

# Understanding the Creation of NIF-Scale Plasmas at the Omega Laser Facility for Laser-Plasma Instability Studies



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## Summary

Hydro-calculated laser-intensity scalings for producing NIF-scale plasmas on OMEGA have been confirmed by experiments and self-similar model predictions



- The hydro-conditions of National Ignition Facility (NIF)-scale plasmas at the quarter-critical density regime are important for understanding laser-plasma instabilities (LPI's)
- A self-similar model for LPI experiments on OMEGA predicts at  $n_c/4$  that
  - $L_n (\mu\text{m}) \propto I^{1/4}$
  - $T_e (\text{keV}) \propto I^{1/2}$
  - $I_{qc} \propto I$
- These predictions are reproduced by 2-D hydro simulations
- DRACO simulations further indicated that scale-length plasmas of  $L_n \sim 500 \mu\text{m}$  can be created with concave spherical half-shells

TC10609

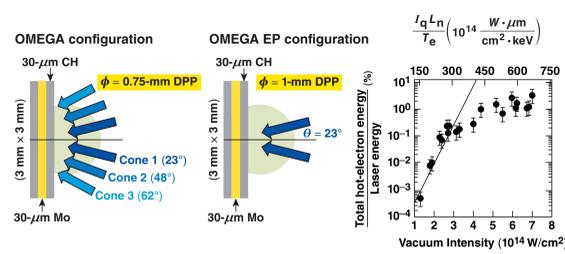
Understanding and mitigating two-plasmon-decay (TPD) instability relies on the accurate knowledge of plasma conditions at  $n_c/4$



- Long-scale-length plasmas ( $L_n > 400 \mu\text{m}$ ), which favor TPD-instability growth, can be encountered in direct-drive-ignition implosions on the NIF
- To understand the laser-intensity scaling of TPD-induced fast electrons, it is crucial to know the exact plasma conditions ( $L_n, T_e, I_{qc}$ ) at the quarter-critical density
- Benchmarking the hydro-simulated plasma conditions at  $n_c/4$  with measurements and model analyses provide more confidence in the TPD-instability studies

TC10610

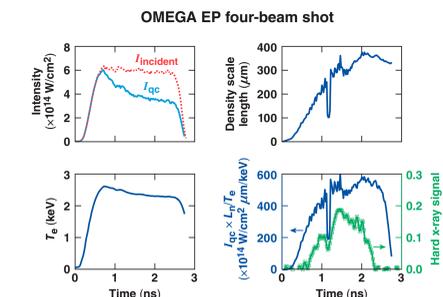
Long-scale-length plasma experiments with planar CH targets have been performed at the Omega Laser Facility using different distributed phase plates (DPP's)\*



TC10617

\*D. H. Froula et al., Phys. Rev. Lett. 108, 165003 (2012).

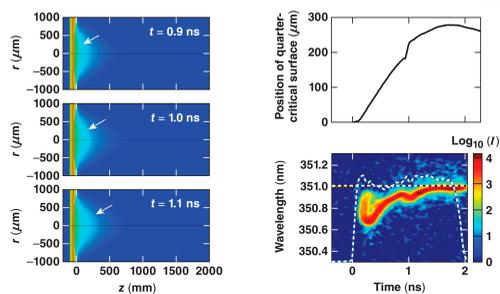
Two-dimensional DRACO simulations\* of these experiments provide the basic plasma conditions at  $n_c/4$  to understand LPI



TC10641

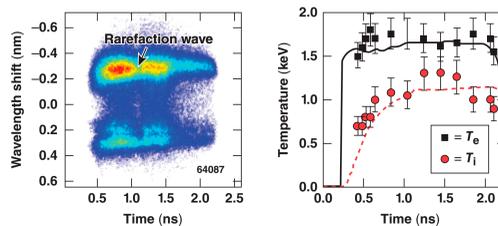
\*S. X. Hu et al., Phys. Plasmas 20, 032704 (2013).

The full-aperture backscatter station (FABS) measurement of light reflection by the rarefaction wave benchmarks the DRACO-predicted hydrodynamics of coronal plasmas



TC10615

The Thomson-scattering measurement of the electron temperature at  $n_c/4$  showed good agreement with hydro simulations



TC10642

The self-similar model\* is used to understand the laser-ablated slab plasma formation



- Solving the hydrodynamic equation with the self-similar dimensionless coordinate  $\xi$

$$\begin{cases} v(m, t) = q_0^{-1/4} (k_0 t)^{3/8} V(\xi) \\ u(m, t) = q_0^{1/4} (k_0 t)^{1/8} U(\xi) \\ \rho(m, t) = q_0^{3/4} (k_0 t)^{-1/8} P(\xi) \\ e(m, t) = q_0^{1/2} (k_0 t)^{1/4} E(\xi) \\ q(m, t) = q_0 Q(\xi) \end{cases}$$

\*S. Atzeni and J. Meyer-ter-Vehn, The Physics of Inertial Fusion: Beam Plasma Interaction, Hydrodynamics, Hot Dense Matter, International Series of Monographs on Physics (Clarendon Press, Oxford, 2004).

TC10615

The laser-intensity scaling of  $L_n$ ,  $T_e$ , and  $I_{qc}$  can be derived from the self-similar model solutions



- The definition of  $L_n$  and the equation-of-state (EOS) relationship of  $e \sim kT_e$

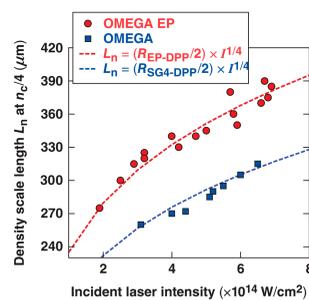
$$L_n = \frac{\rho}{|d\rho/dz|} = \rho \left| \frac{d\rho}{dm} \right|^{-1} \left| \frac{d\rho}{dm} \right| = 1 \left| \frac{d\rho}{d\xi} \cdot \frac{d\xi}{dm} \right|^{-1}$$



$$\begin{cases} L_n \propto I^{1/4} \lambda_0^{3/4} \\ T_e \propto I^{1/2} \lambda_0^2 \\ I_{qc} \propto I \end{cases}$$

TC10616

DRACO-simulated intensity scaling of  $L_n$  is in very good agreement with the self-similar model prediction

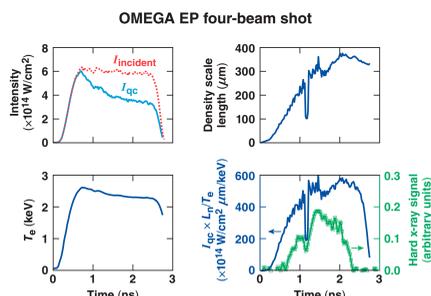


Direct measurements of  $L_n$  are underway.\*

TC10645

\*D. Haberberger, this conference.

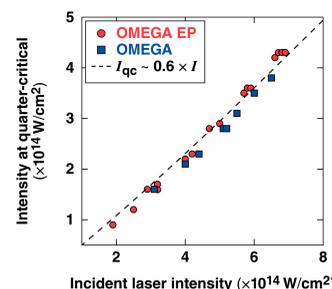
Two-dimensional DRACO simulations\* of these experiments provide the basic plasma conditions at  $n_c/4$  to understand LPI



TC10641

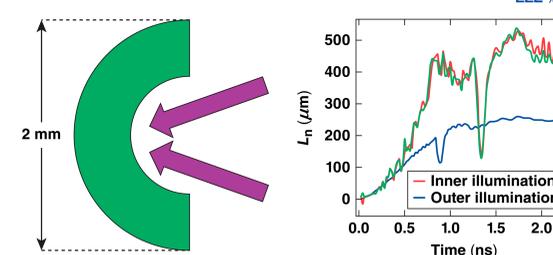
\*S. X. Hu et al., Phys. Plasmas 20, 032704 (2013).

Both DRACO simulations and the self-similar model predict the linear scaling of  $I_{qc}$  with the incident intensity



TC10619

With concave spherical half-shells, DRACO simulations predicted NIF-scale plasmas with even longer density scale lengths ( $L_n \sim 500 \mu\text{m}$ )



Measurements showed ~3x higher hard x-ray signals with inner illumination than outer illumination at the same intensity.

TC10620

## Summary

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- A self-similar model for *LPI* experiments on OMEGA predicts at  $n_c/4$  that
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- These predictions are reproduced by 2-D hydro simulations
- *DRACO* simulations further indicated that scale-length plasmas of  $L_n \sim 500 \mu\text{m}$  can be created with concave spherical half-shells

# Collaborators

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# Understanding and mitigating two-plasmon–decay (TPD) instability relies on the accurate knowledge of plasma conditions at $n_c/4$

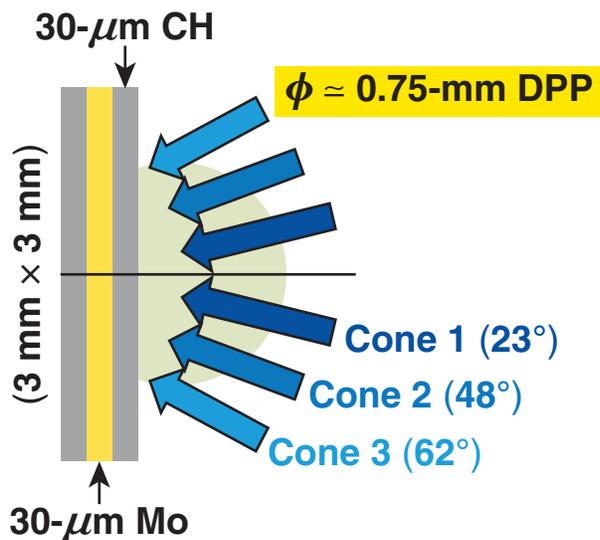


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- To understand the laser-intensity scaling of TPD-induced fast electrons, it is crucial to know the exact plasma conditions ( $L_n, T_e, I_{qc}$ ) at the quarter-critical density
- Benchmarking the hydro-simulated plasma conditions at  $n_c/4$  with measurements and model analyses provide more confidence in the TPD-instability studies

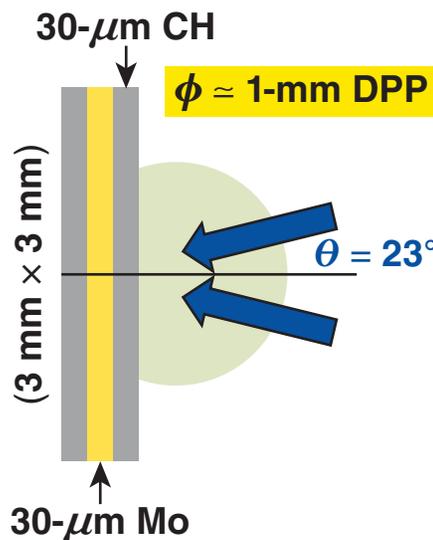
# Long-scale-length plasma experiments with planar CH targets have been performed at the Omega Laser Facility using different distributed phase plates (DPP's)\*



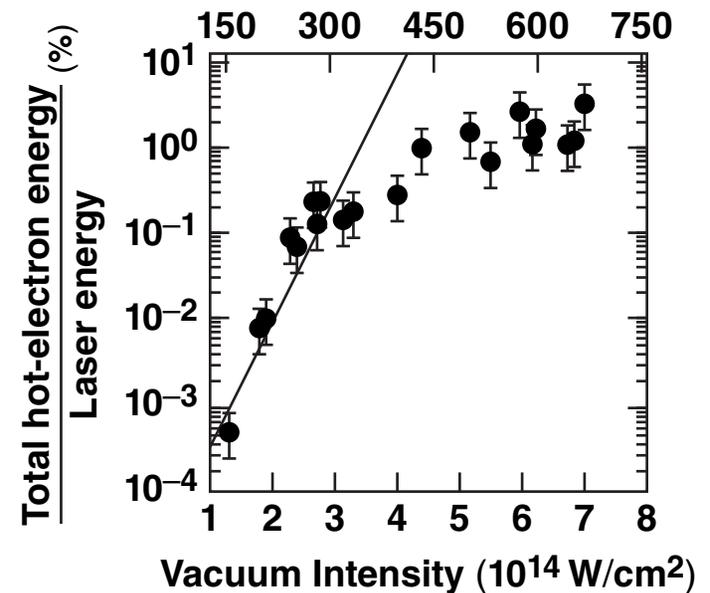
OMEGA configuration



OMEGA EP configuration

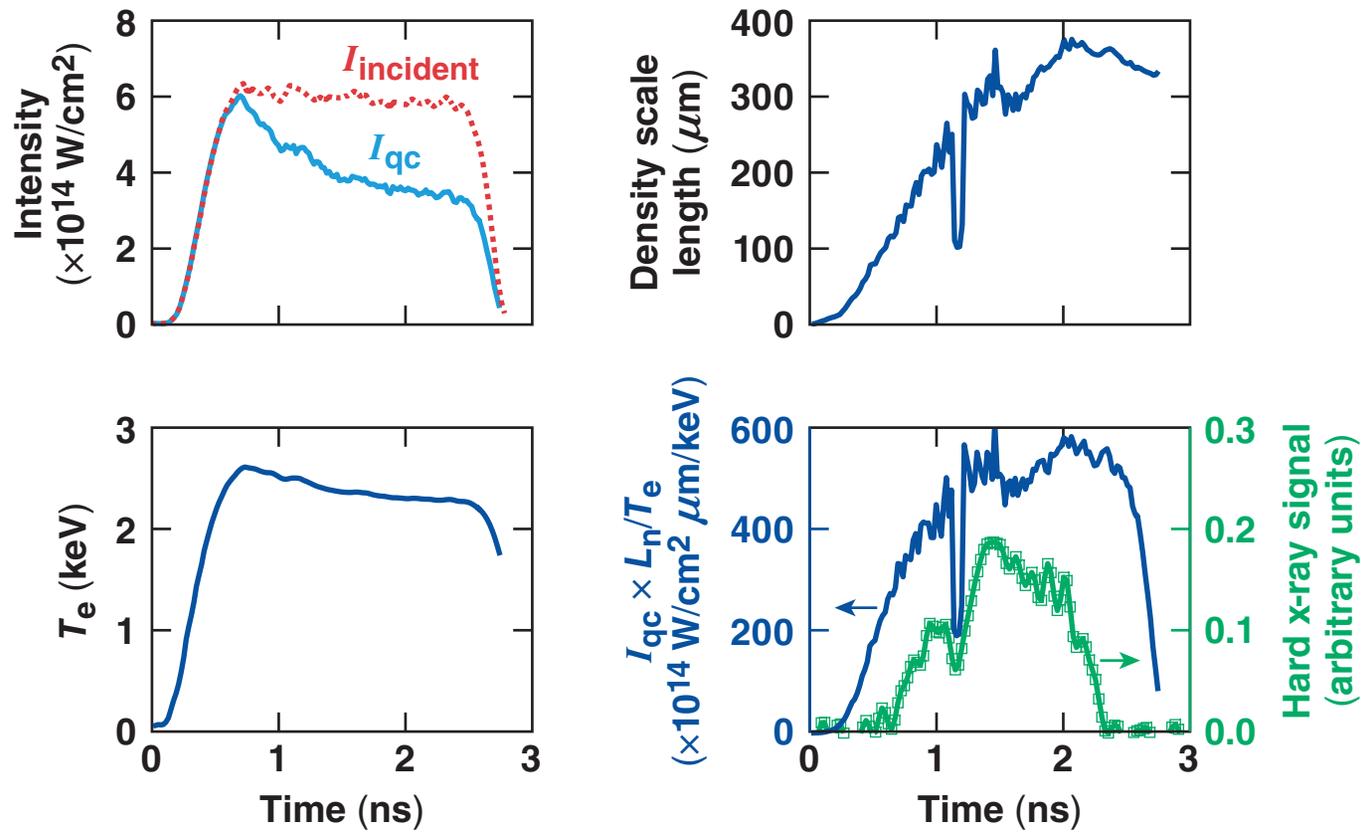


$$\frac{I_q L_n}{T_e} \left( 10^{14} \frac{\text{W} \cdot \mu\text{m}}{\text{cm}^2 \cdot \text{keV}} \right)$$

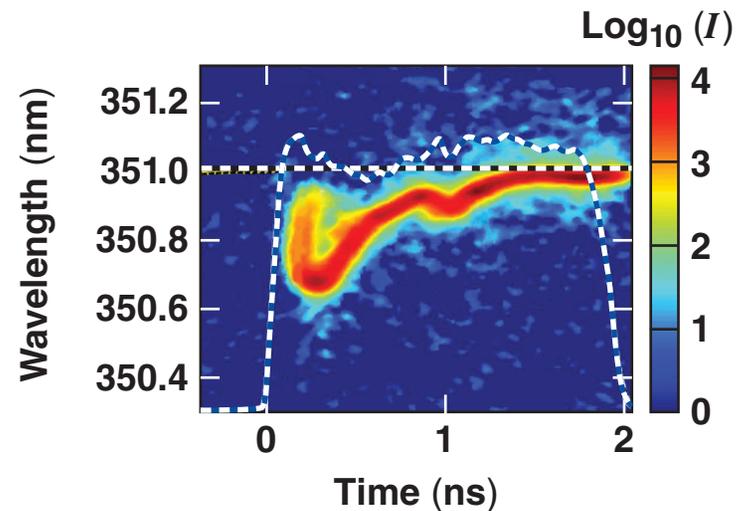
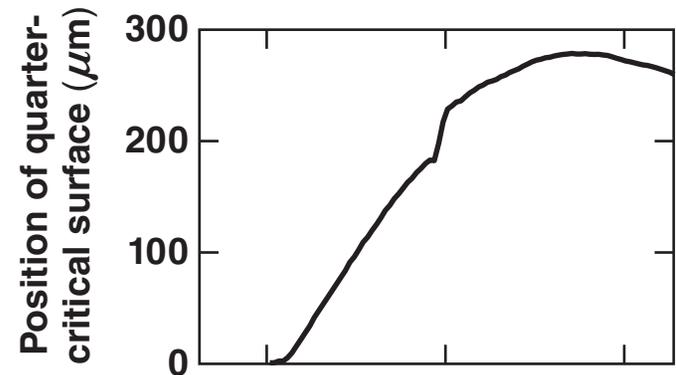
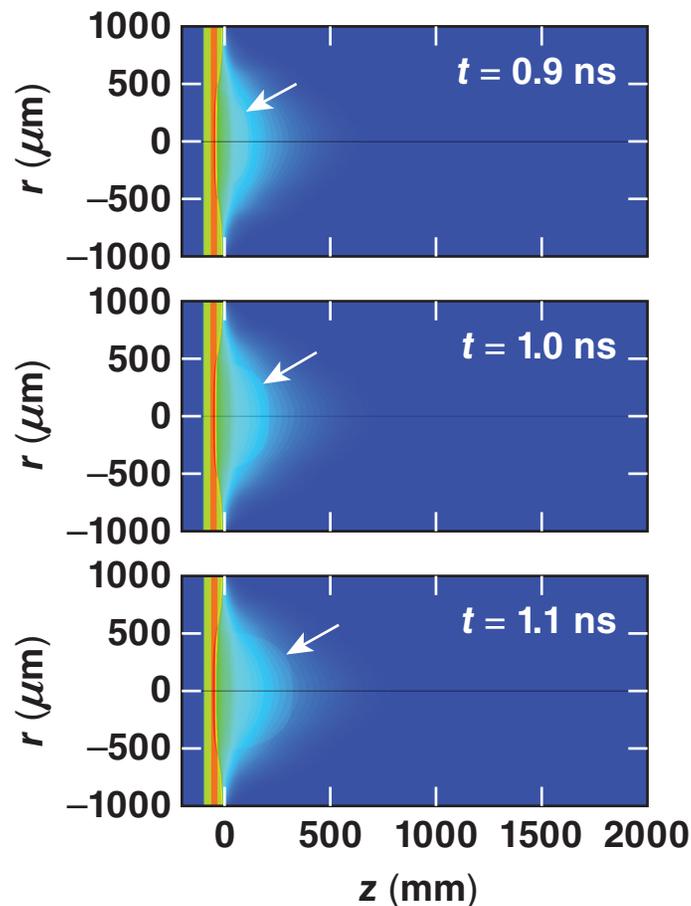


# Two-dimensional *DRACO* simulations\* of these experiments provide the basic plasma conditions at $n_c/4$ to understand LPI

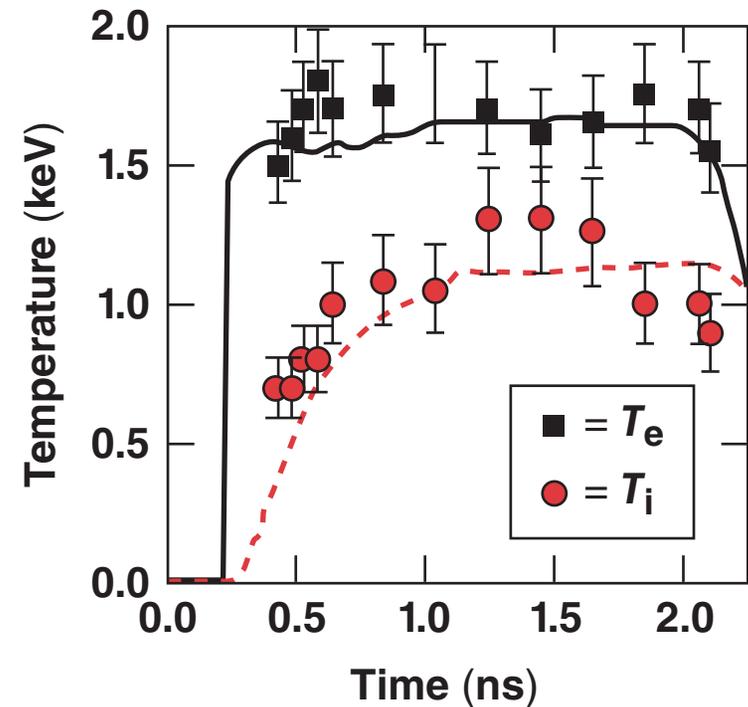
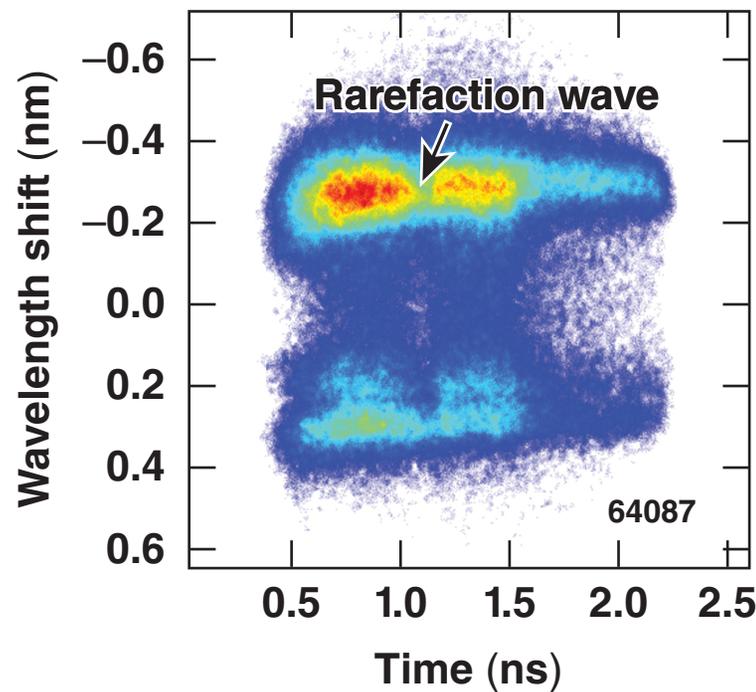
## OMEGA EP four-beam shot



# The full-aperture backscatter station (FABS) measurement of light reflection by the rarefaction wave benchmarks the *DRACO*-predicted hydrodynamics of coronal plasmas



# The Thomson-scattering measurement of the electron temperature at $n_c/4$ showed good agreement with hydro simulations



# The self-similar model\* is used to understand the laser-ablated slab plasma formation

- Solving the hydrodynamic equation with the self-similar dimensionless coordinate  $\xi$

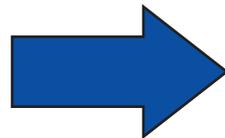
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# The laser-intensity scaling of $L_n$ , $T_e$ , and $I_{qc}$ can be derived from the self-similar model solutions

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$$L_n = \frac{\rho}{|d\rho/dz|} = \rho / \left| \rho \frac{d\rho}{dm} \right| = 1 / \left| \frac{d\rho}{dm} \right| = 1 / \left| \frac{d\rho}{d\xi} \cdot \frac{d\xi}{dm} \right|.$$

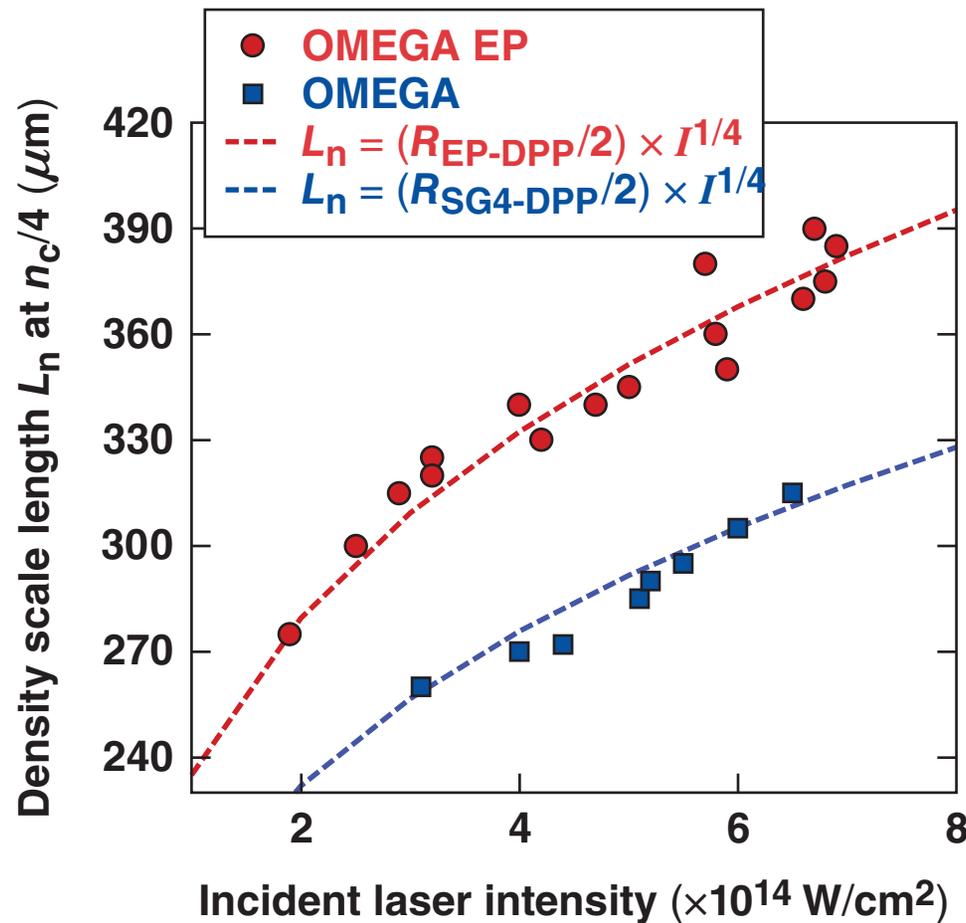


$$L_n \propto I^{1/4} \lambda_0^{1/4}.$$

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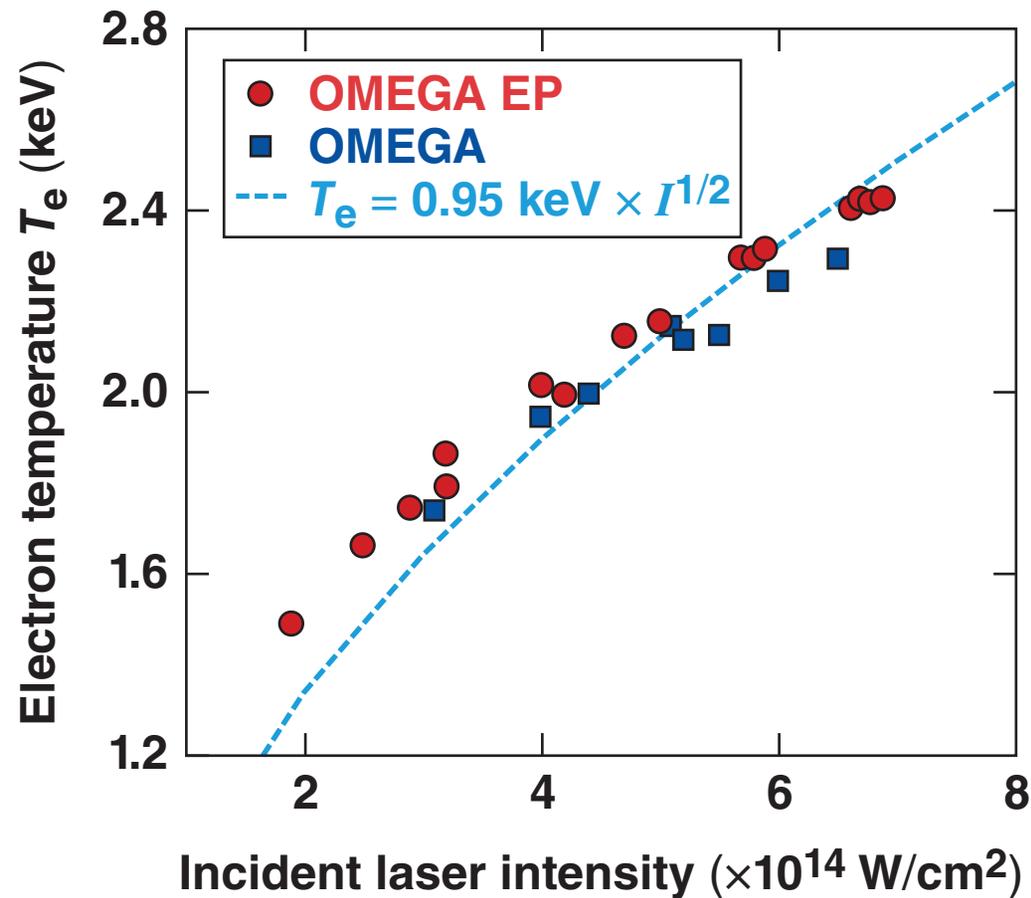
$$I_{qc} \propto I$$

# DRACO-simulated intensity scaling of $L_n$ is in very good agreement with the self-similar model prediction

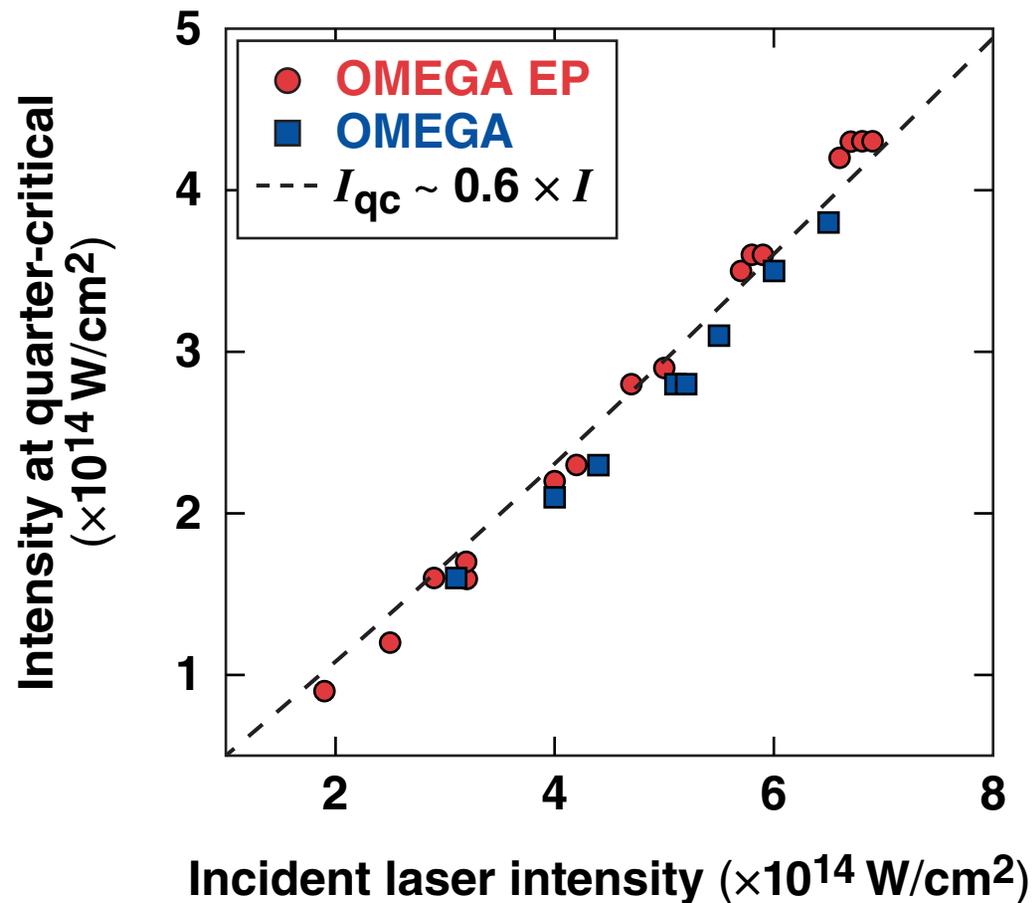


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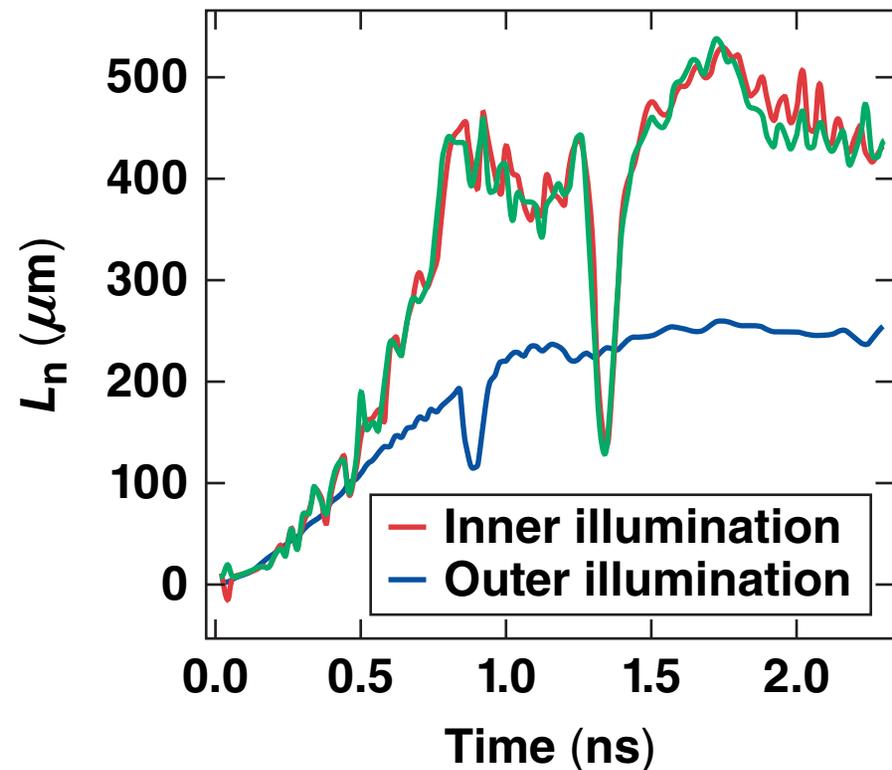
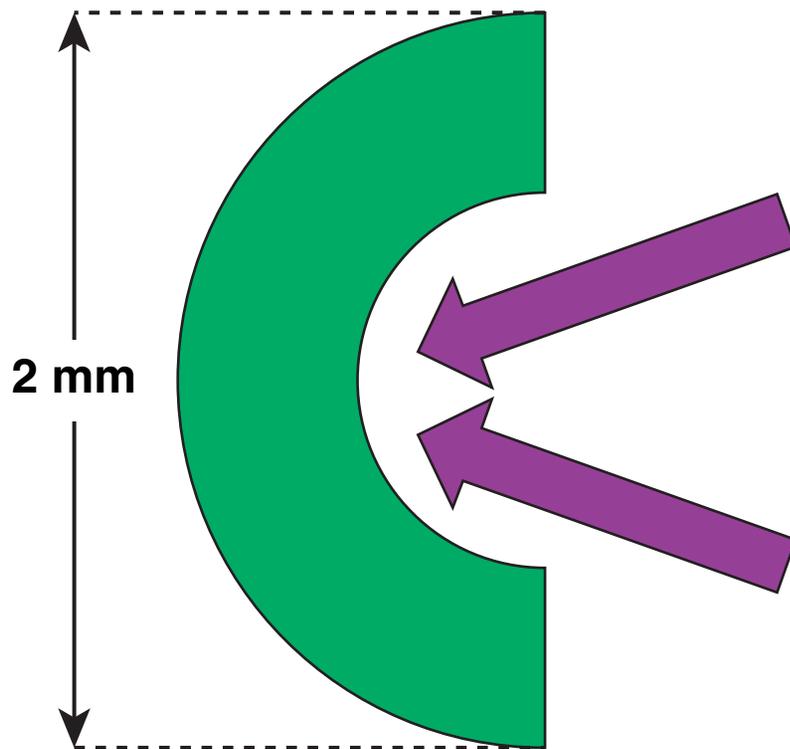
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