The Mitigating Effect of Wave Dissipation on Hot-Electron Generation Caused by Two-Plasmon Decay (TPD) in Inhomogeneous Plasmas



Summary

TPD preheat can be reduced through the manipulation of the collisional and Landau damping of Langmuir and ion-acoustic waves

- Langmuir wave (LW) collisional damping has an impact on the growth rate/thresholds [in addition to the hydrodynamic variables ($\Delta n_e, T_e$)]
 - importance increases with the scale length
- Nonlinear saturation is sensitive to plasma composition ($\langle Z \rangle$, T_i , $\langle Z \rangle^2$) via the ion-acoustic damping rate
 - predictions with the 2-D code ZAK*
 - hot-electron predictions with the 2-D code QZAK**
 - suggestions for planar experiments

^{*}D. F. DuBois et al., Phys. Rev. Lett. <u>74</u>, 3983 (1995);

D. A. Russell and D. F. DuBois, Phys. Rev. Lett. <u>86</u>, 428 (2001).

^{**} D. A. Russell et al., "Two-plasmon-decay turbulence driven by the shared-wave triad of two crossed beams," this conference.



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The ZAK model of TPD is used to predict linear instability and nonlinear saturation caused by density fluctuations

• "Extended" Zakharov equations used in ZAK*

$$\nabla \cdot \left[D_{\text{LW}} - \omega_0^2 (\delta n + \frac{\delta N}{n_0}) \right] E = \left(\frac{e}{4m_c} \nabla \cdot \left[\nabla (E_0 \cdot \overline{E}) - E_0 \nabla \cdot \overline{E} \right] + S_E$$
$$D_{\text{IAW}} \delta n = \nabla^2 \left(|E|^2 + \frac{1}{4} |E_0|^2 \right) / (16\pi M_i) + S_{\delta n}$$
TPD source term

Dispersion relations
for LW and IAW
$$D_{LW} = \left[2i\omega_{p0}\left(D_t + \frac{\nu_e}{*}\right) + 3\nu_e^2\nabla^2\right]$$
$$D_{IAW} = \left(D_t^2 + 2\nu_i * D_t - c_s^2\nabla^2\right)$$

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The collisional TPD threshold can be made to exceed the gradient threshold (for Si at ignition scale)



Two-dimensional ZAK calculations assume two overlapped plane electromagnetic (EM) waves* polarized in their plane of incidence

- This has the feature that TPD is driven unstable convectively even when the single-beam intensities are below threshold
- <|E|²> is computed at saturation



*D. T. Michel et al., "Experimental Demonstration of the Two-Plasmon-Decay Common-Wave Process," this conference.

ZAK Calculations

Increased collisional damping of LW's leads to a lower saturated level of electrostatic fluctuations for $Z_{eff} = 14$



R. Betti, "Advanced Ablator Target Designs for Shock and Hot Spot Ignition on the National Ignition Facility," this conference.

The nonlinear evolution of TPD in the ZAK model depends on the ion-acoustic damping rate v_i

• "Extended" Zakharov equations used in ZAK* $\nabla \cdot \Big[D_{LW} - \omega_0^2 (\delta n + \delta N) / n_0 \Big] E = \frac{(e/4 m_c) \nabla \cdot [\nabla (E_0 \cdot \overline{E}) - E_0 \nabla \cdot \overline{E}]}{P_{LW} \delta n} + S_E$ $D_{IAW} \delta n = \nabla^2 \Big(|E|^2 + \frac{1}{4} |E_0|^2 \Big) / (16 \pi M_i) + S_{\delta n}$ TPD source term

> Dispersion relations for LW and IAW

$$\boldsymbol{D}_{\text{LW}} = \left[2i\omega_{p0} \left(\boldsymbol{D}_{t} + \boldsymbol{v}_{e} \right) + 3\boldsymbol{v}_{e}^{2} \nabla^{2} \right]$$

$$D_{\text{IAW}} = \left(D_t^2 + 2\nu_i * D_t - c_s^2 \nabla^2\right)$$

Only important in the nonlinear stage

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Decreasing the IAW damping leads to a dramatic reduction in the level of electrostatic fluctuations



The reason that decreasing the IAW damping lowers the saturation level is not fully understood

- Langmuir-decay instability
- Ponderomotive response to unstable LW spectrum



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QZAK* is an extension of ZAK to self-consistently include hot-electron generation in the quasilinear approximation

• The diffusion equation
$$\frac{\partial \langle f_e \rangle}{\partial t} + \frac{\partial}{\partial \vec{v}} \cdot \left(D(\vec{v}) \cdot \frac{\partial \langle f_e \rangle}{\partial \vec{v}} \right) = \sigma(\langle f_e \rangle - f_M)$$



*D. A. Russell *et al.*, "Two-Plasmon-Decay Turbulence Driven by the Shared-Wave Triad of Two Crossed Beams"; H. X. Vu *et al.*, "Hot Electron Generation by 'Cavitating' Langmuir Turbulence in the Nonlinear Stage of the Two Plasmon Decay Instability," this conference.

The LW Landau damping rate $v_e(k)$ is recalculated based on the evolving distribution function

• The damping is the only mechanism where the particles act back on the waves

$$\nu_{e}(\vec{k},t) = \frac{\pi}{2} \frac{\omega_{e}^{3}}{k^{2}} \int d\vec{v} \frac{\vec{k} \cdot \partial \langle f_{e}(\vec{v},t) \rangle}{\partial \vec{v}} \delta(\omega_{pe} - \vec{k} \cdot \vec{v})$$



QZAK predicts less hot-electron production for a plasma with weakly damped ion-acoustic waves



Ion-acoustic damping can be manipulated by modifying the plasma composition

- Ion Landau damping decreases with $ZT_e / T_i (c_s \gg v_{ti})$
- Electron Landau damping in IAW's is always weak ($v_{te} \gg c_s$)
- Light ions (e.g., hydrogen) can increase the damping rate
- The multi-ion dispersion relation is solved by finding the most weakly damped mode*
- Ion-ion collisions complicate the matter
- Part of the spectrum will be collisionally damped because of ion viscosity ($k\lambda_{ii} < 10$)

^{*}E. A. Williams et al., Phys. Plasmas 2, 129 (1995).

These effects can be investigated experimentally in planar targets on OMEGA/OMEGA EP[†]



[†]S. X. Hu *et al.*, "Analyses of Long-Scale-Length Plasma Experiments with Different Ablator Materials on the OMEGA EP Laser System," this conference.
^{*}R. K. Follet *et al.*, "Thomson-Scattering Measurements of Ion-Acoustic Wave Amplitudes Driven by the Two-Plasmon Decay Instability," this conference.

Summary/Conclusions

TPD preheat can be reduced through the manipulation of the collisional and Landau damping of Langmuir and ion-acoustic waves

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