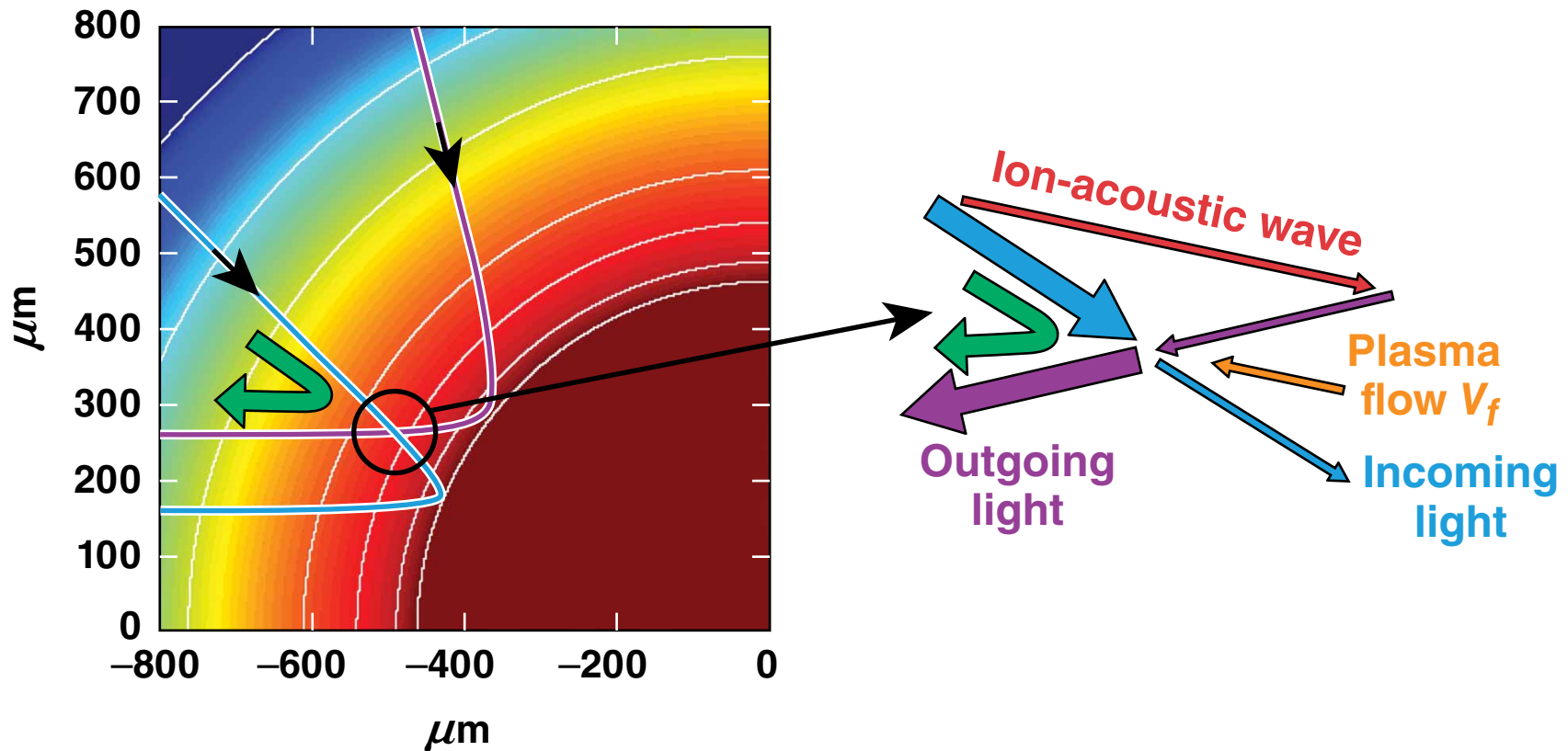


Time-Resolved Scattered-Light Spectroscopy in Direct-Drive-Implosion Experiments



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Summary

Scattered-light spectrum simulations indicate that anomalous absorption affects the latter part of implosions



- Time-dependent scattered-laser-light spectra in the SBS range (351 ± 1 nm) are modeled by a combination of hydrodynamic and ray-tracing codes
- Most features observed in the scattered-light spectra are well reproduced by the modeling
- The largest discrepancy in the modeling suggests that absorption is over-estimated in the later part of the pulse, but scaling the total absorption to match observations still does not accurately reproduce the spectra
- Cross-beam transfer of energy out of the beam-profile center might be the physical process behind the discrepancy

Collaborators



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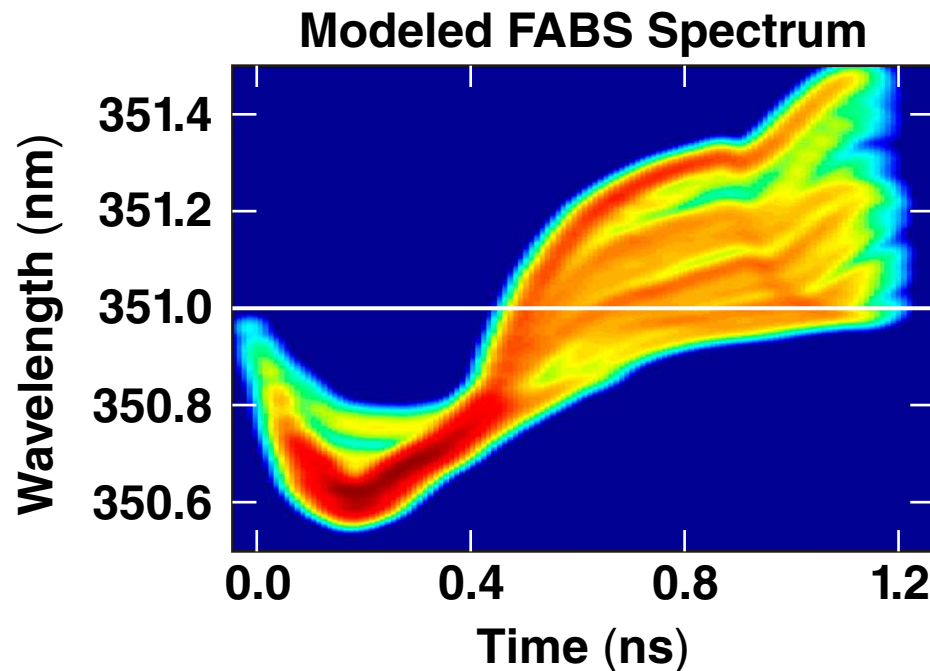
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Modeled Spectra

Time-dependent scattered-laser-light spectra in the SBS range (351 ± 1 nm) are modeled for OMEGA implosions



20- μ m plastic shell
1-ns square pulse

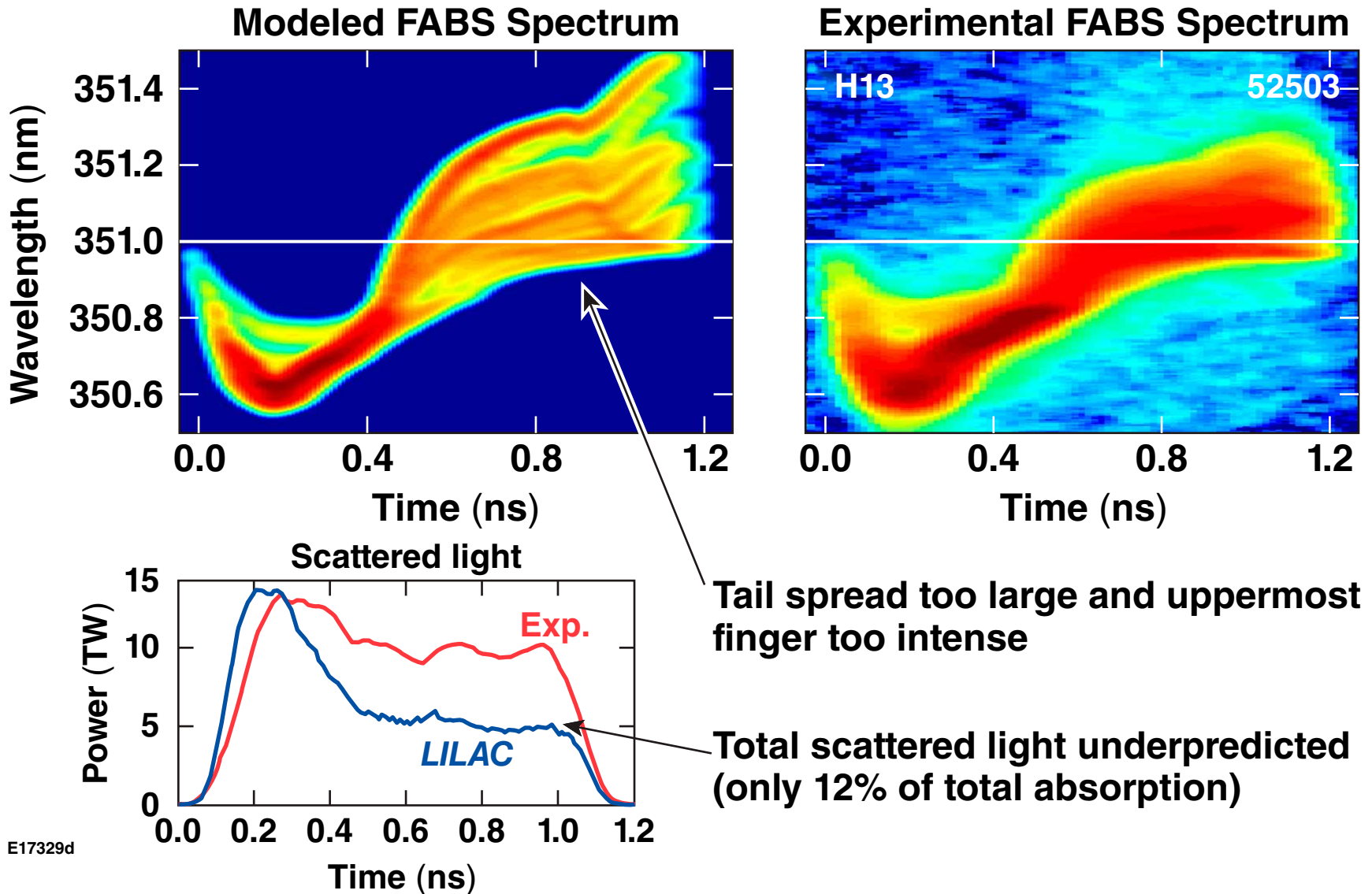
- A combination of codes is used
 - *LILAC*¹: 1-D hydrodynamic code predicts time-dependent implosion profiles
 - *SAGERAYS*²: Ray traces laser light through the corona and calculates spectral shift³
 - *MATLAB* code calculates total spectrum collected from all 60 OMEGA beams

¹J. A. Delettrez *et al.*, Phys. Rev. A **36**, 3926 (1987).

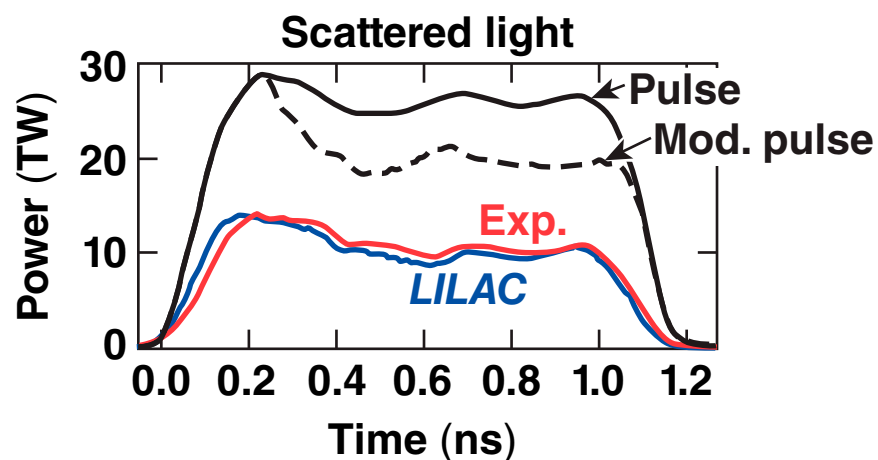
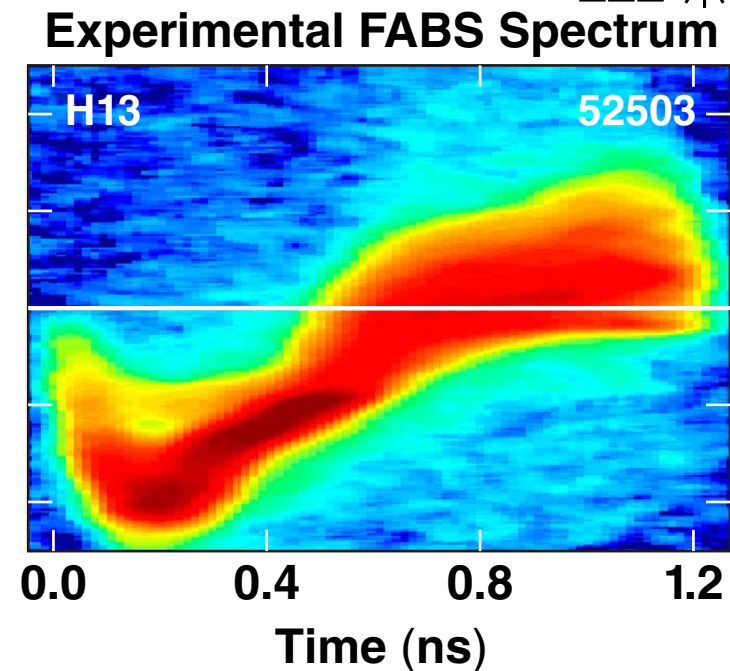
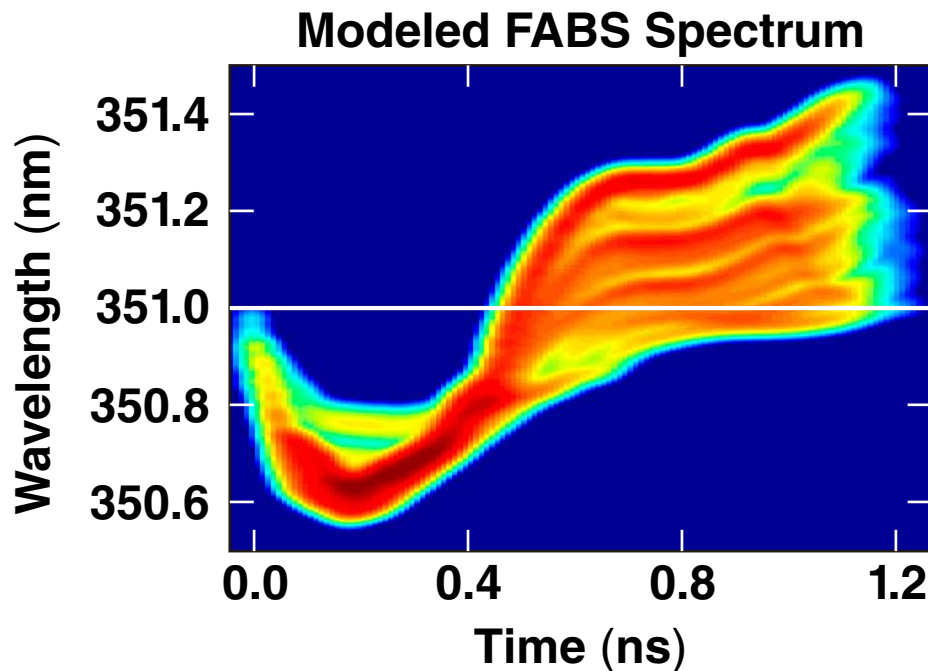
²R. S. Craxton and R. L. McCrory, J. Appl. Phys. **56**, 108 (1984).

³T. Dewandre, J. R. Albritton, and E. A. Williams, Phys. Fluids **24**, 528 (1981).

Modeled spectra show all the basic structures of the experimental spectra but differ in some details



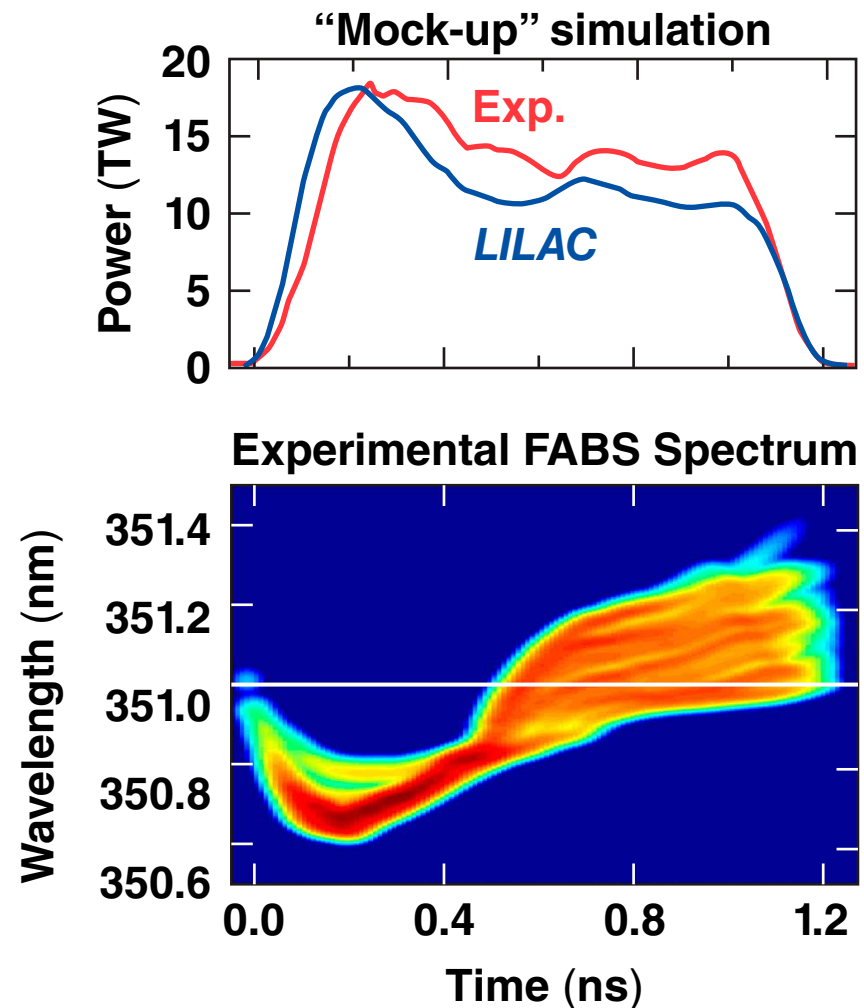
Modeling with the pulse power scaled to reproduce the observed time-dependent absorption does not significantly improve the spectral shift predictions



Where the energy is absorbed seems important, not just how much is absorbed.

Cross-beam transfer of energy from the beam profile center toward the profile edge might be consistent with the observations

- Removes energy from rays closest to center of beam profile that penetrate furthest towards the critical surface and are responsible for the uppermost finger of the spectrum tail
- Redistributes that energy to rays farther out in the beam profile where absorption is less
- Should result is a spectrum that better matches observations
 - removes energy from the uppermost finger
 - decreases total absorption/ increases total scattered energy



Cross-Beam Power Transfer

EM-seeded SBS cross-beam power transfer might cause some laser energy to “bypass” the high-absorption zone



- Ion-acoustic wave (IAW) transfers energy from a “pump” EM wave to a “seed” EM wave

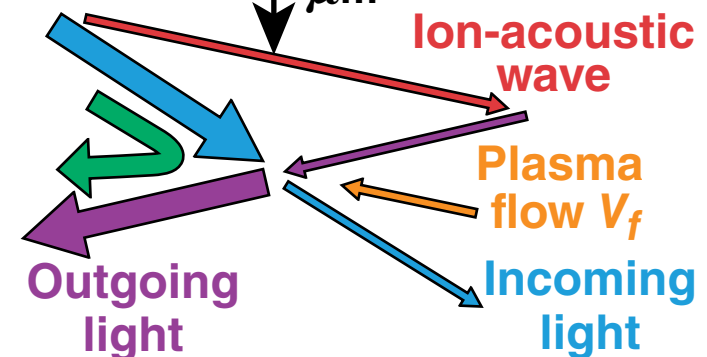
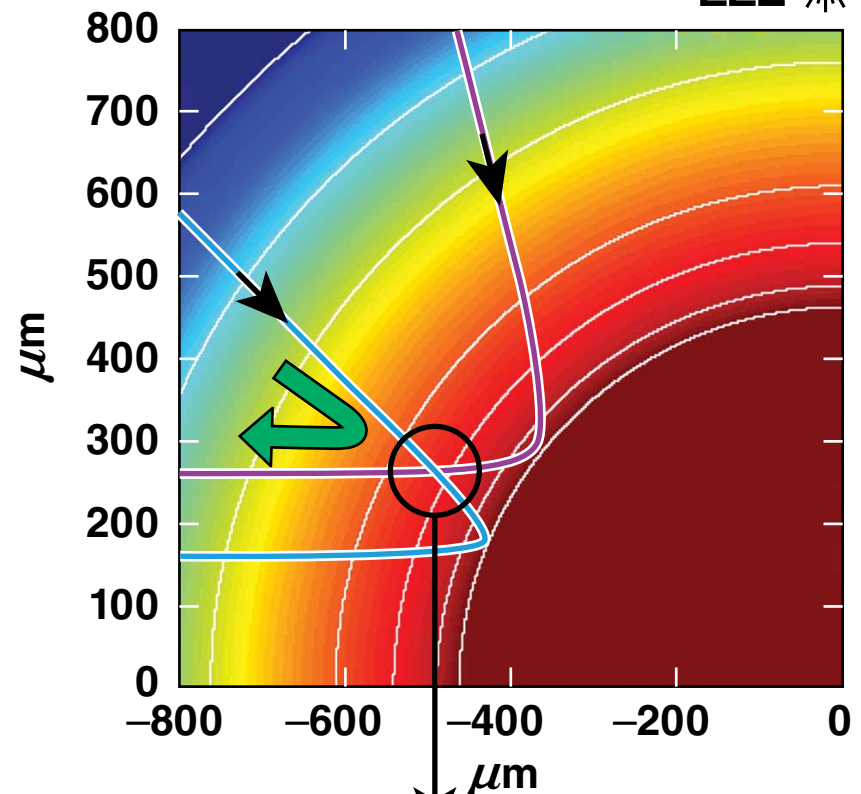
$$\omega_{\text{pump}} = \omega_{\text{seed}} + \omega_{\text{IAW}}$$

$$\vec{k}_{\text{pump}} = \vec{k}_{\text{seed}} + \vec{k}_{\text{IAW}}$$

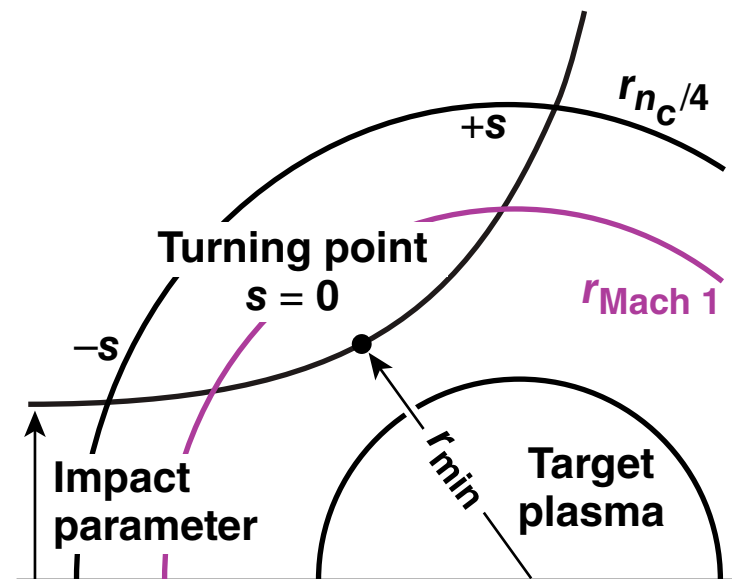
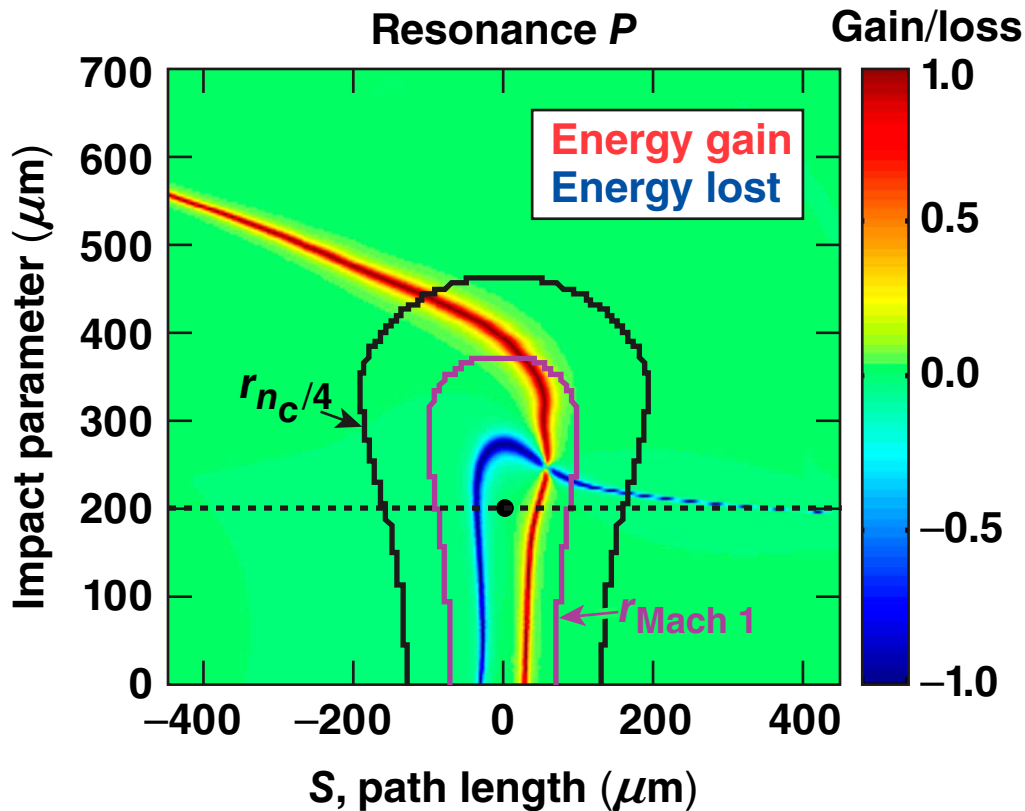
$$0 = \pm c_s |k_{\text{IAW}}| + \vec{v}_f \cdot \vec{k}_{\text{IAW}} - \omega_{\text{IAW}}$$

- Light entering the plasma can transfer energy to light that is leaving the plasma so that some laser energy “bypasses” the high-absorption region, reducing the total absorbed power

Because the EM seed amplitude is of the same order as the pump, very small gains of only a few percent could significantly affect the absorbed energy.



Beamlet crossings calculated from ray-trace and OMEGA beam geometry indicate that energy is typically lost by incoming beamlets



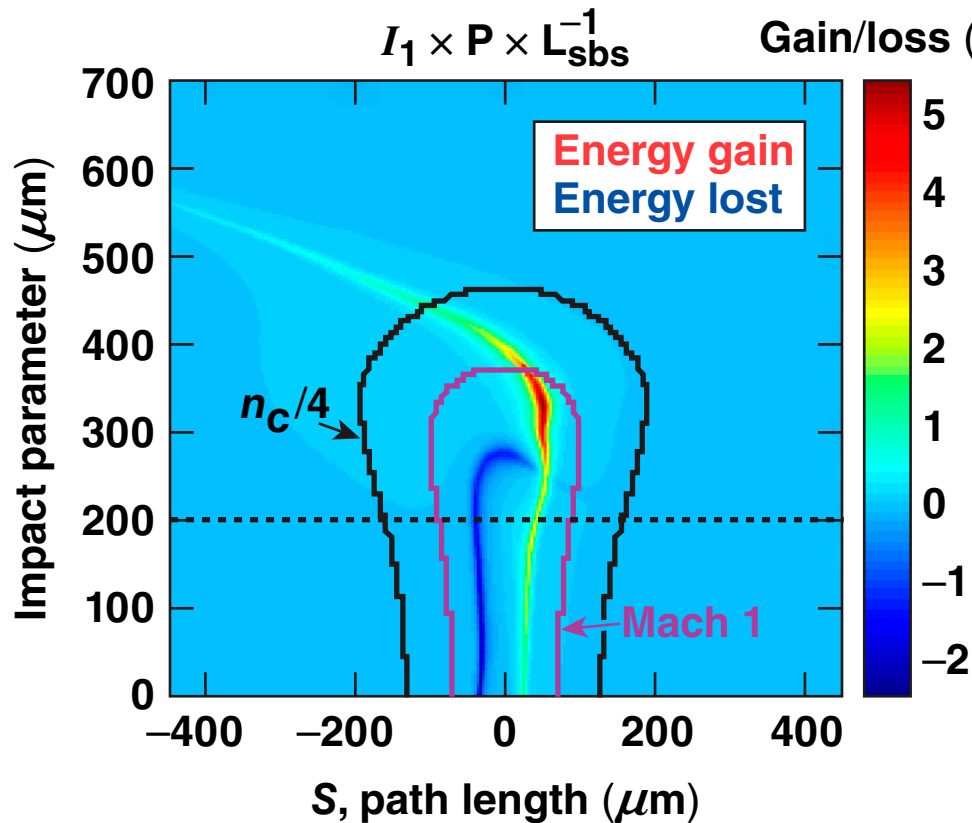
For one set of beamlets from one beam crossing the reference beam is at 40°

Resonance function* (P) is a measure of how close the conditions are to resonance for SBS cross-beam transfer

$$0 = \pm c_s |k_{IAW}| + \vec{v}_f \cdot \vec{k}_{IAW} - \omega_{IAW}$$

*C. J. Randall, J. R. Albritton, and J. J. Thomson, Phys. Fluids 24, 1474 (1981).

The strength of the transfer is estimated using the spatial gain length* L_{SBS} for crossing planar waves

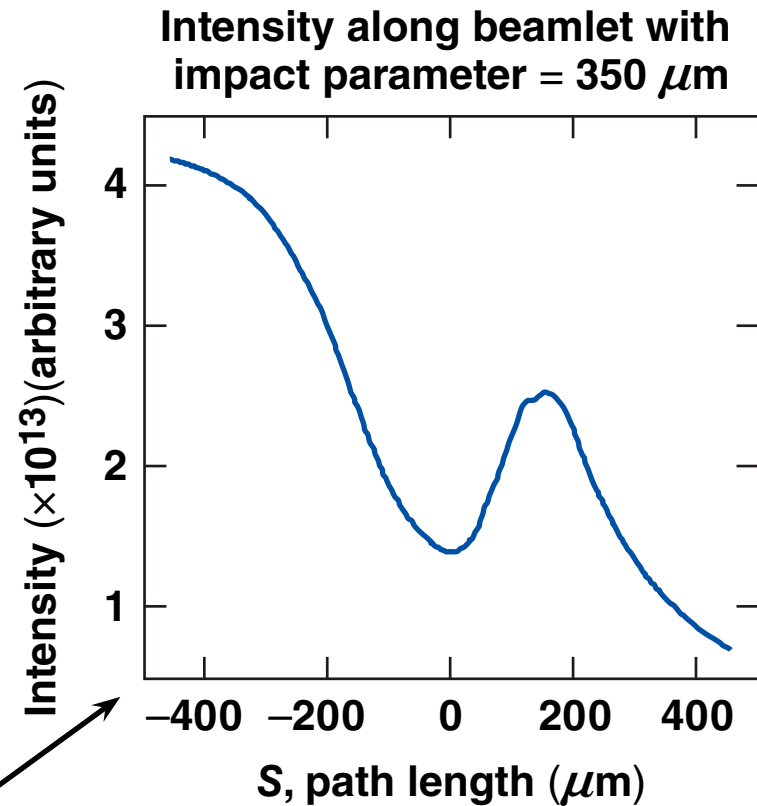
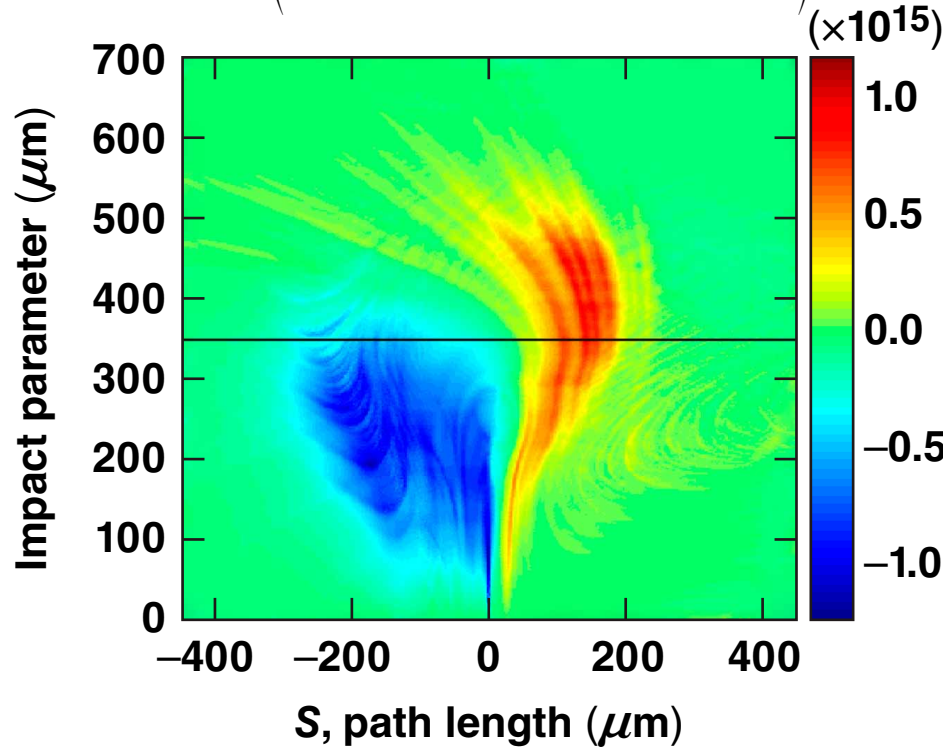


For one set of beamlets from one beam crossing, the reference beam is at 40°

$$L_{\text{SBS}}^{-1} = 2.8 \times 10^{-2} \frac{1}{\nu_i \lambda_{0,\mu\text{m}}} \frac{n_e/n_c}{\sqrt{1-n_e/n_c}} \frac{I_{14} \lambda_{0,\mu\text{m}}^2}{T_{e,\text{keV}} (1+3T_i/ZT_e)} P(\eta) (\mu\text{m}^{-1})$$

Calculating the energy lost/gained along each beamlet supports the transfer of energy out of beam center

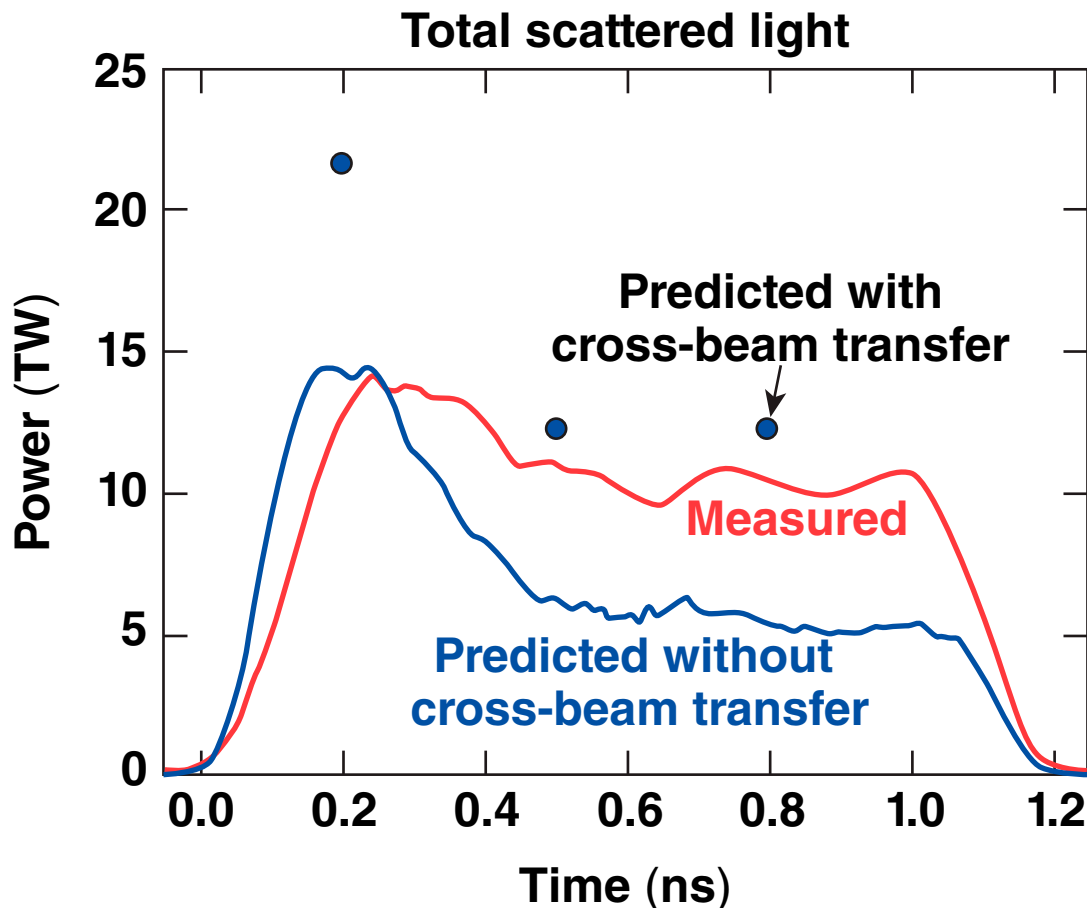
$$d(IA) = -IA \left(\frac{1}{L_{\text{abs}}} + \sum_{\text{all beams}} \sum_{\varphi} \frac{1}{L_{\text{SBS}}} P \right) ds$$



The rate of change in intensity caused by cross-beam transfer and absorption can be integrated along each path to determine the intensity

Summed over all sets of beamlets from all beams crossing the reference beam

Cross-beam transfer scattered-light modeling improves the match to experimental data later in the implosion



- Early in the implosion modeling now shows too much scattered light
- Integrating cross-beam transfer into the hydrocode may improve the agreement

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