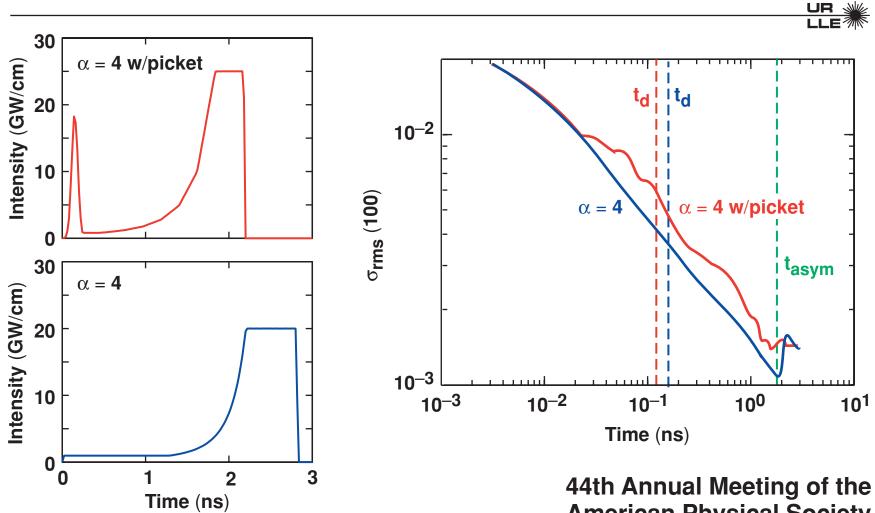
2-D SSD Model for Arbitrary Pulse Shapes Used in the Multidimensional Hydrodynamic Code DRACO



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Summary

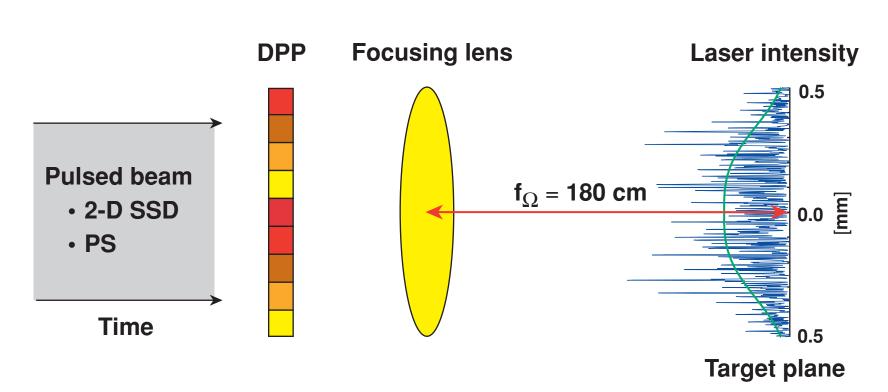
The pulse shape affects the laser nonuniformity on target

- A continuous model for 2-D SSD that accurately calculates the nonuniformity for arbitrary pulse shapes was developed.
- The model is deterministic; no need to average many runs.
- The nonuniformity predicted by the new model closely matches far-field simulations.



- Derivation of the new 2-D SSD model
- Comparison of the model with far-field simulations
- Comparison of the model with a stochastic model

SSD smoothes laser nonuniformities because the hydrodynamic response time is longer than the coherence time of the laser



- At every instant in time a speckle pattern is produced in the far-field.
- The speckle evolves continuously in time, but the pattern is statistically unique every coherence time.
- The long hydrodynamic response time causes the laser speckle field to be integrated and results in a smooth profile on target over time.

The instantaneous nonuniformity of the laser intensity does not change in time

• The intensity perturbations have a fixed amplitude, but with modes that continuously change phase due to 2-D SSD:

$$I(\theta, t) = I_0(t) + \sum_{\ell} a_{\ell} \cos\left[\theta + \phi_{\ell}(t)\right]$$
$$\sigma_{rms}^2(\ell, t) = f\left(\cos\left[\theta + \phi_{\ell}(t)\right]\right)$$

 The boundary conditions in 2-D ALE codes do not readily permit this behavior.

An auxillary model for laser intensity perturbations is used for *DRACO*

 The phases are fixed to satisfy boundary conditions while the amplitudes vary: UR 🔌

$$\hat{\mathbf{I}}(\boldsymbol{\theta}, \mathbf{t}) = \mathbf{I}_{\mathbf{0}}(\mathbf{t}) + \sum_{\ell} \hat{\mathbf{a}}_{\ell}(\mathbf{t}) \cos\left[\boldsymbol{\theta} + \boldsymbol{\phi}_{\ell}\right]$$
$$\hat{\sigma}_{rms}^{2}(\ell, \mathbf{t}) \propto \left| \int_{\mathbf{0}}^{\mathbf{t}} \hat{\mathbf{a}}_{\ell}(\mathbf{t}) d\mathbf{t} \right|^{2}$$

• A physical connection between the two laser intensity perturbation models is made by forcing the time evolution of the power spectra to match:

$$\sigma_{rms}^{2}(\ell,t) \equiv \hat{\sigma}_{rms}^{2}(\ell,t)$$

The time evolution of the power spectrum for any $\ell\text{-mode}$ is related to the near-field autocorrelation

 $\sigma_{rms}^{2}(\ell,t) \propto \left| \int_{0}^{t} (NF^{*} \times NF) dt \right|^{2}$

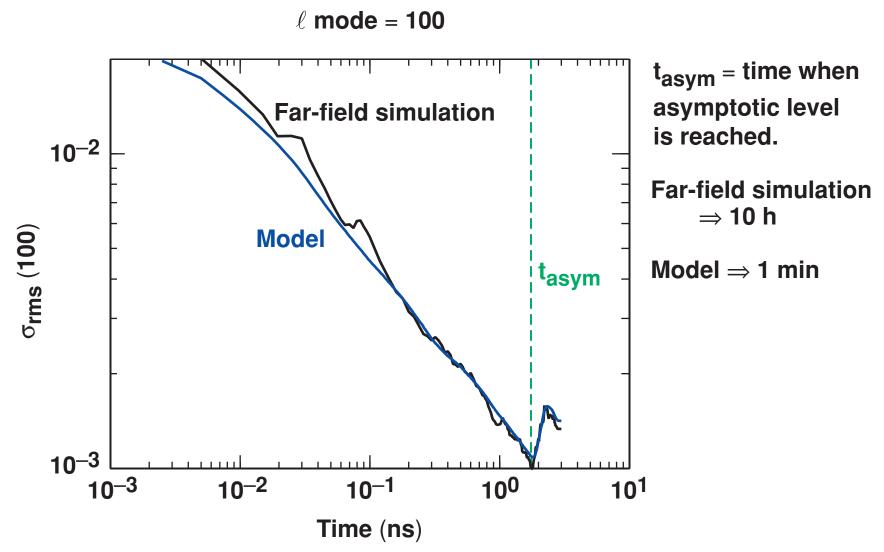
- The near-field term NF includes a random phase plate and 2-D SSD + PS.
- Bessel function expansions are used to simplify the correlation integral.
- The pulse is linearly interpolated.
 - An exact evaluation of the correlation integral within each linear section is possible.
- The exact solution of σ_{rms}^2 is then used to solve for the amplitudes of the auxillary model $\hat{a}_{\ell}(t)$.

There are many advantages to the new 2-D SSD model

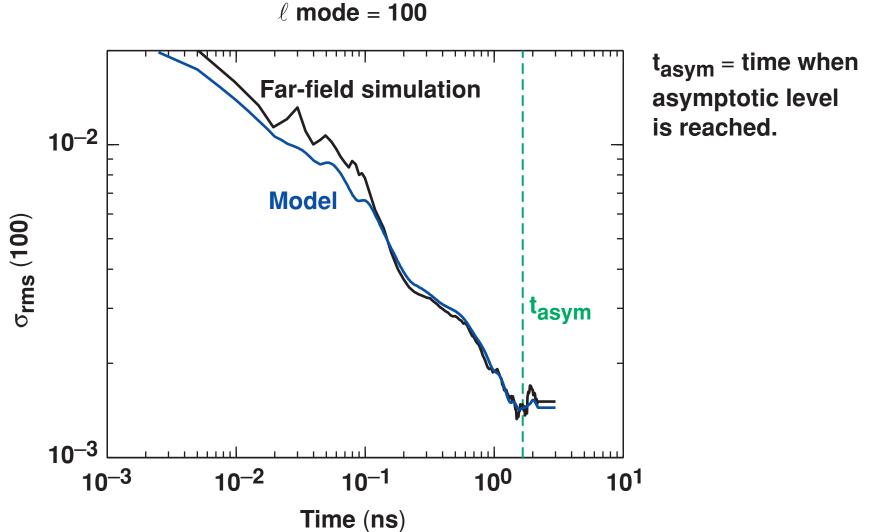
 Avoids the requirement of averaging multiple runs when using a flipping model. UR 🚽

- Continuously describes the mode amplitudes.
- Can model arbitrary pulse shapes.
- Explicitly uses the 2-D SSD parameters; not just phenomenological behavior.
- Accounts for the 2-D distribution of bandwidth.
- Much faster than calculating the mode amplitudes with a full far-field simulation.

The new 2-D SSD model matches the results of a far-field simulation of an α = 4 pulse



The new 2-D SSD model matches the results of a far-field simulation of an α = 4 with a picket pulse*



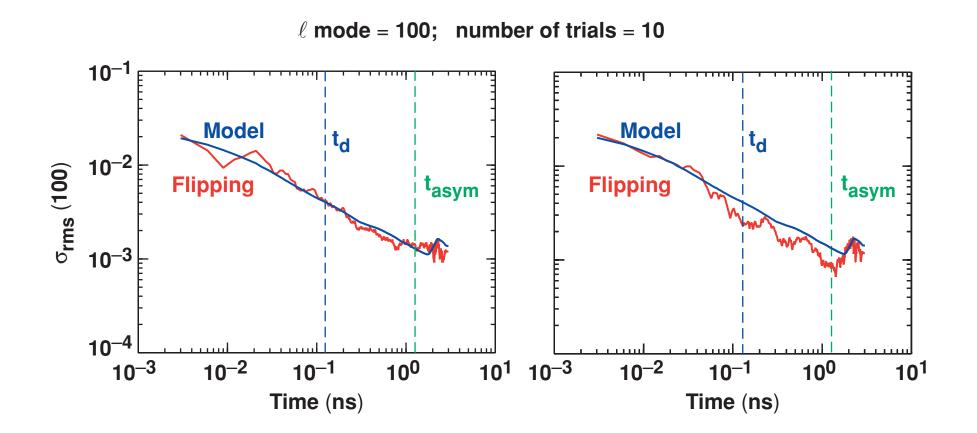
The pulse shape's effect on nonuniformity can have a negative impact near the decoupling time

ℓ mode = 100 The fast-rising Гd edge of the picket 10-2 competes with the smoothing effect of SSD. σ_{rms} (100) $\alpha = 4 \text{ w/picket}$ $\alpha = 4$ t_{asym} = time when asymptotic level lasym is reached t_d = laser decoupling time 10⁻³ 10⁻³ 10⁰ 10⁻² 10⁻¹ 10¹ Time (ns)

A random flipping model in *DRACO* compares with the new model if enough trials are averaged

 ℓ mode = 100; number of trials = 500 **Stochastic flipping** model: 10^{-1} • The phase is td flipped every coherence time. The phase has Flipping 10⁻² flipped ~ 40 times **Model** σ_{rms} (100) by the decoupling time t_d. t_{asym} = time when 10⁻³ asymptotic level is reached lasym t_d = laser decoupling time 10^{-4} 10⁻³ 10⁻² 10⁻¹ **10⁰** 10¹ Time (ns)

Unpredictable results occur if only ten trials are used



 Some sets of random trials can yield erroneous nonuniformity levels near the decoupling time t_d.

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