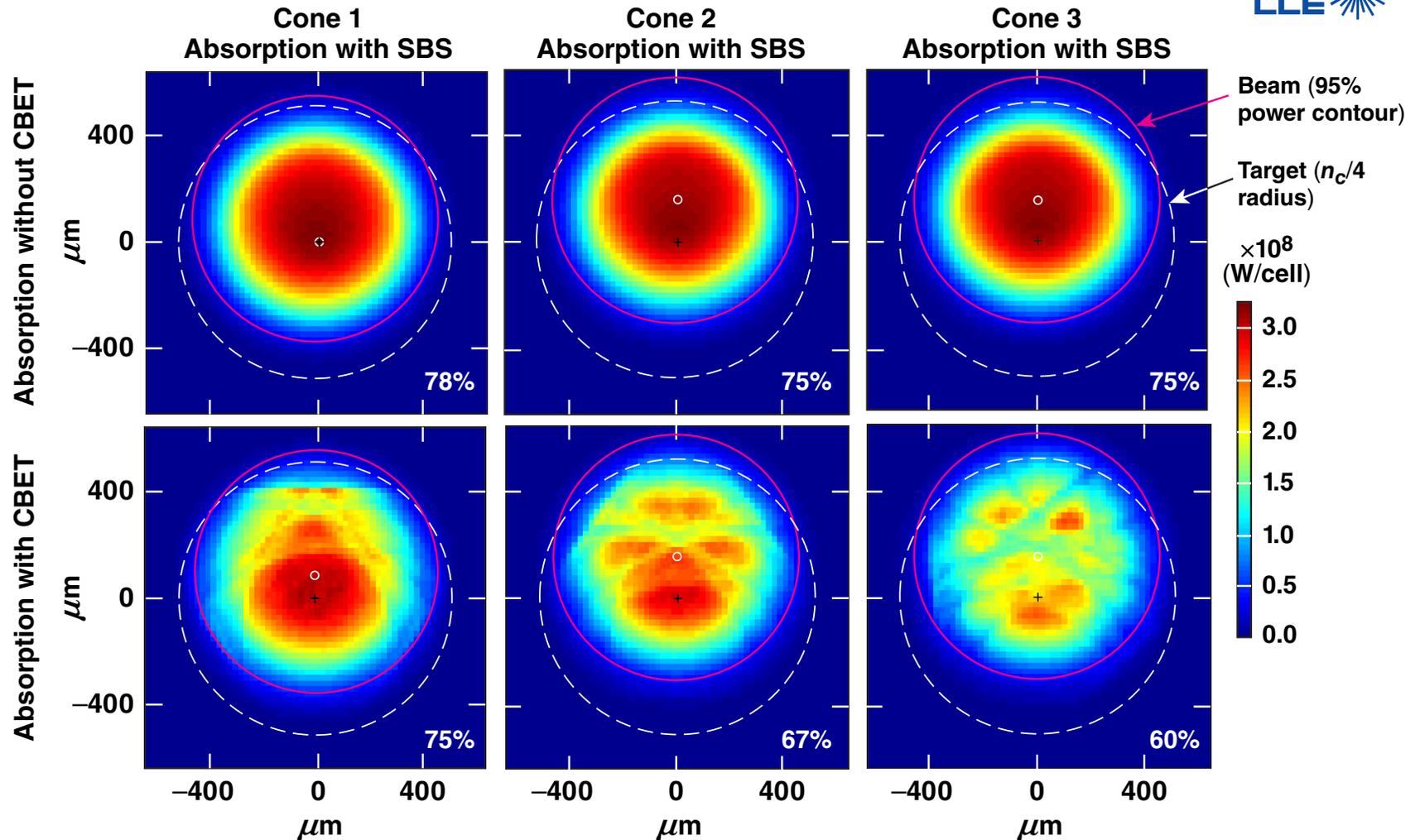


Cross-Beam Energy Transfer in Polar-Drive Implosions



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41st Annual Anomalous
Absorption Conference
San Diego, CA
19–24 June 2011

Cross-beam energy transfer (CBET) predictions show significant absorption profile modifications in Polar Drive (PD)

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 - preventing energy from reaching the high absorption region
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- **CBET must be included in PD implosion modeling and its mitigation through phase plate design will be studied**

Collaborators



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Cross-Beam Energy Transfer

EM-seeded SBS cross-beam energy transfer causes some laser energy to “bypass” the high-absorption zone



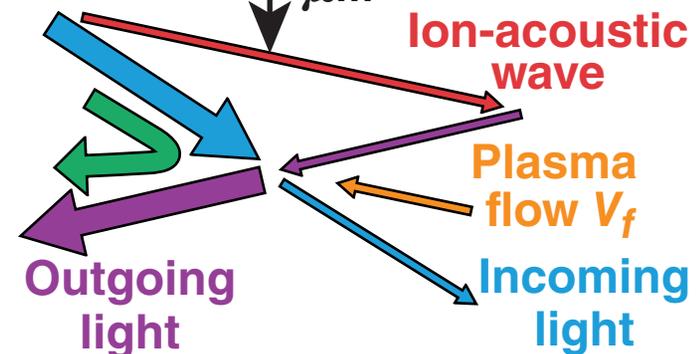
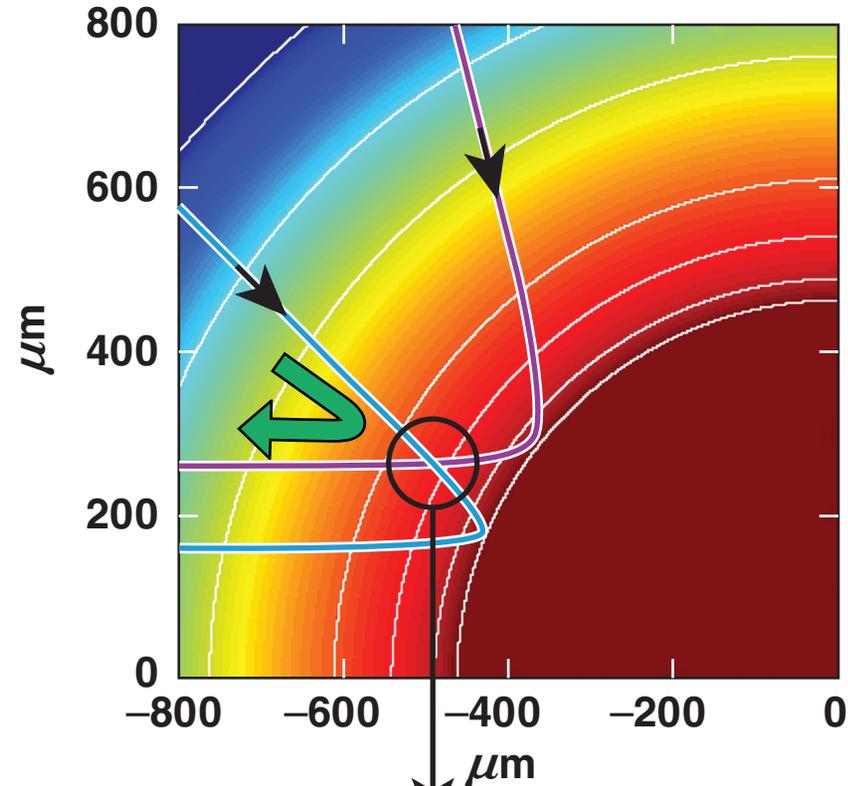
- Ion-acoustic wave (IAW) transfers energy from a “pump” EM wave (light entering plasma) to a “seed” EM wave (light leaving plasma)

$$\omega_{\text{pump}} = \omega_{\text{seed}} + \omega_{\text{IAW}}$$

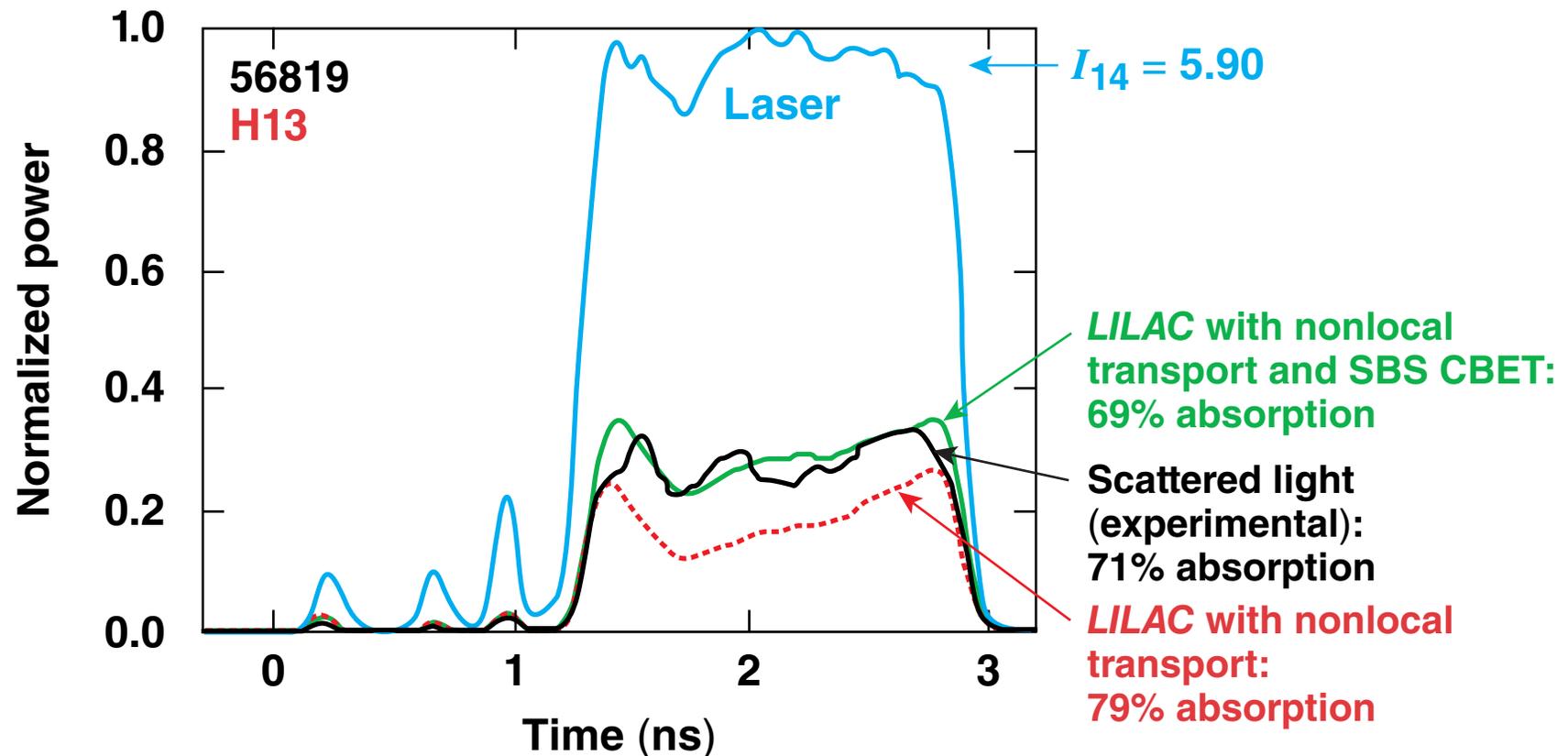
$$\vec{k}_{\text{pump}} = \vec{k}_{\text{seed}} + \vec{k}_{\text{IAW}}$$

$$0 = \pm c_s |k_{\text{IAW}}| + \vec{v}_f \cdot \vec{k}_{\text{IAW}} - \omega_{\text{IAW}}$$

Because the EM-seed amplitude is of the same order as the pump, very small gains of only a few percent can significantly reduce the absorbed energy.



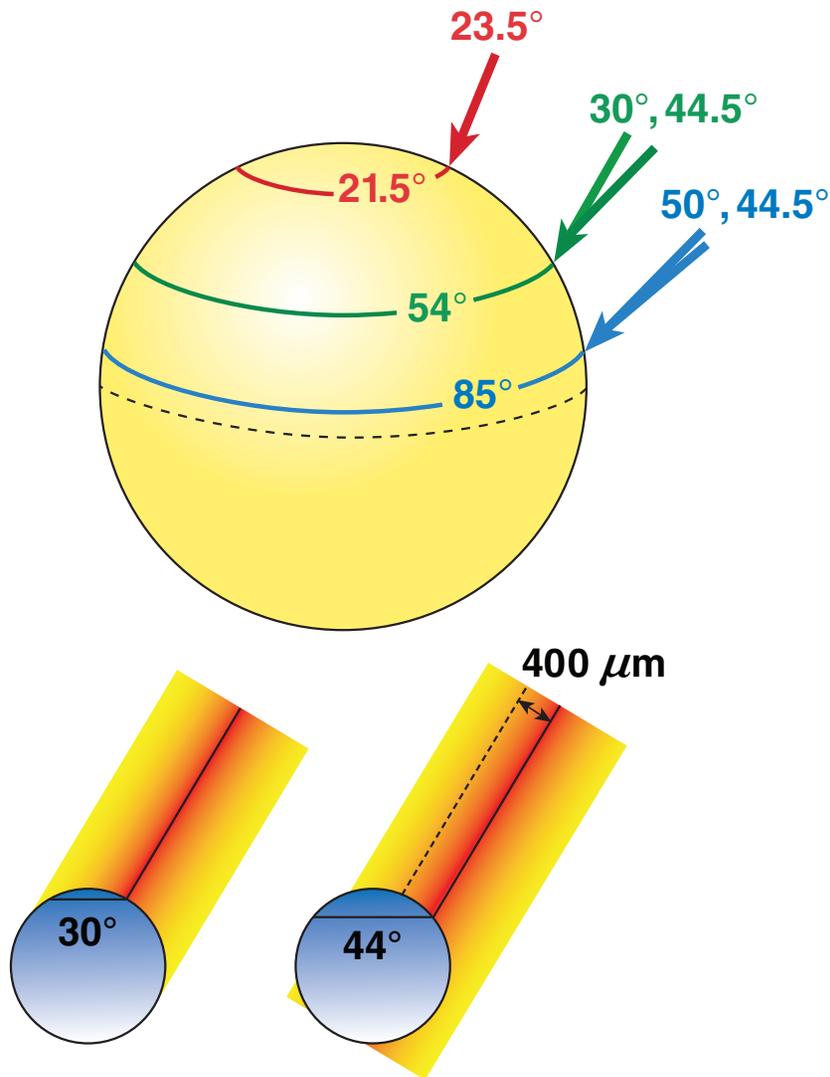
Accurate simulations of direct-drive implosions on OMEGA require CBET in the hydrocode*



*W. Seka, "Reducing the Cross-Beam Energy Transfer in Direct-Drive Implosions Through Laser-Irradiation Control," this conference.

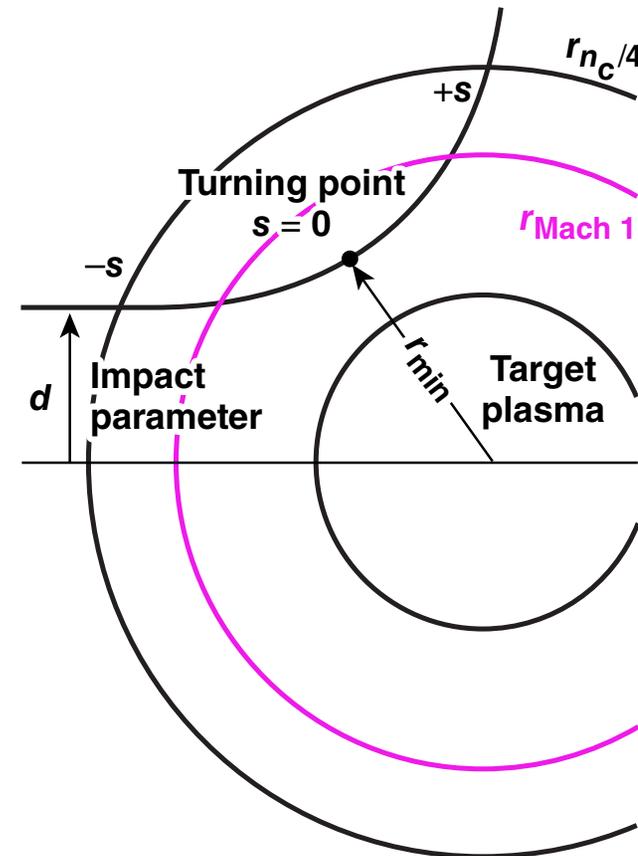
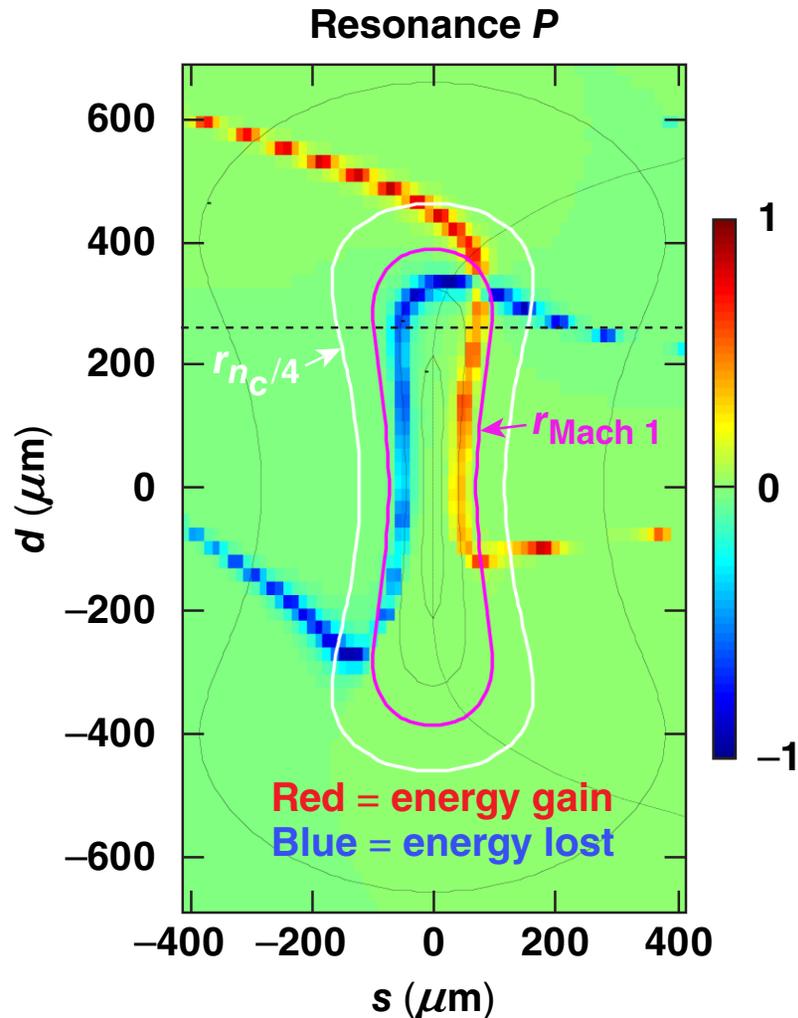
Polar Drive

Direct-drive experiments on the NIF require the nonsymmetric polar-drive geometry



- Quasi symmetric intensity on target is achieved through a combination of spot shape, pulse shape, and beam-pointing control.
- Repointing beams from the x-ray-drive pointing leads to higher angles of incidence at the equator relative to the pole.
- The effects of CBET on beam intensity and uniformity with repointed PD beams is now modeled.

Resonance function* (P) is a measure of how close the conditions are to resonance for SBS cross-beam transfer



CBET resonance:

$$0 = \pm c_s |k_{IAW}| + \vec{v}_f \cdot \vec{k}_{IAW} - \omega_{IAW}$$

For one set of beamlets from one beam crossing one other beam.

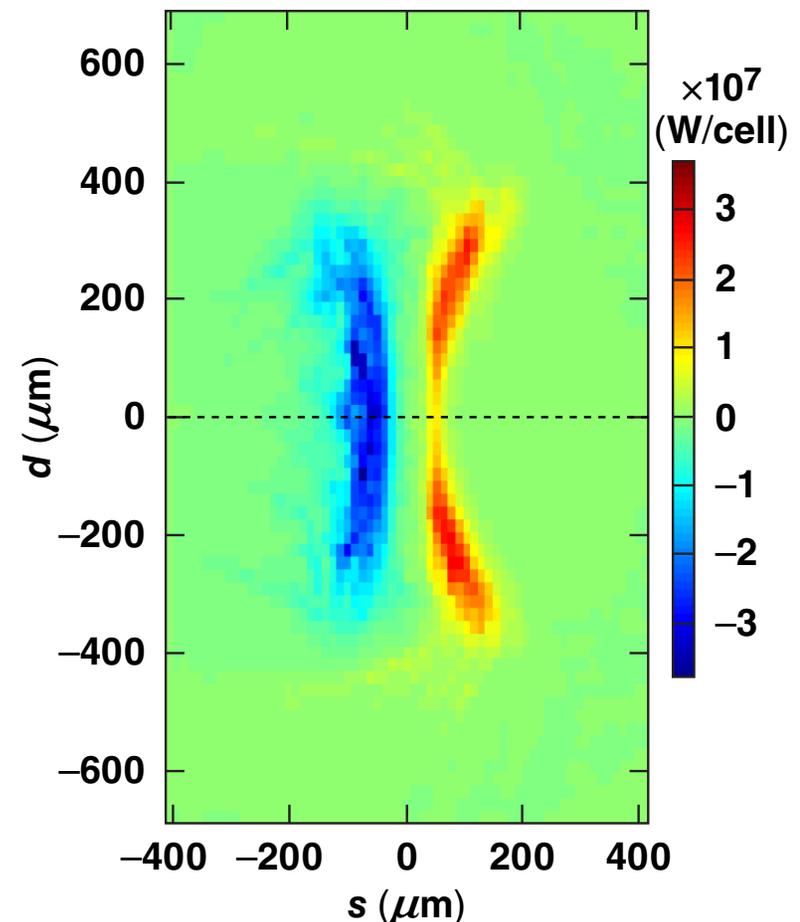
Iterative calculations find the energy lost/gained along each beamlet for a beam

- The strength of the transfer is estimated using the spatial gain length* L_{SBS} for crossing planar waves

$$L_{SBS}^{-1} = 2.8 \times 10^{-2} \frac{1}{v_i \lambda_{0, \mu m}} \frac{n_e/n_c}{\sqrt{1-n_e/n_c}} \frac{I_{14} \lambda_{0, \mu m}^2}{T_{e, \text{keV}} (1+3T_i/ZT_e)} P(\eta) (\mu m^{-1})$$

- The rate of change in intensity due to cross-beam transfer and absorption can be integrated along each path to determine the intensity

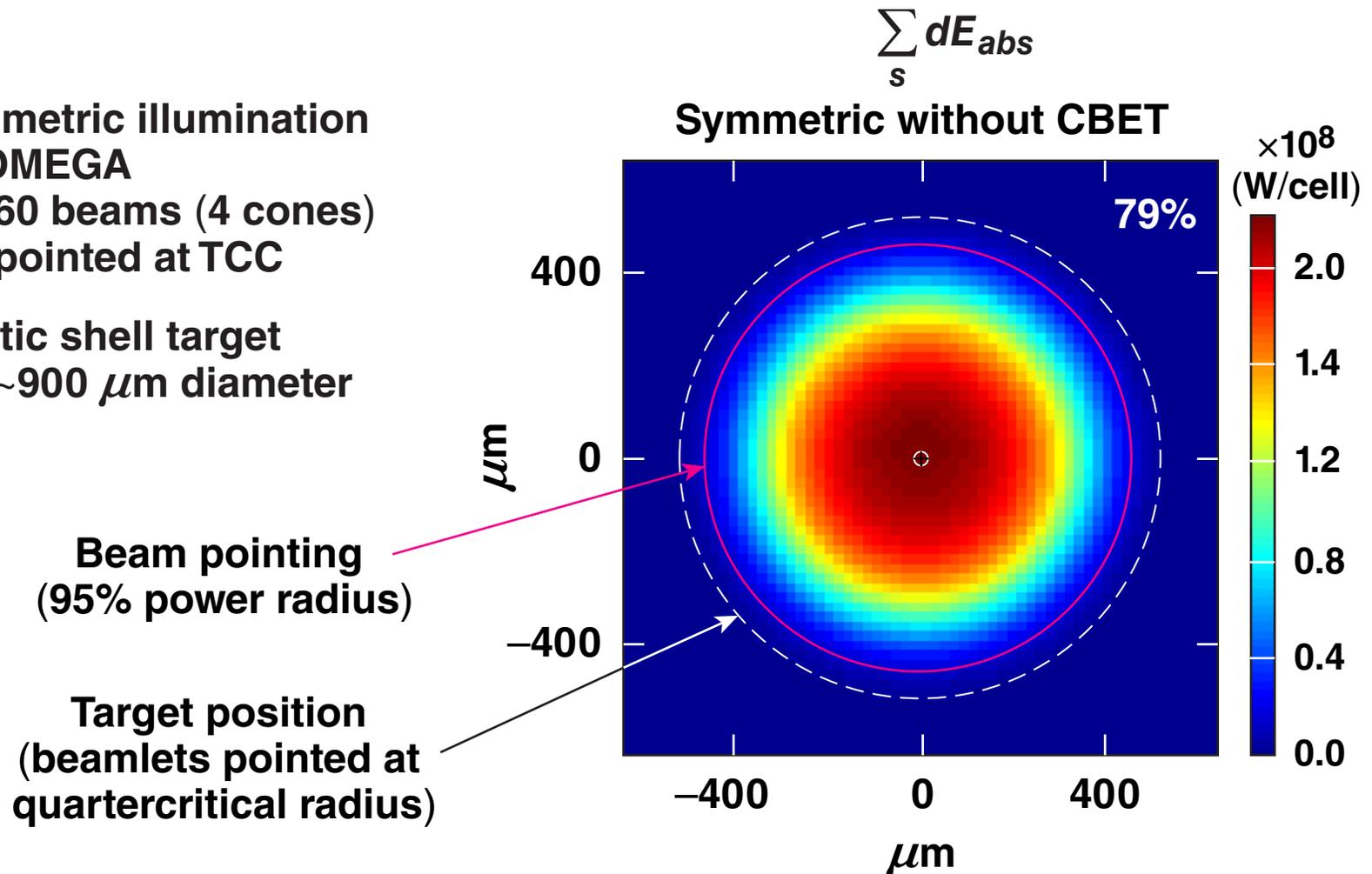
$$d(IA) = -IA \left(\frac{1}{L_{Abs}} + \sum_{\text{all beams}} \frac{1}{L_{SBS}} P \right) ds$$



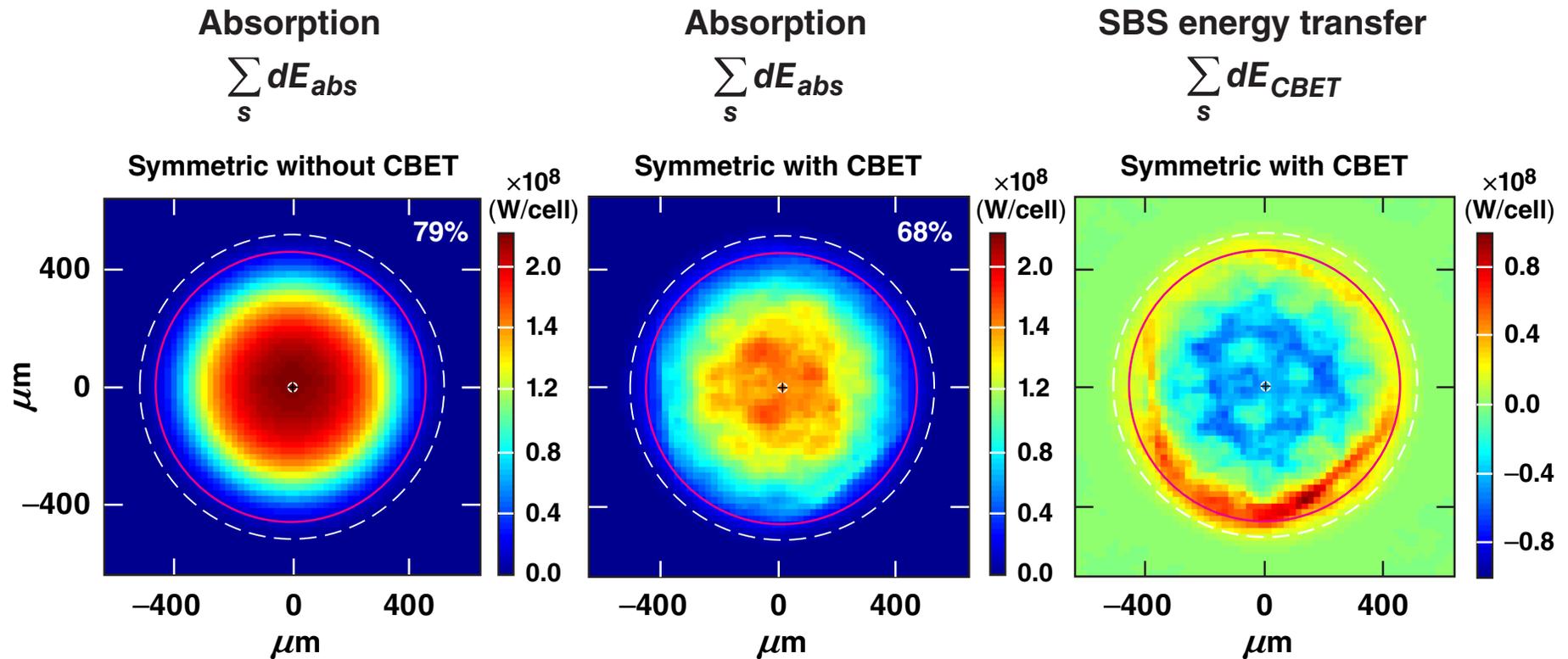
*J. F. Myatt et al., Phys. Plasmas 11, 3394 (2004).
E17997b

The integral of absorbed energy along the beamlet path shows no fine structure with symmetric illumination and without CBET

- Symmetric illumination on OMEGA
 - 60 beams (4 cones)
 - pointed at TCC
- Plastic shell target
 - $\sim 900 \mu\text{m}$ diameter

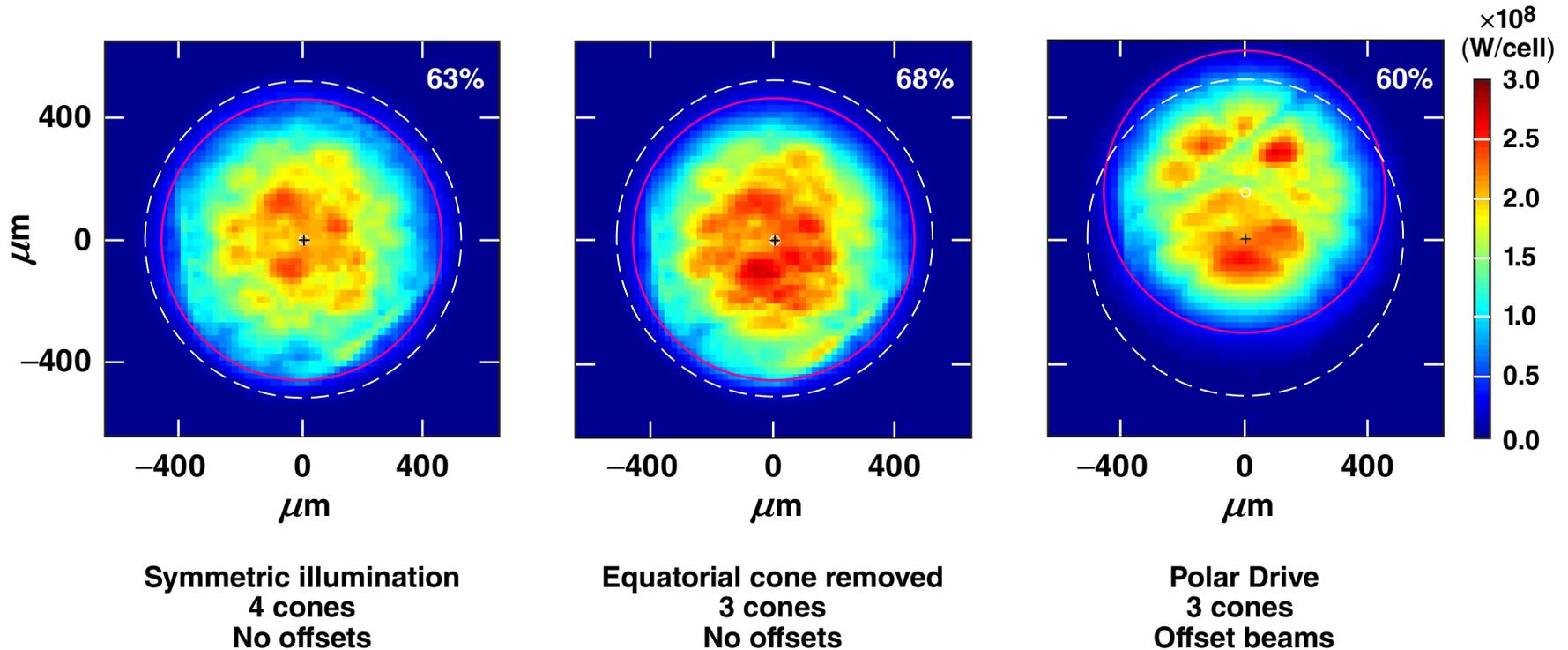


Including CBET with symmetric illumination transfers energy from inner beam and reduces the absorption

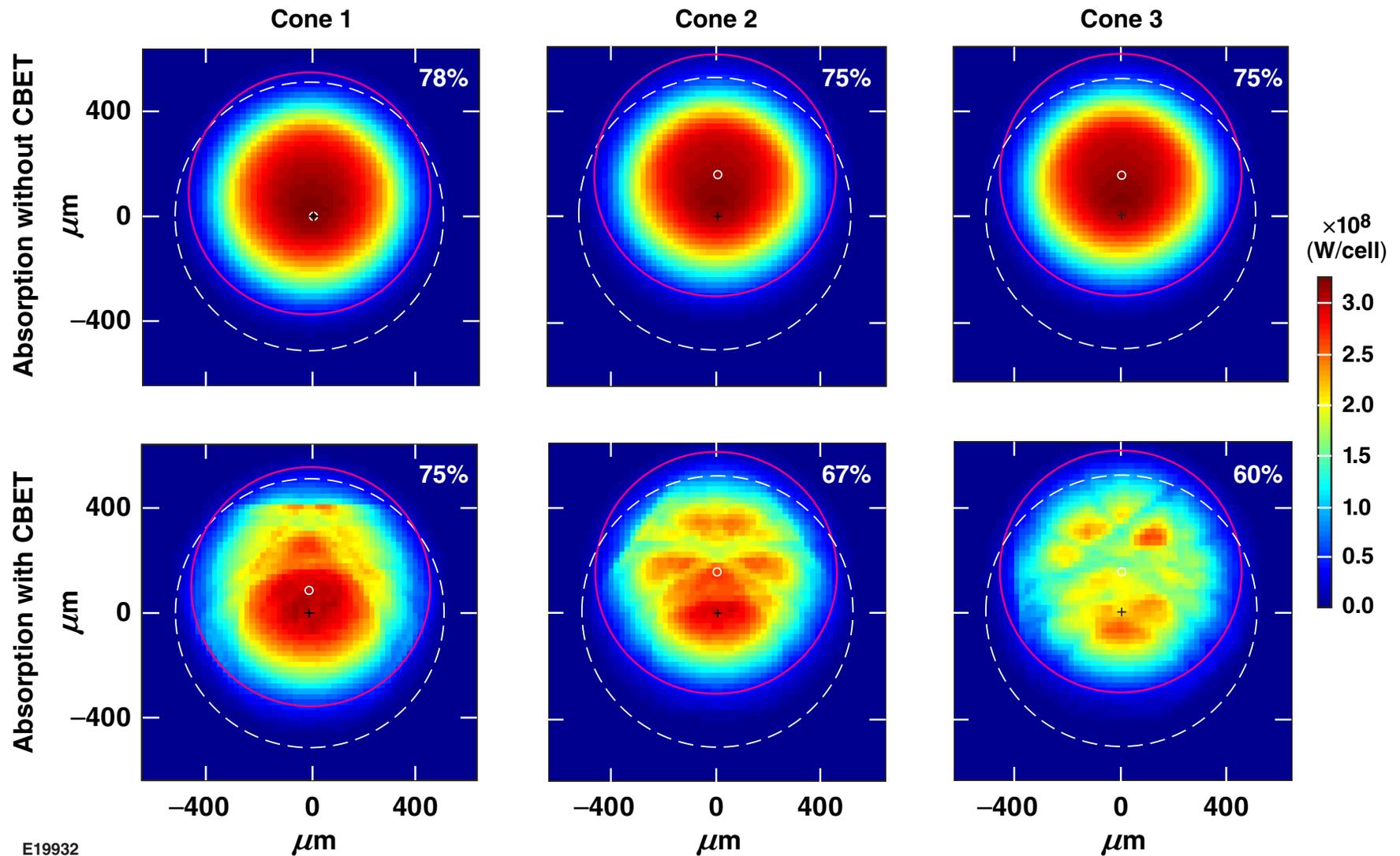


PD affects the absorption primarily through the beam pointing offset

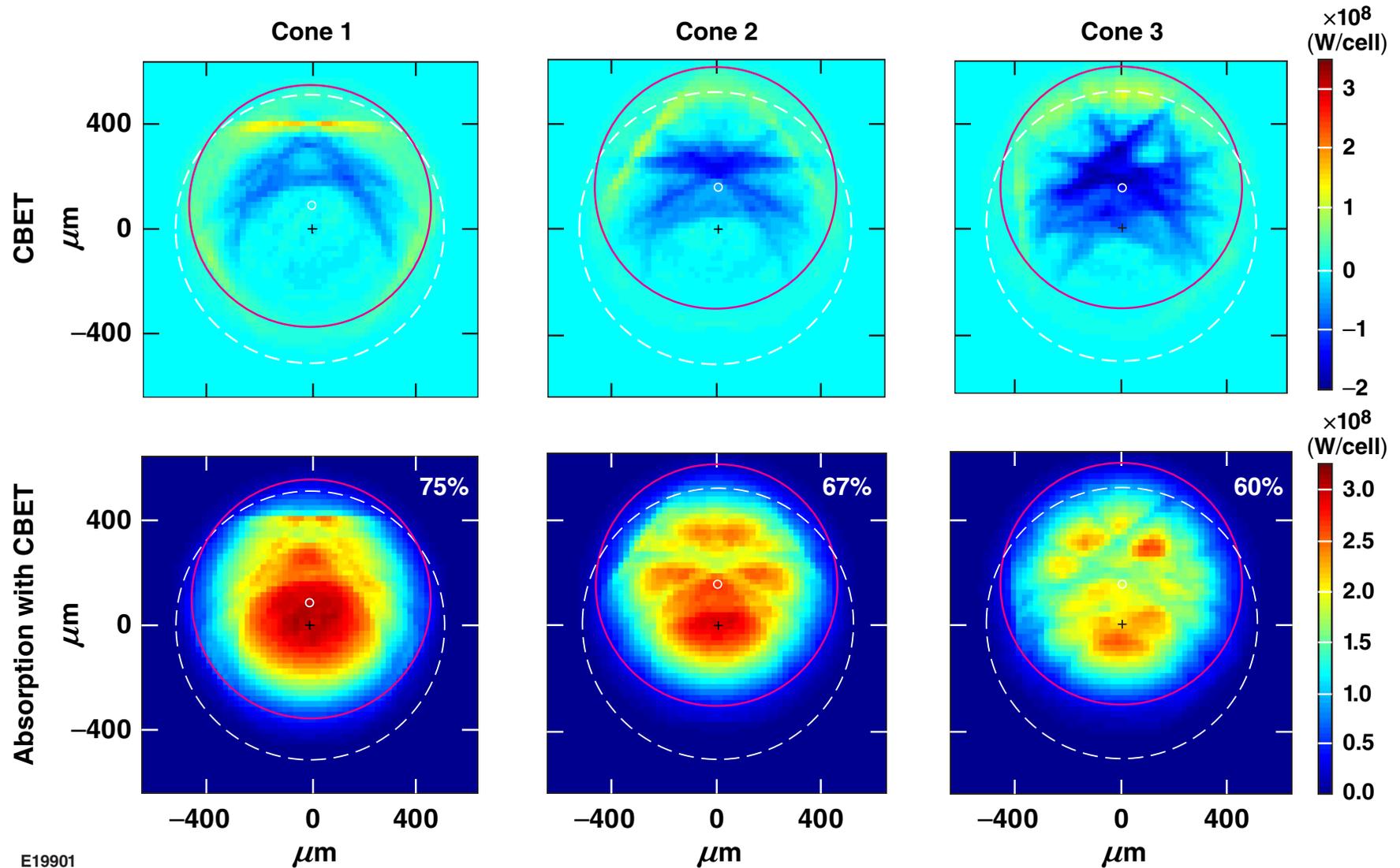
Absorbed energy $\sum_s dE_{abs}$



CBET has a measured effect on the cones closer to the equator



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