

A Wave-Based Model for Cross-Beam Energy Transfer in Inhomogeneous Plasmas



J. F. MYATT, J. G. SHAW, R. K. FOLLETT, D. H. EDGELL, V. N. GONCHAROV, A. V. MAXIMOV, R. W. SHORT, W. SEKA, and D. H. FROULA

University of Rochester, Laboratory for Laser Energetics

Summary

A 3-D wave-based model of CBET* has been successfully developed in LPSE



- This model solves the time-enveloped Maxwell equations in 3-D coupled to a fluid equation for the low-frequency ion-acoustic response
 - radiative boundary conditions, arbitrary incident beams
- A solid understanding of CBET has been obtained on OMEGA
 - coordinated program of theory, numerical simulations, and experiments
- CBET mitigation is a crucial part of LLE's 100-Gbar plan

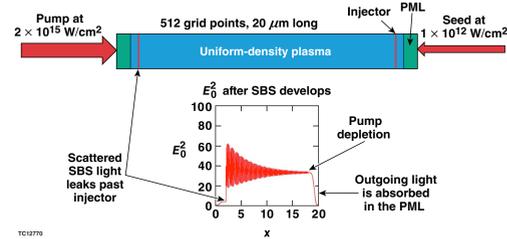
TC12767

*cross-beam energy transfer

The boundary conditions use a "total-field/scattered-field" formulation together with a perfectly matched absorbing layer (PML)

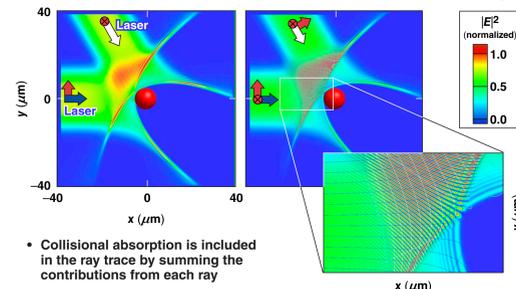


- Inject a pump wave on the left-hand side and a (weak) seed wave on the right-hand side
- Match the stimulated Brillouin scattering (SBS) resonance condition for the seed by changing its frequency and/or adding a flow velocity to the plasma



TC12770

The laser light is partially coherent; the intensity is not the sum of the intensity of individual beams



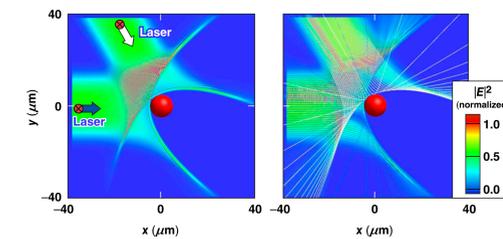
- Collisional absorption is included in the ray trace by summing the contributions from each ray

TC12773

This effect can have a dramatic impact on laser absorption and the drive of an ICF target



- It appears to operate in the regime of a convective amplifier for direct-drive ICF, which may be a tractable problem to describe with rays

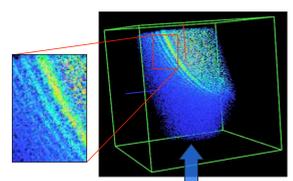


TC12776

A wave-based CBET model is required for several important reasons



- There are uncertainties associated with ray-based CBET models that are hard to quantify without comparison with a more-fundamental model
- The model's correctness is empirically determined; however, experimental tests of CBET are integrated experiments (indirect)
- Caustic surfaces/turning points (field swelling, Airy-like patterns)
- Beam speckle (spatial and temporal incoherence)
- Polarization effects
- The IAW response is approximate in ray-based CBET (steady state, strong damping)

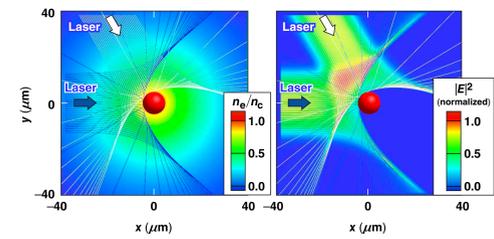


TC12777

Laser-beam propagation and energy deposition is computed in ICF* design codes using ray tracing



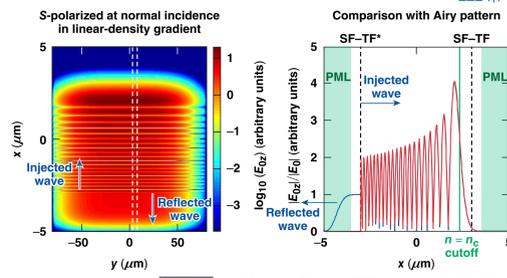
- Absent nonlinearity, the geometric-optics approximation is well justified based on the long plasma scale lengths
- Power is deposited based on collisional absorption of laser light



TC12768

*inertial confinement fusion

The LPSE electromagnetic (EM) wave solver reproduces analytical results for propagation in an inhomogeneous plasma



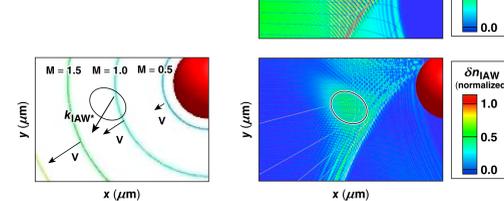
TC12771

D. E. Merewether, R. Fisher, and F. W. Smith, IEEE Trans. Nucl. Sci. 27, 1829 (1980). *scattered field-total field

The electric-field grating resonantly excites ion-acoustic waves because of the plasma-flow Doppler shift



- If the two EM waves have equal frequencies, the ion-acoustic perturbation will be large if $k \cdot v = c_s$
- Mach number $M = |v|/c_s$



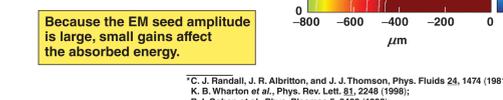
TC12774

*ion-acoustic wave

All direct-drive CBET calculations have been performed using a 1-D description* that has been adapted to geometric ray tracing



- Unlike x-ray drive, the presence of supersonic plasma flow enables the process to be resonant*
- Three-wave SBS equations are computed (pairwise) for each beam crossing using a generalization of Randall et al.* and are implemented in-line in 1-D LILAC

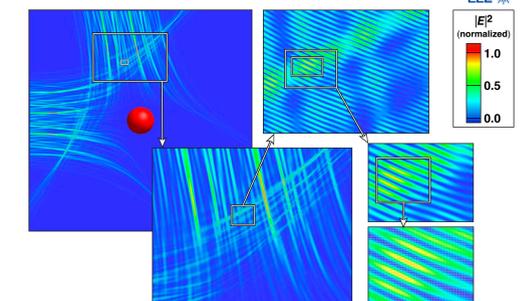


E117944

Because the EM seed amplitude is large, small gains affect the absorbed energy.

*C. J. Randall, J. R. Albritton, and J. J. Thomson, Phys. Fluids 24, 1474 (1981); K. B. Wharton et al., Phys. Rev. Lett. 81, 2248 (1998); B. J. Cohen et al., Phys. Plasmas 5, 3408 (1998); H. A. Rose and S. Ghosal, Phys. Plasmas 5, 1461 (1998).

Spatial incoherence can be modeled with no difficulty; temporal incoherence is only slightly harder



TC12778

1/10,000 original

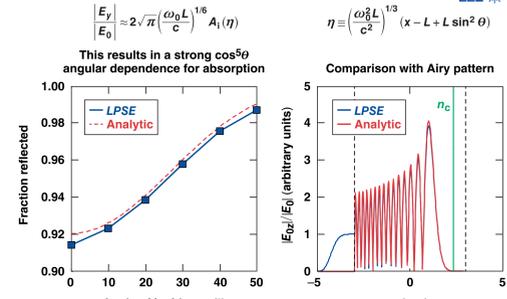
Wave-based models face several challenges and complications, but these can be overcome



- We are computing the motion of a semi-classical object by solving the (vector) Schrödinger equation
- Several key technical challenges were solved
 - time-enveloped (vector) wave equation in 3-D in strongly inhomogeneous plasma of a useful size
 - efficient algorithm is required (i.e., not Crank-Nicholson)
 - very complicated boundary conditions (in 3-D)
 - coupled to a plasma model
 - parallel efficiency [scalable solver for $O(10^9)$ computational cells]
- The resulting wave solver is practical to run in 3-D
 - 100 Intel cores, could scale to 1000's
 - pioneering use/visualization of large datasets/fast disks at LLE [O(100 GB) sets]

TC12769

Obliquely incident light turns at a lower density (shifted by $\cos^2\theta$ in a linearly varying density profile)

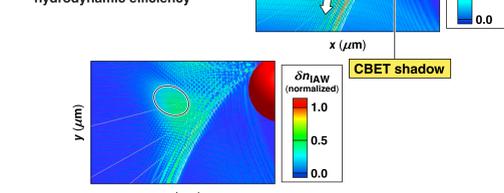


TC12772

This becomes an induced SBS process; laser energy changes direction



- Laser energy can be redirected before it has reached its turning point
- This leads to a reduction of absorption and hydrodynamic efficiency

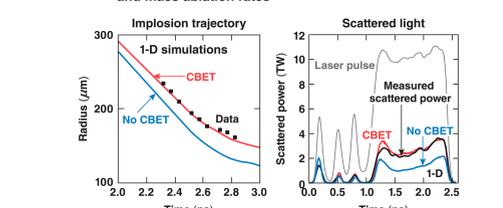


TC12775

A nonlinear CBET model is required to obtain agreement between 1-D predictions and OMEGA experimental data



- The CBET model used to obtain agreement with $\alpha > 3.5$ data (not compromised by mix) is ray-based
 - scattered-light power and spectrum, shell trajectories, and mass ablation rates



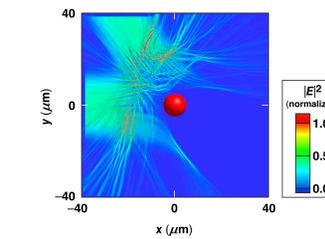
E81916K

At high enough laser intensities, CBET may not act as a simple spatial amplifier



- Shock-ignition experiments exceed filamentation thresholds

TC12779



Summary

A 3-D wave-based model of CBET* has been successfully developed in *LPSE*

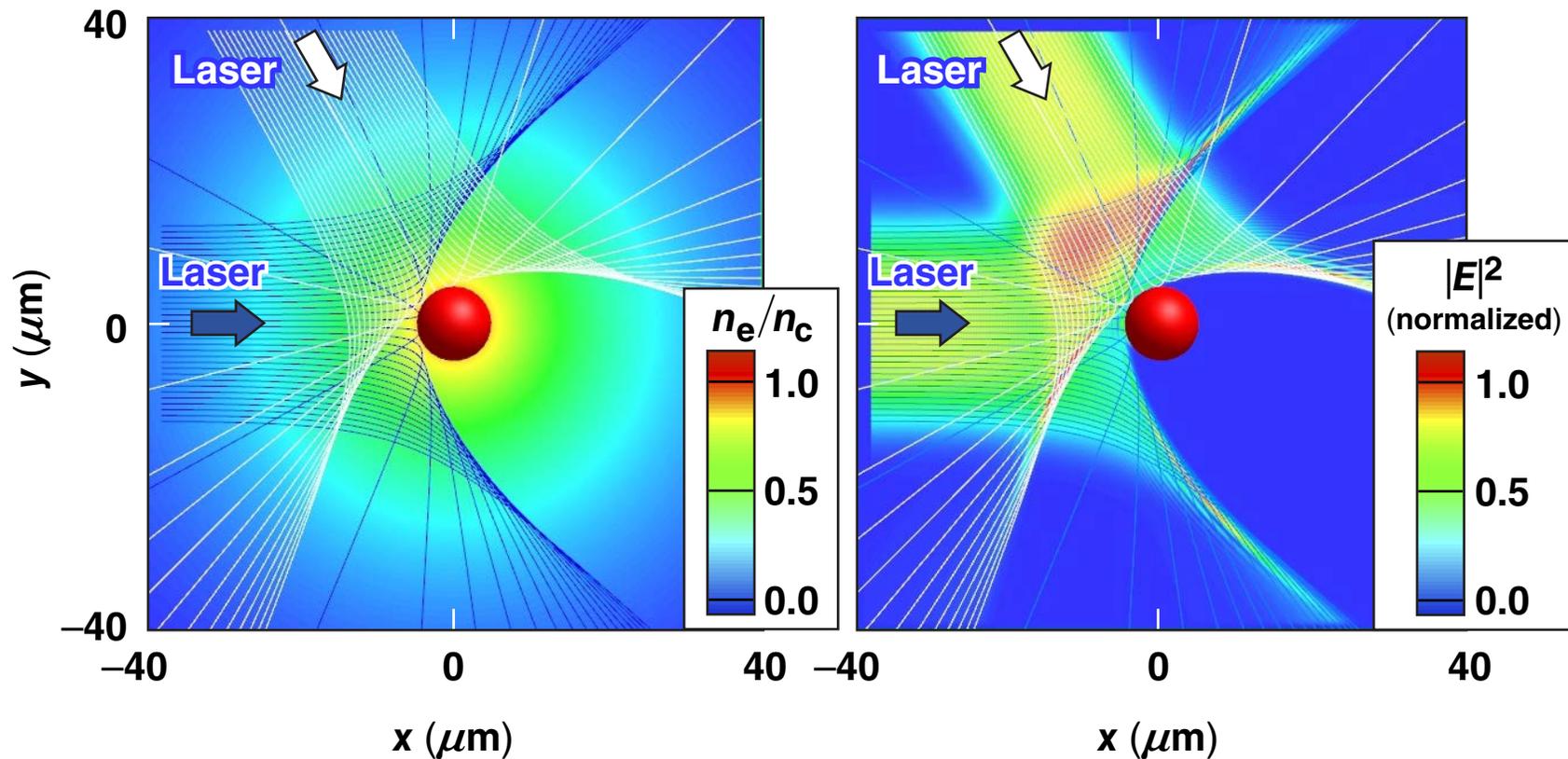


- This model solves the time-enveloped Maxwell equations in 3-D coupled to a fluid equation for the low-frequency ion-acoustic response
 - radiative boundary conditions, arbitrary incident beams
- A solid understanding of CBET has been obtained on OMEGA
 - coordinated program of theory, numerical simulations, and experiments
- CBET mitigation is a crucial part of LLE's 100-Gbar plan

Laser-beam propagation and energy deposition is computed in ICF* design codes using ray tracing



- Absent nonlinearity, the geometric-optics approximation is well justified based on the long plasma scale lengths
- Power is deposited based on collisional absorption of laser light



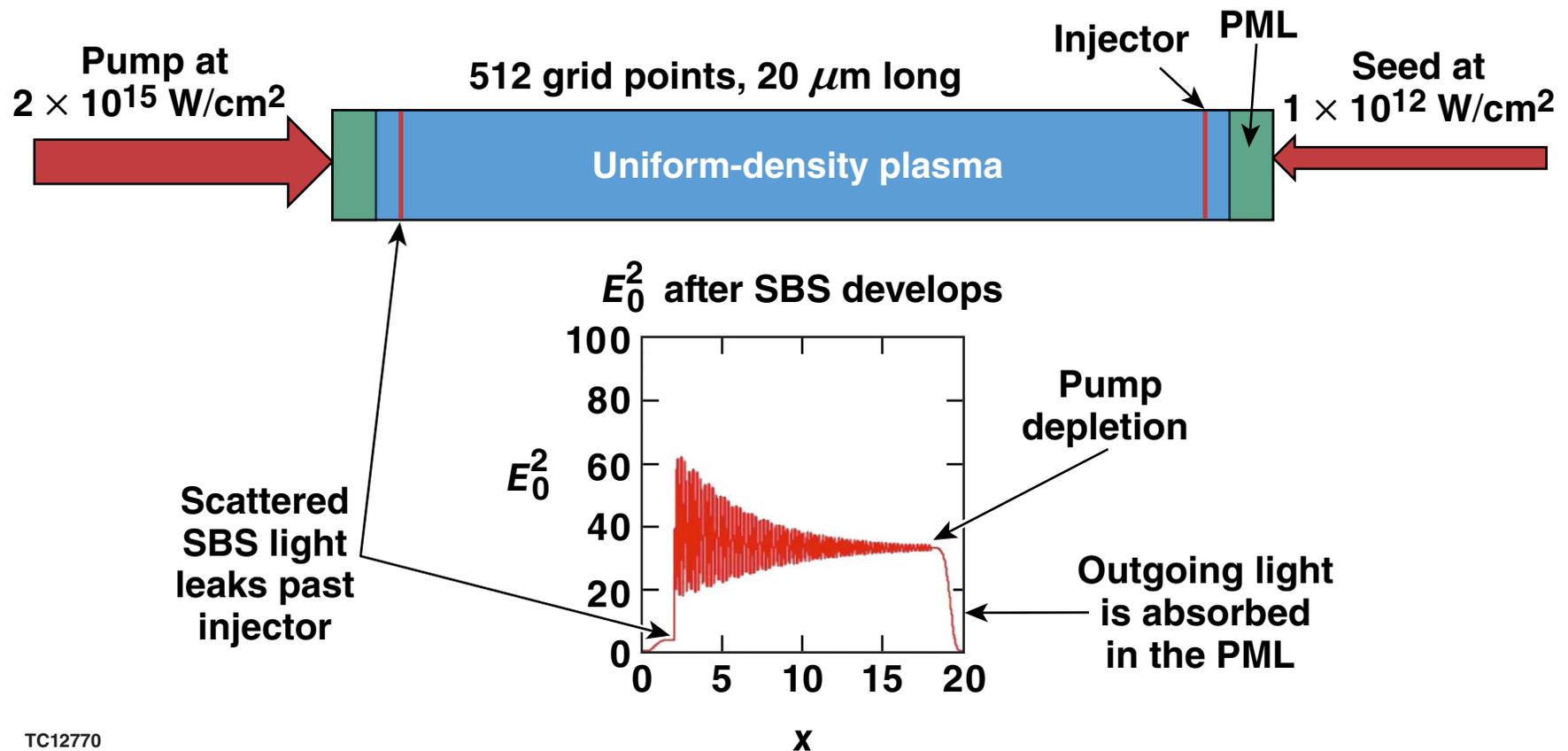
Wave-based models face several challenges and complications, but these can be overcome



- We are computing the motion of a semi-classical object by solving the (vector) Schrödinger equation
- Several key technical challenges were solved
 - time-enveloped (vector) wave equation in 3-D in strongly inhomogeneous plasma of a useful size
 - efficient algorithm is required (i.e., not Crank–Nicholson)
 - very complicated boundary conditions (in 3-D)
 - coupled to a plasma model
 - parallel efficiency [scalable solver for $O(10^9)$ computational cells]
- The resulting wave solver is practical to run in 3-D
 - 100 Intel cores, could scale to 1000's
 - pioneering use/visualization of large datasets/fast disks at LLE [$O(100\text{ GB})$ sets]

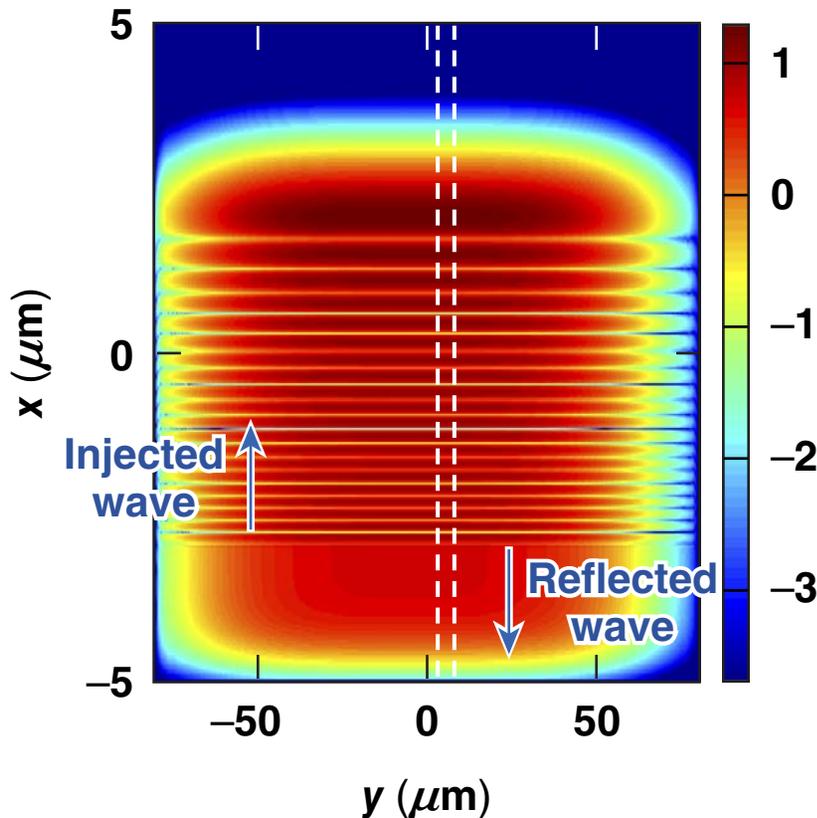
The boundary conditions use a “total-field/scattered-field” formulation together with a perfectly matched absorbing layer (PML)

- Inject a pump wave on the left-hand side and a (weak) seed wave on the right-hand side
- Match the stimulated Brillouin scattering (SBS) resonance condition for the seed by changing its frequency and/or adding a flow velocity to the plasma

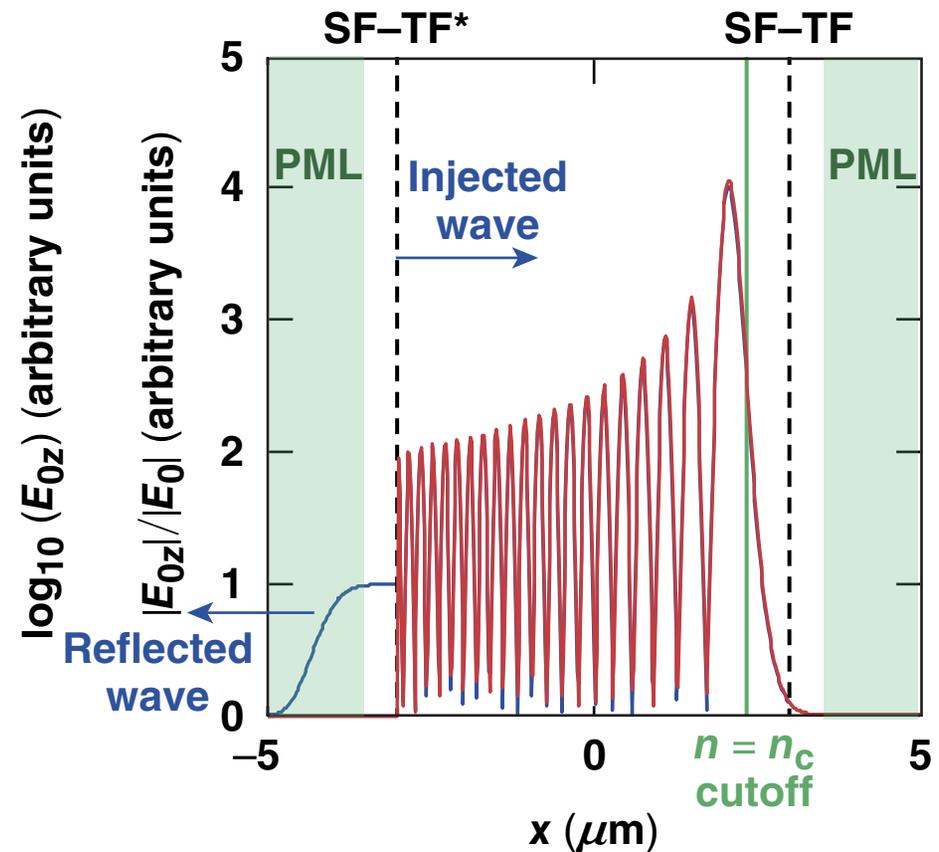


The *LPSE* electromagnetic (EM) wave solver reproduces analytical results for propagation in an inhomogeneous plasma

S-polarized at normal incidence
in linear-density gradient



Comparison with Airy pattern



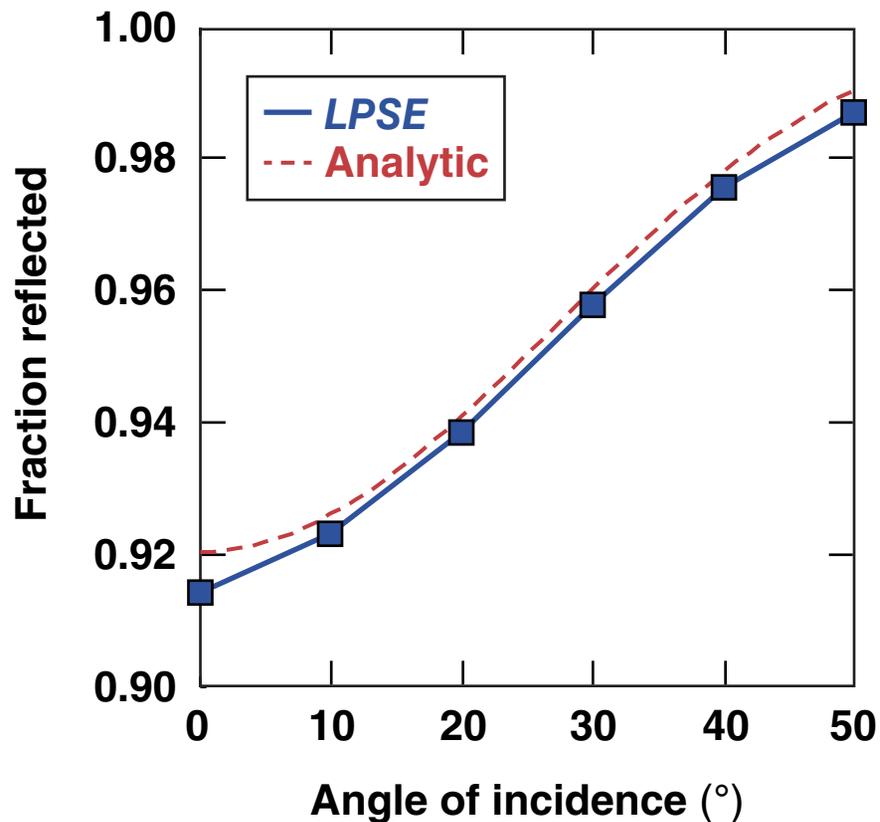
D. E. Merewether, R. Fisher, and F. W. Smith, IEEE Trans. Nucl. Sci. 27, 1829 (1980).
*scattered field–total field

Obliquely incident light turns at a lower density (shifted by $\cos^2\theta$ in a linearly varying density profile)

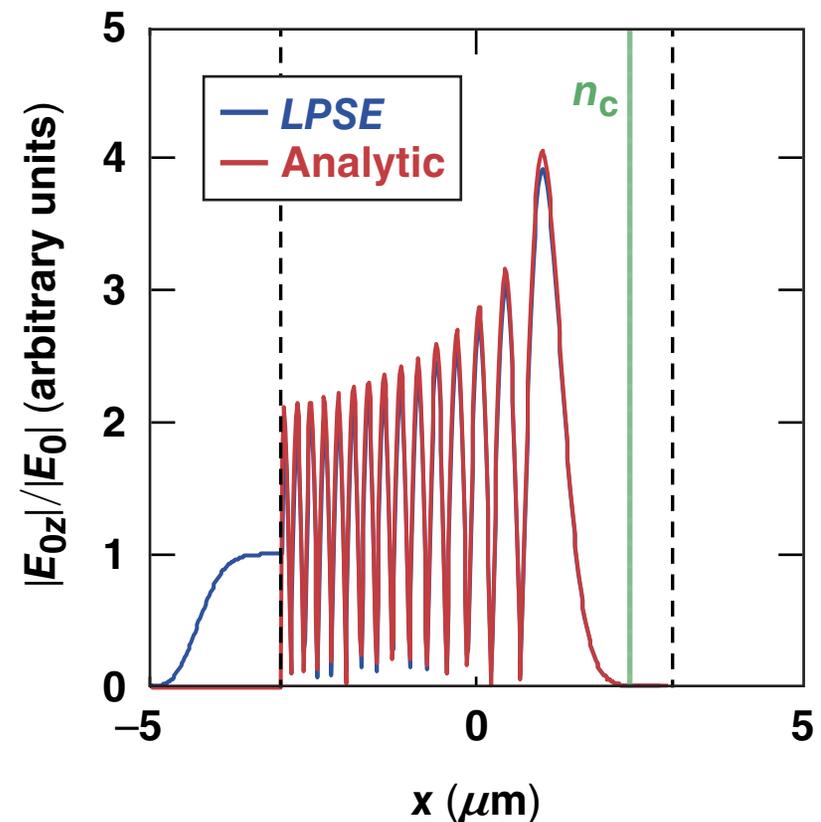
$$\frac{|E_y|}{|E_0|} \approx 2\sqrt{\pi} \left(\frac{\omega_0 L}{c}\right)^{1/6} A_i(\eta)$$

$$\eta \equiv \left(\frac{\omega_0^2 L}{c^2}\right)^{1/3} (x - L + L \sin^2 \theta)$$

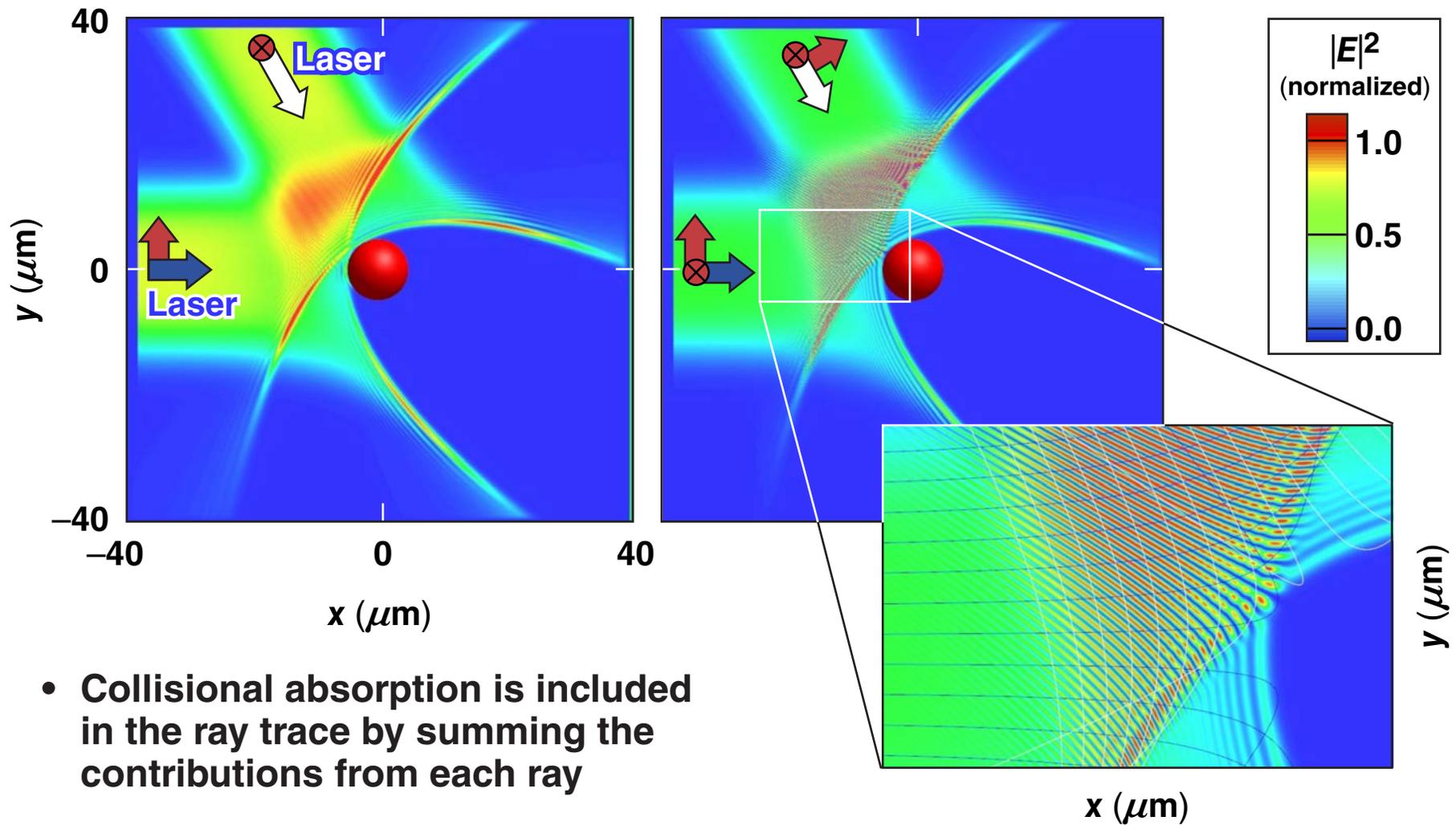
This results in a strong $\cos^5\theta$ angular dependence for absorption



Comparison with Airy pattern



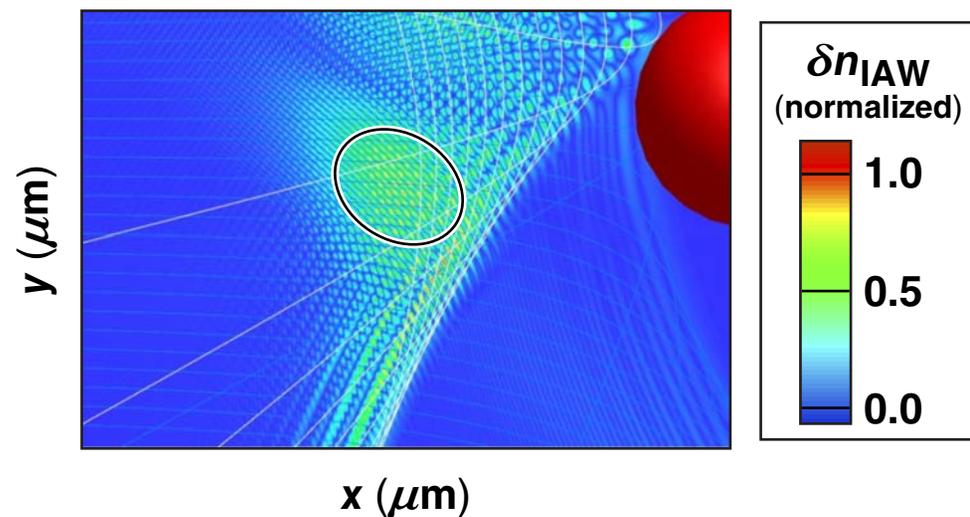
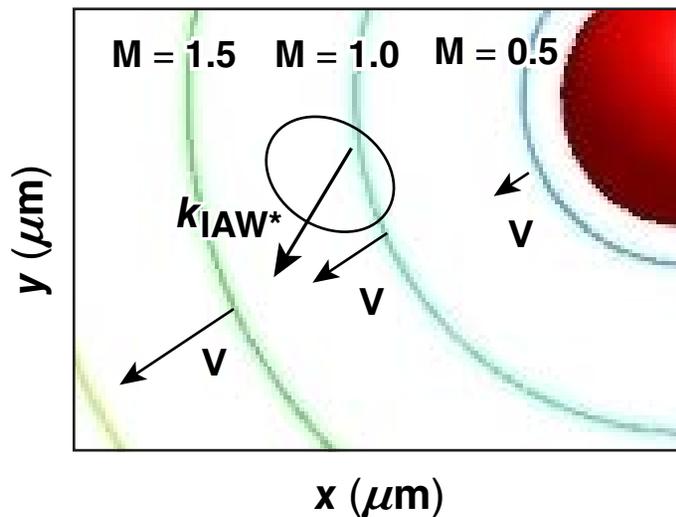
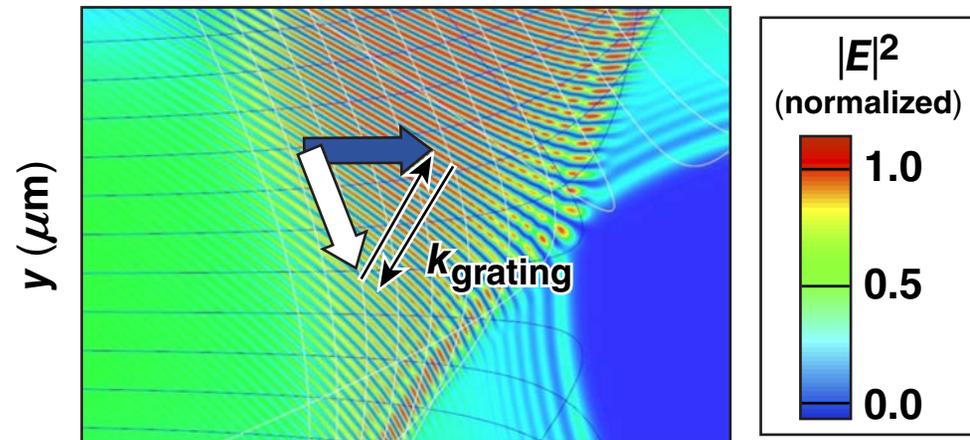
The laser light is partially coherent; the intensity is not the sum of the intensity of individual beams



- Collisional absorption is included in the ray trace by summing the contributions from each ray

The electric-field grating resonantly excites ion-acoustic waves because of the plasma-flow Doppler shift

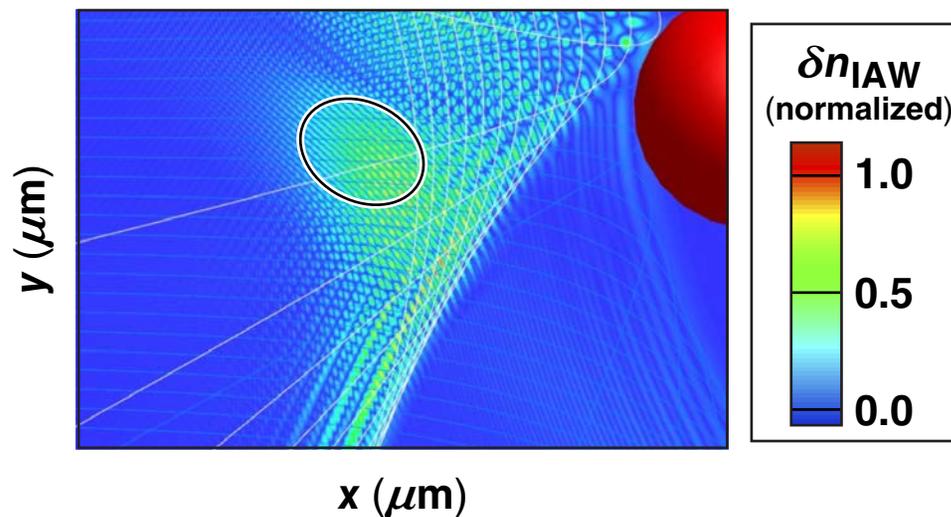
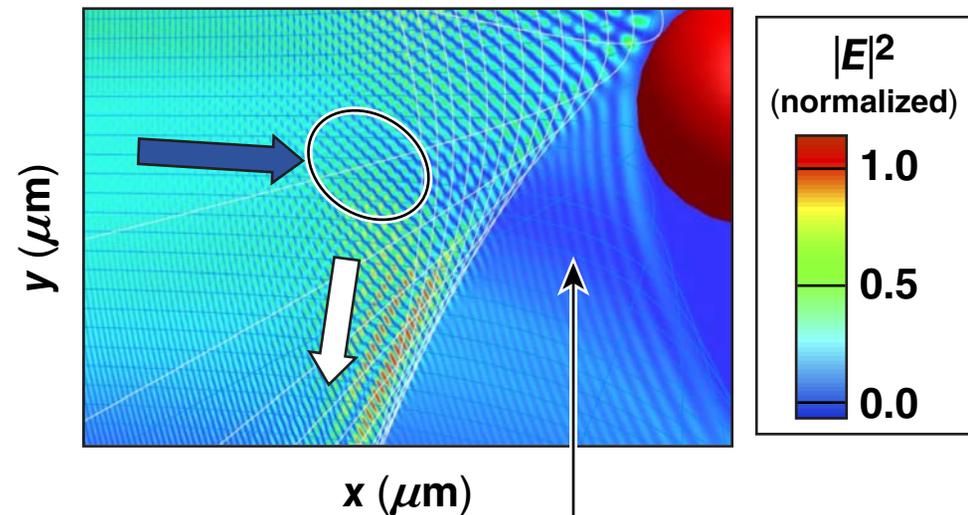
- If the two EM waves have equal frequencies, the ion-acoustic perturbation will be large if $k \cdot v = c_s$
- Mach number $M = |v|/c_s$



*ion-acoustic wave

This becomes an induced SBS process; laser energy changes direction

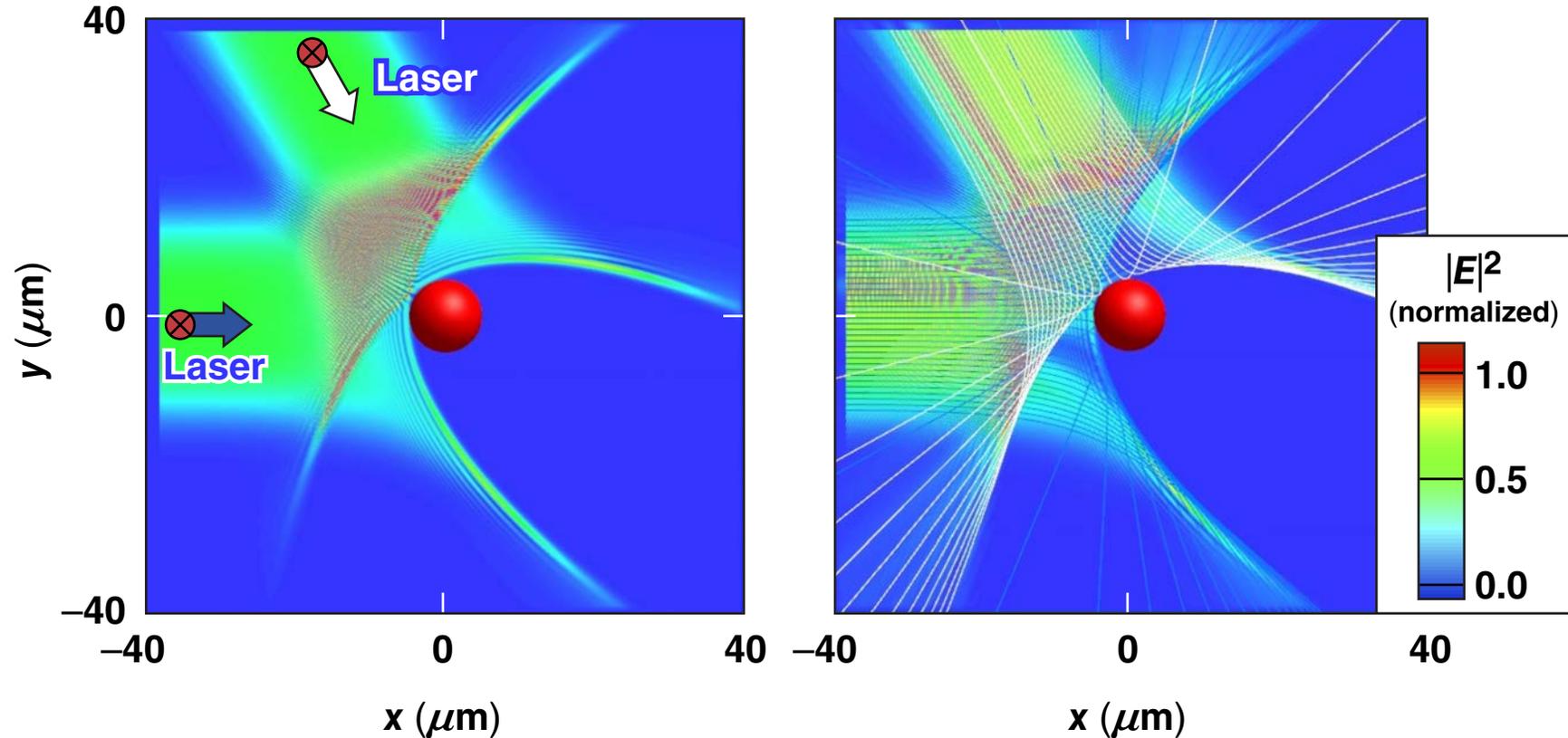
- Laser energy can be redirected before it has reached its turning point
- This leads to a reduction of absorption and hydrodynamic efficiency



CBET shadow

This effect can have a dramatic impact on laser absorption and the drive of an ICF target

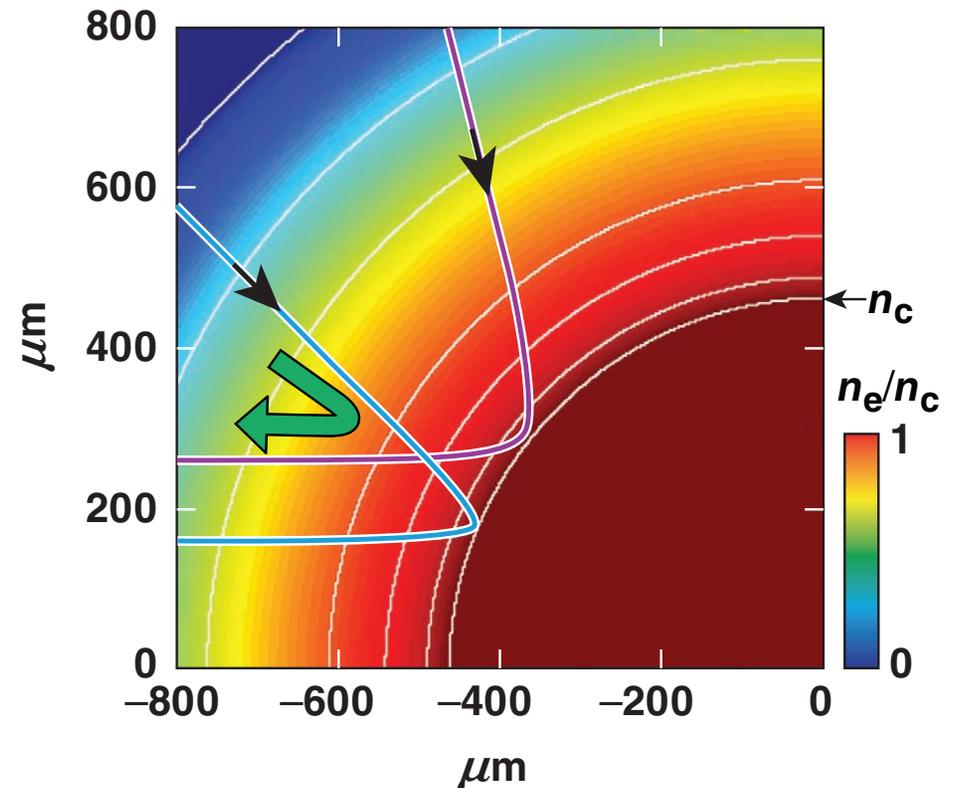
- It appears to operate in the regime of a convective amplifier for direct-drive ICF, which may be a tractable problem to describe with rays



All direct-drive CBET calculations have been performed using a 1-D description* that has been adapted to geometric ray tracing

- Unlike x-ray drive, the presence of supersonic plasma flow enables the process to be resonant*
- Three-wave SBS equations are computed (pairwise) for each beam crossing using a generalization of Randall *et al.** and are implemented in-line in 1-D *LILAC*

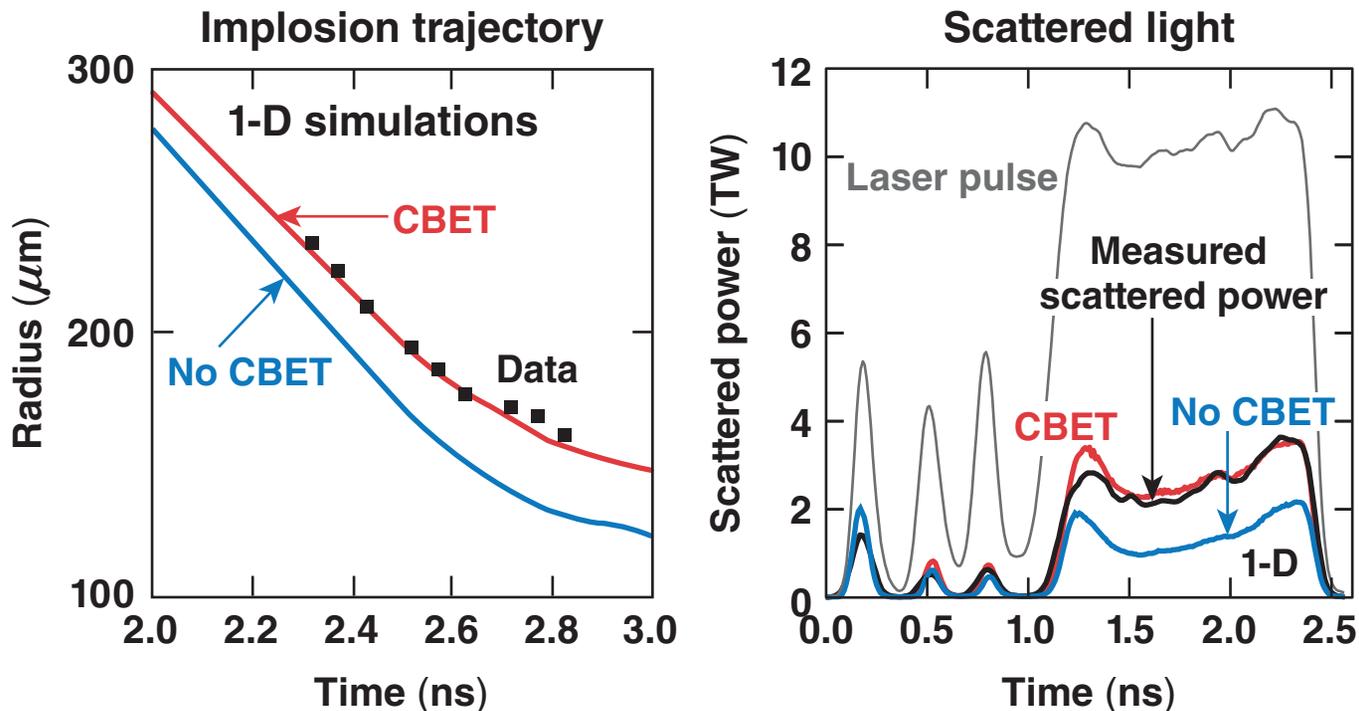
Because the EM seed amplitude is large, small gains affect the absorbed energy.



*C. J. Randall, J. R. Albritton, and J. J. Thomson, *Phys. Fluids* **24**, 1474 (1981);
K. B. Wharton *et al.*, *Phys. Rev. Lett.* **81**, 2248 (1998);
B. I. Cohen *et al.*, *Phys. Plasmas* **5**, 3408 (1998);
H. A. Rose and S. Ghosal, *Phys. Plasmas* **5**, 1461 (1998).

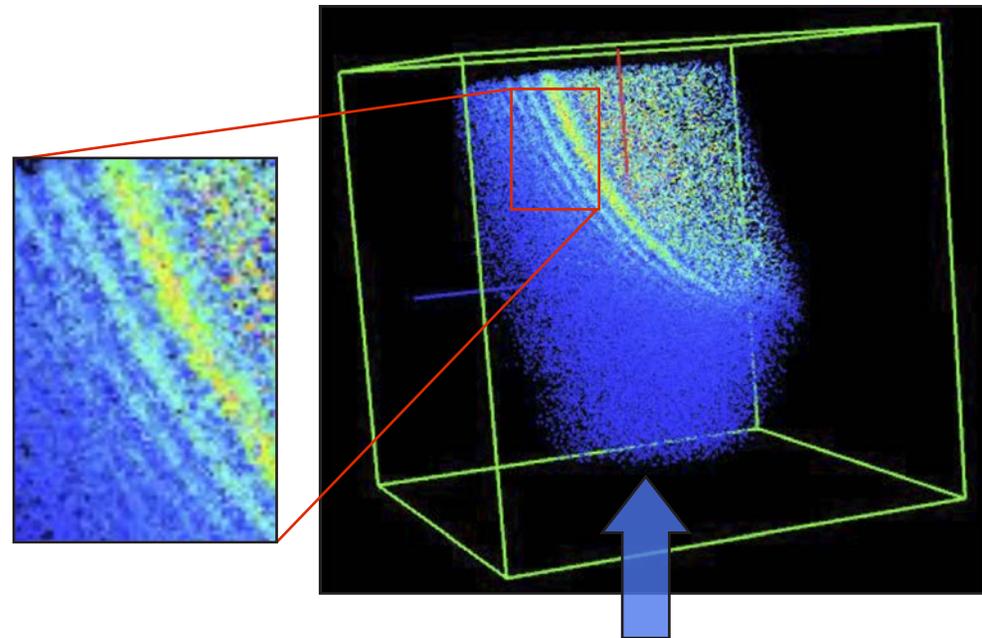
A nonlinear CBET model is required to obtain agreement between 1-D predictions and OMEGA experimental data

- The CBET model used to obtain agreement with $\alpha > 3.5$ data (not compromised by mix) is ray-based
 - scattered-light power and spectrum, shell trajectories, and mass ablation rates

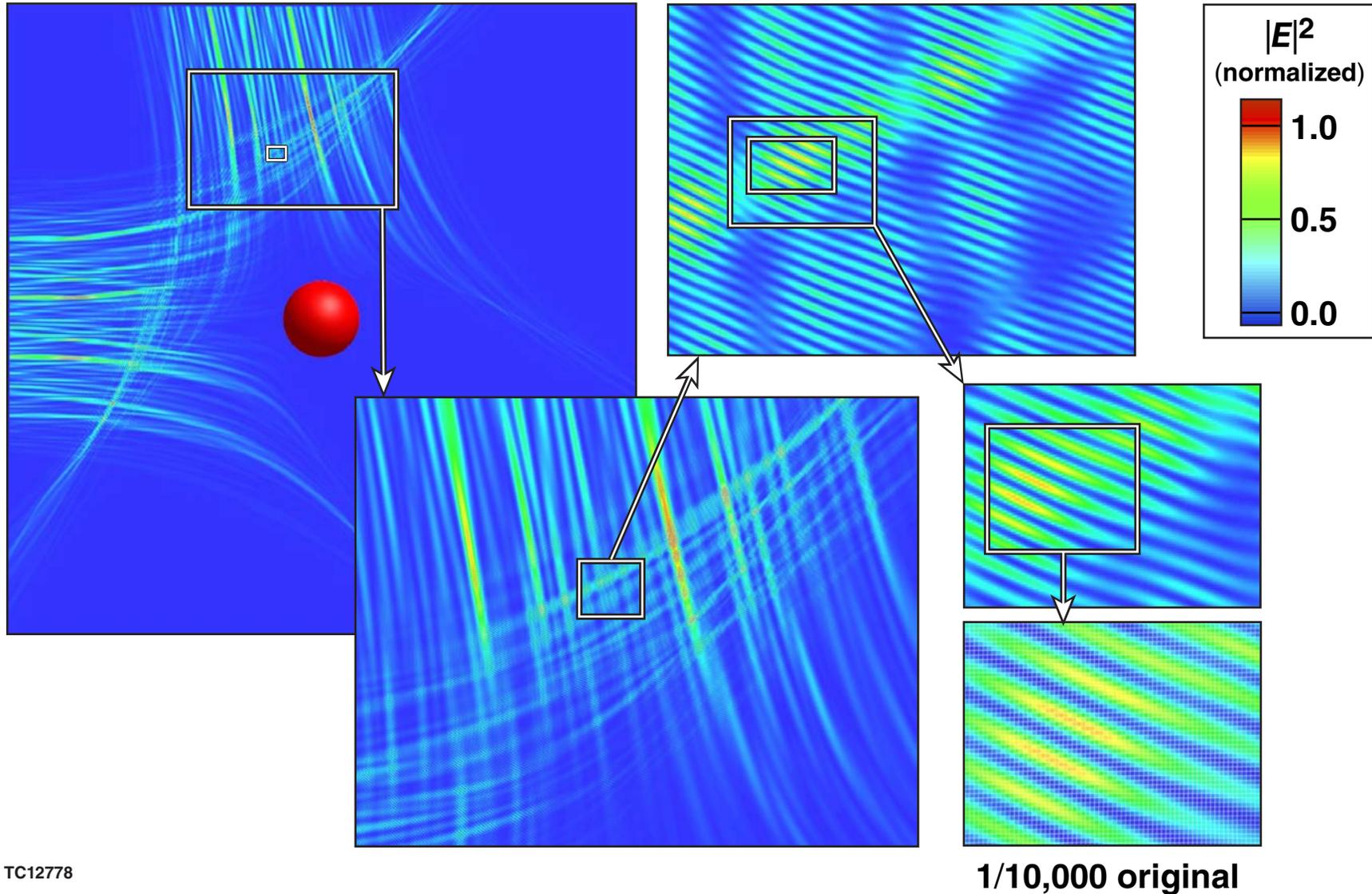


A wave-based CBET model is required for several important reasons

- There are uncertainties associated with ray-based CBET models that are hard to quantify without comparison with a more-fundamental model
- The model's correctness is empirically determined; however, experimental tests of CBET are integrated experiments (indirect)
- Caustic surfaces/turning points (field swelling, Airy-like patterns)
- Beam speckle (spatial and temporal incoherence)
- Polarization effects
- The IAW response is approximate in ray-based CBET (steady state, strong damping)



Spatial incoherence can be modeled with no difficulty; temporal incoherence is only slightly harder



At high enough laser intensities, CBET may not act as a simple spatial amplifier

- Shock-ignition experiments exceed filamentation thresholds

