

# SBS in Multiple-Species Plasmas

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## Summary

# We developed a model to study fast- and slow-wave SBS and how nonlinearities affect them

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- **Fast-wave SBS reflectivities are**
  - significantly reduced by nonlinear saturation for  $T_e/T_i \gtrsim 10$
  - not affected by nonlinear saturation for  $T_e/T_i \lesssim 5$
- **In CH plasmas, slow-wave SBS reflectivity is much lower than fast-wave SBS reflectivity.**
- **Slow-wave SBS reflectivities are significantly increased by nonlinear effects, but remain much smaller than fast-wave SBS.**

# Outline

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- **Description of the model**
- **Fast-wave SBS at low and high ion temperatures**
- **Slow-wave SBS**
- **Conclusions**

# We modeled SBS in CH plasmas at OMEGA-like parameters

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- We modeled backscattering using a one-dimensional fluid model with phenomenological Landau-damping terms.
- Light wavelength:  $0.35 \mu\text{m}$
- Pump intensity:  $I = 2.45 \times 10^{15} \frac{\text{W}}{\text{cm}^2}$
- Electron density:  $n_e/n_{\text{cr}} = 0.46$
- Electron temperature:  $T_e = 1 \text{ keV}$

# Model consists of electron density, ponderomotive potential, ion-fluid equations, and Poisson equations



- **Electron density:**  $n_e = \exp(\varphi) - P = 1 + \delta n_e$
- **Ponderomotive potential:**  $\partial_x P = \alpha \mathfrak{S}(n_{e1}) \sqrt{P^2 + C^2}$ ,  
 $\alpha = (v_e/c)^2 / ks$ ,  $C = |a_0|^2 + |a_1|^2$

- **Ion-fluid equations for species “s”:**

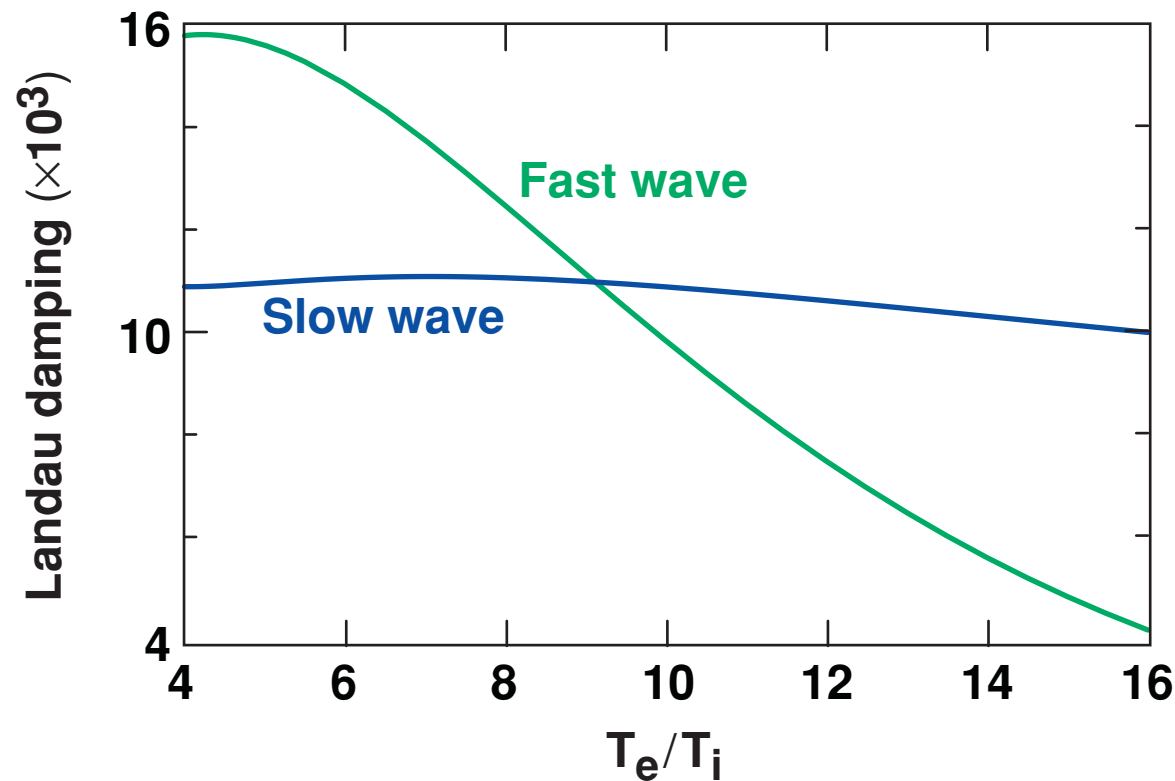
$$\partial_t n_s + \partial_x (n_s u_s) + \lambda n_s = 0$$

$$\partial_t u_s + u_s \partial_x u_s + \beta_s \partial_x \varphi + \gamma_s v_{ths}^2 n_s^{(\gamma_s - 2)} \partial_x n_s + \lambda u_s = 0$$

- $\lambda$ –Landau damping coefficients;  $\gamma_s$ –polytropic indices;  $\beta_s = \frac{z_s M_r}{z_r M_s}$
- **Poisson equation:**  $\partial_{xx}^2 \varphi - n_e + \sum_{s=l,h} \alpha_s n_s = 0$ ;  $\alpha_s = Z_s n_{0s} / n_{0e}$

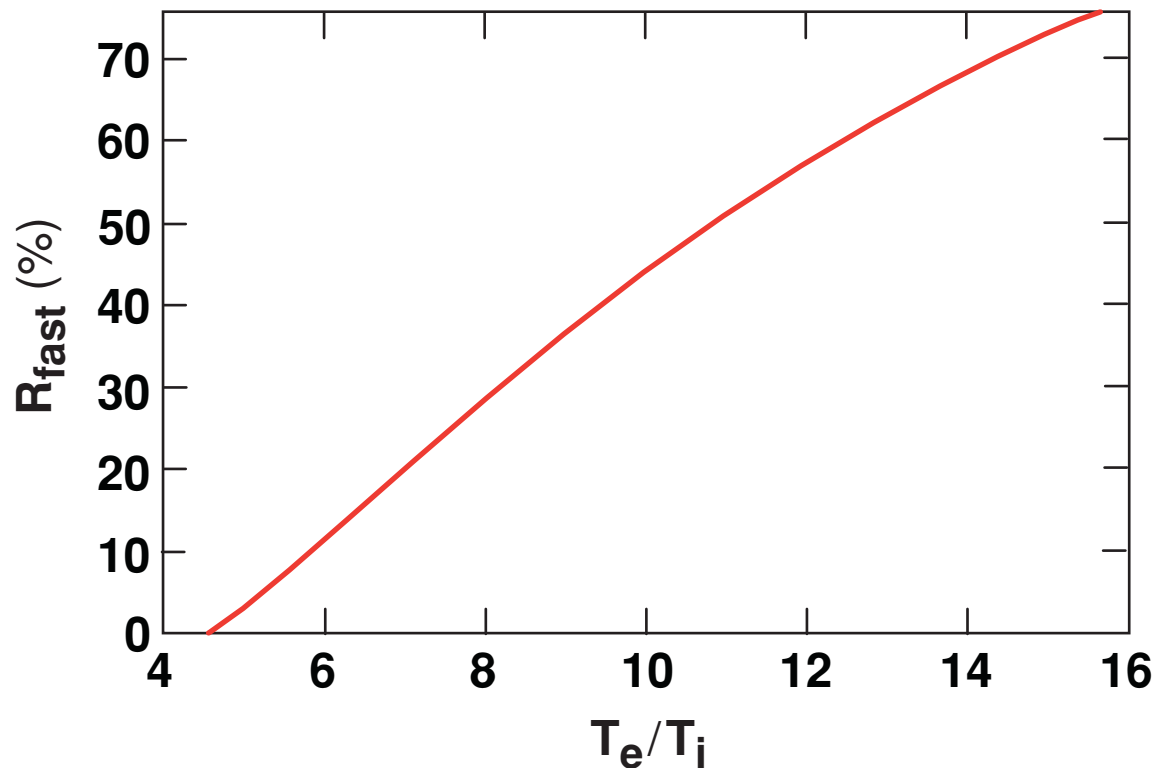
# Results of kinetic model<sup>1</sup> are used to close the system of fluid Equations

- $\gamma = 3$  gives best agreement between fluid and kinetic phase velocities.
- Phenomenological damping terms correspond to imaginary parts of kinetic frequencies.

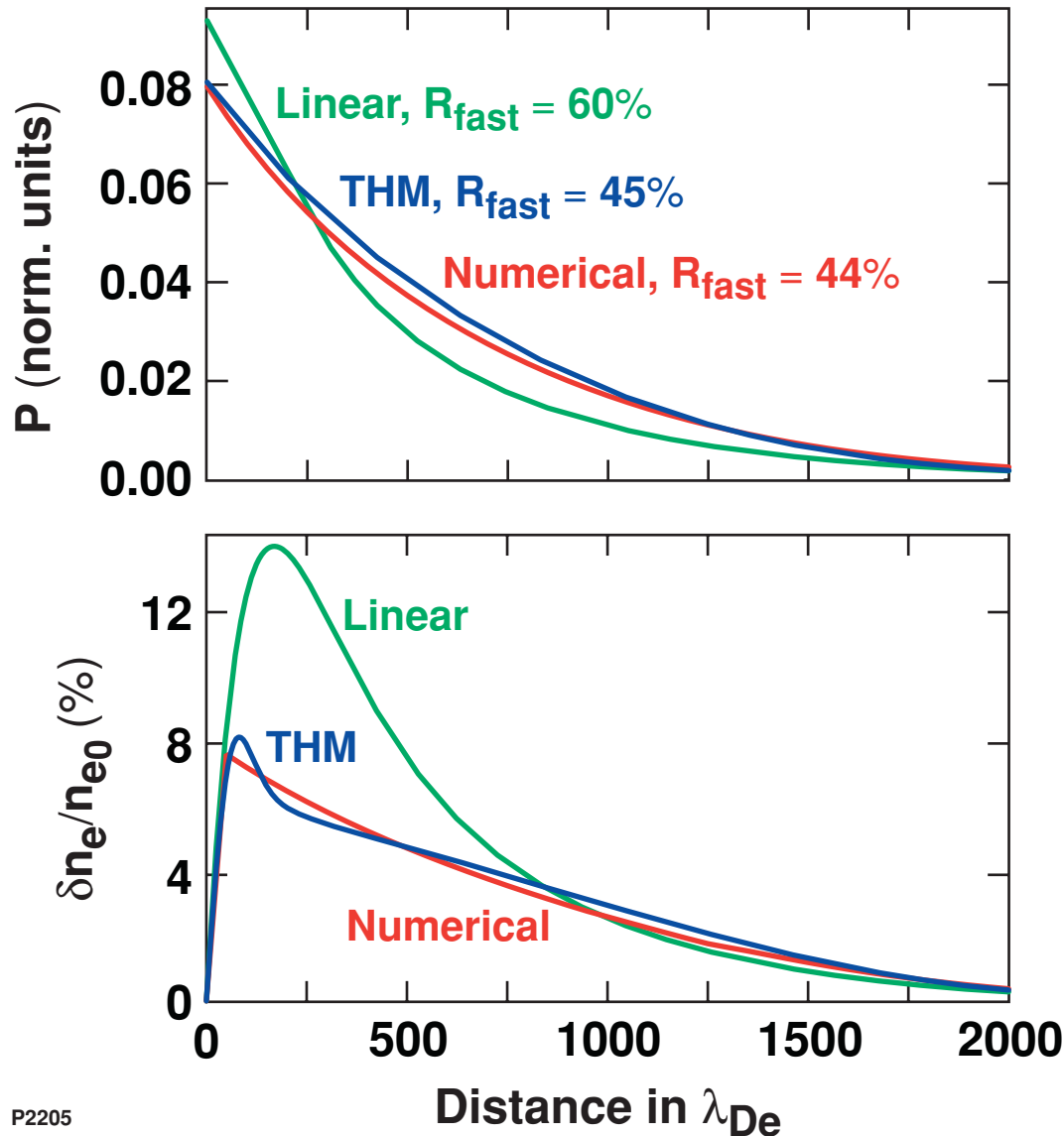


# Damping of fast wave increases with ion temperature

- Due to exponential dependence on damping, fast-wave SBS reflectivity decreases fast with increasing ion temperature.
- For  $T_e/T_i = 5$ ,  $\gamma/\omega_{pi} = 1.56 \times 10^{-2}$ ,  $R_{fast} = 3\%$  → IA wave amplitude is small → nonlinearities are weak.



For  $T_e/T_i = 10$ , damping is small and fast-wave SBS reflectivity is large



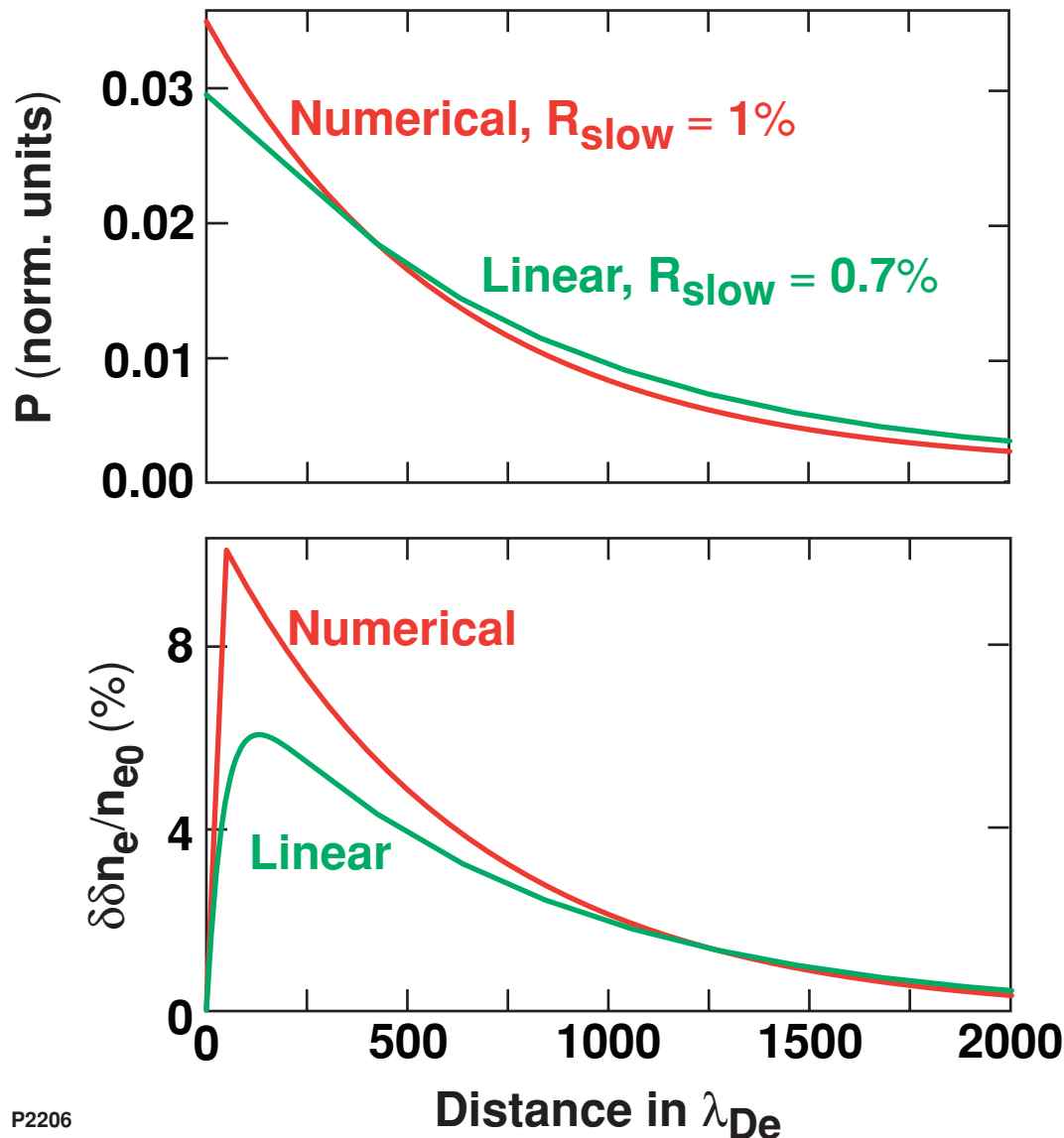
- $\gamma/\omega_{pi} = 10^{-2} \Rightarrow R_{fast} = 44\%$   
IA wave amplitude is large enough to be affected by nonlinear saturation.

- The two-harmonics model<sup>1</sup> (THM) has been extended for two-ion plasmas.

<sup>1</sup> J. A. Heikkinen *et al.*,  
Phys. Plasmas 27, 707 (1984).



# In CH plasmas slow-wave SBS reflectivity is much smaller than fast-wave SBS reflectivity



- In slow-wave protons are pushed away from carbon ions shielding  $\phi$  and thus reducing  $\delta n_e$  and  $R_{slow}$ .
- For  $I = 2 \times 10^{15} \text{ W/cm}^2$ , slow-wave SBS is below threshold.
- For  $I = 7 \times 10^{15} \text{ W/cm}^2$ ,  $T_e/T_i = 10$   $\gamma/\omega_{pi} = 10^{-2}$  and  $R_{slow} = 1\%$ .
- Nonlinear saturation for  $\delta n_h \gg \delta n_c \rightarrow$  increase in  $\delta n_e$ .

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- **Slow-wave SBS reflectivities are significantly increased by nonlinear effects, but remain much smaller than fast-wave SBS.**
- **Application of our model to different component plasmas is underway**