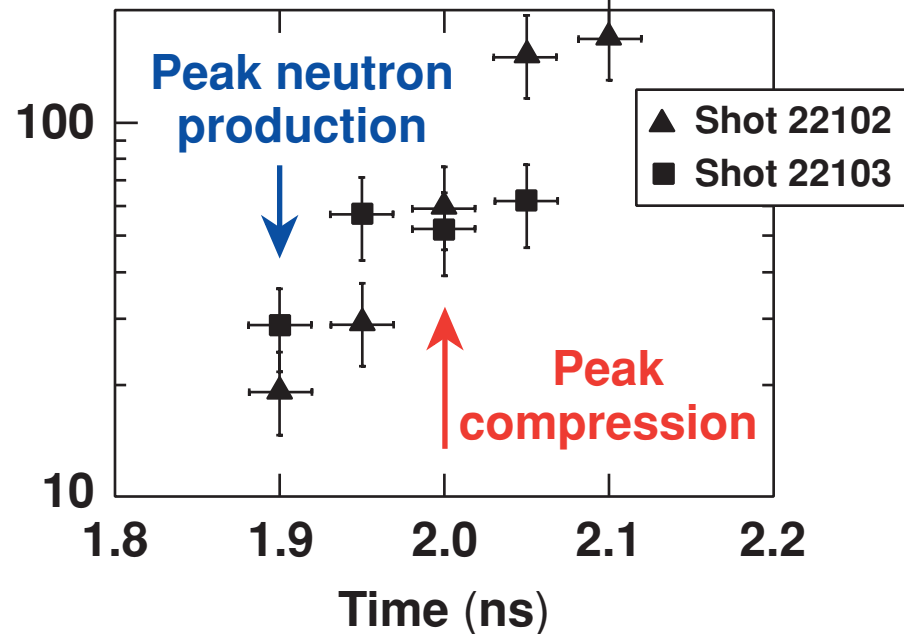


# Hydrodynamic Growth of Shell Modulations in the Deceleration Phase of Spherical Direct-Drive Implosions

2- $\mu\text{m}$ ,  
Ti-doped layer



Modulation  $\sigma_{\text{rms}}$  of  $\frac{\delta(\rho r)}{\rho r}$  (%)



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44th Annual Meeting of the  
American Physical Society  
Division of Plasma Physics  
Orlando, FL  
11–15 November 2002

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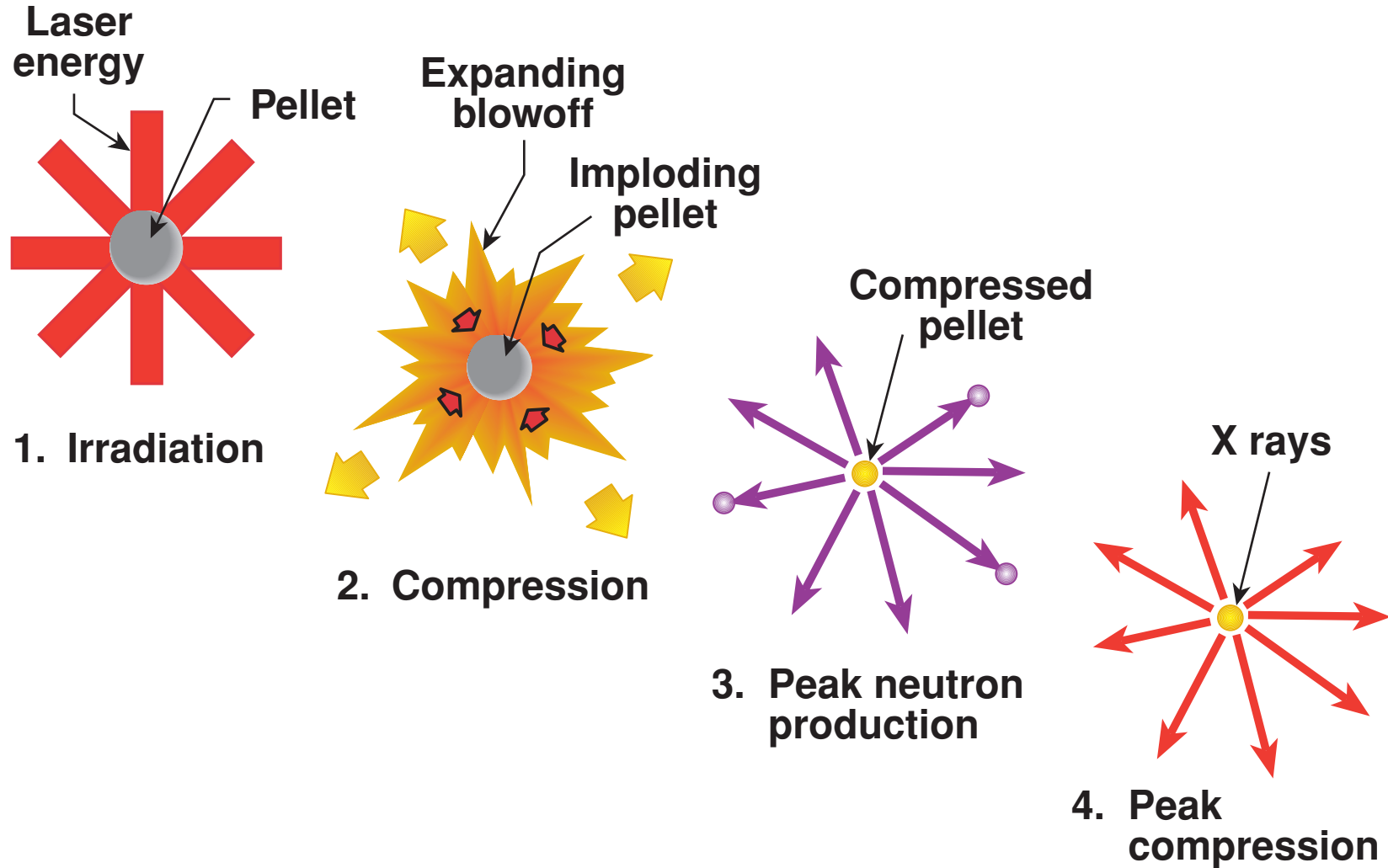
**J. Koch**

## Shell modulation growth has been measured near peak compression of spherical implosions

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- **20- $\mu\text{m}$ -thick CH shells with CH Ti layers, filled with 4 and 18 atm  $\text{D}^3\text{He}$ , were imploded with  $\sim 23\text{-kJ}$ , 1-ns square laser pulses.**
- **Measured perturbations have highest amplitudes at wavelengths of about 40 to 50  $\mu\text{m}$  (corresponding to mode number  $l \sim 6$ ).**
- **At peak neutron production, inner-shell areal-density modulation level,  $\delta(\rho r)/\rho r$ , is 20% and grows to  $\sim 50\%$  at peak compression 100 ps later due to Rayleigh–Taylor instability.**
- **For the same period, shell modulations grow up to about 1.5 times due to Bell–Plesset effects.**
- **At peak compression the inner part of the shell has a higher modulation level than the bulk of the shell.**

# Laser-induced ablation is used to generate ultra-high pressures on a fusion capsule to compress it



## Related talks

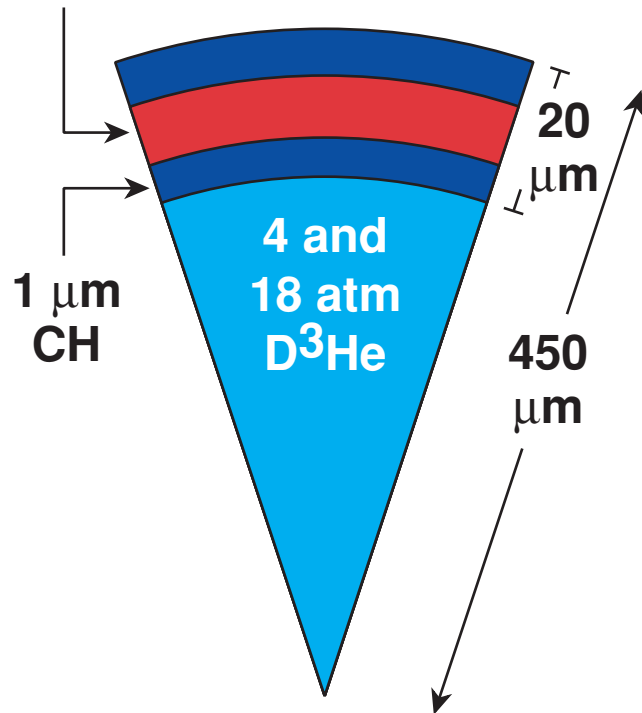
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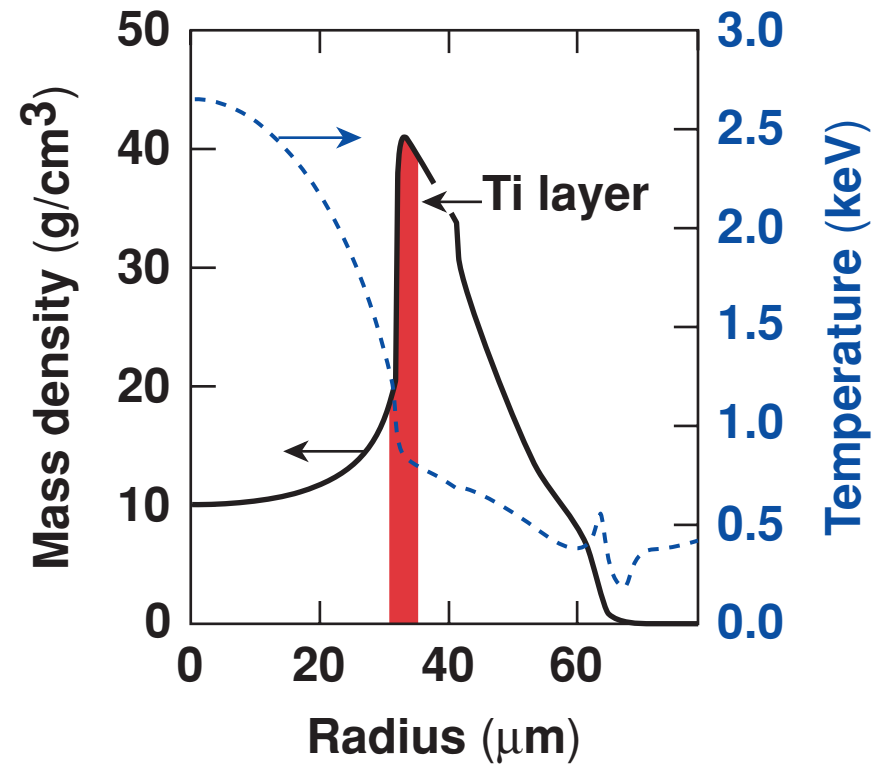
- BO2.001** R. Epstein *et al.*, Modeling of Fuel–Pusher Mix Effects in 1-D Simulations of Cryogenic, All-DT Ignition Capsule Implosions
- BO2.002** S. P. Regan *et al.*, Experimental Investigation of Fuel–Pusher Mix in Direct-Drive Implosions on OMEGA
- BO3.003** J. A. Frenje *et al.*, Effects of Fuel–Pusher Mix on Direct-Drive Implosions of  $^3\text{He}$ -Gas-Filled, CD-Layered Plastic Capsules on OMEGA
- FO2.001** P. B. Rahda *et al.*, The Effect of Laser Nonuniformities on Plastic-Shell Direct-Drive Implosions on OMEGA
- RI1.005** C. K. Li *et al.*, Capsule Areal-Density Asymmetries and Time Evolution Inferred from 14.7-MeV Proton Line Structure in OMEGA  $\text{D}^3\text{He}$  Implosions

# Targets with titanium-doped layers were used to measure shell modulations

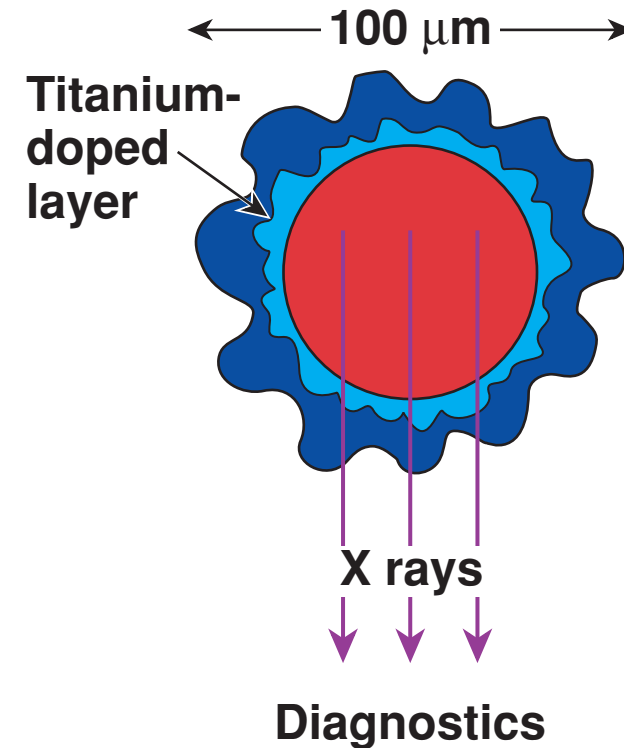
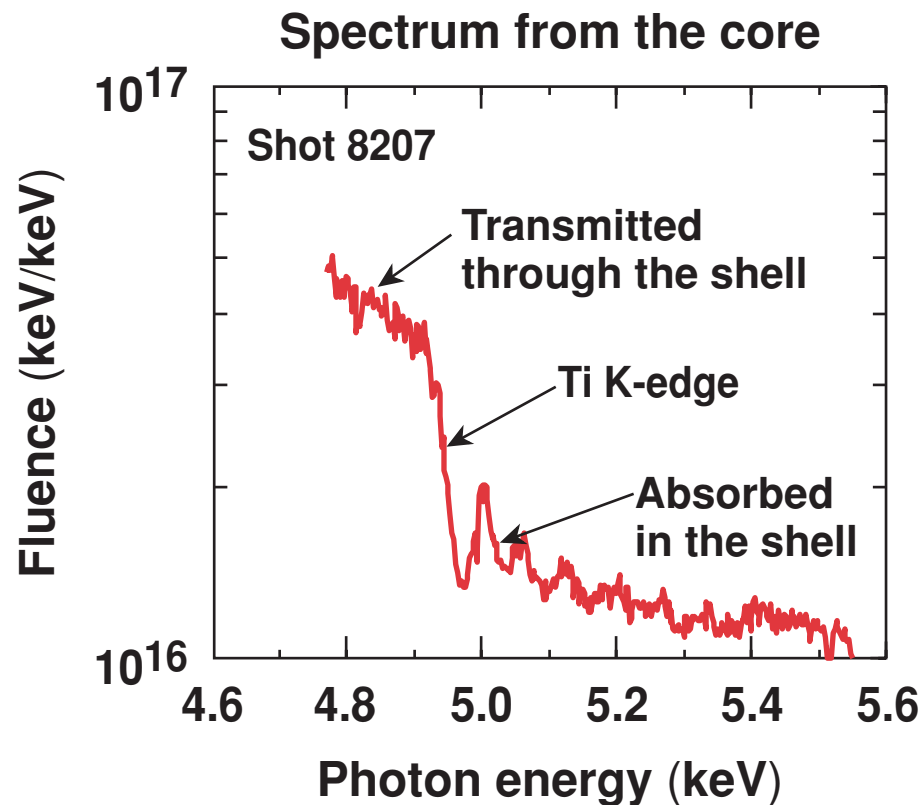
2- $\mu\text{m}$ ,  
Ti-doped layer (6% by atom)



1-D simulations  
at peak compression

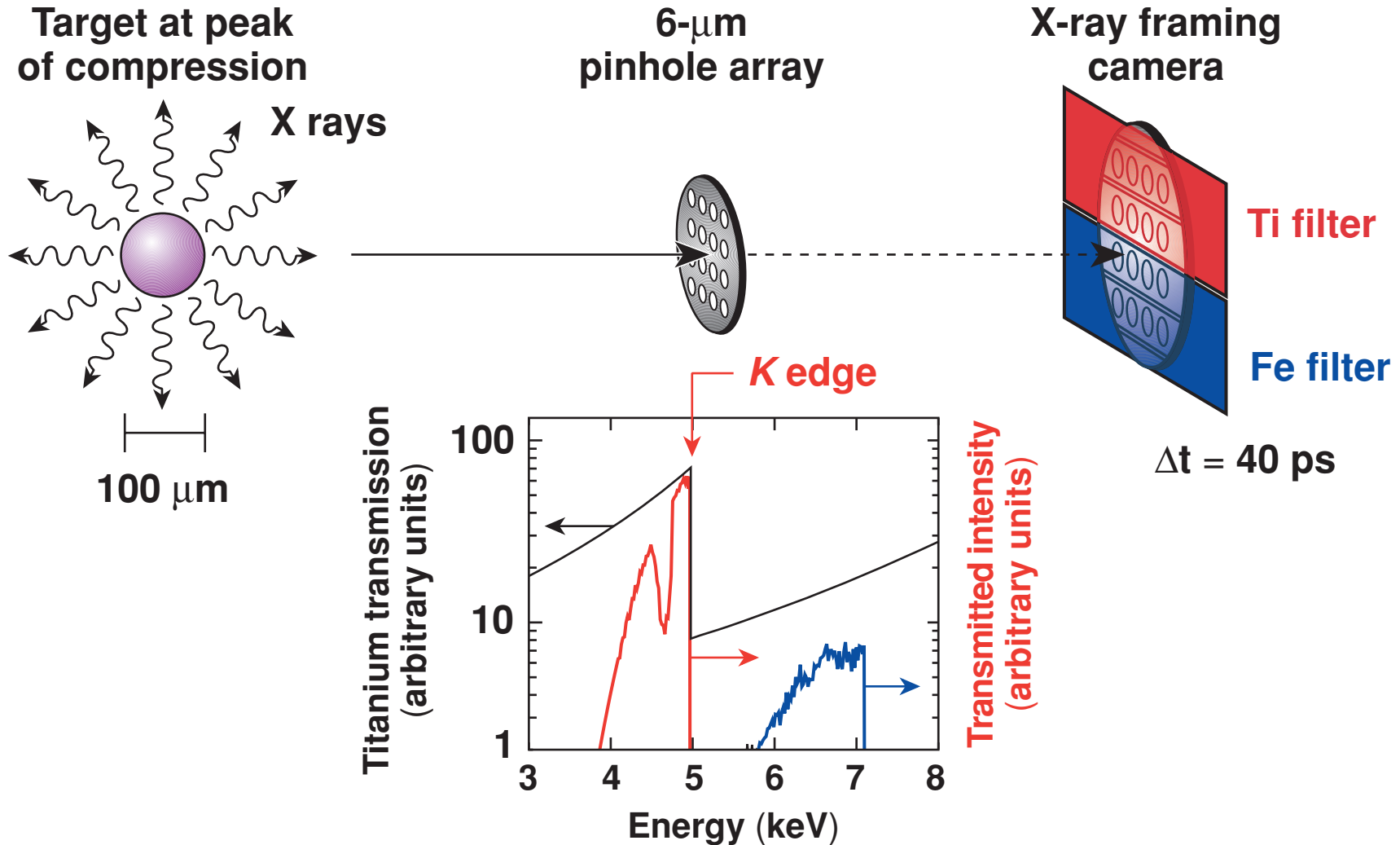


Core images at two x-ray energies, highly absorbed and nonabsorbed by the shell, are used to measure the integrity of the inner portion of the shell



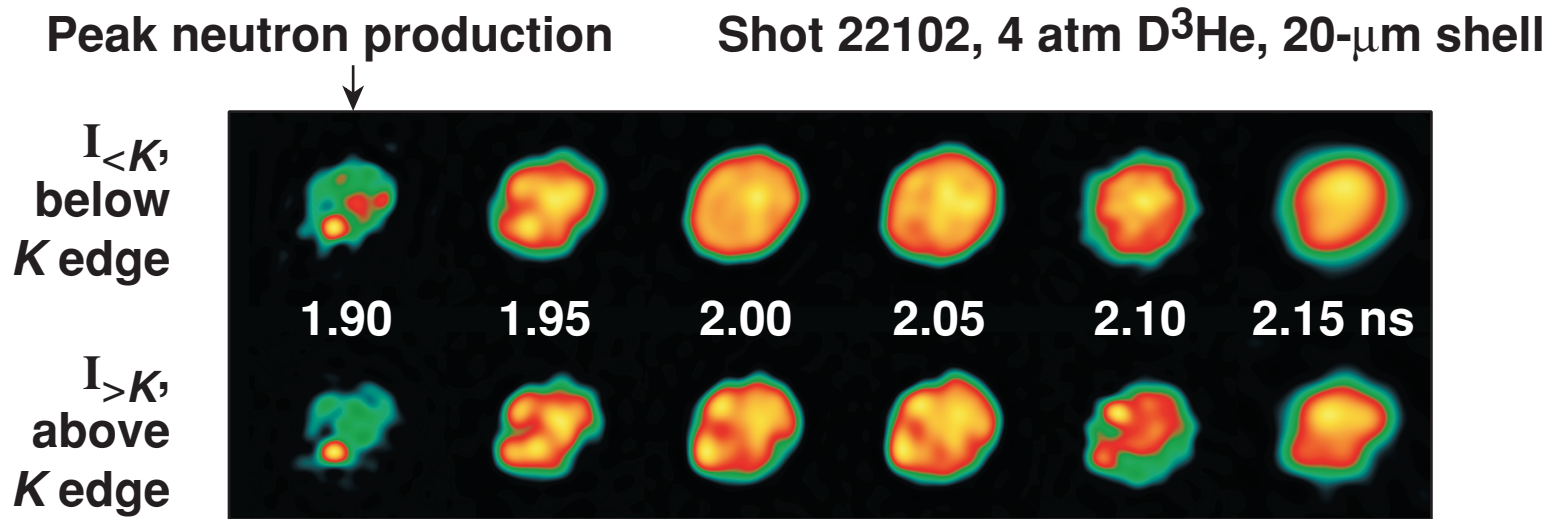
# Shell-Integrity Measurements

## X-ray framing cameras are the primary diagnostics of shell nonuniformity\*





# The ratio of images above and below the $K$ edge is related to areal-density modulations in the shell\*

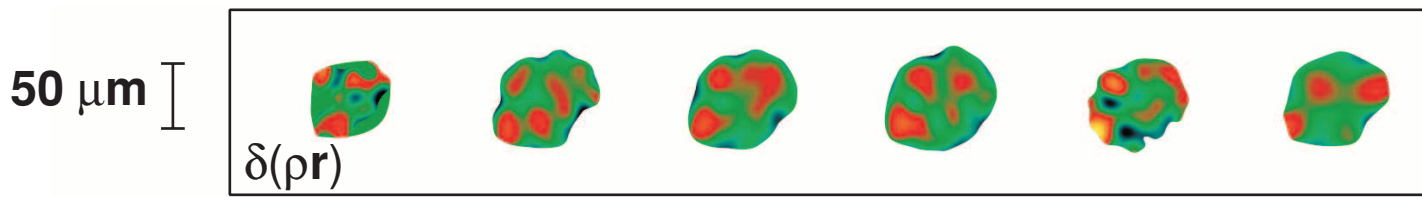


50  $\mu m$

Peak compression

$$\delta(\rho r) = \frac{\delta \left\{ \ln \left( \frac{I_{<K}}{I_{>K}} \right) \right\}}{\mu_{>K} - \mu_{<K}}$$

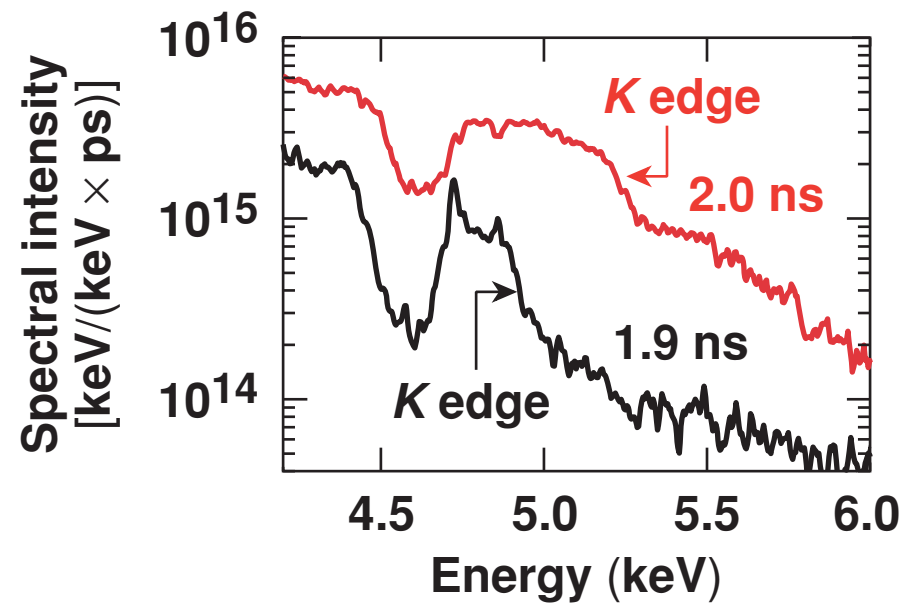
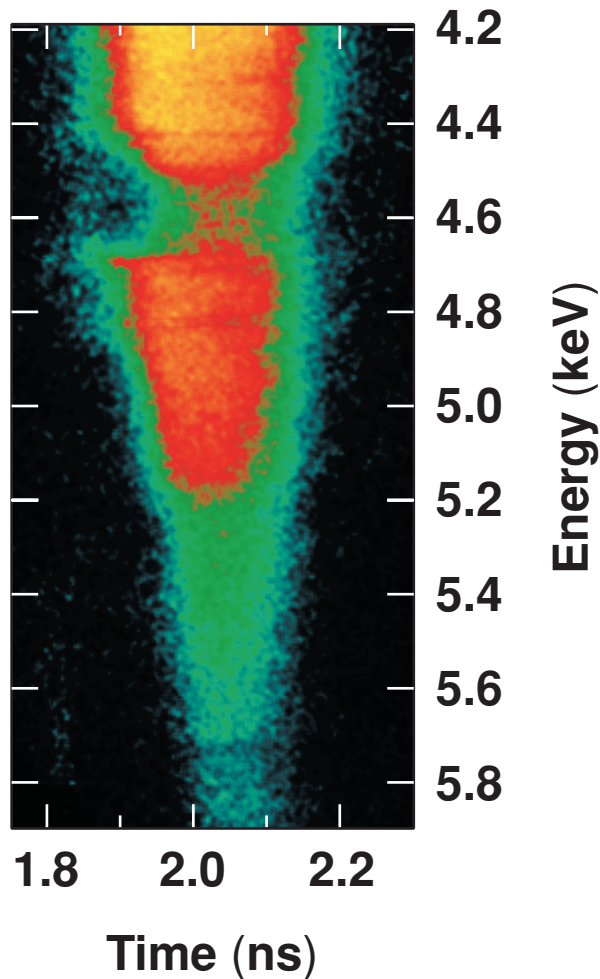
$\mu_{>K}, \mu_{<K}$ : titanium absorption rates above and below the  $K$  edge



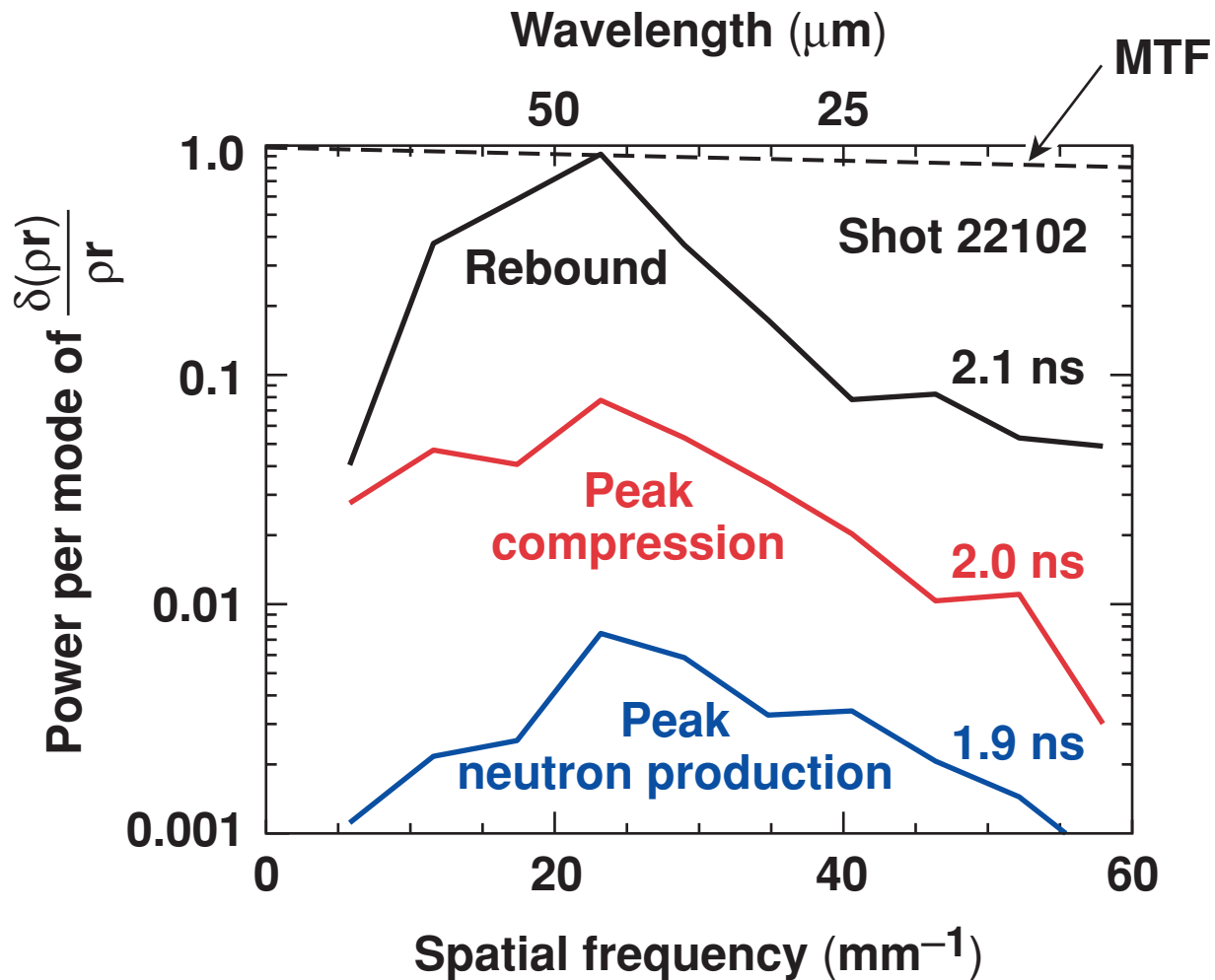
\* V. A. Smalyuk *et al.*, Phys. Rev. Lett. 87, 155002 (2001).

# Evolution of average titanium areal density is measured using streaked spectrometers

Streak 22102

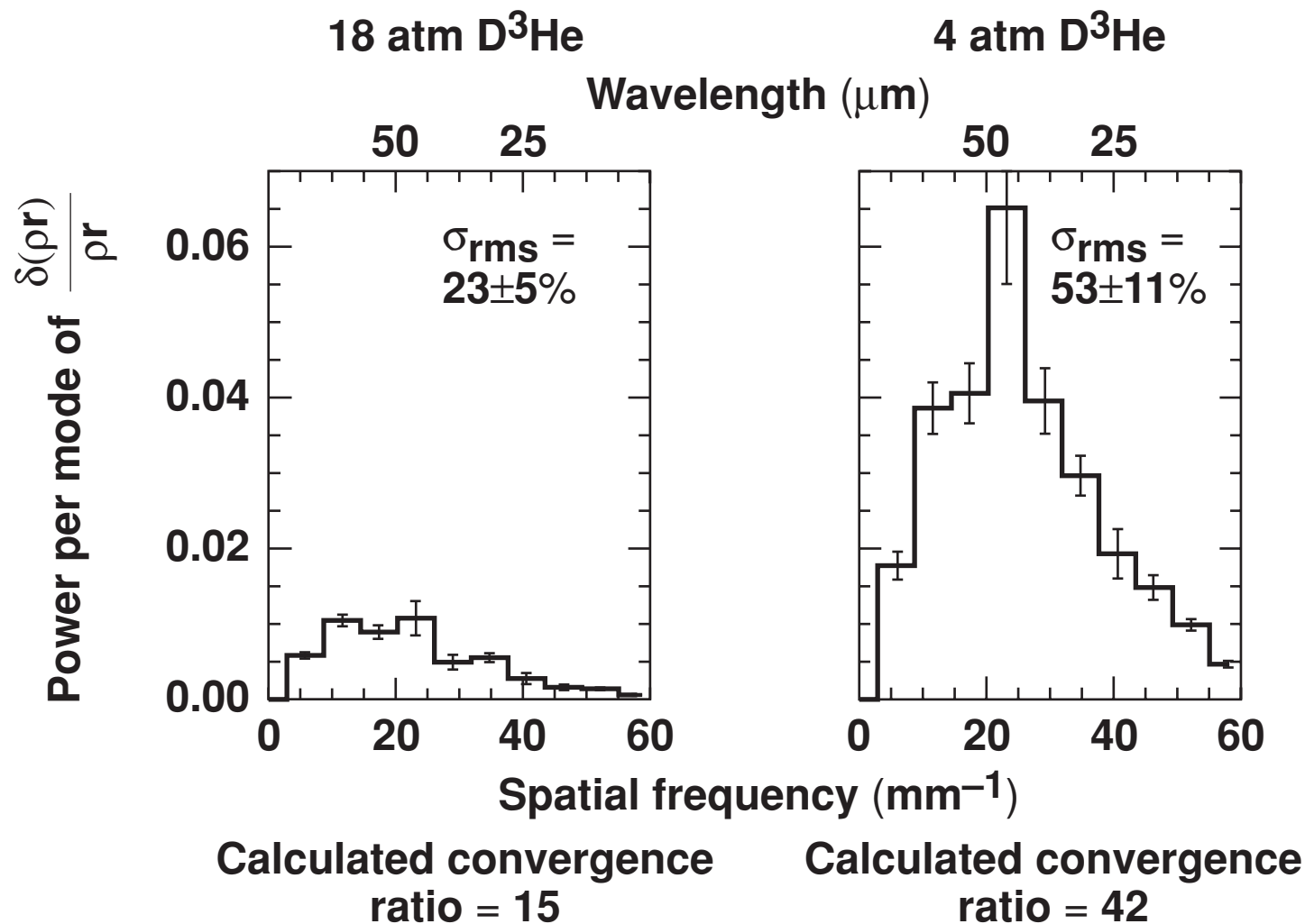


# Inner-shell modulations grow throughout the implosion's deceleration phase

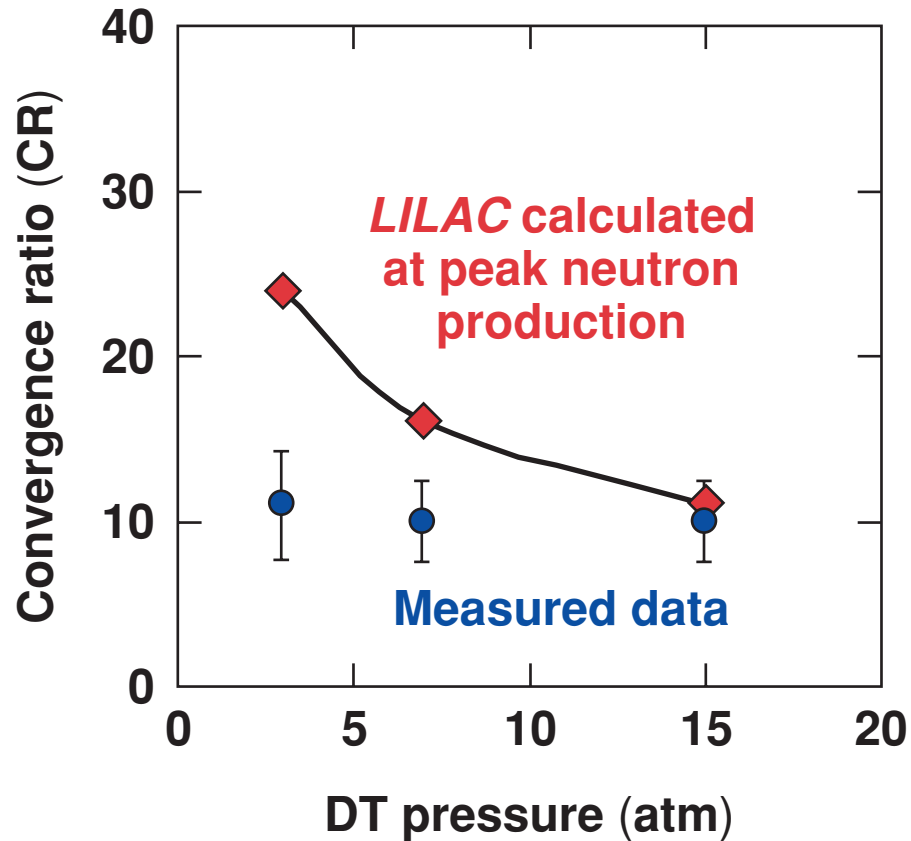


- The spectrum peaks at  $\ell$ -mode  $\sim 6$

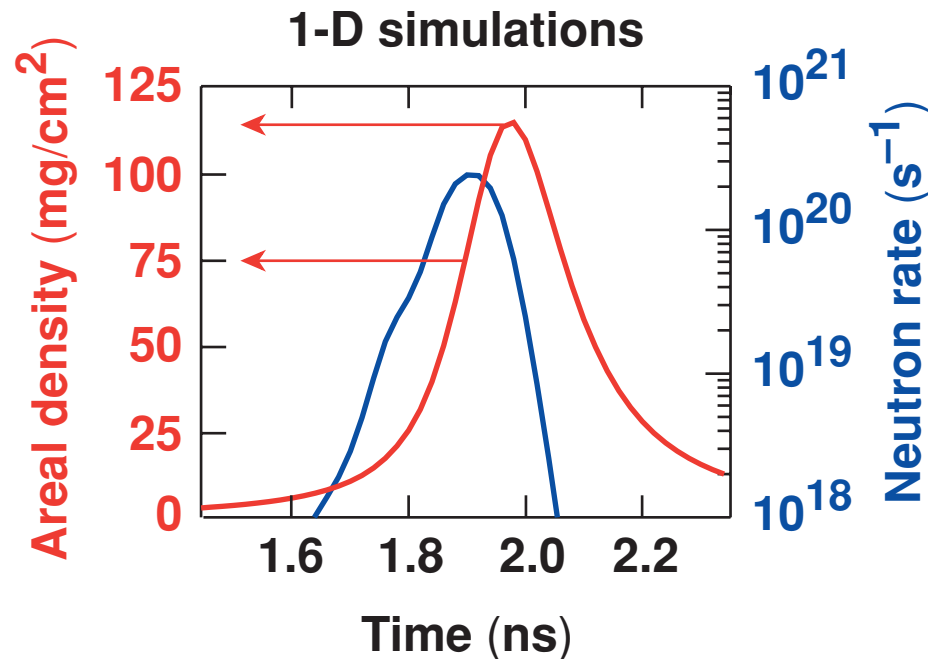
# At peak compression more-unstable implosions have a higher level of inner-shell of modulations\*



# Measured target convergencies are similar for implosions with 3 and 15 atm of DT

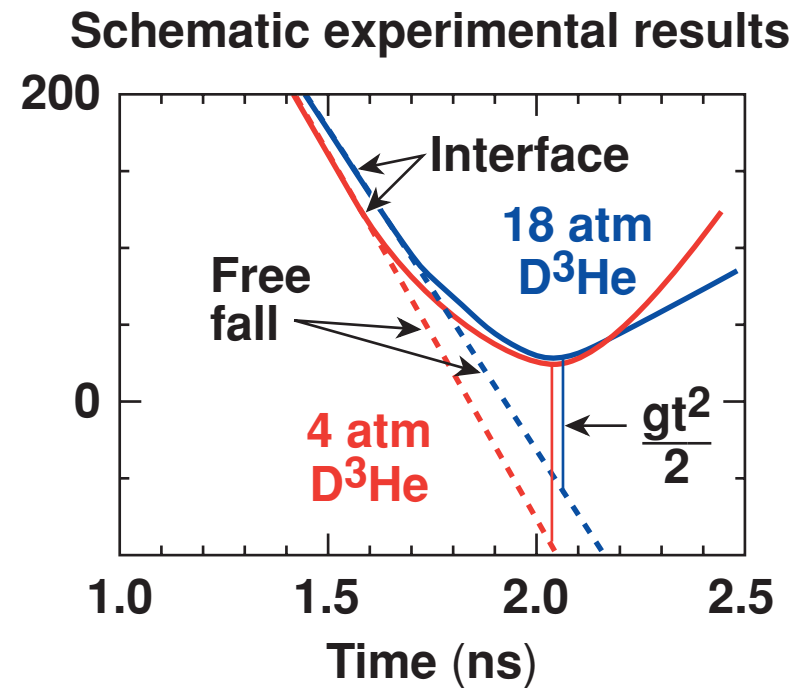
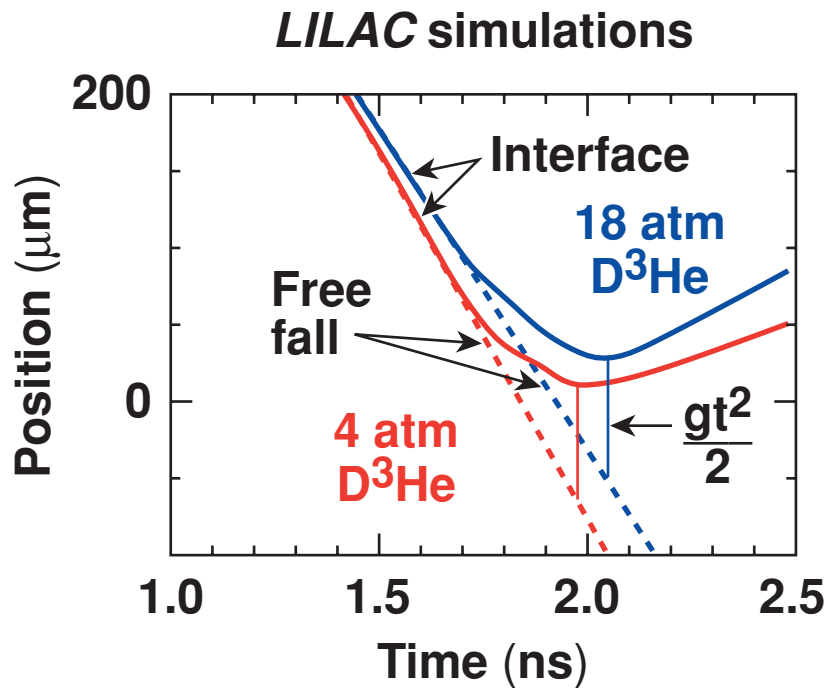


The target areal density grows by a factor of 1.5 from time of peak neutron production ( $\sim 1.9$  ns) to time of peak compression ( $\sim 2.0$  ns)



- Shell modulation due to Bell-Plesset (BP) convergent growth is proportional to shell thickness:  $\delta r \sim d$ .
- Measured shell integrity growth  $\delta \rho r / \rho d$  does not include Bell-Plesset effects.
- Areal density  $\rho d$  increases 1.5 times for 100 ps due to both thickness  $d$  and density  $\rho$ ; therefore Bell-Plesset growth is up to 1.5 times for 100 ps.

# Targets with 18 atm are more stable than targets with 4 atm D<sup>3</sup>He in the deceleration phase of spherical implosions



Classical Rayleigh–Taylor growth factor

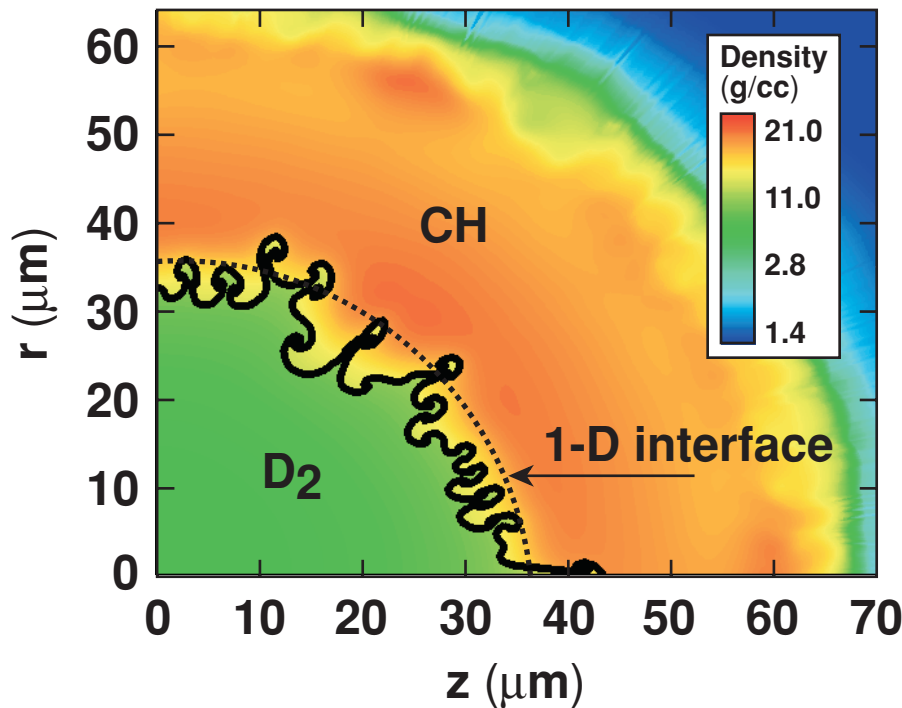
$$GF = \exp \left[ \sqrt{A_T \frac{\ell gt^2}{R}} \right], \text{ where}$$

$\ell$  – mode number,  $g$  – acceleration,  $R$  – inner surface radius,  $t$  – time, and  $A_T$  – Atwood number.

# Small-scale perturbations are in their nonlinear phase during deceleration

20- $\mu\text{m}$ -CH shell; 15-atm-fill laser imprint at peak neutron production

$$\ell = 2-80$$

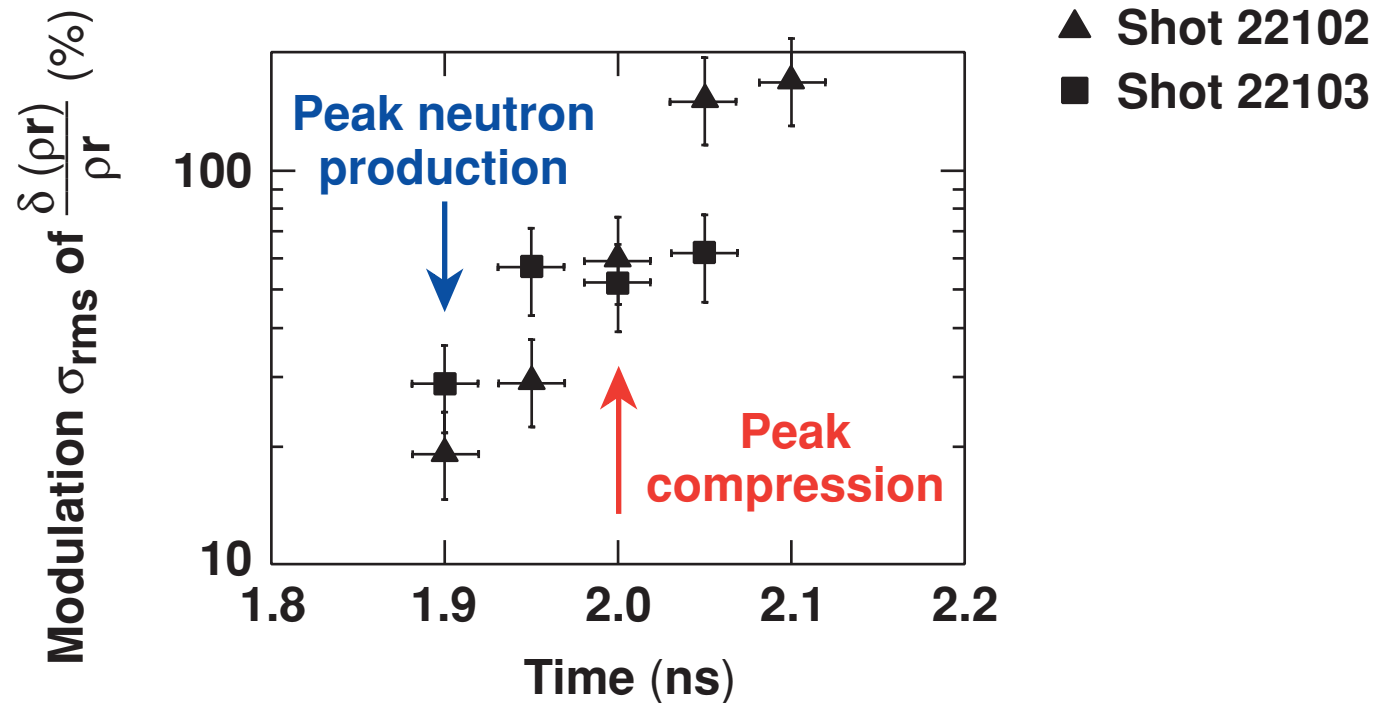


- Fuel-shell mix has been measured with neutron, particle, and spectroscopic diagnosis.

D. D. Meyerhofer *et al.*, Phys. Plasmas **8**, 2251 (2001).  
P. B. Radha *et al.*, Phys. Plasmas **9**, 2208 (2002).  
S. P. Regan *et al.*, Phys. Rev. Lett. **89**, 085003-1 (2002).  
C. K. Li *et al.*, Phys. Rev. Lett. **89**, 165002-1 (2002).

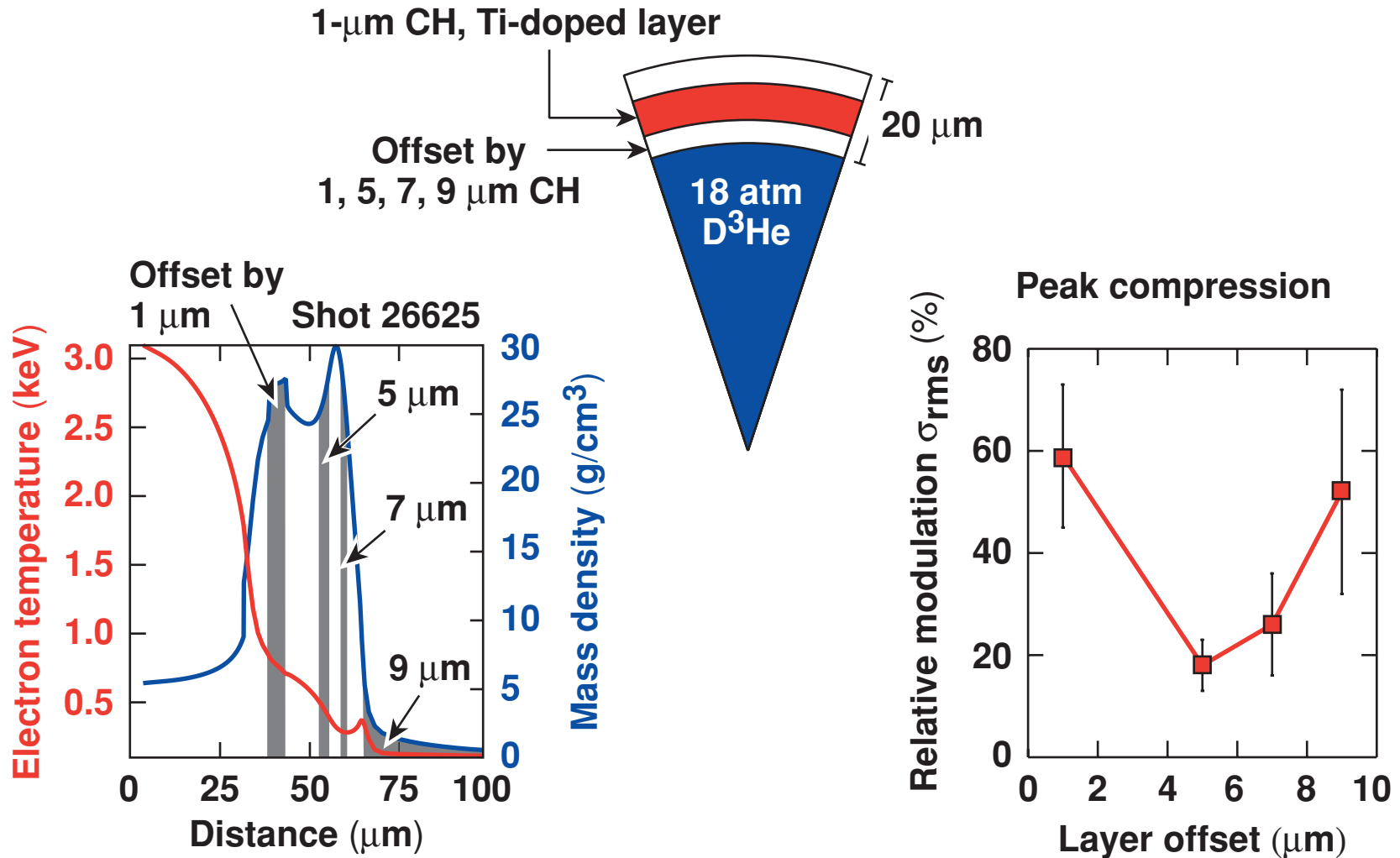


Inner-shell modulation grows from 20% at peak neutron production to 50% at peak compression for  $\sim 100$  ps due to Rayleigh–Taylor instability.



- In addition, the shell modulations grow by up to a factor of  $\sim 1.5$  due to Bell-Plesset effects.

# The central part of the shell is more uniform than the inner and outer surfaces\*



\* V. A. Smalyuk *et al.*, submitted to Phys. of Plasmas (2002).

## Future Work

# Future work will combine time-resolved shell areal-density with fuel–shell mix measurements



- Differential imaging with titanium 1s–2p absorption will provide much more sensitive shell-integrity  $\delta(\rho r)/\rho r$  measurements.
- Absorption spectroscopy of titanium 1s–2p region will provide shell-density, temperature, and areal-density measurements.
- Time-resolved fuel–shell mix measurements will be performed using CD shells filled with T<sub>2</sub> fuel.
- Time-resolved spectroscopic measurements of mix will be performed with Cl dopants in the shell and Ar dopants in the fuel.

## Shell modulation growth has been measured near peak compression of spherical implosions

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