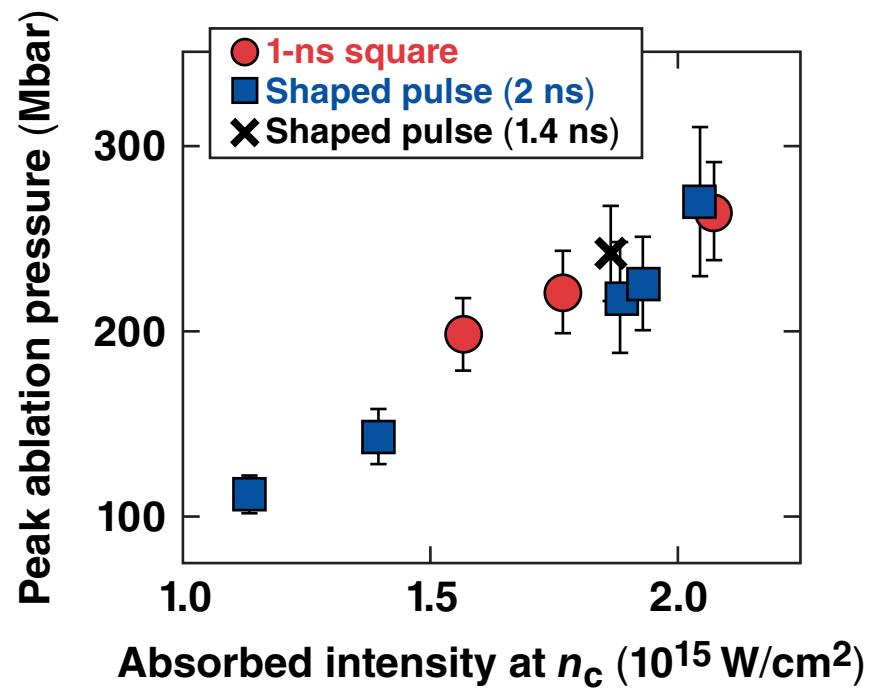
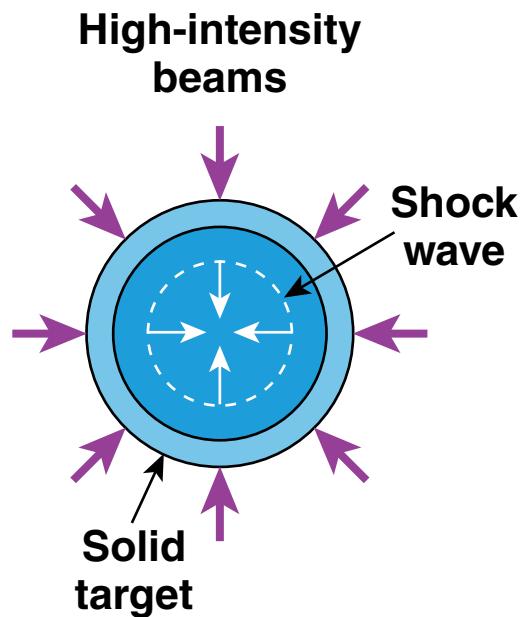


Strong-Shock Generation and Laser–Plasma Interactions for Shock-Ignition Inertial Fusion



W. Theobald
University of Rochester
Laboratory for Laser Energetics

44th Annual Anomalous
Absorption Conference
Estes Park, CO
8–13 June 2014

High-intensity shock-ignition (SI) experiments on OMEGA provide valuable data on hot-electron production and shock ablation pressure



- Multiple overlapping high-intensity beams generate an energetic hot-electron distribution (<100 keV) with a conversion efficiency of up to ~10%
- Turning off smoothing by spectral dispersion (SSD) in overlapping high-intensity beam experiments increased the hot-electron fraction and shock strength
- Stimulated Raman scattering (SRS) increases significantly (~5×) when SSD is turned off, while two-plasmon–decay (TPD) is unaffected
- Moderate hot electrons may be beneficial to shock ignition by coupling energy to the outer layer of the imploding capsule

The inferred ablation pressures of 270 Mbar at $\sim 3 \times 10^{15} \text{ W/cm}^2$ approaches the minimum SI requirement of 300 Mbar.

Collaborators



R. Nora,* M. Lafon, K. S. Anderson, M. Hohenberger, F. J. Marshall, D. T. Michel,
T. C. Sangster, W. Seka, A. A. Solodov, C. Stoeckl, B. Yaakobi, and R. Betti*

University of Rochester
Laboratory for Laser Energetics

*also Departments of Mechanical Engineering and Physics

A. Casner and C. Reverdin

CEA, DAM, DIF
Arpajon, France

X. Ribeyre and A. Vallet

CELIA
University of Bordeaux, France

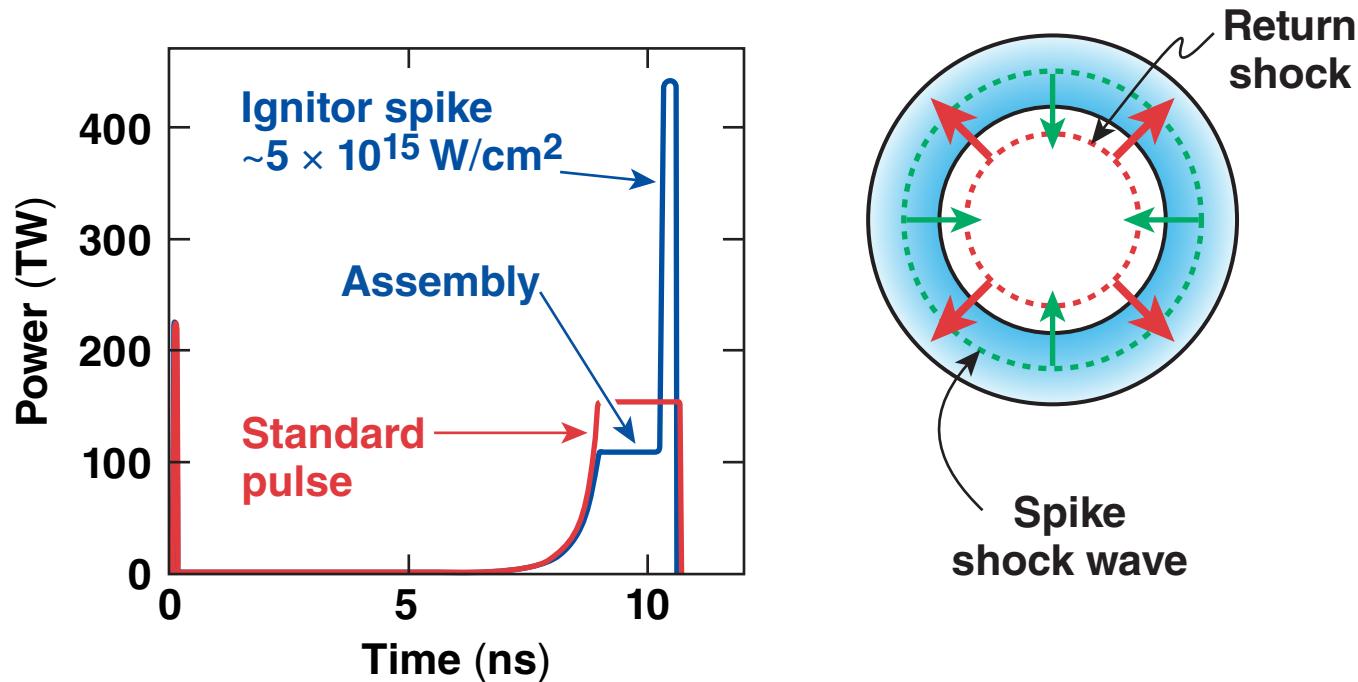
J. Peebles

University of California, San Diego
La Jolla, CA

M. S. Wei

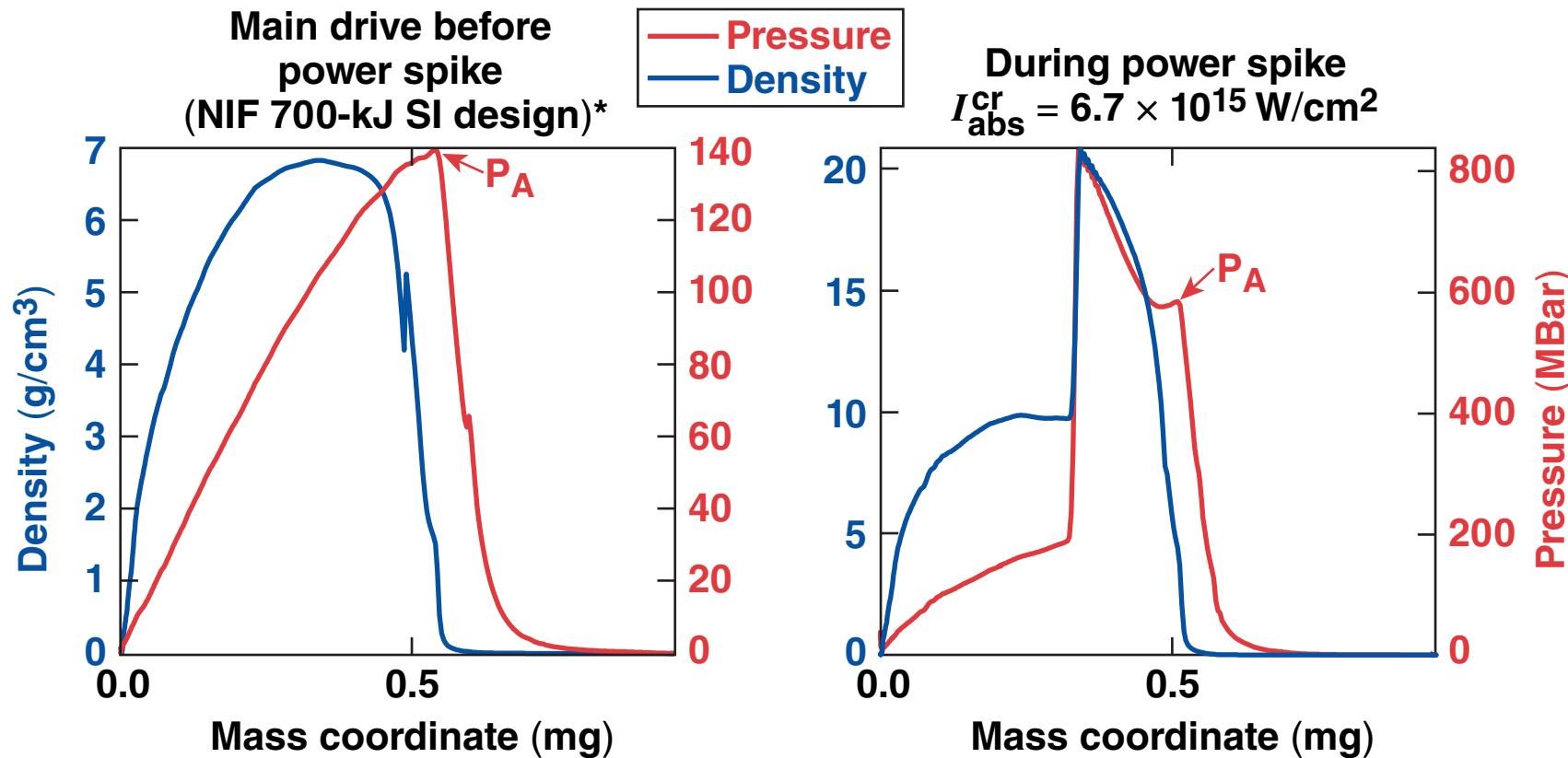
General Atomics
San Diego, CA

Sufficient ablation pressure and $T_{\text{hot}} < 150 \text{ keV}$ must be demonstrated for shock ignition during the ignitor spike



- Critical issues for shock ignition
 - demonstrate ~300- to 400-Mbar spike-generated ablation pressure
 - demonstrate hot-electron temperatures of $\leq 150 \text{ keV}$ generated by spike

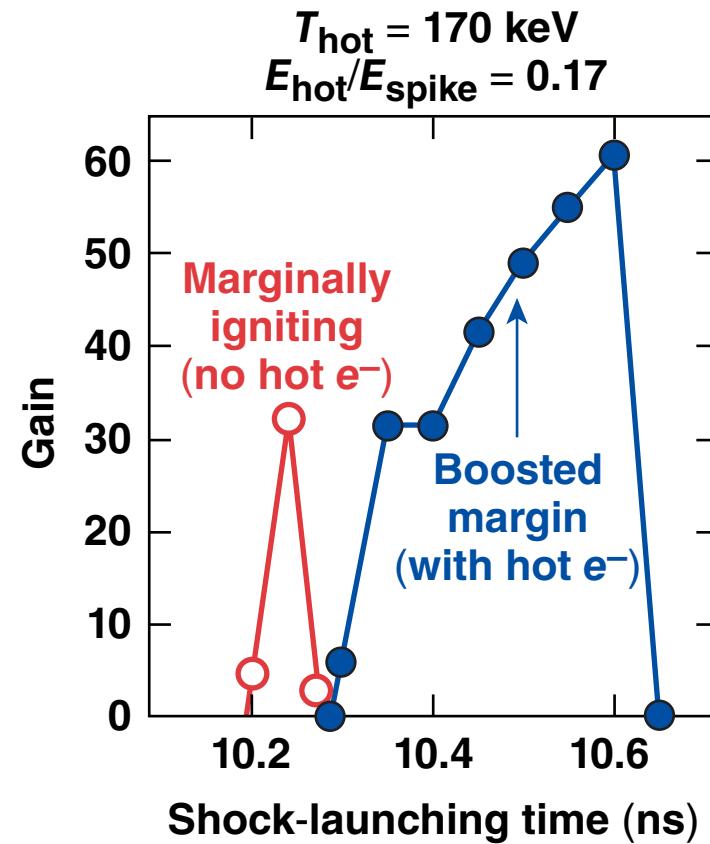
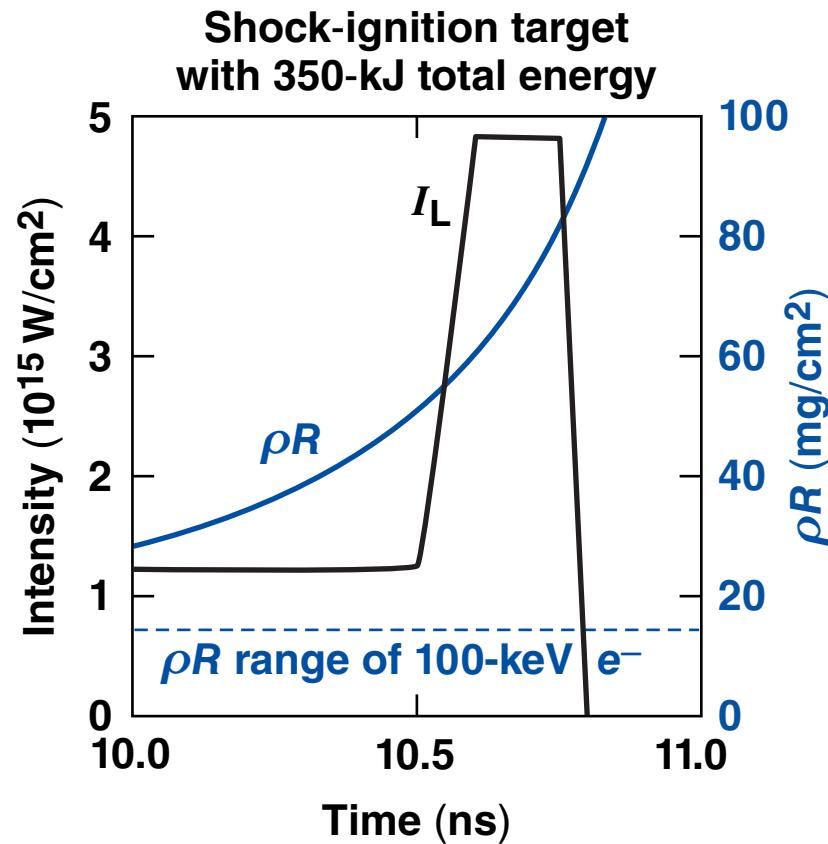
SI requires minimum ablation pressures of ~300 Mbar; up to ~600 Mbar is required for large ignition margins



Minimum ablation pressure for SI $\sim 2 \times (P_A \text{ main drive}) \sim 300 \text{ Mbar}$.

Ablation pressure for the National Ignition Facility (NIF) 700-kJ polar-drive (PD) SI design with an ignition threshold factor (ITF) = 4 $\sim 600 \text{ Mbar}$.

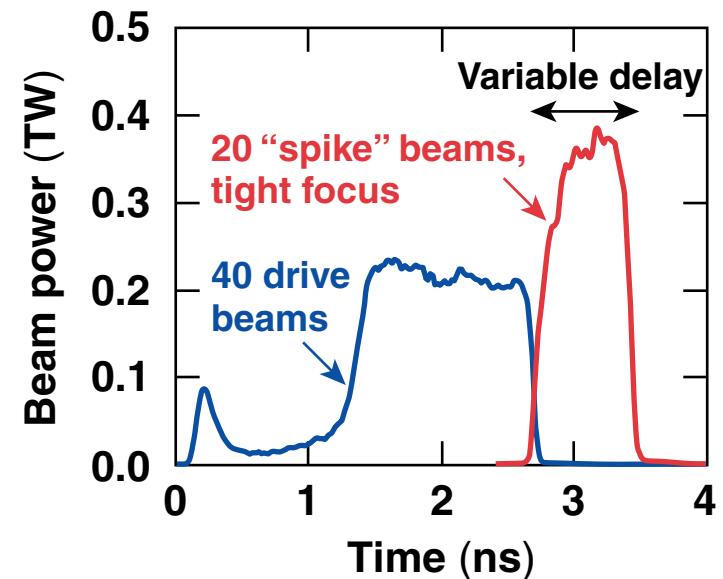
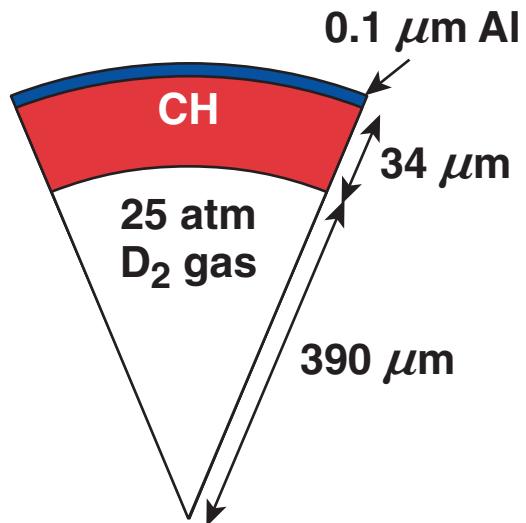
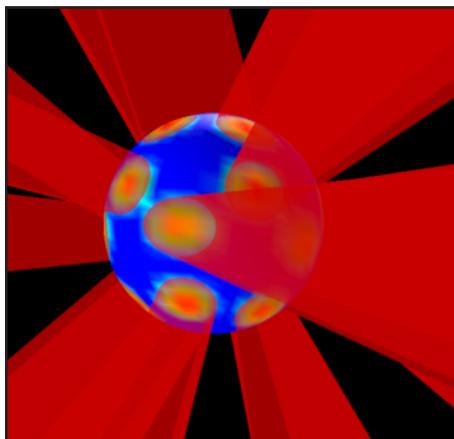
Laser–plasma interaction during the spike pulse and hot-electron generation are important issues for shock ignition



If the ρR is high enough when hot electrons are produced, they will be stopped in the outer regions of the shell.

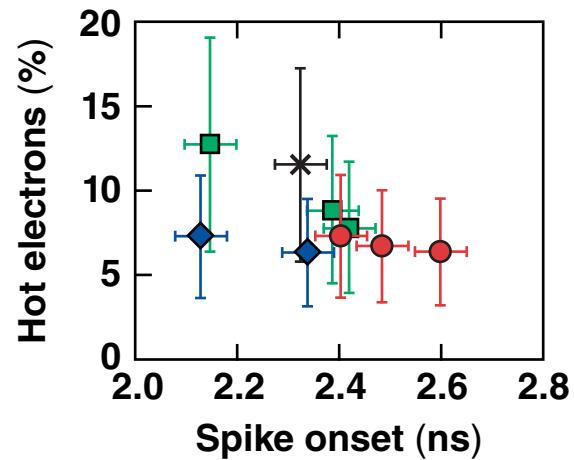
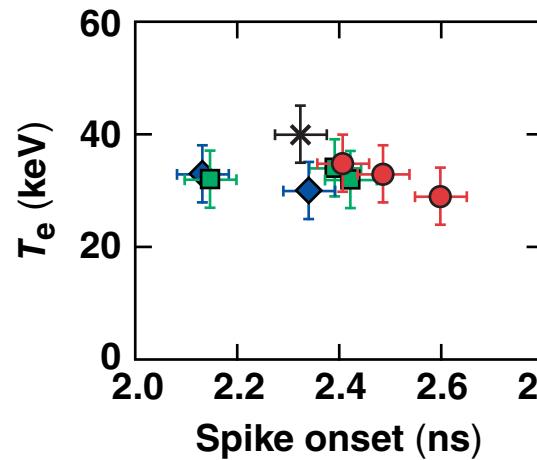
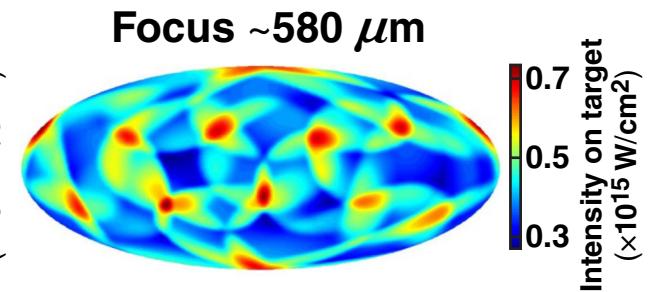
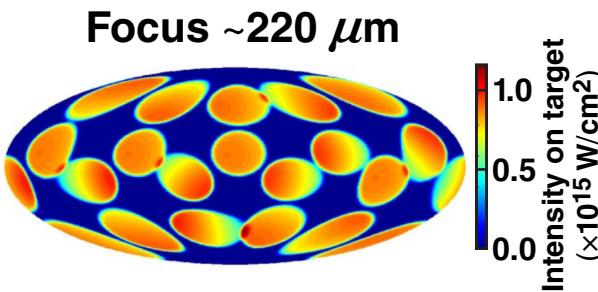
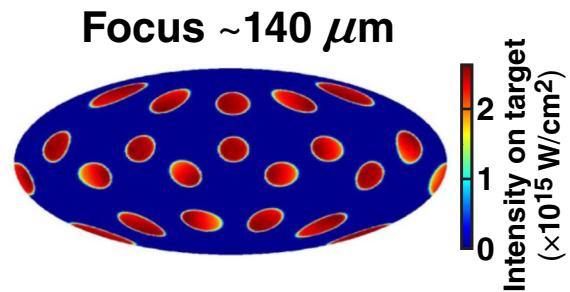
R. Betti et al., J. Phys., Conf. Ser. **112**, 022024 (2008);
S. Gus'kov et al., Phys. Rev. Lett. **109**, 255004 (2012).

The 40 + 20 configuration was used to study the interaction with single high-intensity beams on OMEGA



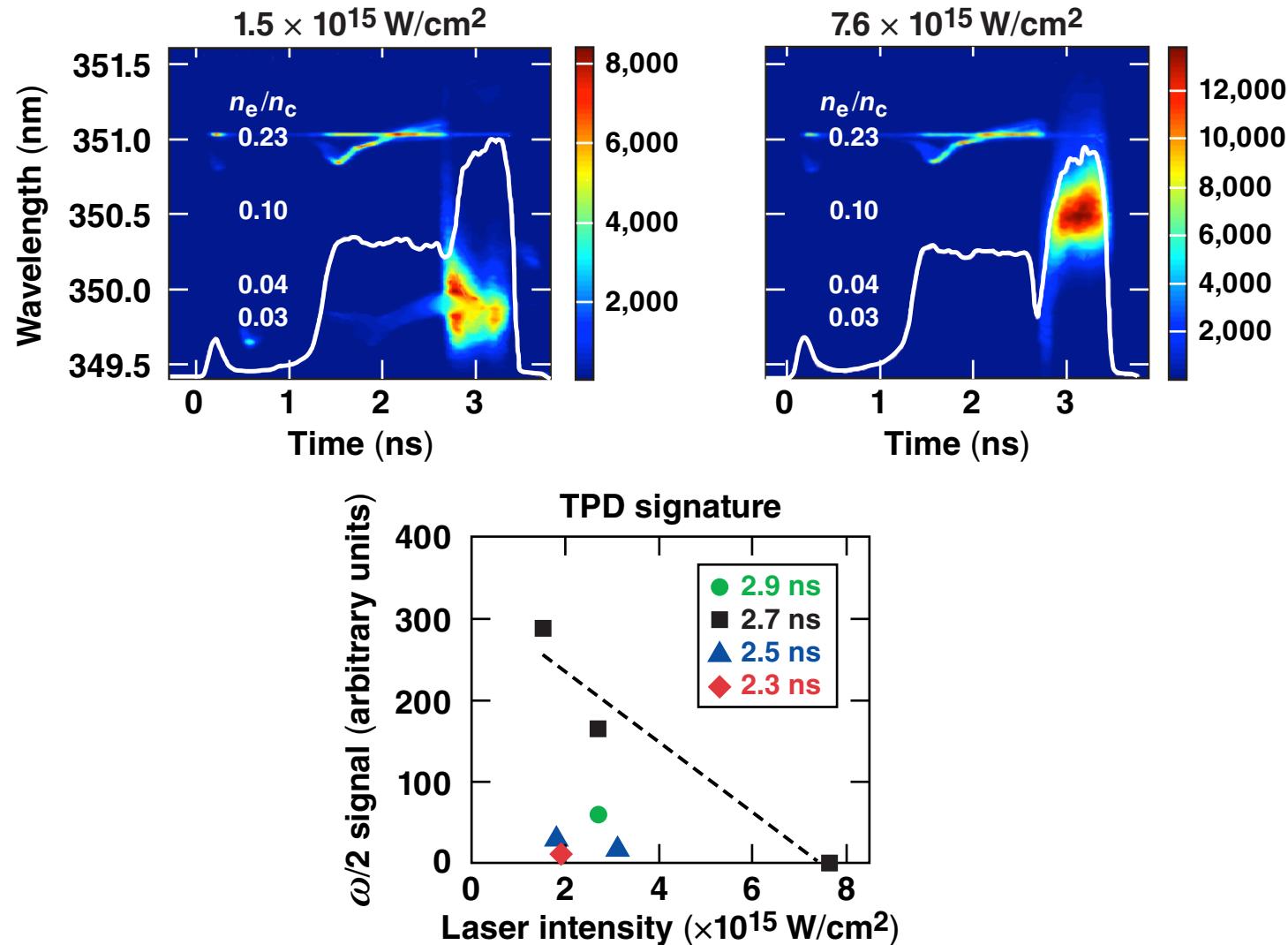
- Single-beam shock-ignition studies in spherical geometry up to $\sim 8 \times 10^{15} \text{ W/cm}^2$
- 40 drive beams: distributed phase plate (DPP), distributed phase rotator (DPR) (no SSD), $\sim 14 \text{ kJ}$
- 20 spike beams: no DPP, $\sim 5 \text{ kJ}$

Up to ~12% of the spike energy was converted into a moderate (~30-keV) hot-electron distribution



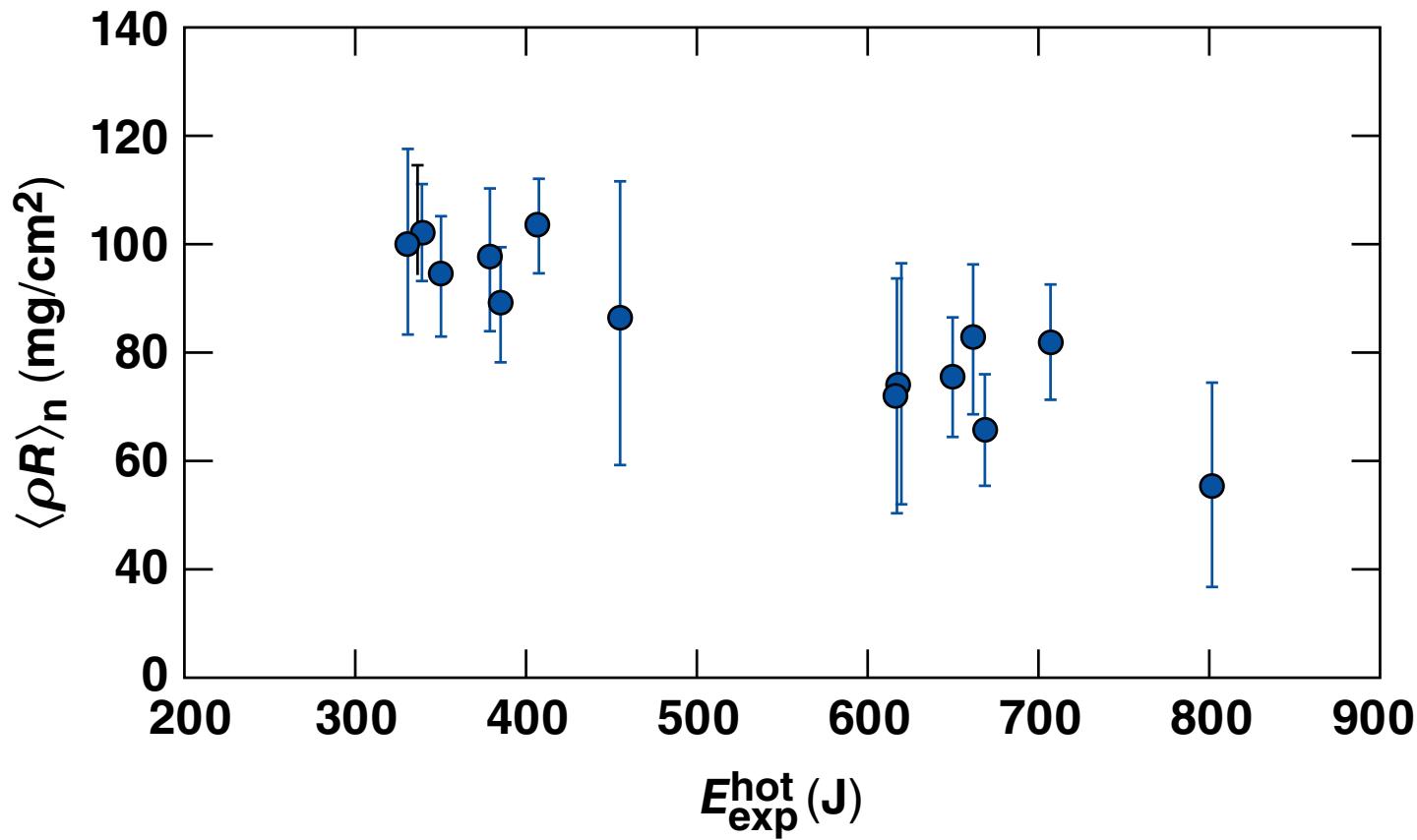
- Hot-electron generation in the 40 + 20 experiment is dominated by SRS
- The TPD instability is strongly suppressed

At high laser intensity, stimulated Brillouin scattering (SBS) signal is emitted near $n_c/4$ and TPD is suppressed



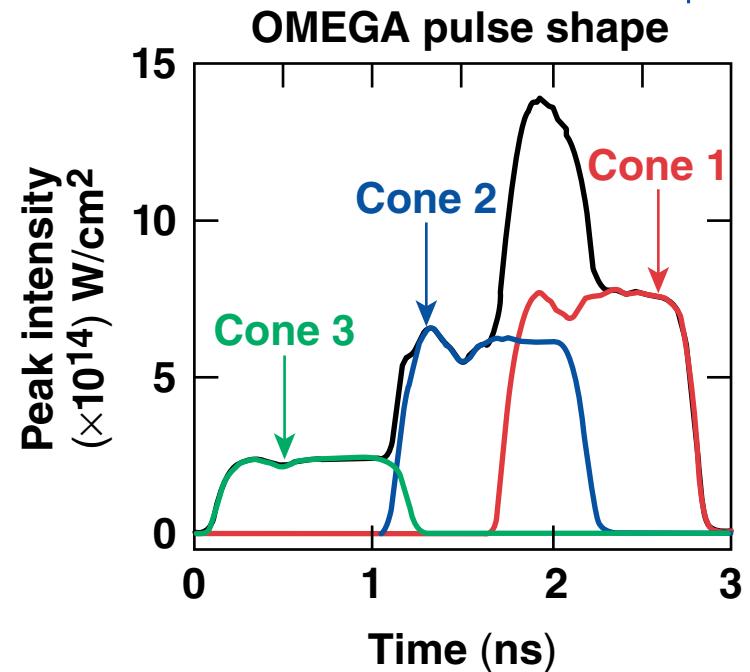
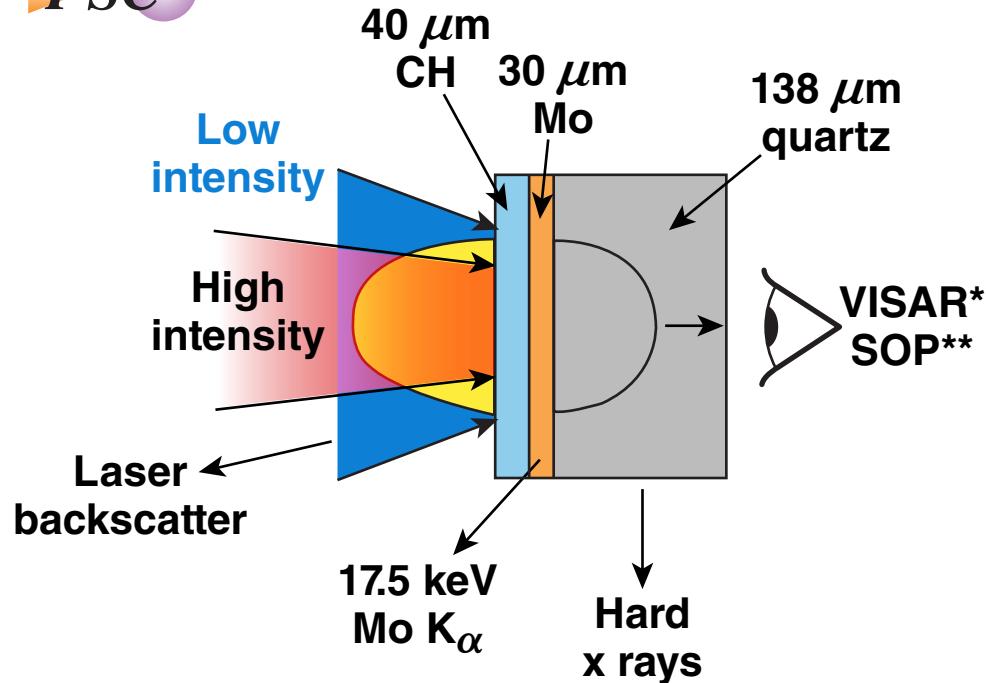
E20064a

Areal-density measurements show implosion degradation with increasing hot-electron production



The areal density of OMEGA targets is too low to stop hot electrons on the ablation surface.

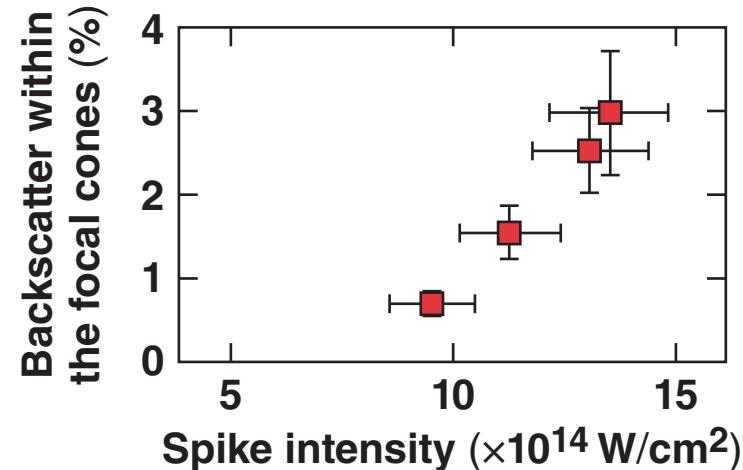
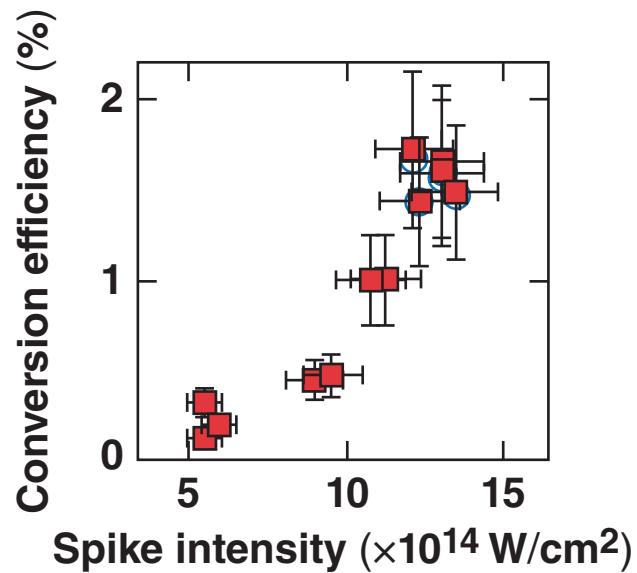
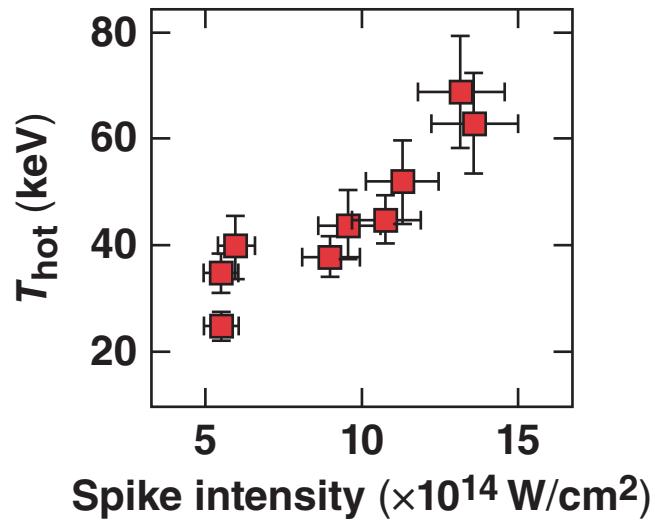
A laser-plasma interaction experiment was performed in planar geometry with overlapping beams to infer ablation pressure and to study hot-electron generation



- Phase plates and DPR's with $\sim 900\text{-}\mu\text{m}$ focal spots were used in plasma-generating beams (cone 2 and cone 3)
- Phase plates with an $\sim 600\text{-}\mu\text{m}$ focal spot were used in six high-intensity beams (cone 1)
- $L_n \sim 350\text{ }\mu\text{m}$

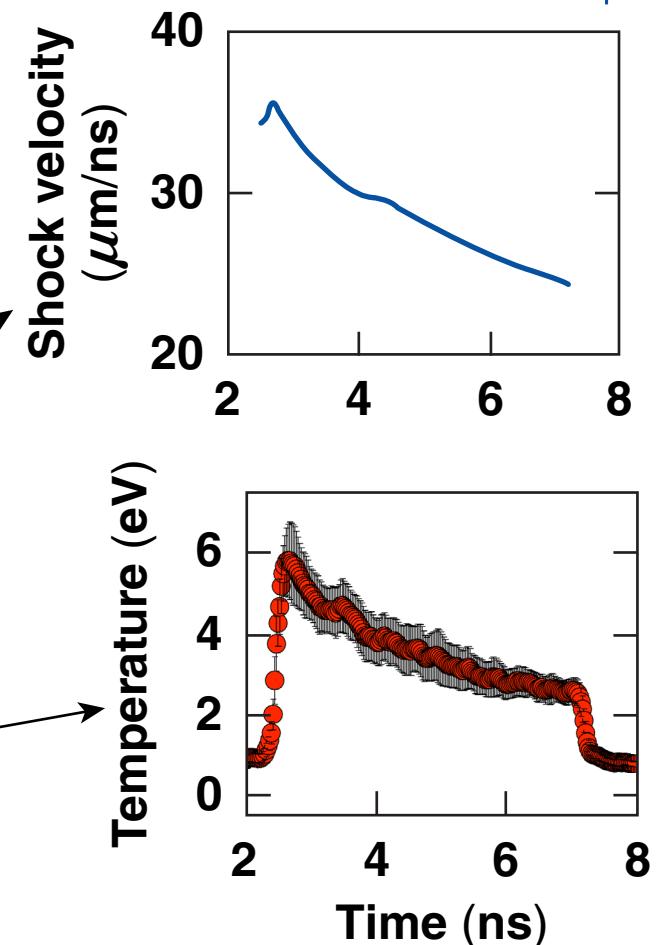
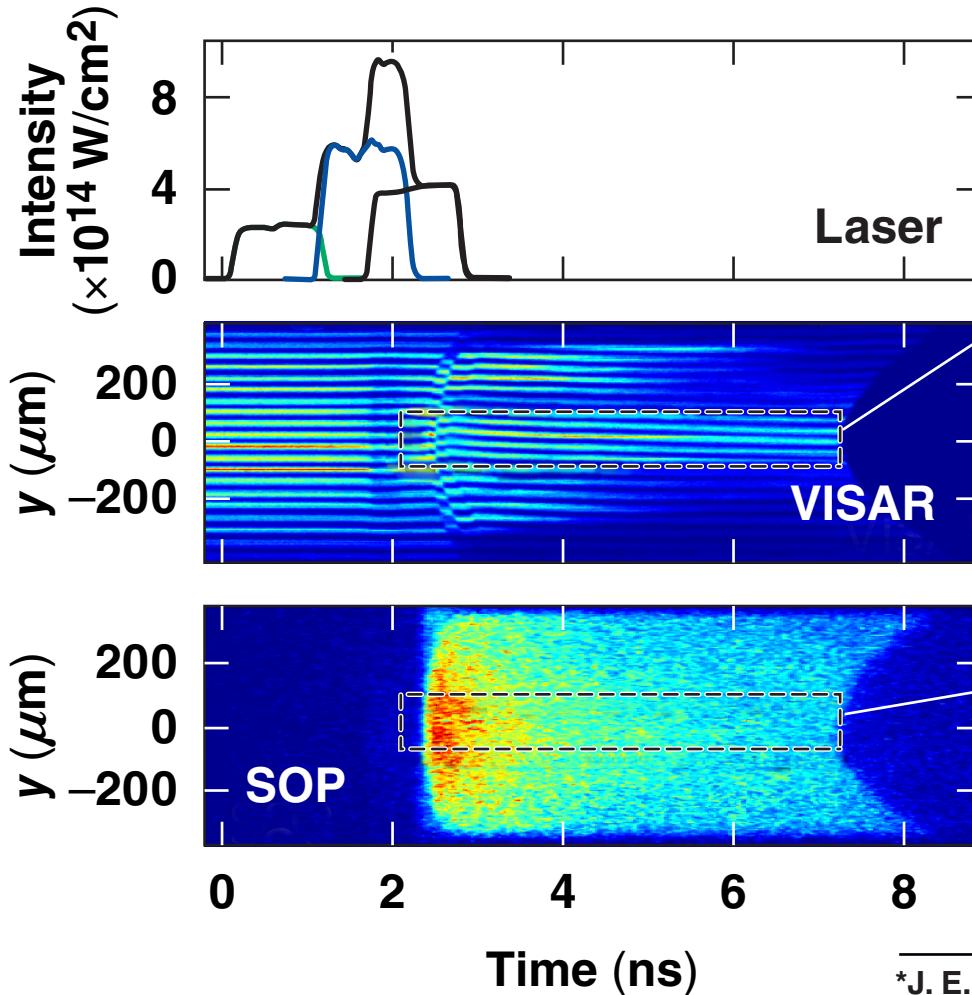
M. Hohenberger et al., Phys. Plasmas 21, 022702 (2014).
* Velocity interferometer system for any reflector
** Streaked optical pyrometer

The hot-electron temperature and hot-electron fraction increase with spike laser intensity



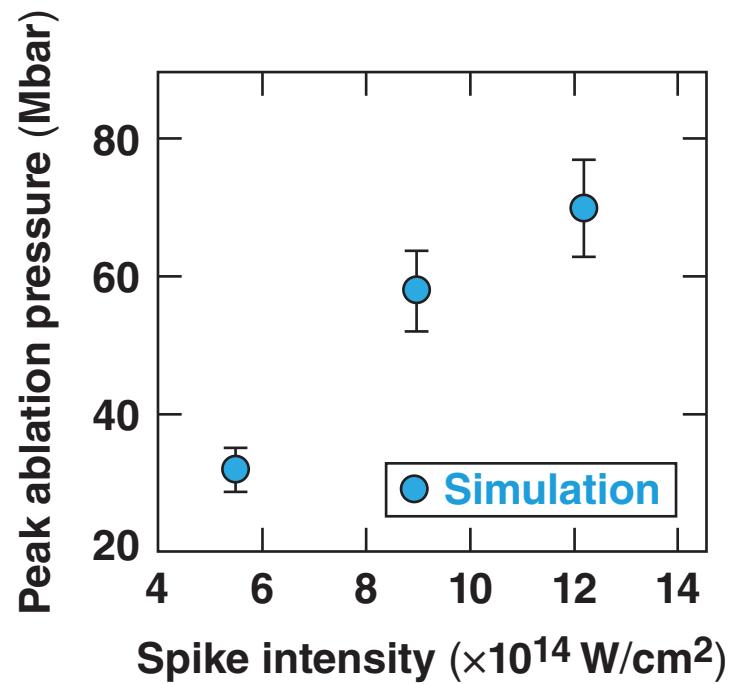
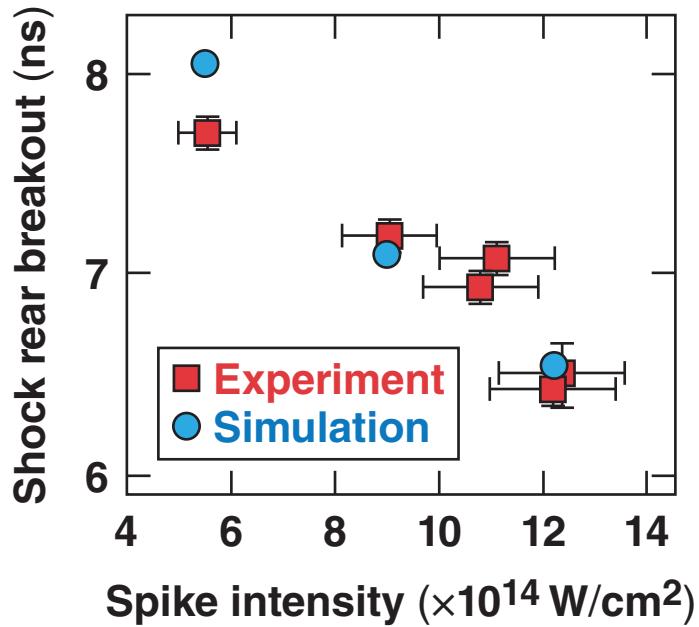
E20783

The shock propagation in quartz was observed with SOP and VISAR*



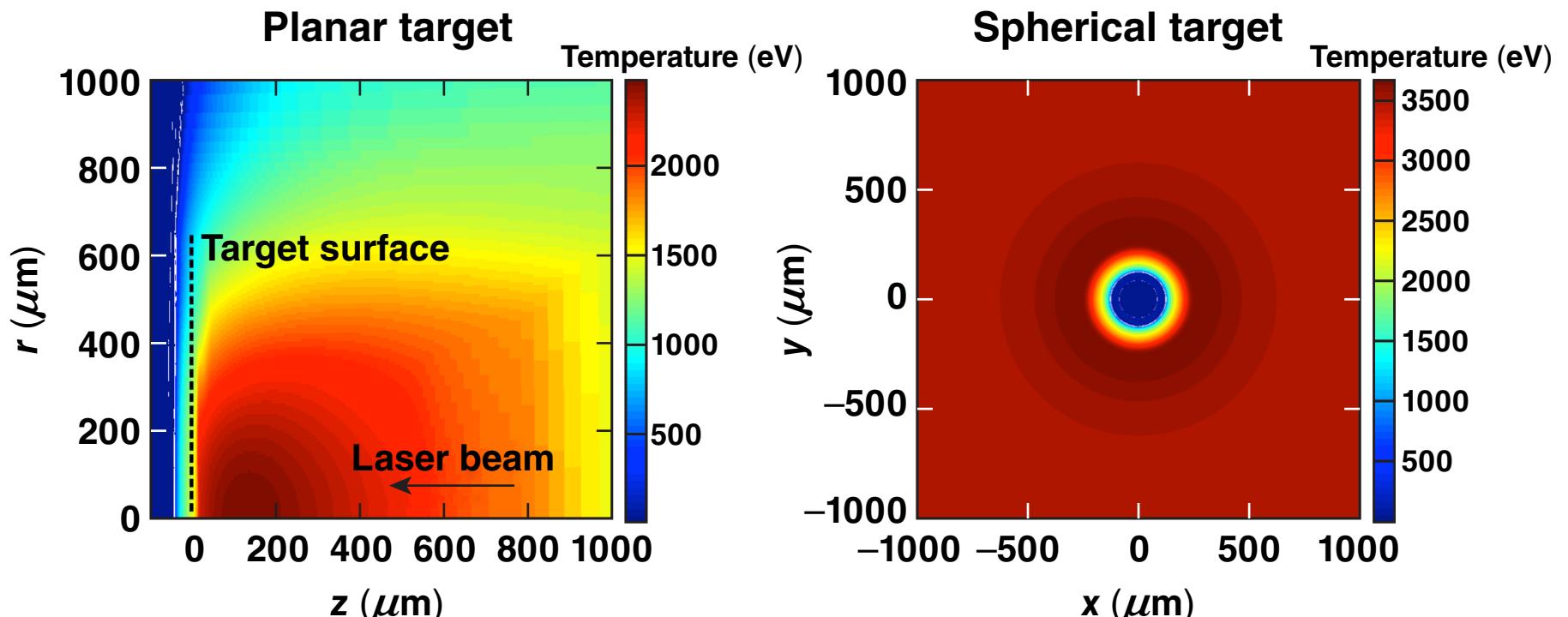
*J. E. Miller et al., Rev. Sci. Instrum. **78**, 034903 (2007);
P. M. Celliers et al., Rev. Sci. Instrum. **75**, 4916 (2004);
M. Hohenberger et al., Phys. Plasmas **21**, 022702 (2014).

Two-dimensional DRACO simulations reproduce well the shock dynamics over a range of spike intensities



- The ablation pressure reaches ~ 75 Mbar at 1.2×10^{15} W/cm 2

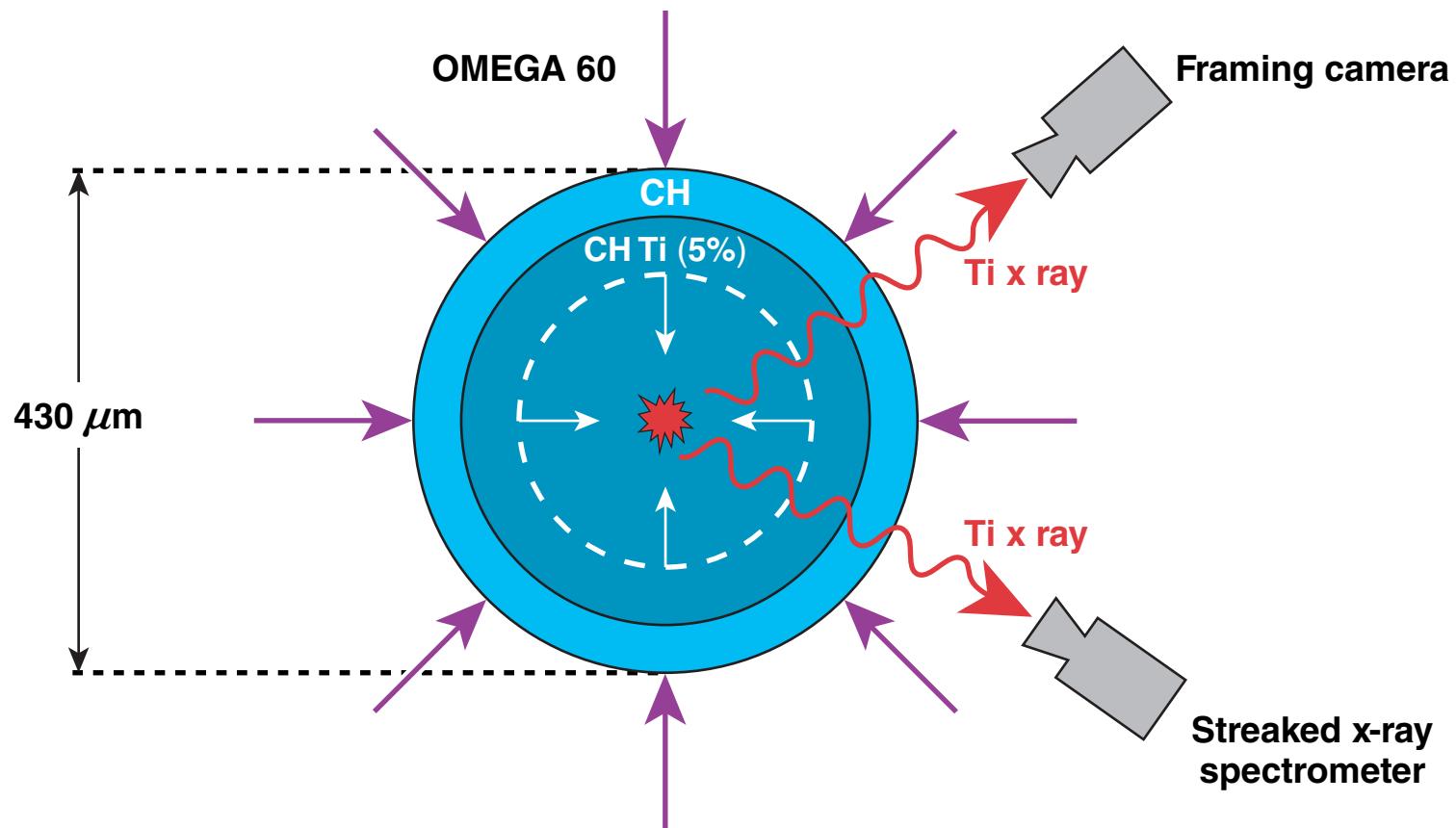
Ablation pressures in planar target experiments are limited by lateral heat losses



$$\text{Heat flux: } \mathbf{Q} = -K_{\text{th}} \nabla T$$

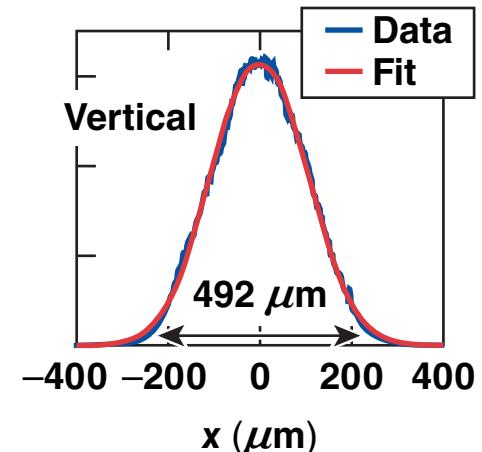
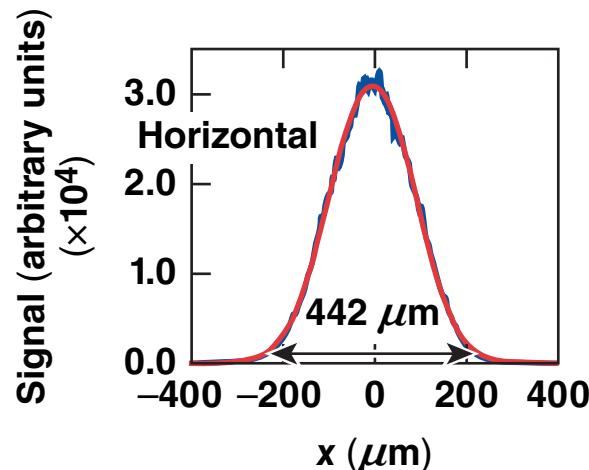
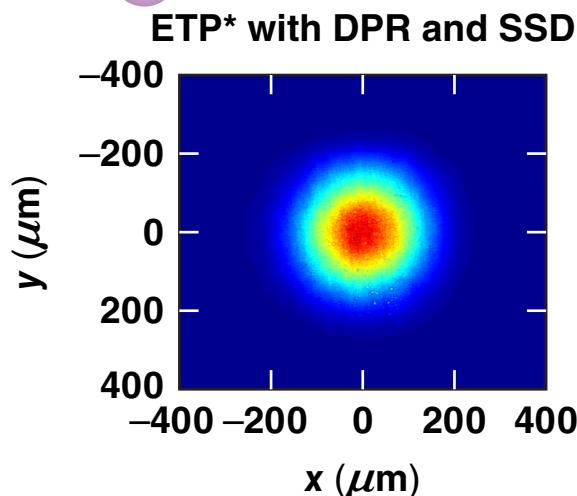
Higher ablation pressures can be achieved in spherical geometry.

A new OMEGA platform has been developed to study the generation of strong shocks in spherical targets

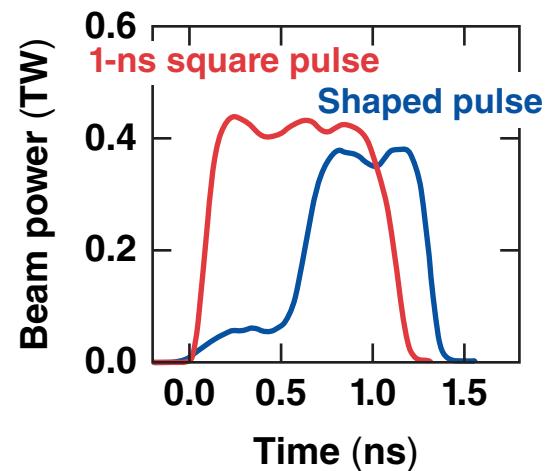


E22457a

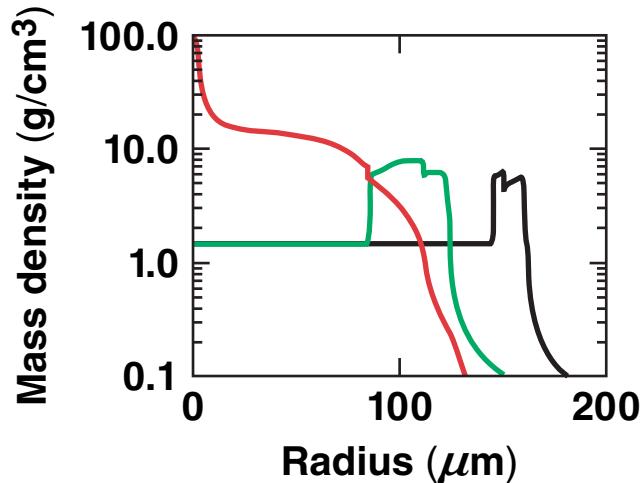
The small solid target was irradiated by 60 high-intensity beams equipped with small-spot phase plates



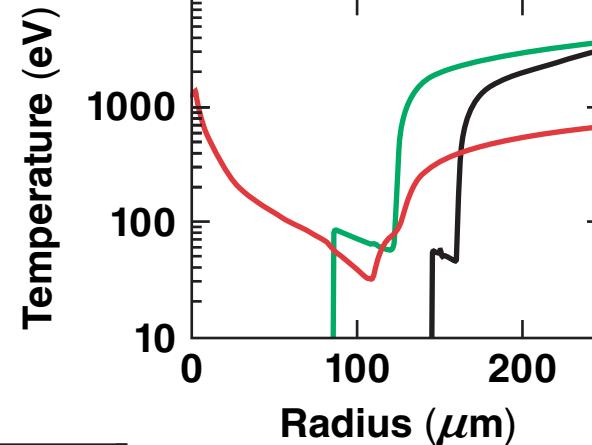
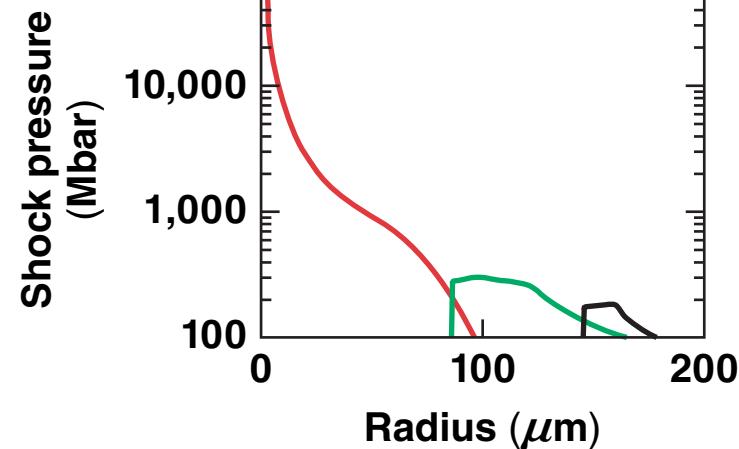
- Small-spot phase plates
- DPR's
- SSD
- $\langle I \rangle \sim 3 \times 10^{15} \text{ W/cm}^2$
- Density scale length $L_{n_c/4} \sim 120 \mu\text{m}$



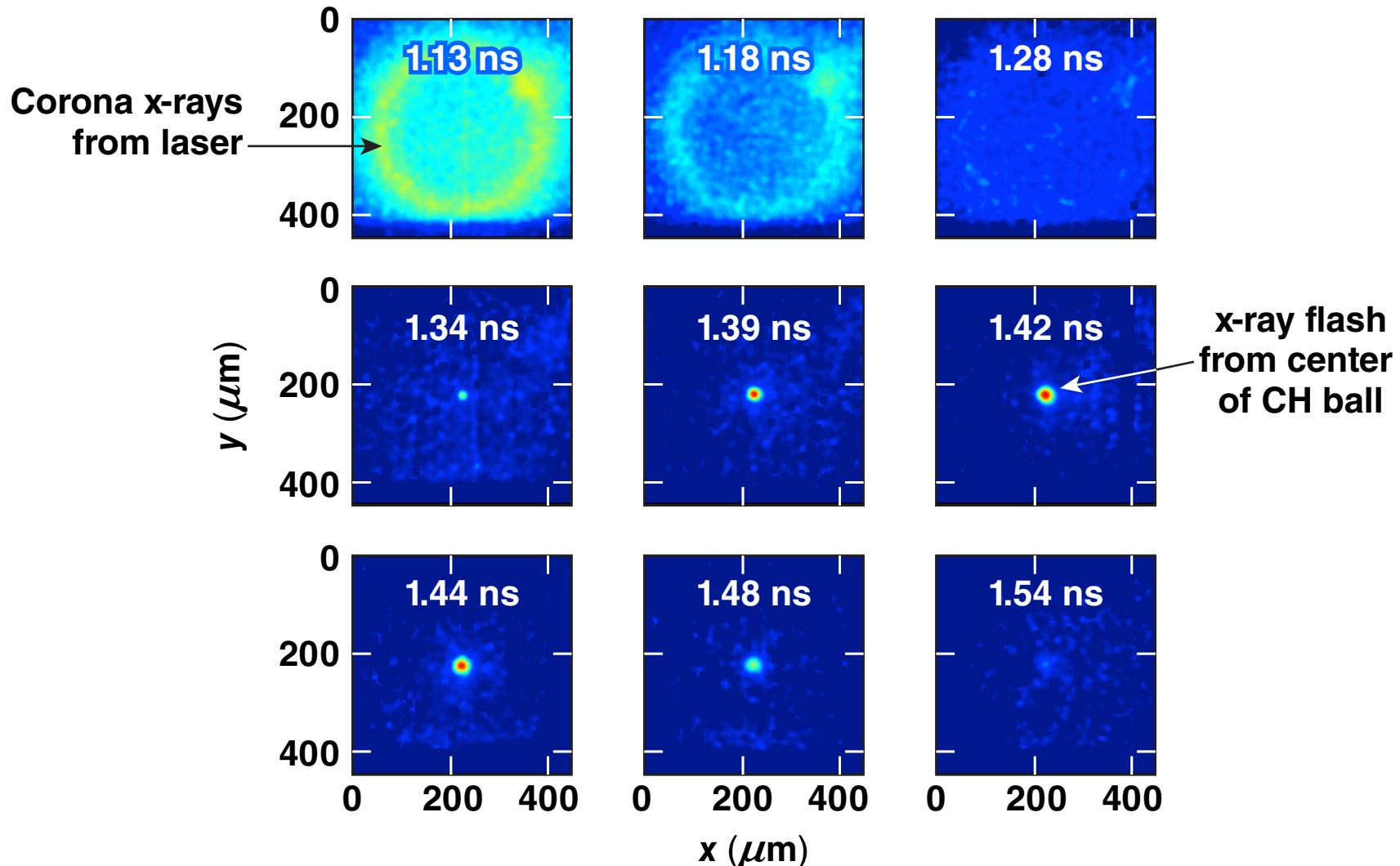
One-dimensional *LILAC* simulations predict a strong spherical shock wave that converges in the center of the solid target



ns
— 1.44
— 1.06
— 0.64

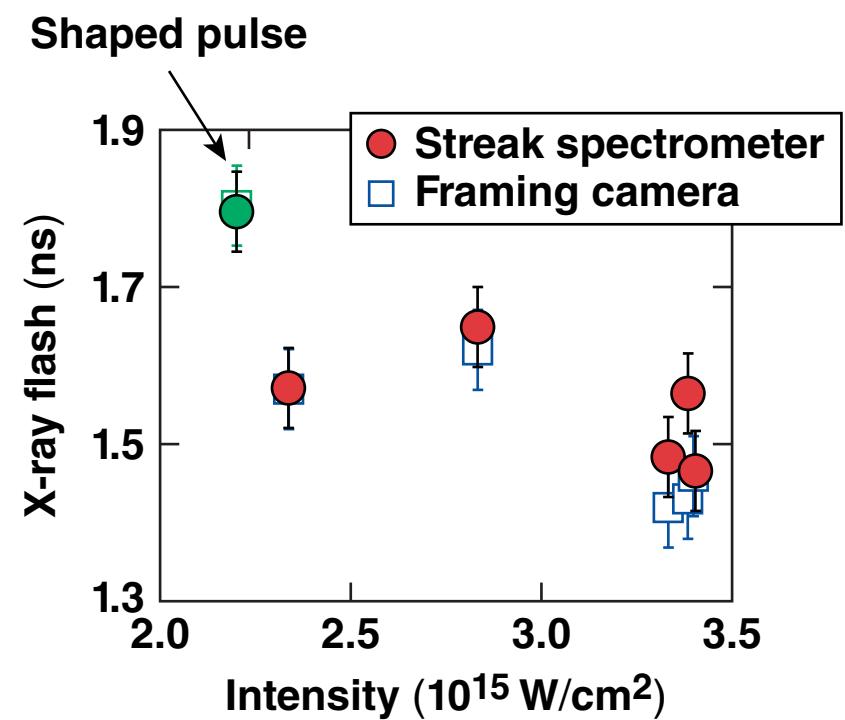
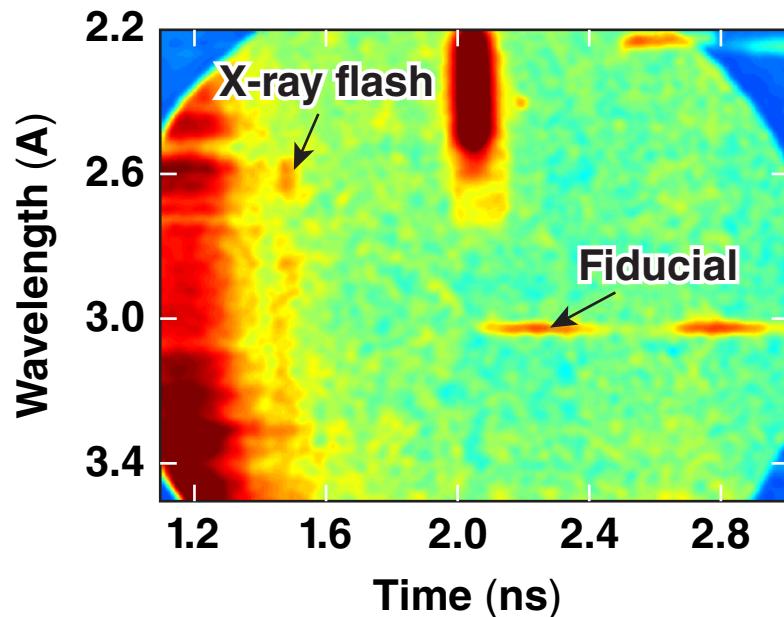


An x-ray framing camera captured a short x-ray flash at the time when the shock converged in the center



E23233

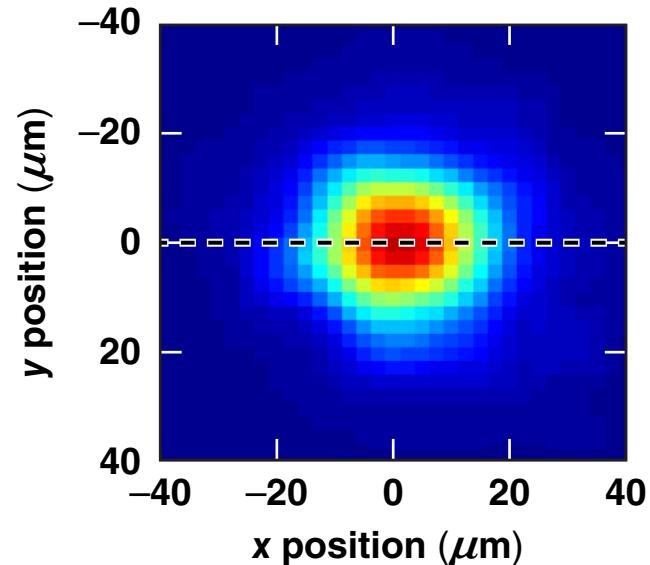
The x-ray flash was measured with a streaked x-ray spectrometer



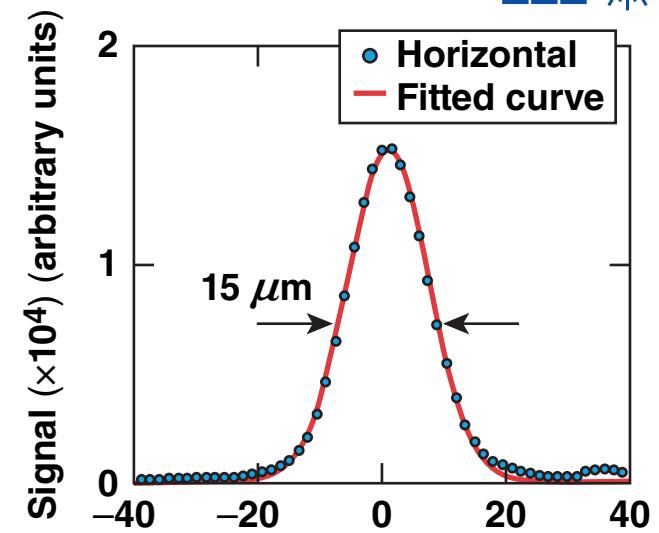
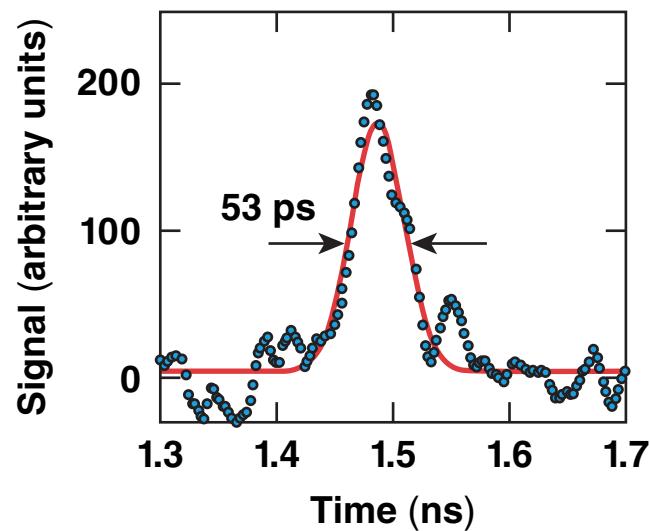
The x-ray flash was emitted from a small volume of $\sim 10^3 \mu\text{m}^3$ in less than 50 ps



Framing-camera images from target center



Ti He α flash from streak camera

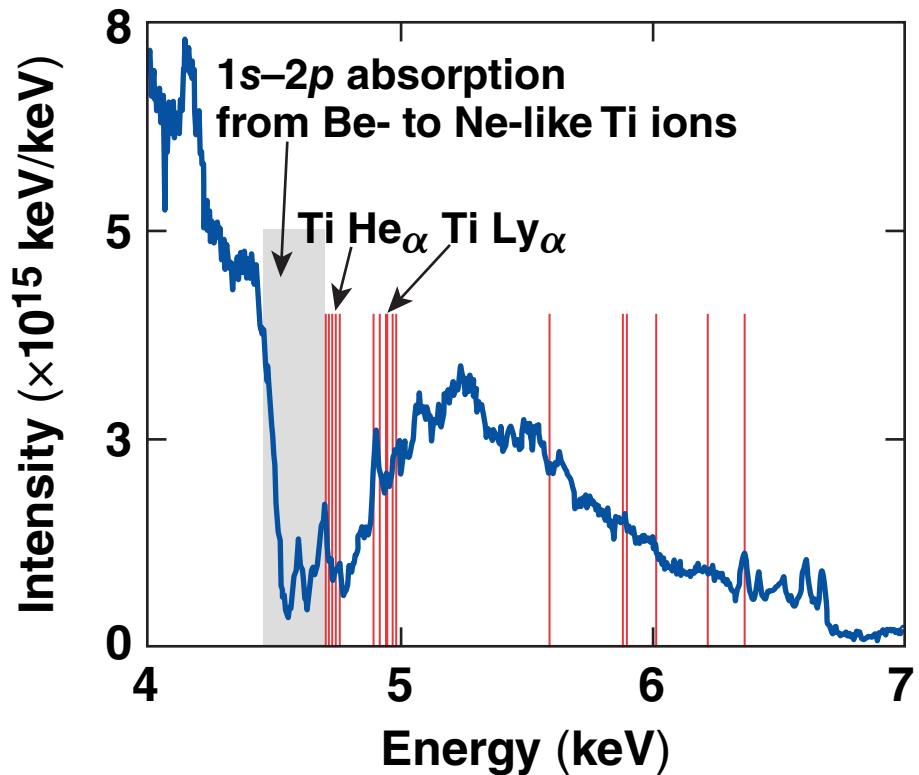


$$\Delta x \approx \sqrt{(15 \mu\text{m})^2 - (12 \mu\text{m})^2} \approx 9 \mu\text{m}$$

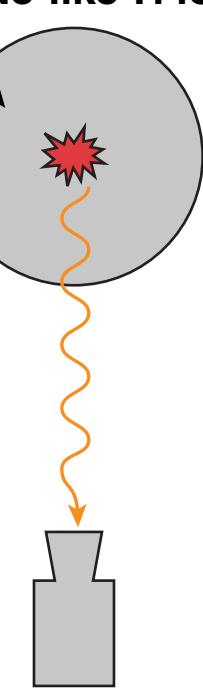
$$\Delta t \approx \sqrt{(53 \text{ ps})^2 - (40 \text{ ps})^2} \approx 35 \text{ ps}$$

E22462a

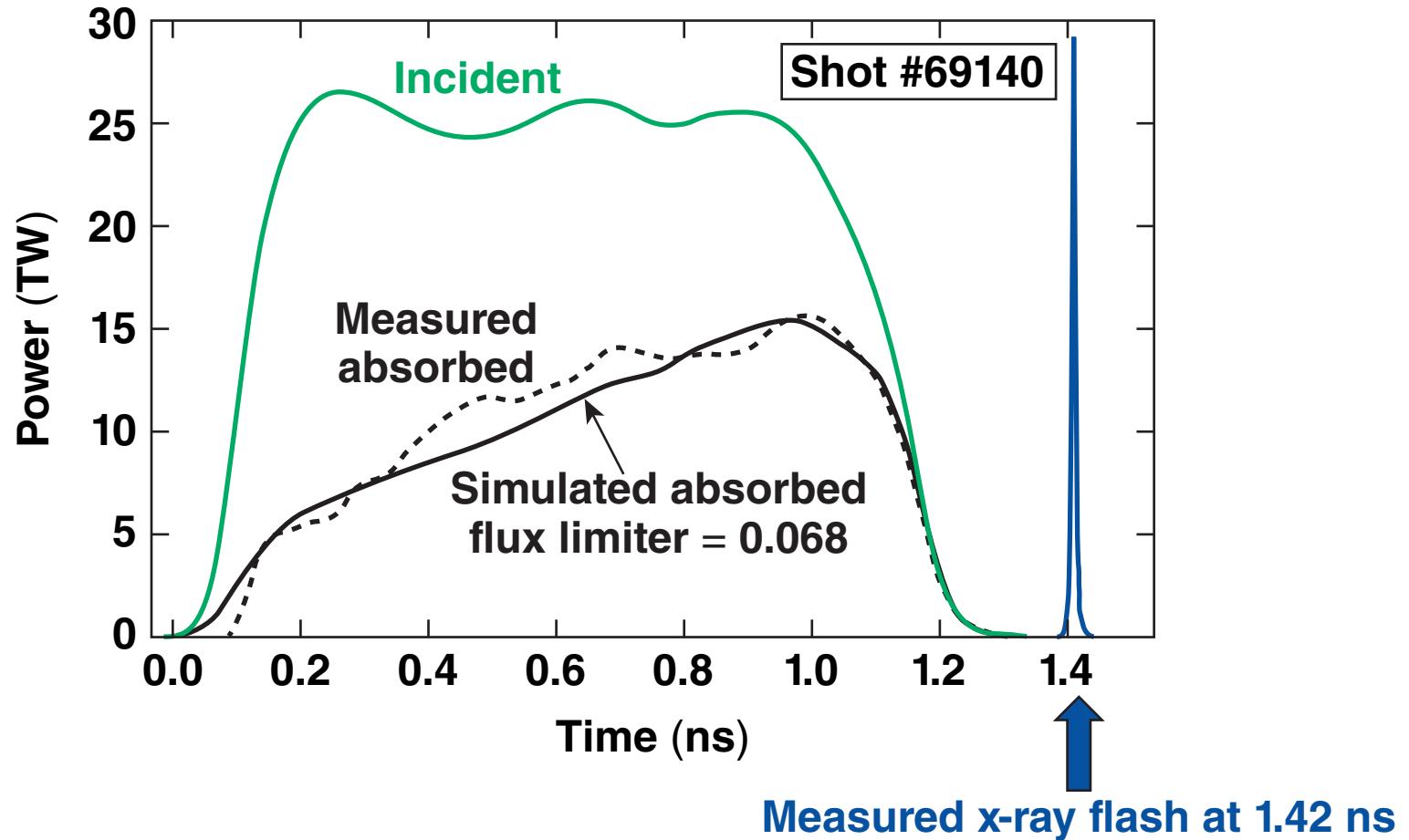
Time integrated x-ray spectra reveal strong 1s–2p absorption features from Ti ions of various charge states



$T_e \sim 10$ to 100 eV
Solid density
Be- to Ne-like Ti ions

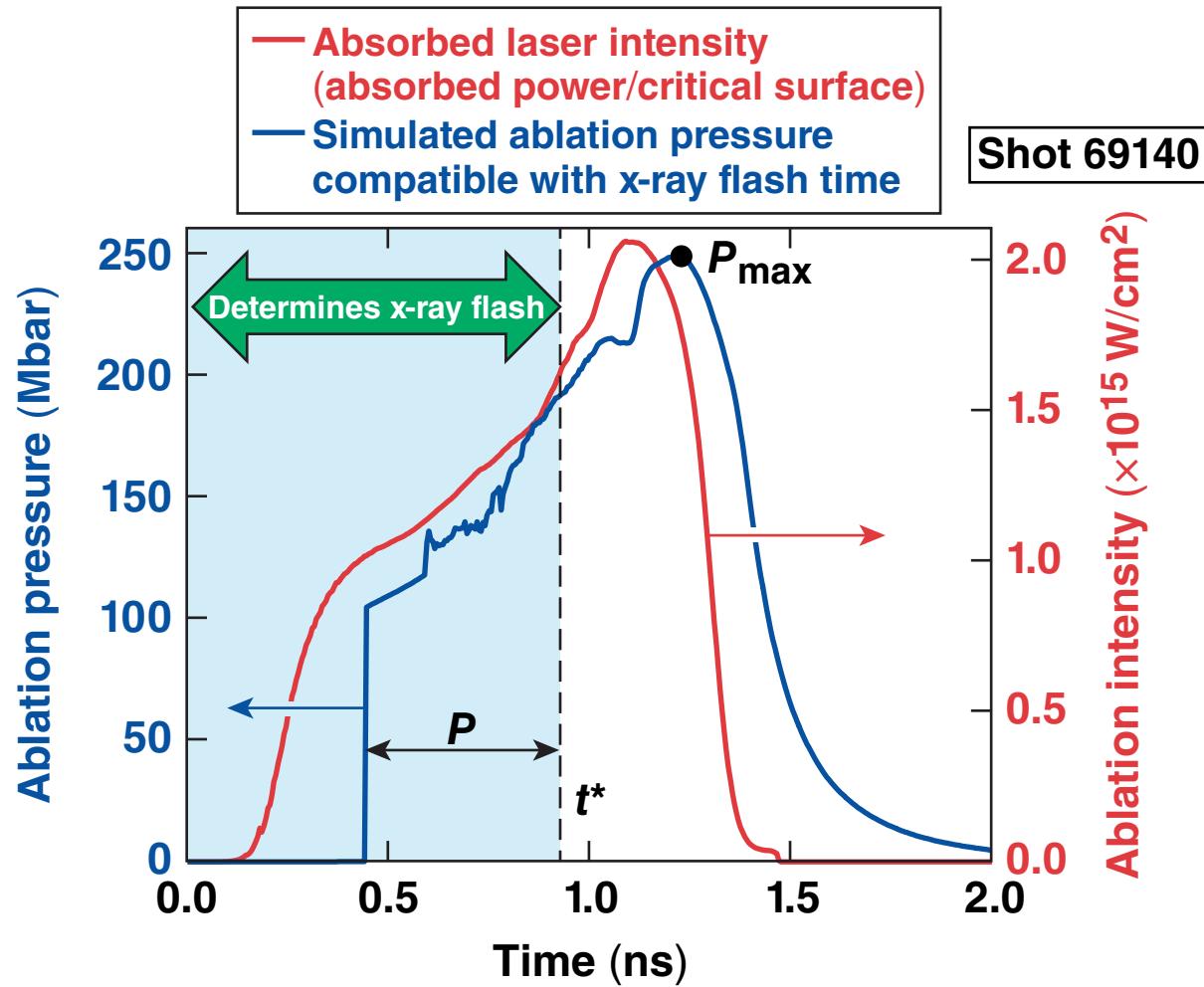


One-dimensional simulations reproduce the measured x-ray flash time when the flux limiter is adjusted to match the absorbed laser-power time history



E23184

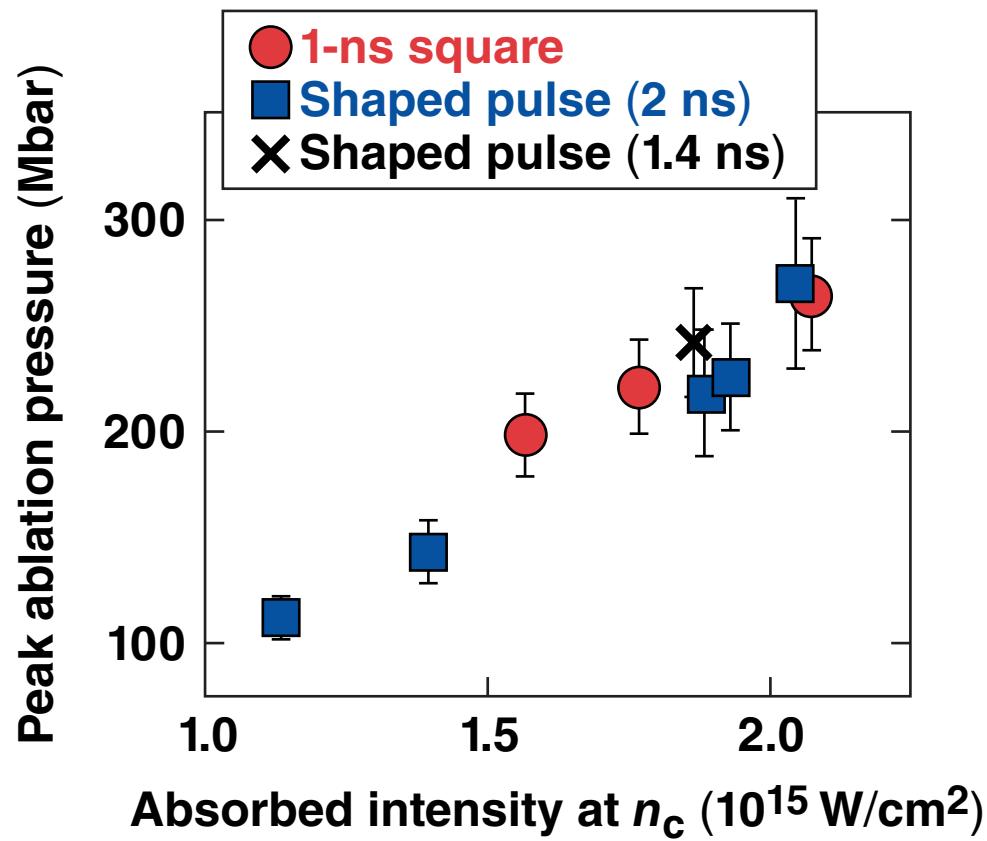
The last 400 ps of the laser pulse does not affect the x-ray flash time



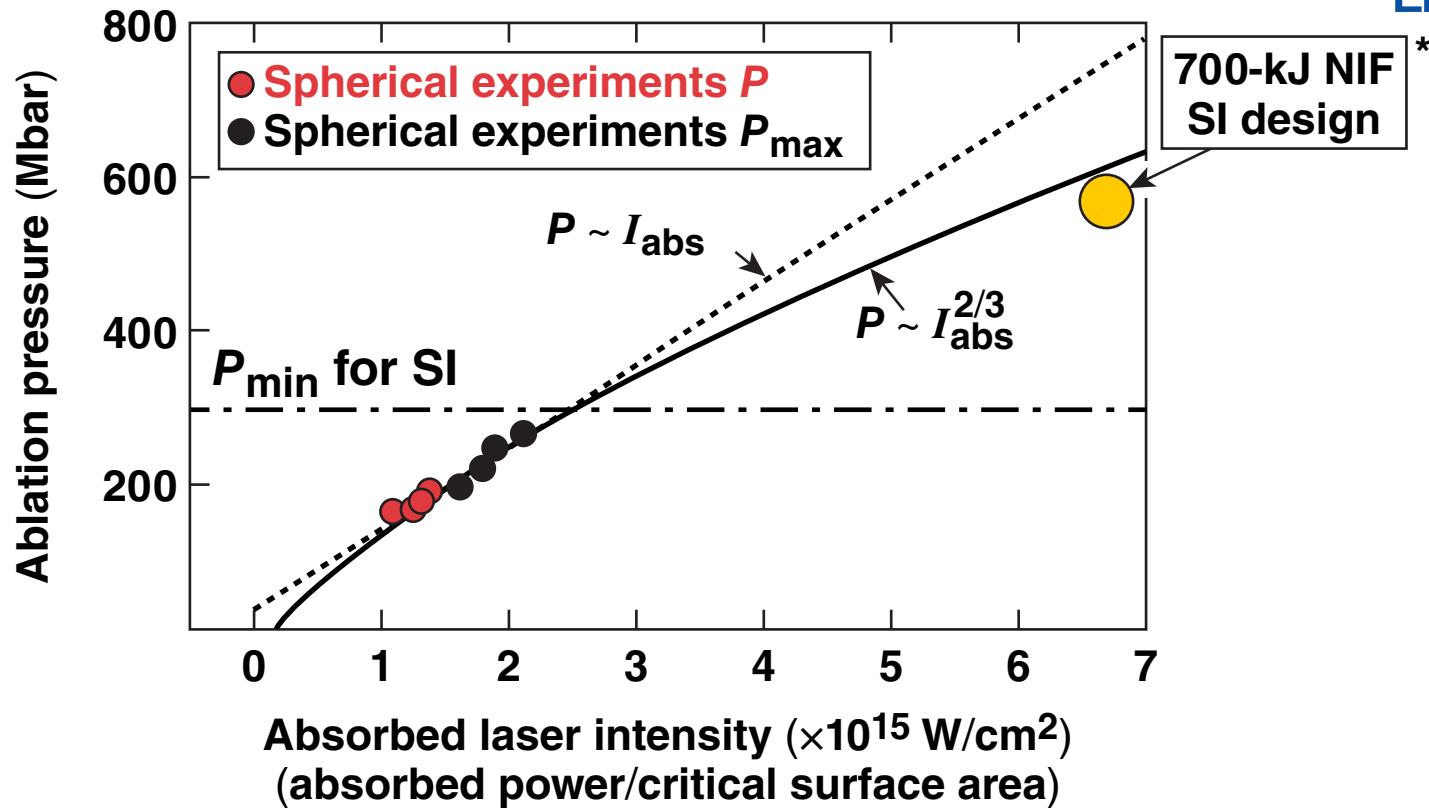
Simulated $P(t < t^*)$ constrained by the measured laser absorption and x-ray flash.
Simulated P_{\max} is constrained only by measured absorption.

E23187

Peak ablation pressures of up to 270 Mbar were inferred from simulations constrained by the observables



A plot of the ablation pressure versus the absorbed intensity shows the extrapolation required for the SI 700-kJ NIF point design



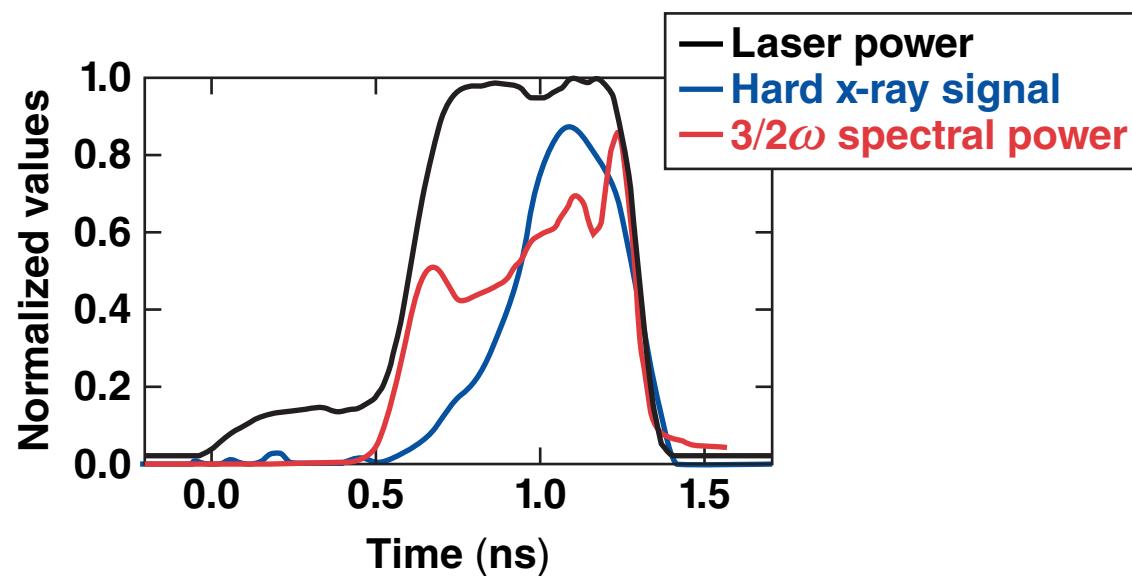
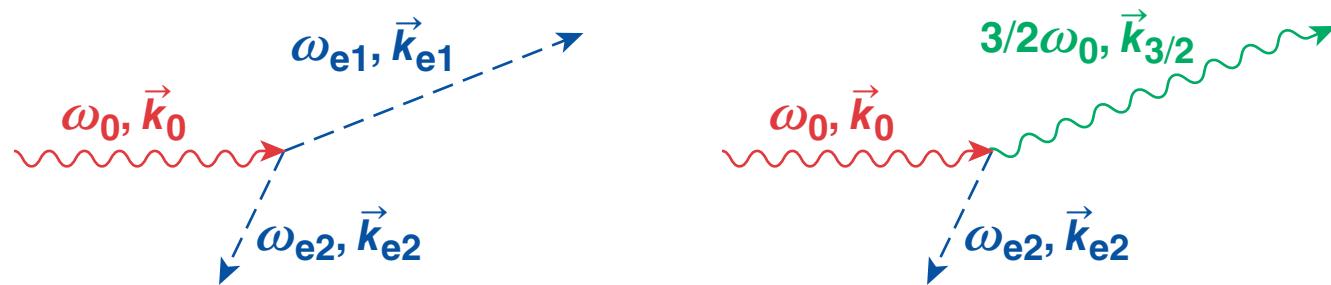
The ablation pressure in the OMEGA experiments approaches the minimum requirements for SI of 300 Mbar. Demonstration of ~ 600 Mbar for the SI NIF design may require experiments on the NIF.

The two-plasmon-decay (TPD) instability is important for the hot electron generation in the spherical strong-shock experiments



$$n_e \sim n_c/4$$

$$\omega_{e1} \sim \omega_{e2} \sim \omega_0/2$$

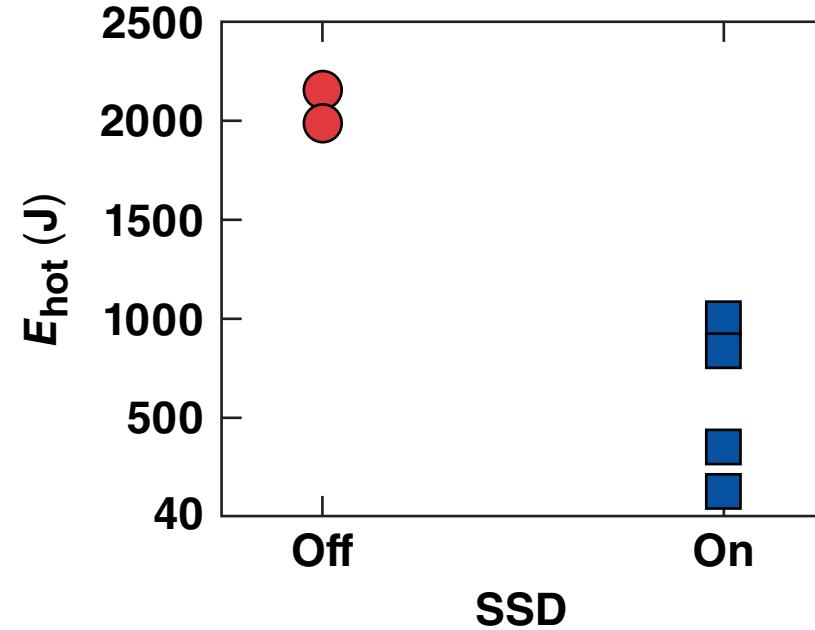
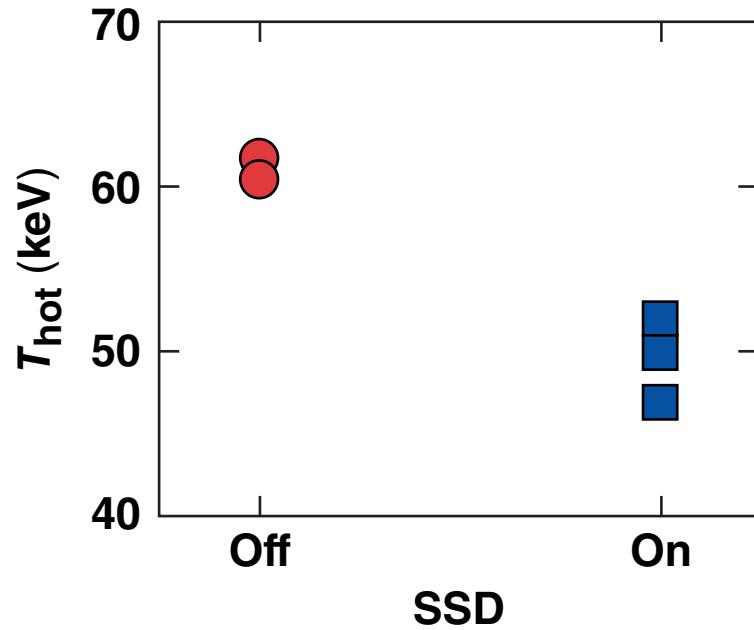


E23189

Switching the SSD bandwidth off has a significant impact on the laser–plasma interaction

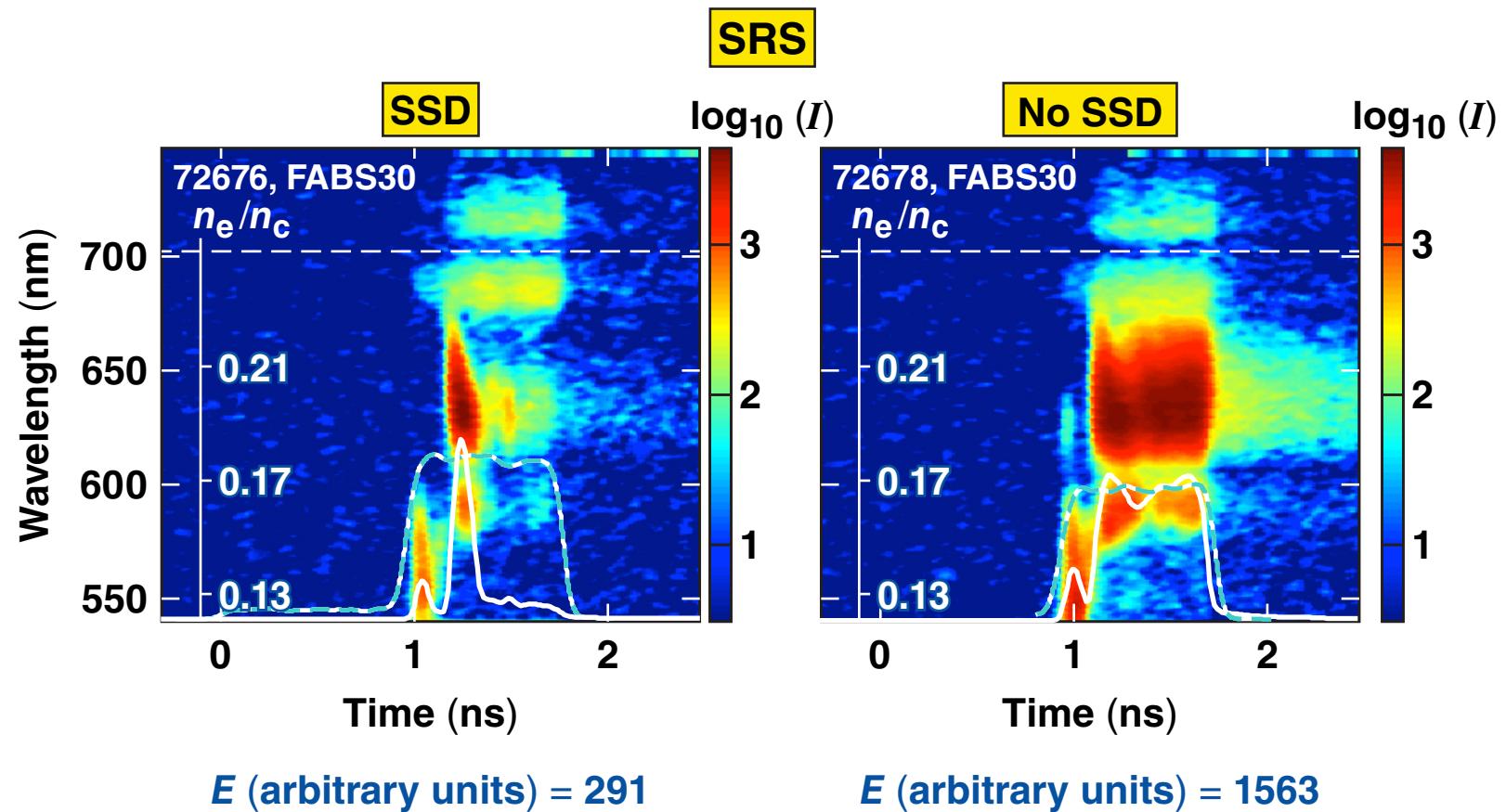


$\sim 3 \times 10^{15} \text{ W/cm}^2$



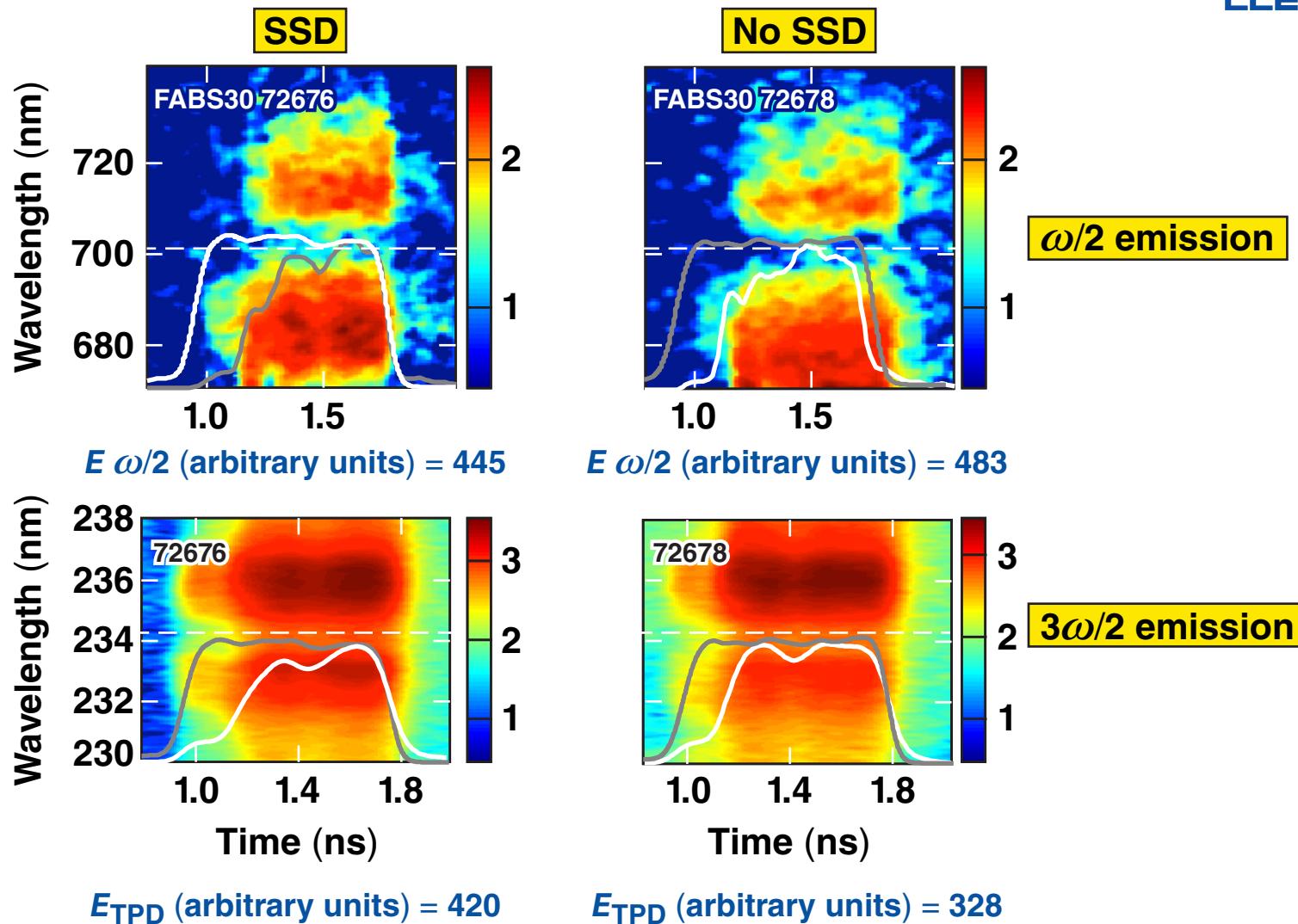
- Measured with nine channel imaging-plate diagnostic
- A hotter electron distribution and more hot electrons were produced when SSD was turned off
- Up to ~10% of the laser energy was converted into hot electrons

SRS increases significantly (~5× in FABS) when SSD is turned off*

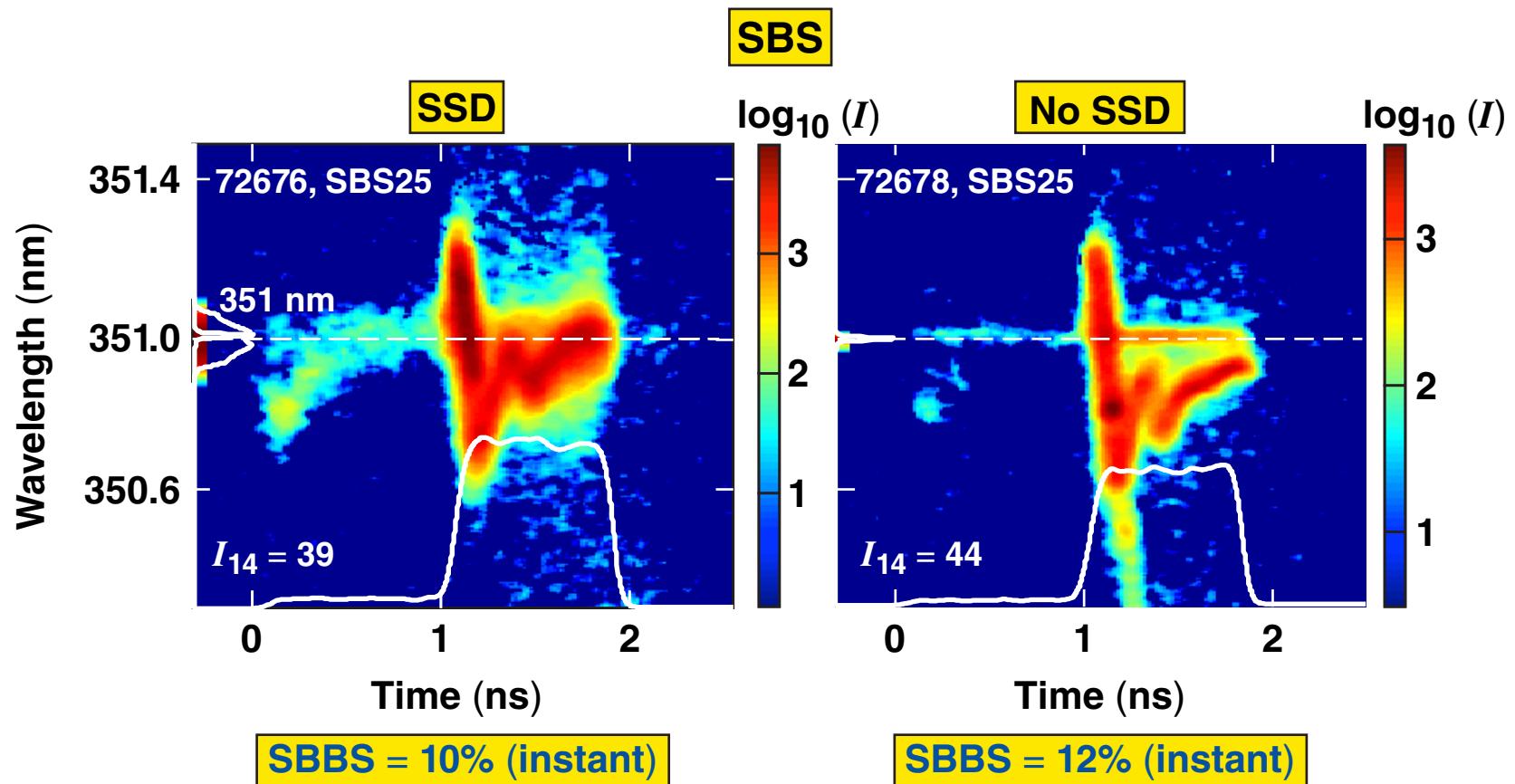


*W. Seka et al., this conference

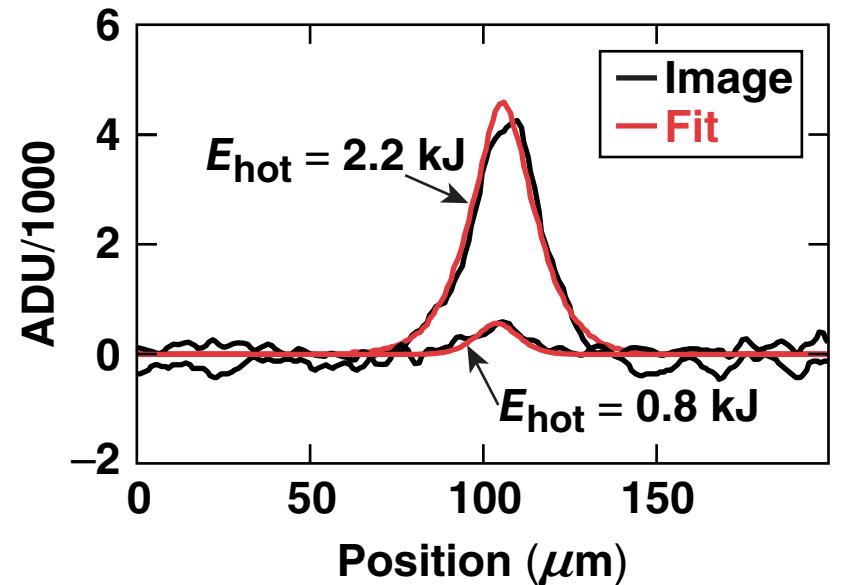
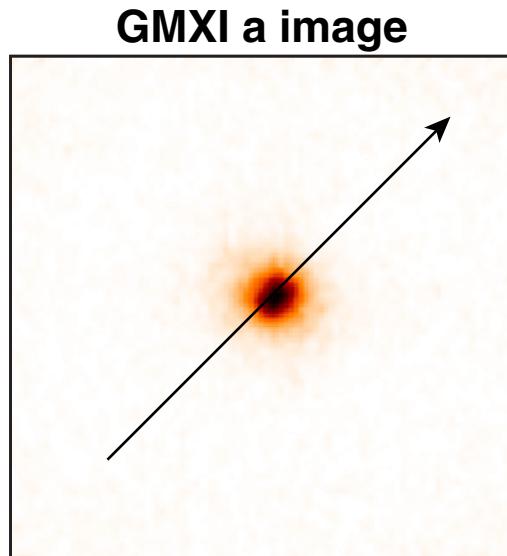
TPD is largely unaffected by SSD in contrast to backscattered SRS



The SBS backscatter signal is insensitive to SSD

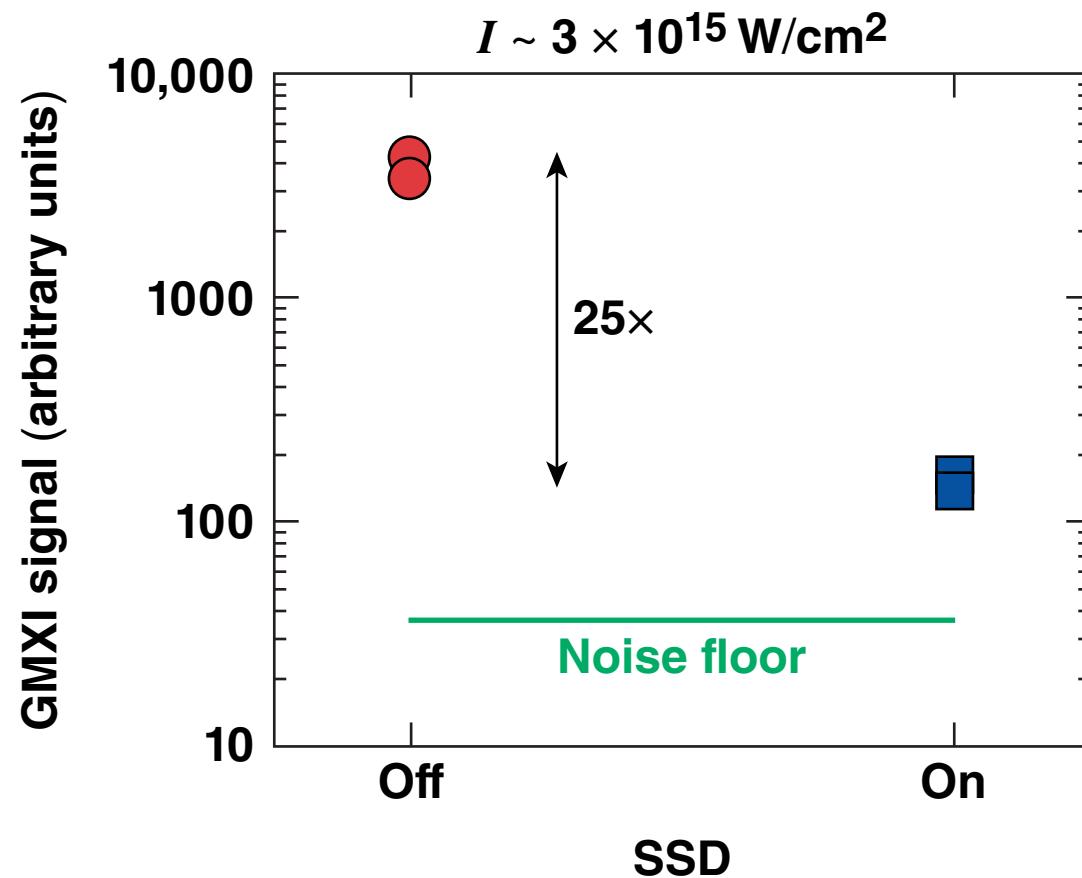


The increase in electron production correlates with an increase in the x-ray emission from the target center



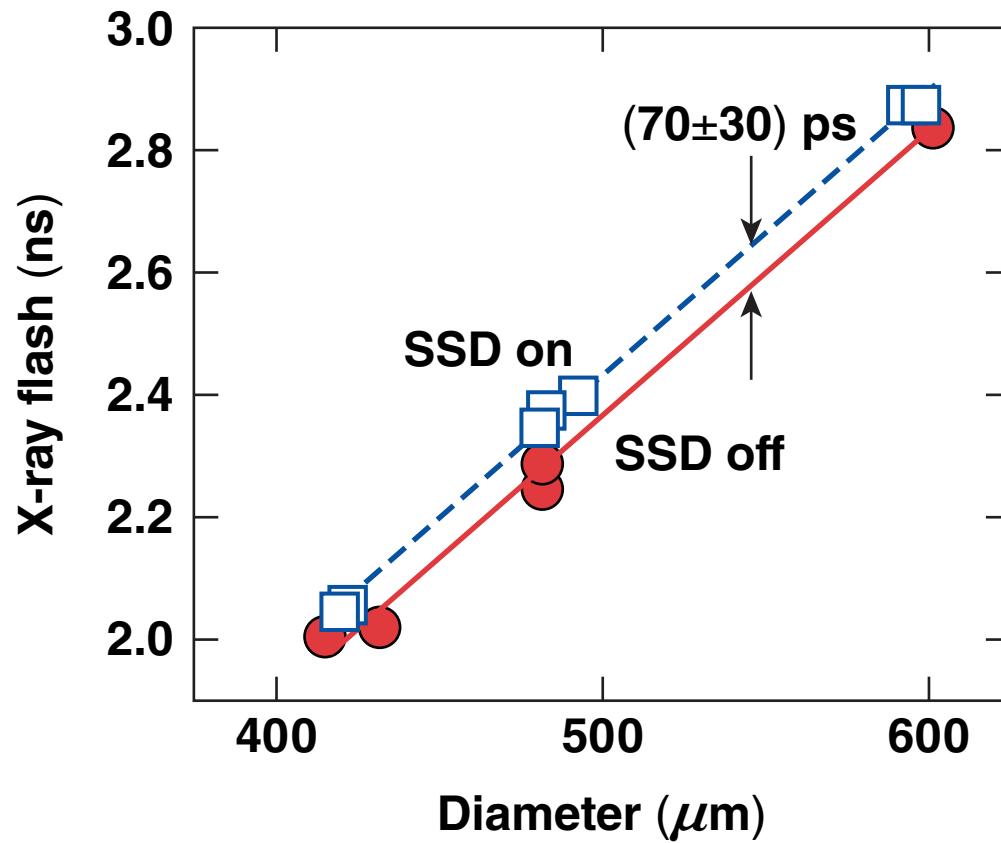
Time integrated x-ray microscope* data from the core center

The x-ray emission strongly increased when SSD was turned off



E23192

An earlier flash time was measured
when SSD was turned off



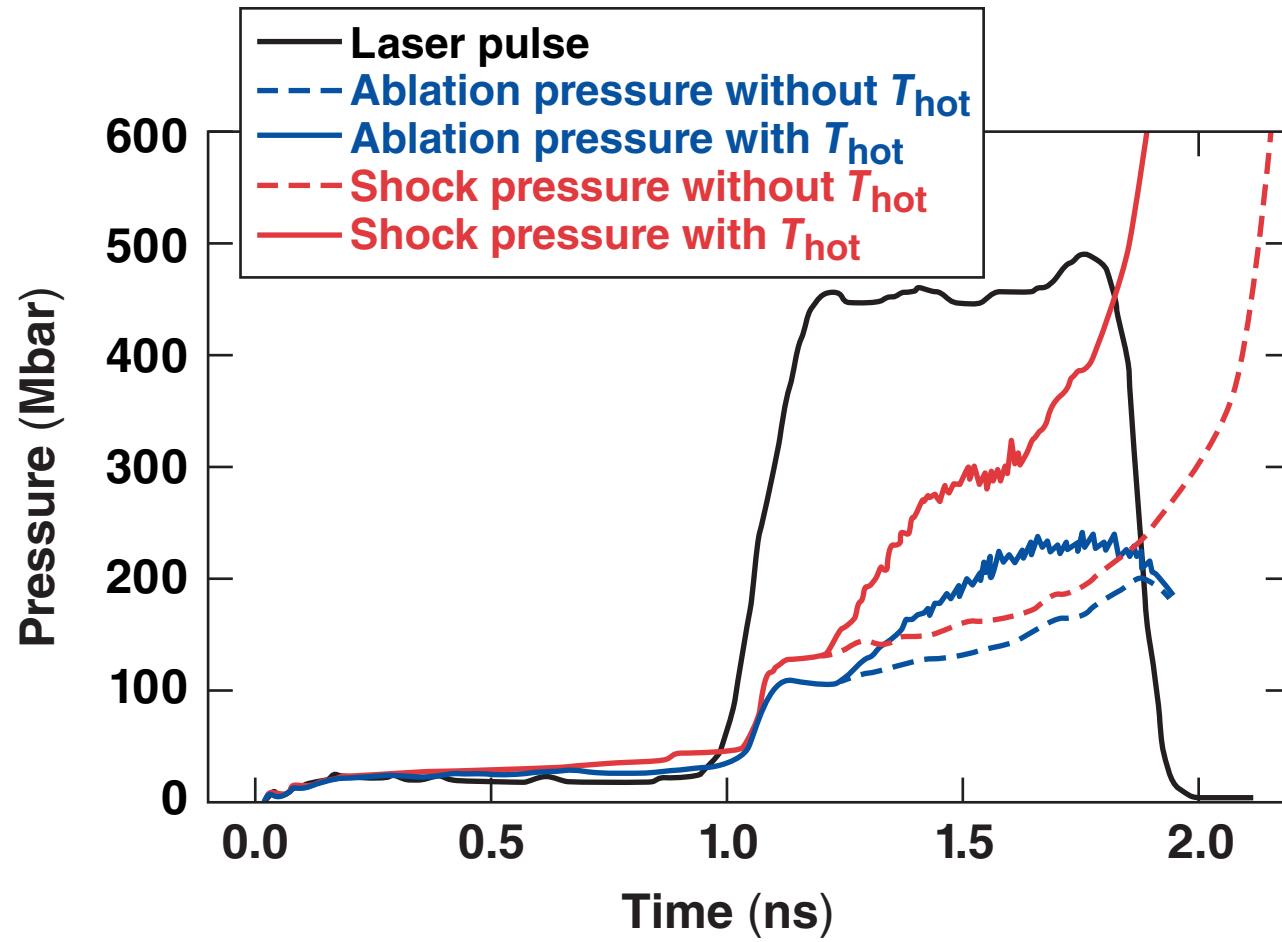
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- Turning off smoothing by spectral dispersion (SSD) in overlapping high-intensity beam experiments increased the hot-electron fraction and shock strength
- Stimulated Raman scattering (SRS) increases significantly (~5×) when SSD is turned off, while two-plasmon–decay (TPD) is unaffected
- Moderate hot electrons may be beneficial to shock ignition by coupling energy to the outer layer of the imploding capsule

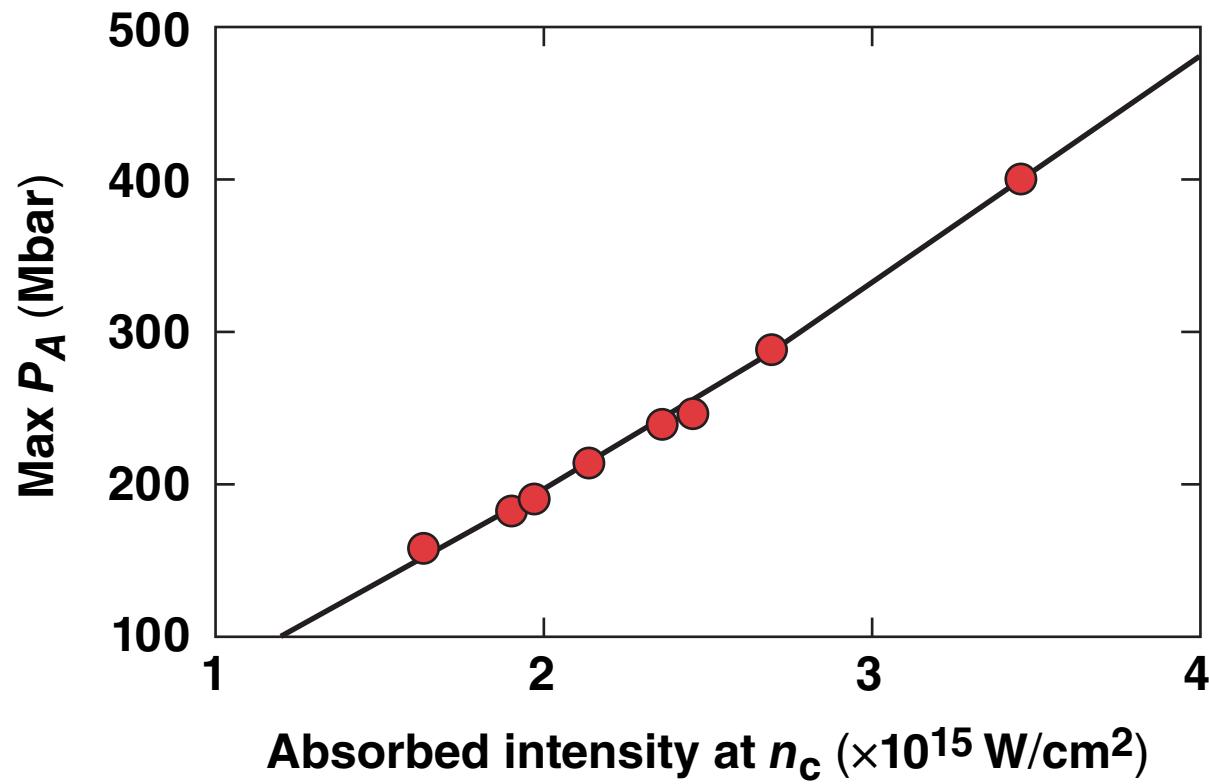
The inferred ablation pressures of 270 Mbar at $\sim 3 \times 10^{15} \text{ W/cm}^2$ approaches the minimum SI requirement of 300 Mbar.

LILAC simulations indicate that hot electrons greatly enhance the shock pressure



E23271

The simulated maximum ablation pressure with hot-electron deposition increases more than linear with the laser intensity



$$P_A^{\max} \sim (80 \text{ Mbar}) \cdot I_{\text{abs}}^{1.3}$$

E23272