Relativistically Transparent Magnetic Filaments:

a short-pulse path to megaTesla fields and efficient gamma radiation



Schematic of a magnetic filament Plasma: $n_e \le a_0 n_c$ r_{mb} r_{mb} r_{mb} $M_{agnetic filament:}$ $B \propto R$ γ radiation

E29601J1

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Rinderknecht et al., New J. Phys 23, 095009 (2021)

Hans Rinderknecht University of Rochester Laboratory for Laser Energetics

Anomalous Absorption Conference 11:00 AM, Tuesday June 7th 2022

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UNIVERSITY OF ROCHESTER, LABORATORY FOR LASER ENERGETICS

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Photon energy Efficiency 10 100% (d) η (8) (a) 8 6 3D-PIC MeV scaling laws 10% 2 1% 0 2 0 2 0 (R/λ) (R/λ) E29602J1

Predicted radiation (I = 5×10²² W/cm²)

Anomalous Absorption Conference 11:00 AM, Tuesday June 7th 2022

Magnetic filaments promise a repeatable and efficient laser-driven source of MT fields, relativistic electrons, and MeV photons

- Intense lasers in relativistically-transparent plasmas generate ultra-strong magnetic filaments
 - Electrons are trapped and accelerated, efficiently radiating MeV-scale photons
- Scaling laws were derived for magnetic filament radiation, and validated with 3-D PIC simulations
 - Efficiency of >10% is predicted for intensity above 6×10^{21} W/cm²
- Experiments have been performed on Texas Petawatt and OMEGA-EP to test these predictions
 - The predicted electron and photon signatures were observed in a subset of experiments



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Collaborators



LLE/UR:

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- Jarrod Williams
- Alex Haid

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Dan Stutman













Relativistic transparency allows an intense laser to interact with a plasma above the classical critical density (n_{cr})



$$a_0 \equiv \frac{|e|E_0}{m_e \omega c} \propto \sqrt{Intensity}$$

> 1 \rightarrow relativistic electrons

 $n_{cr} = \frac{\epsilon_0 m_e}{e^2} \omega_L^2$ = critical plasma density limit for a laser of frequency ω_L

 $a_0 > 1 \rightarrow$ Relativistic Transparency: $n_e \le n_{cr} \langle \gamma \rangle \approx n_{cr} a_0$

¹Z. Gong, et al., Phys. Rev. E 102, 013206 (2020)



The ponderomotive force drives a relativistic current, producing a strong azimuthal magnetic field: a magnetic filament



3-D PIC simulations $(a_0 = 50)^1$:

Magnetic field of current normalized to laser field:



Quasi-static magnetic fields of the order of the oscillating laser field are produced.

¹Z. Gong, et al., Phys. Rev. E 102, 013206 (2020)



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Electrons oscillate within the filament: those in phase with the laser are rapidly accelerated,



¹Z. Gong, et al., Phys. Rev. E 102, 013206 (2020);

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Electrons oscillate within the filament: those in phase with the laser are rapidly accelerated, and radiate in the strong magnetic field.



The maximum magnetic field seen by electrons depends on the laser focal radius and the initial transverse momentum.



¹Z. Gong, et al., Phys. Rev. E 102, 013206 (2020);

20); ²D. J. Stark, et al., Phys. Rev. L

Analytical scaling laws and 3-D PIC simulations of magnetic filaments predict > 10% radiation efficiency of MeV photons from sufficiently intense lasers







Initial experiments to study relativistically transparent magnetic filaments were performed at the Texas Petawatt Laser (TPW)¹









- Wavelength:
- Energy:
 - **Duration**:
- Power:

- 1057 nm
- 98.8 ± 6.0 Joules 140 fs
 - 694 ± 38 TW

- Intensity:
- Radius:
- $[1.09 \pm 0.07] \times 10^{21} \text{ W/cm}^2 (a_0 = 29.9 \pm 1.0)$
- 2.6 ± 0.12 μm (at 50% peak intensity)
- Pointing:
- 8-µrad rms \rightarrow 5-µm rms on target
- Primary diagnostics: EPPS electron spectrometer² Gamma calorimeter



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Microchannel targets filled with low-density foam ($n_e = 5 \text{ or } 10 n_{cr}$) were developed for this campaign

Kapton wedge



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The 'hot' electron temperature was elevated on 2 of 8 microchannel shots





ROCHESTER H.G. Rinderknecht, et al., New J. Phys. 23, 095009 (2021)



ROCHESTER H.G. Rinderknecht, et al., New J. Phys. 23, 095009 (2021)

12 $6-\mu m ID$ $6-\mu m ID$ \diamond 30 mg/cm^3 15 mg/cm^3 Electron T_1 and T_2 (MeV) 10 $\overline{\mathbf{\Phi}}$ \diamond 8 ∮ \diamond -6 $\overline{\mathbf{\Phi}}$ $\overline{\mathbf{Q}}$ 4 2 0 3-D PIC 2 8 Shot number

Electron temperatures

Given the pointing stability and channel size, the probability of observing N interactions is:

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Ν	Probability
0	0.21
1	0.36
2	0.27
3	0.12
4	0.03
5+	< 0.01

We conclude that the predicted electron acceleration was observed in a subset of these experiments.



H.G. Rinderknecht, et al., New J. Phys. 23, 095009 (2021)

A follow-on campaign on OMEGA-EP (Dec 2022) used a new target engineering technique to improve laser-target coupling: 3-D printing *inside* foam



Target Geometry



Electron radiograph (empty channels)



Foam-filled channel targets:





Shot #	Beam	Peak I ×10 ²⁰	a₀ max	Target ρ
35891	BL	3.49	16.8	Empty channels
35892	SL	1.42	10.7	Empty channels
35893	BL	10.1	28.6	3 mg/cc (1 n _{cr})
35894	SL	3.84	17.6	3 mg/cc (1 n _{cr})
35895	BL	7.69	25.0	3 mg/cc (1 n _{cr})
35896	SL	3.01	15.6	3 mg/cc (1 n _{cr})
35897	BL	9.64	28.0	5 mg/cc (1.6 n _{cr})
35898	SL	3.08	15.8	3 mg/cc (1 n _{cr})
35899	BL	6.76	23.4	5 mg/cc (1.6 n _{cr})
35900	SL	2.74	14.9	5 mg/cc (1.6 n _{cr})
35901	BL	6.56	23.1	5 mg/cc (1.6 n _{cr})
35902	SL	4.49	19.1	5 mg/cc (1.6 n _{cr})
35903	BL	6.80	23.5	Empty channels
35904	SL	1.55	11.2	5 mg/cc (1.6 n _{cr})

This technique increases the density of channel openings and guarantees the foam fully fills the channel.



The first OMEGA-EP experiments show electron acceleration on all target types: No significant difference was observed between target types**



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Subsequent 3-D PIC simulations predict that, for 0.7 ps pulses, the magnetic field amplitude and electron and x-ray spectra are independent of fill density



All cases essentially behave the same as the empty channel

${ m (M/MeV)}_{10_{10}} \ (\#/MeV)_{10_{10}} \ 10_{8}$ 10^{12} Empty Ejected $1 n_c$ electrons $-5 n_c$ 2080 100 120 140 4060 ε (MeV) $\frac{10}{10} (\#/\mathrm{MeV})^{-10} \mathrm{MeV}$ 10^{15} Empty **Photons** $1 n_c$ $-5 n_c$ 0 2 4 ε_{γ} (MeV)

Slide courtesy of: Kathleen Weichman



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For future EP experiments, we will continue to test out advanced targets for improved laser-channel coupling and "long-pulse" coupling physics

Improved target arrays (June 15, 2022)



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Contentrator targets (Matt Van-Dusen Gross (GS), FY23)



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With 10-PW lasers now becoming available, magnetic filaments promise exciting opportunities for high-flux gamma-ray sources

Laser	ELI-NP [†]		ELI-Beamlines L4 [‡]	
λ	0.8 µm		1.057 μm	
т	23 fs		150) fs
Peak power	10 PW		10 PW	
Intensity (a_0)	5×10 ²² W/cm ² (153)		5×10 ²² W/cm ² (202)	
Design choice:	S _α = 0.01	S _α = 0.05	S _α = 0.01	S _α = 0.05
Photon energy <ε∗>	68 MeV	9.2 MeV	96 MeV	19 MeV
Total energy $E_{\gamma,tot}$	111 J	51 J	797 J	727 J
# photons N _γ	1.0×10 ¹³	3.5×10 ¹³	5.2×10 ¹³	2.5×10 ¹⁴
Efficiency η	48%*	22%*	53%*	48%*

By varying the channel design, the photon spectrum and flux may be optimized.

[†] D. Ursescu, et al., Romanian Reports in Physics 68, S11 (2016) [‡] S. Weber, et al., Matter and Radiation at Extremes 2, 149 (2017)

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Magnetic filaments promise a repeatable and efficient laser-driven source for MT fields, relativistic electrons and MeV photons



- Electrons are trapped and accelerated, efficiently radiating MeV-scale photons
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Appendix

* First reference ** Second reference † Third reference ‡ Fourth reference

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Electrons orbit within a magnetic boundary. Those in phase with the laser are accelerated, and radiate by deflecting in the strong magnetic field.



Using simple assumptions for the electron acceleration and orbits, we derive scaling laws for the radiation from magnetic filaments

0: Electrons are thermal

$$f_e(\acute{\mathbf{U}}_e, t) = \frac{N_e}{T_e} \exp\left[-\frac{\acute{\mathbf{U}}_e}{T_e}\right], \text{ where } N_e = n_e\left(\pi R^2\right)(c\tau)$$

1: Electron acceleration is linear in time

$$T_{e}(t) = C_{T} a_{0}\left(\frac{ct}{\lambda}\right) mc^{2} \equiv C_{T} a_{0} t_{v} mc^{2}$$

2: Radiation is synchrotron-like

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$$\frac{dP}{d\dot{\mathsf{U}}} = f_r \frac{4}{9} \alpha_{fsc} \frac{mc^2}{\hbar} \left(\frac{B}{B_{cr}}\right) F\left[\frac{\dot{\mathsf{U}}}{\dot{\mathsf{U}}_c}\right], \quad \text{where} \quad \dot{\mathsf{U}}_c = \frac{3}{2} \chi \gamma mc^2, \quad F[x] \equiv \frac{9\sqrt{3}}{8\pi} x \int_x^\infty K_{5/3}(z) dz \quad \left[\int_0^\infty F(y) dy = 1\right]$$

3: The laser depletes by heating electrons

$$\frac{E_e}{E_{Laser}} \le 1 \quad \to \quad t_{\nu,\max} \le \frac{\sqrt{\pi}}{4(\ln 2)^{3/2}} \frac{1}{C_T S_\alpha}.$$

We define:
$$t_{v,cut} \equiv f_t t_{v,max} \approx 0.768 \frac{f_t}{C_T S_{\alpha}}$$

These assumptions have four constants: f_i , f_t , f_r , C_T and four design parameters: a_0 , S_{α} , R/λ , c_T/λ

H.G. Rinderknecht, et al., New J. Phys. **23**, 095009 (2021) doi: <u>10.1088/1367-2630/ac22e7</u>



Using simple assumptions for the electron acceleration and orbits, we derived scaling laws for the radiation from magnetic filaments

Radiation property:	<u>if focal radius R < r_{mb}:</u>	if focal radius R > r _{mb} :
Photon energy	$\left\langle \acute{\mathbf{U}} \right\rangle_{tot} \approx 1.38 \times 10^{-6} f_t^2 a_0^3 S_\alpha^{-1} \mathbf{R}_\lambda \lambda_{\mu m}^{-1} m_e c^2$	$\langle \hat{\mathbf{U}}_{i} \rangle_{tot} \approx 4.40 \times 10^{-7} \sqrt{f_i} f_i^2 a_0^3 S_{\alpha}^{-3/2} \lambda_{\mu m}^{-1} m_e c^2$
Radiated energy	$E_{\gamma,tot} \approx 7.74 \times 10^2 f_r f_t^3 C_T^{-1} a_0^5 R_{\lambda}^4 \tau_{\nu} m_e c^2$	$E_{\gamma,tot} \approx 7.84 \times 10^{1} f_{i}f_{r}f_{t}^{3}C_{T}^{-1}a_{0}^{5}S_{\alpha}^{-1}R_{\lambda}^{2}\tau_{v} m_{e}c^{2}$
# photons	$N_{\gamma,tot} = 5.59 \times 10^8 f_r f_t C_T^{-1} a_0^2 S_\alpha R_\lambda^3 \tau_\nu \lambda_{\mu m}$	$N_{\gamma,tot} = 1.78 \times 10^8 \sqrt{f_i} f_r f_t C_T^{-1} a_0^2 S_{\alpha}^{1/2} R_{\lambda}^2 \tau_{\nu} \lambda_{\mu m}$
Radiation efficiency	$\eta_{\gamma} = 2.88 \times 10^{-7} f_r f_t^3 C_T^{-1} a_0^3 R_{\lambda}^2 \lambda_{\mu m}^{-1}$	$\eta_{\gamma} = 2.92 \times 10^{-8} f_r f_t^3 f_i C_T^{-1} a_0^3 S_{\alpha}^{-1} \lambda_{\mu m}^{-1}$

→ 4 Parameters: Intensity (a_0), Relativistic transparency ($S_{\alpha} = n_e/n_{cr}a_0$), Focal radius (R_{λ}), Pulse duration (τ_v) → 3 Constants: Initial e- momentum scalar ($f_i \sim 1$), cutoff scalar ($f_t < 1$), radiation duty cycle ($f_r < 1$)

> *For more details, please see:* H.G. Rinderknecht, et al., New J. Phys. 23, 095009 (2021) doi:10.1088/1367-2630/ac22e7



The scaling laws show good agreement with 3-D PIC simulations that varied the focal radius, with reasonable constants



doi:10.1088/1367-2630/ac22e7



Primary diagnostics were an electron spectrometer (EPPS) and a gamma calorimeter (GCAL) in the expected radiation direction





A factor of 5 difference in photon brightness is predicted between microchannel and 'solid' targets.

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The 'hot' electron temperature was elevated on 2 of 8 microchannel shots







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ROCHESTER H.G. Rinderknecht, et al., New J. Phys. 23, 095009 (2021) doi: 10.1088/1367-2630/ac22e7 LABORATORY FOR LASER ENERGETICS

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We conclude that the predicted electron acceleration was observed in a subset of these experiments.





The number of photons > 10 keV also scaled with hot electron temperature as expected.

