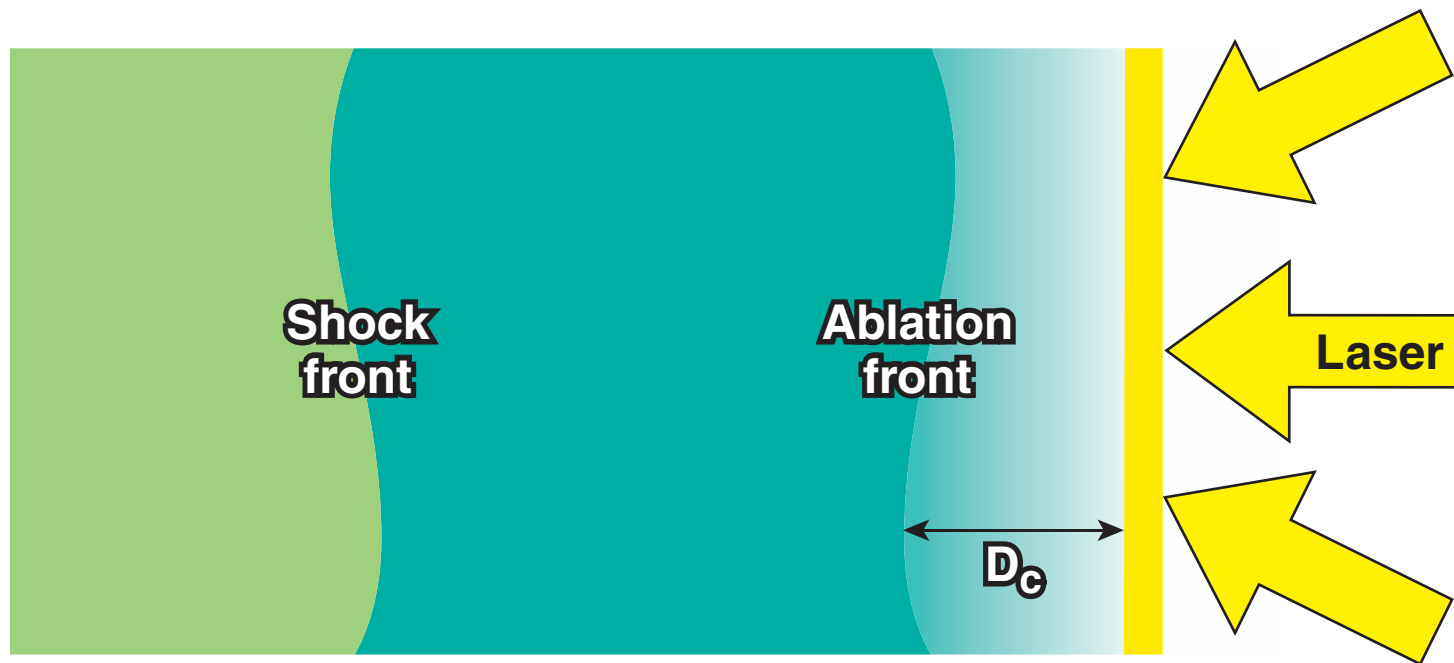


Ablative Richtmyer–Meshkov Instability as a Test of Thermal Conduction Models Used in Hydrosimulations of ICF Experiments



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Summary

Ablative Richtmyer–Meshkov (RM) evolution is sensitive to coronal conditions

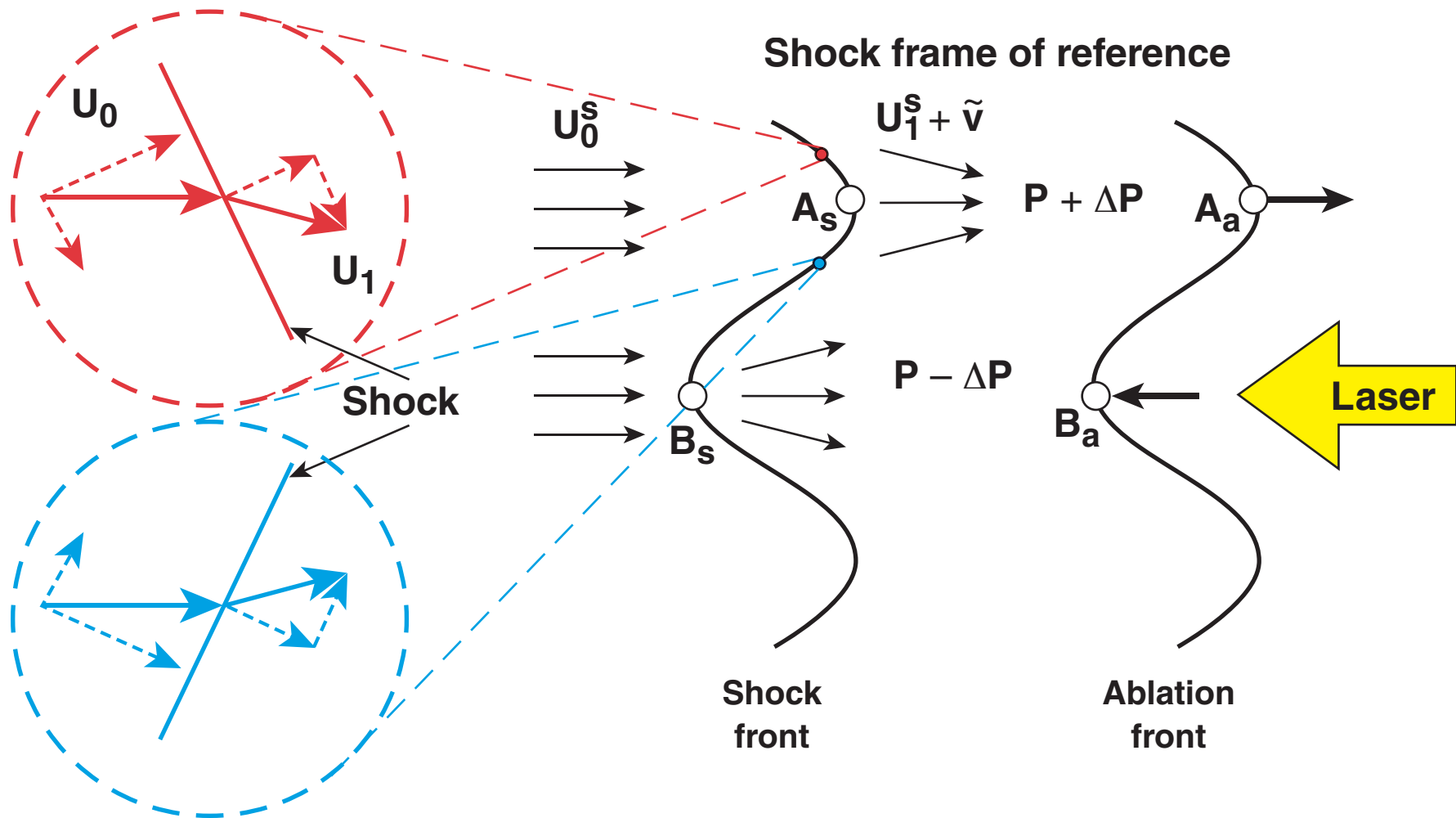


- In the presence of mass ablation, the RM evolution depends on the size of conduction zone D_c
 - $kD_c \ll 1$, ablation front is unstable (Landau-Darrieus instability)
 - $kD_c > 1$, ablation front is stable, perturbations oscillate* with $\omega = k\sqrt{V_a V_{bl}}$
- V_{bl} and D_c depend on the thermal transport models
- Shock velocity and RM measurements are consistent with nonlocal and time-dependent flux-limiter models

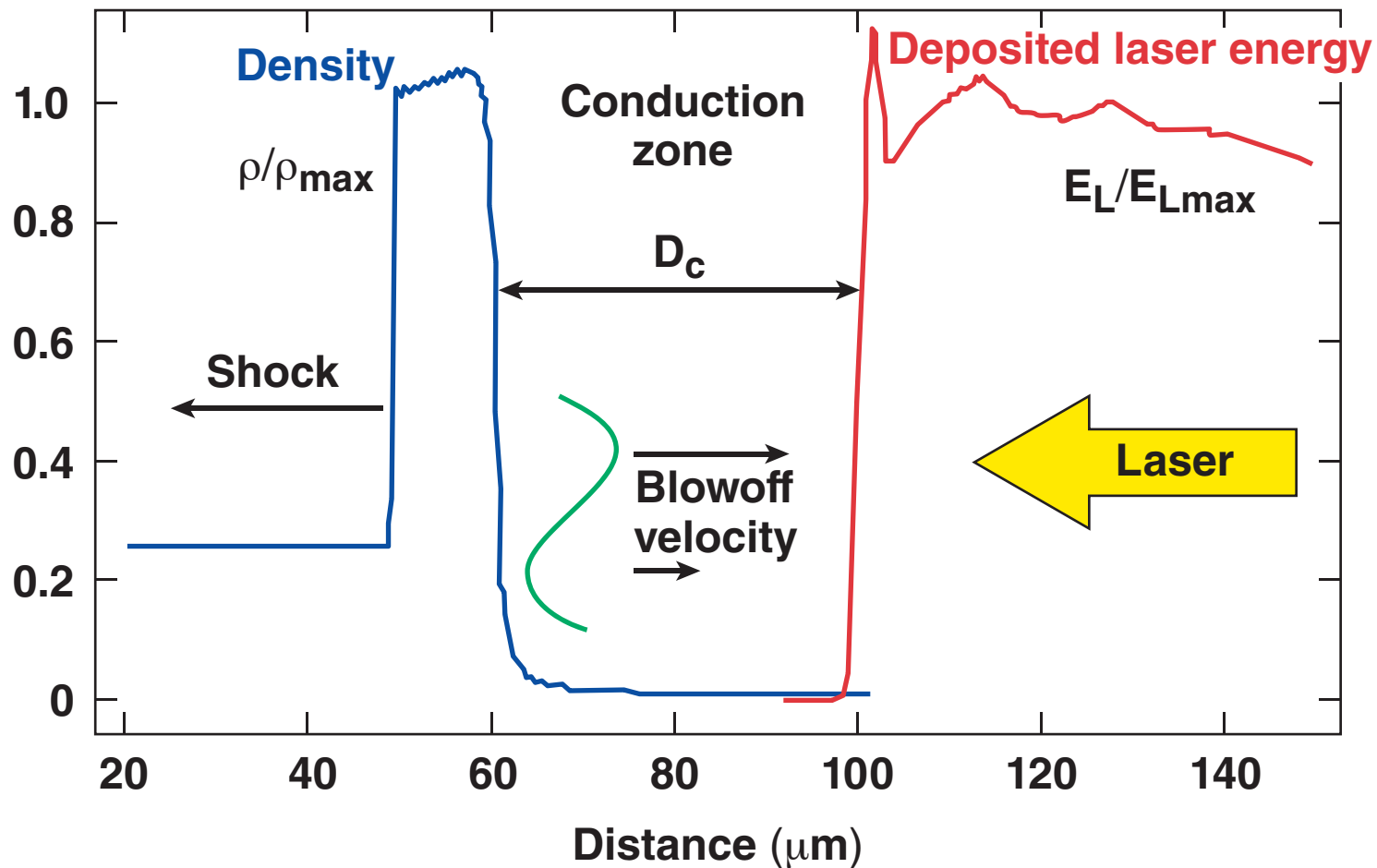
RM evolution is a good test of hydrocodes

*A. L. Velikovich *et al.*, Phys. Plasmas **5**, 1491 (1998); V. N. Goncharov, Phys. Rev. Lett. **82**, 2091 (1999); Y. Aglitskiy *et al.*, Phys. Rev. Lett. **87**, 265001 (2001).

During the shock propagation, conditions arise for Richtmyer–Meshkov-like instability



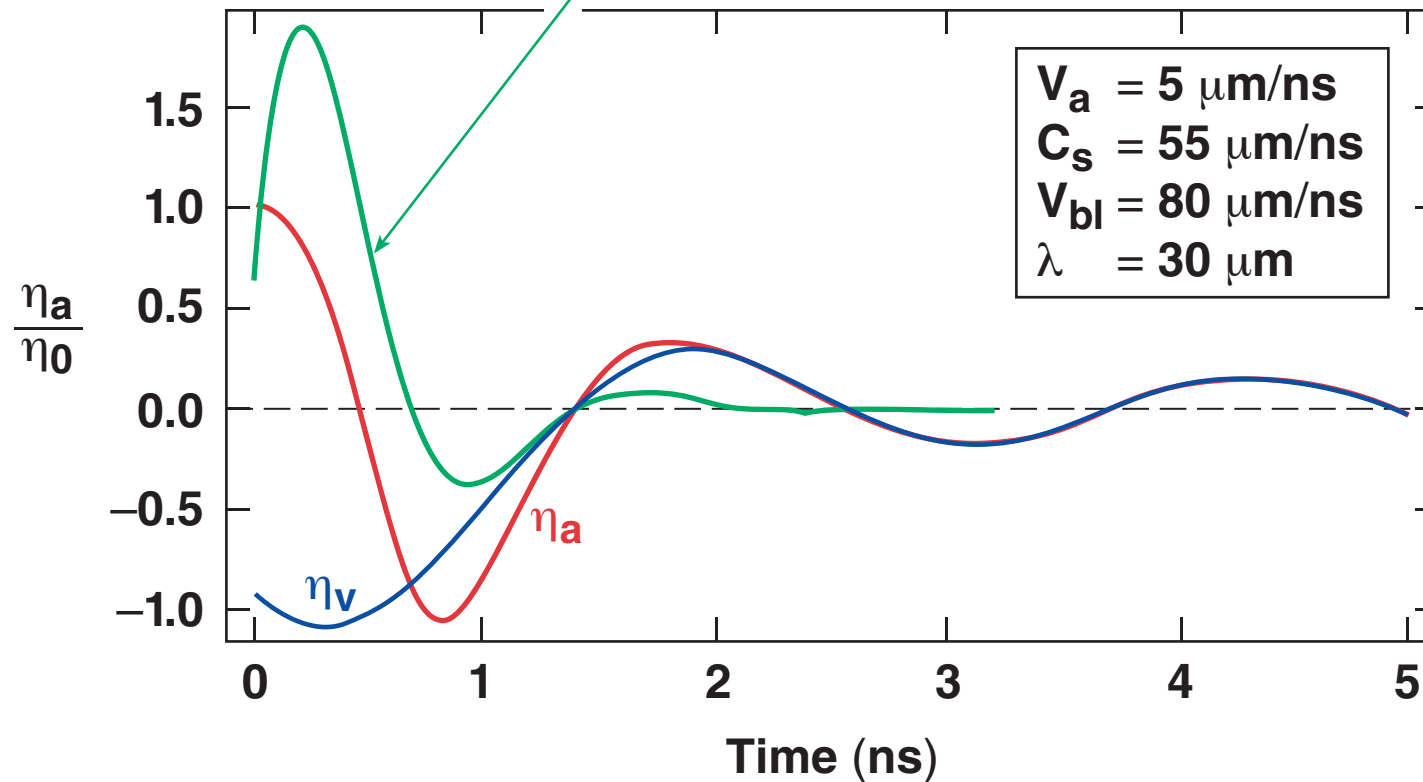
The sharp-boundary model* assumes that perturbation is localized inside the conduction zone



Ablation-front perturbations asymptotically oscillate in time

$$\eta_a = \eta^{cl}(t) - \eta_{\infty}^{cl}(t) + e^{-2kV_a t} (\alpha \cos \omega t + \beta \sin \omega t) + \eta_v(t)$$

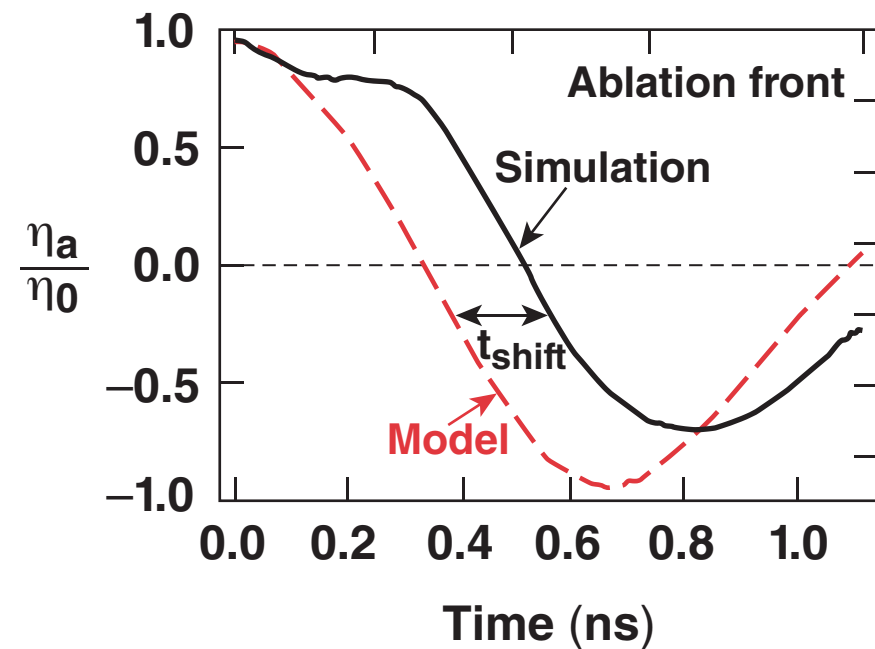
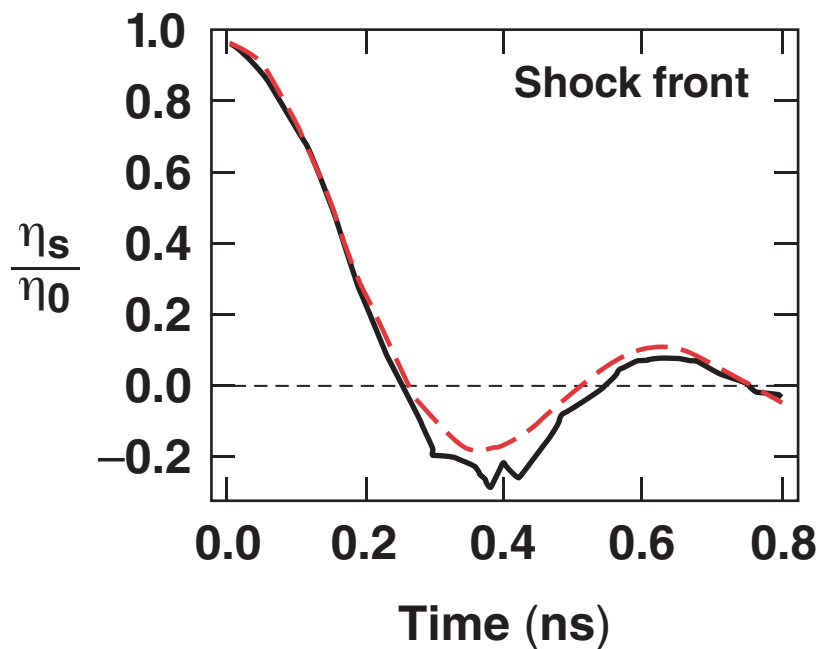
Vorticity set by shock



$$\omega = k \sqrt{V_a V_{bl}}, V_{bl} \propto V(1/k)$$

A detailed comparison between the model and simulation results shows a discrepancy at the beginning of the pulse

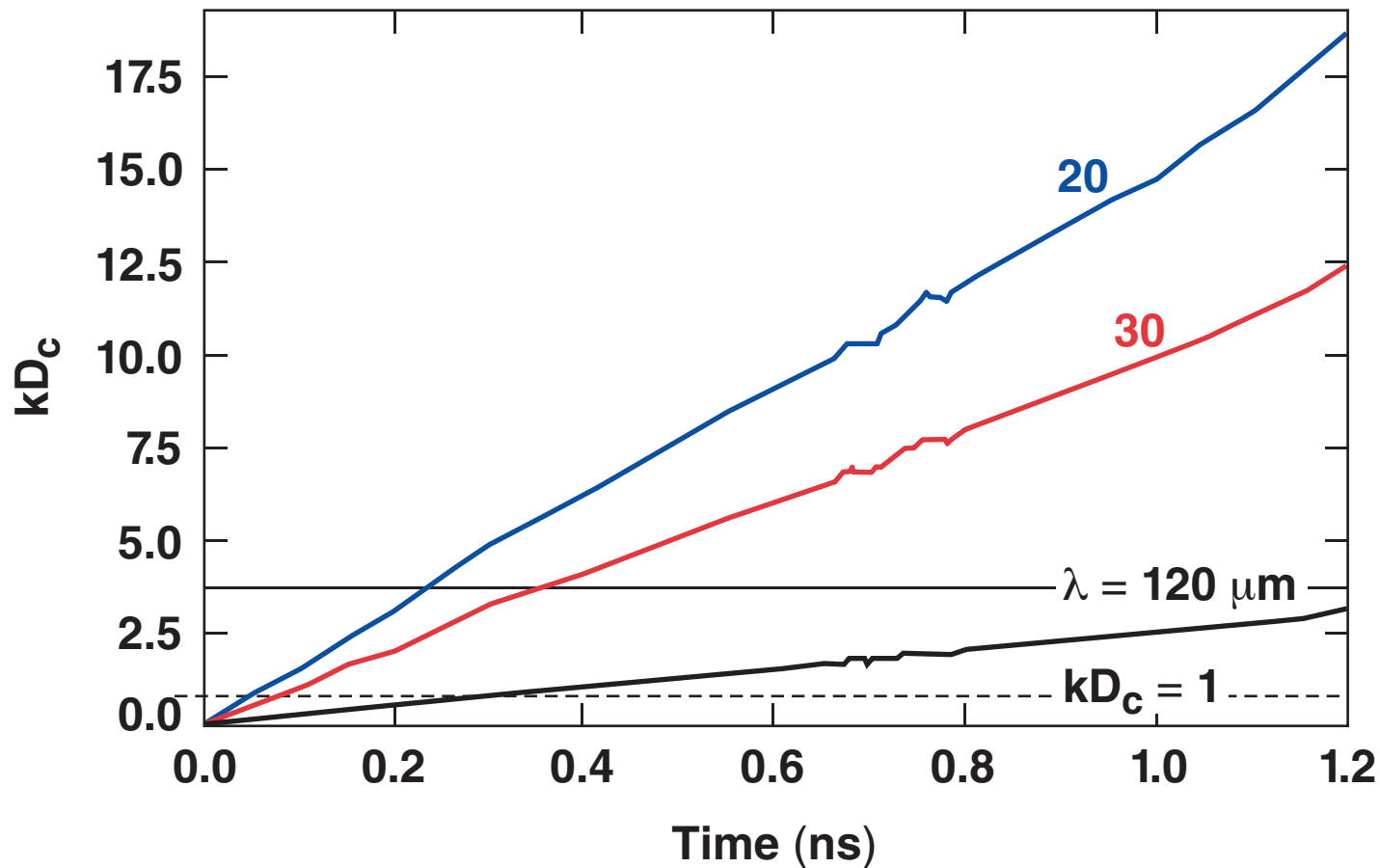
- DD foil, $I = 4 \times 10^{14} \text{ W/cm}^2$, $\lambda = 30 \text{ }\mu\text{m}$



The discrepancy between model and simulation is due to a small conduction zone at the beginning of the pulse



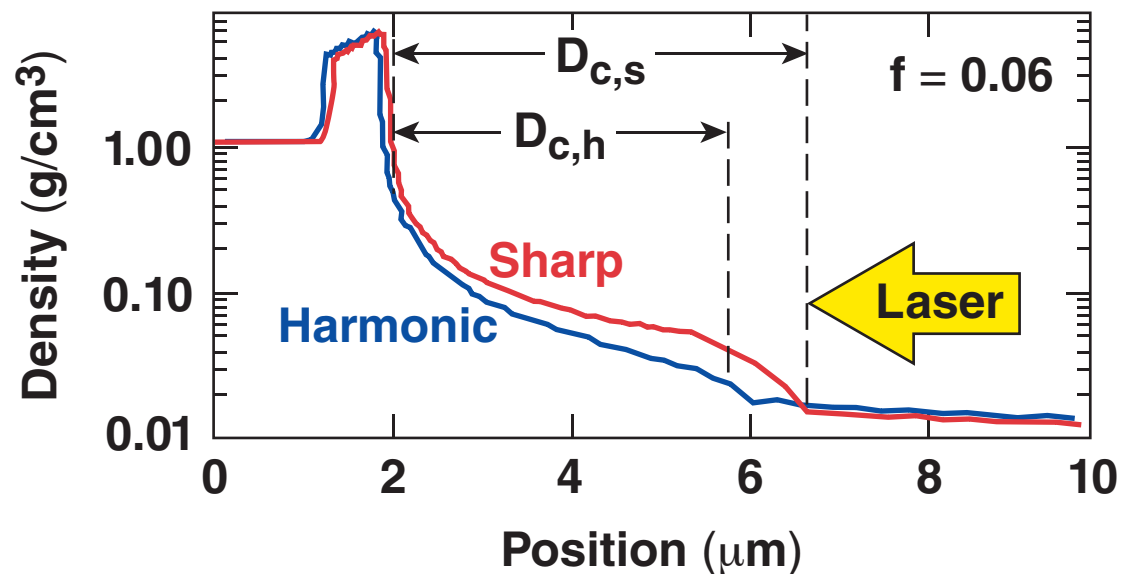
- DD foil, $I = 4 \times 10^{14} \text{ W/cm}^2$



The location of the phase reverse depends on the size of the conduction zone.

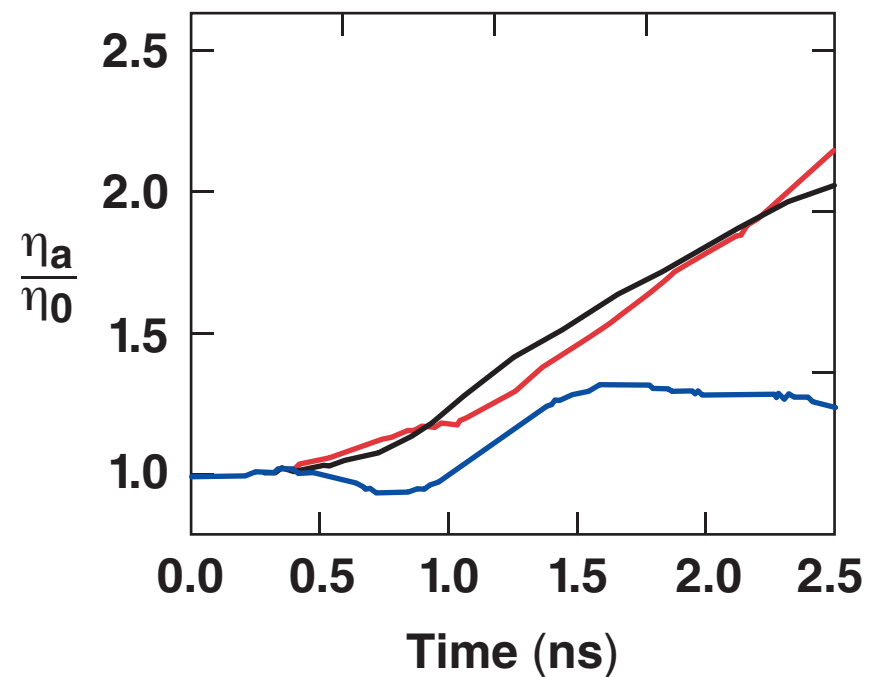
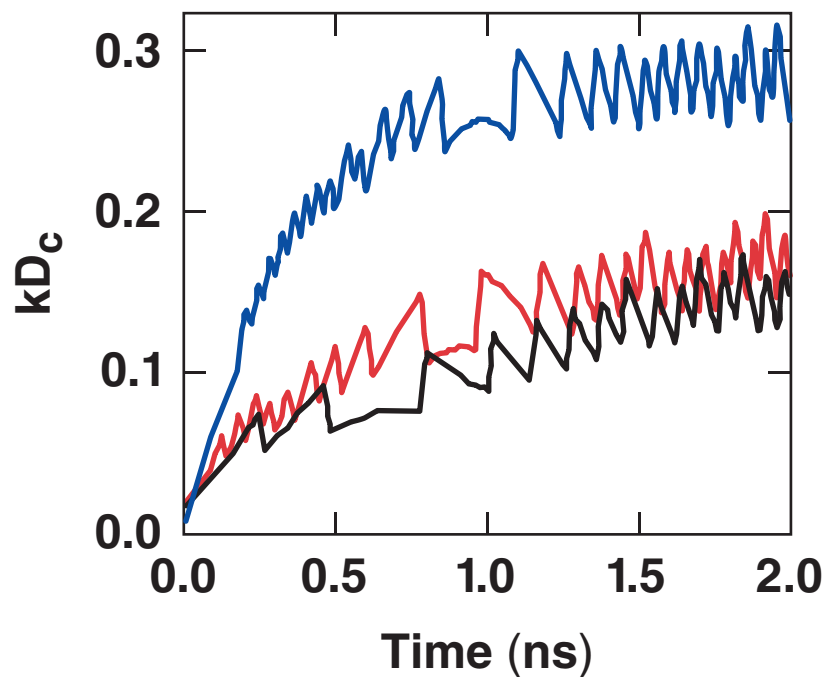
The size of the conduction zone depends on the thermal transport model

- Flux-limited thermal transport* has traditionally been used in hydrocodes
- $q_{SH} = -\kappa \nabla T$ $q_{FS} = nTV_T$
- Sharp cutoff $q_{eff} = \min(q_{SH}, fq_{FS})$ $0.04 < f < 0.1$
- Harmonic $q_{eff} = \frac{q_{SH} fq_{FS}}{q_{SH} + fq_{FS}}$



The ablation front is unstable in the case of a reduced conduction zone

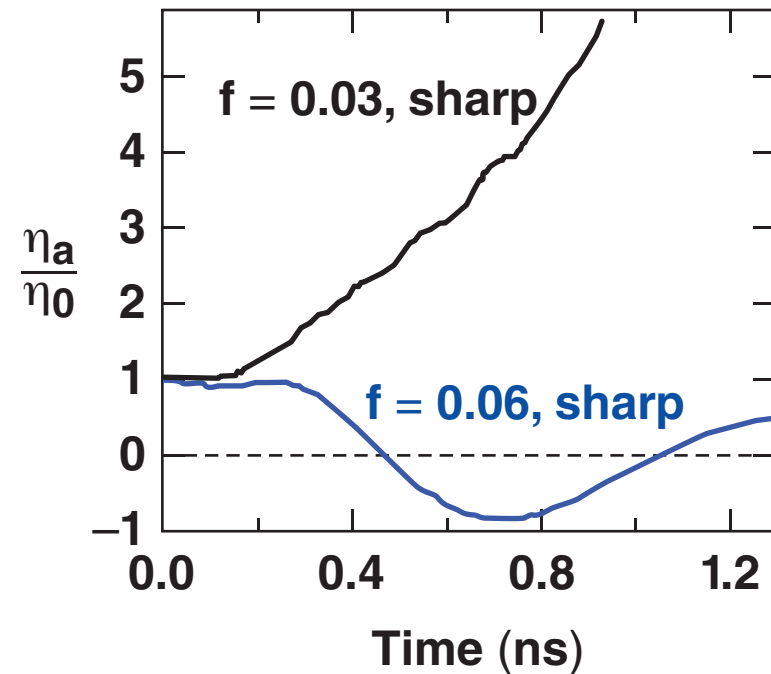
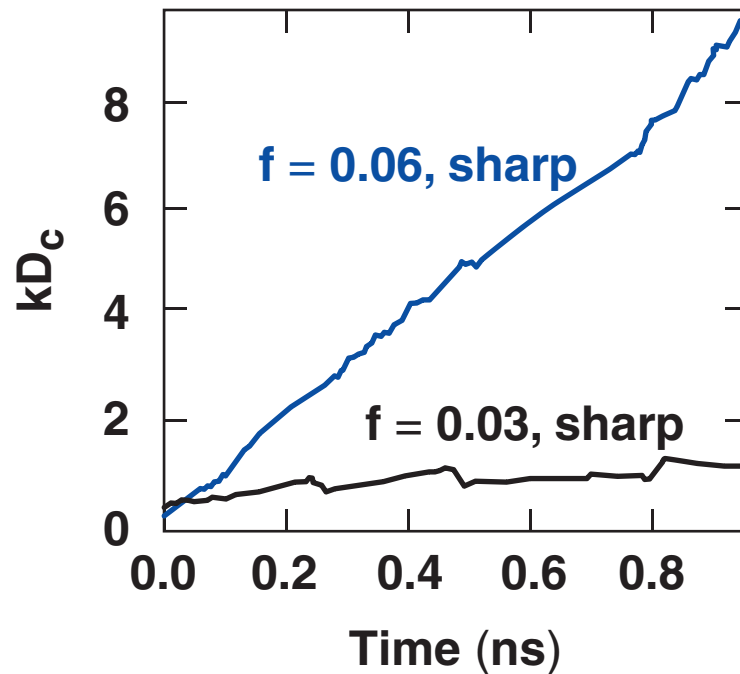
$\lambda = 200 \mu\text{m}$



- $f = 0.06$, sharp
- $f = 0.03$, sharp
- $f = 0.05$, harmonic

Shorter wavelengths are also sensitive to the transport model parameters

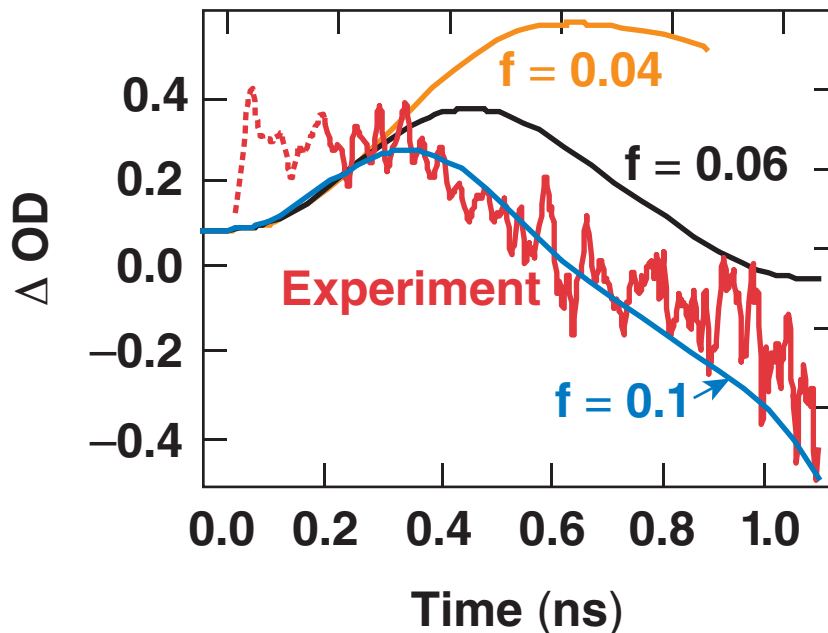
- DD foil, $I = 2 \times 10^{14} \text{ W/cm}^2$, $\lambda = 20 \mu\text{m}$



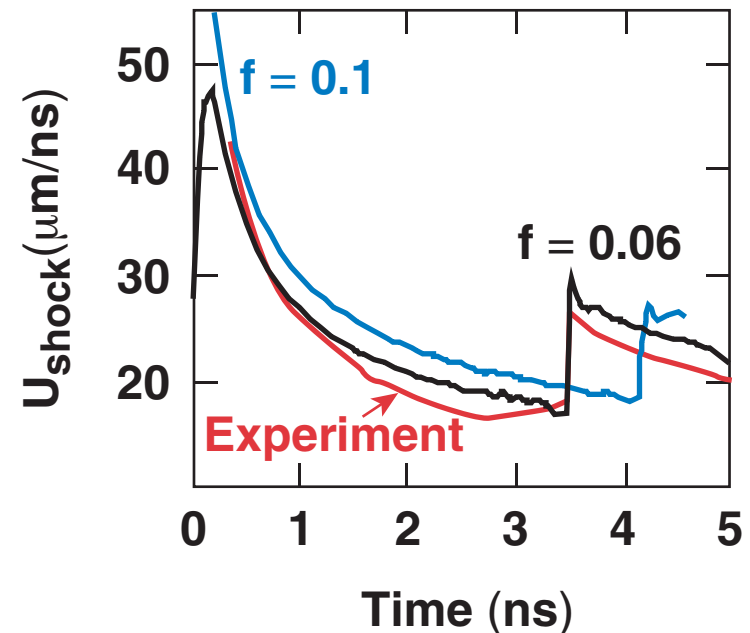
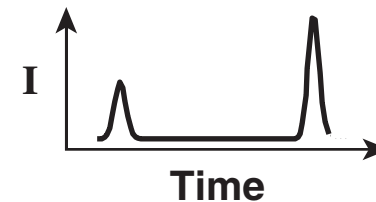
A single-valued flux limiter is not consistent with the experimental results

RM Experiment*

- 2-ns square pulse, $I = 4 \times 10^{14} \text{ W/cm}^2$
- CH foil $d = 40 \text{ }\mu\text{m}$, $\lambda = 20 \text{ }\mu\text{m}$



Dual Shock Timing**



* O. Gotchev, Ph.D. thesis, University of Rochester, 2004.

** E. Vianello *et al.*, Bull. Am. Phys. Soc 48, 205 (2003).

A nonlocal transport model has been developed to test the results of flux-limited approximation

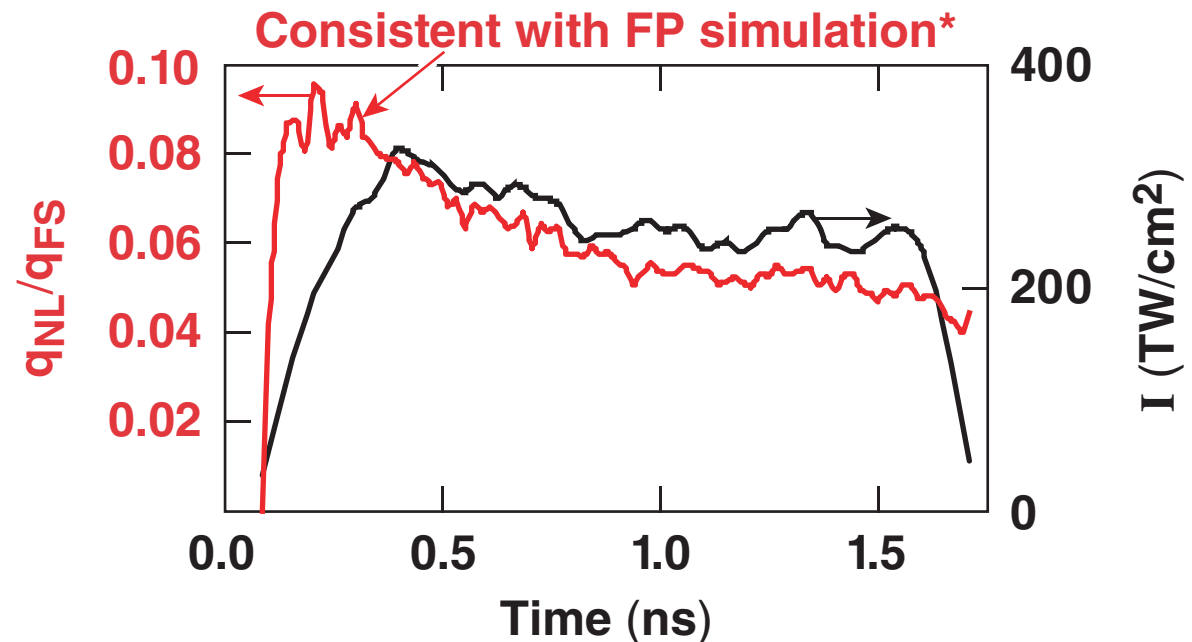
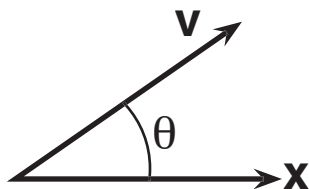
Electric field Collisional frequency

$$\bullet \quad \mathbf{v} \frac{\partial \mathbf{f}}{\partial \mathbf{x}} + \frac{\mathbf{eE}}{m} \frac{\partial \mathbf{f}_0}{\partial \mathbf{v}_x} = -\nu (\mathbf{f} - \mathbf{f}_0) \Rightarrow \mathbf{f} = \int^{\mathbf{x}} \left(\mathbf{f}_0 - \frac{\mathbf{eE}}{m\nu} \frac{\partial \mathbf{f}_0}{\partial \mathbf{v}_x} \right) e^{-\xi} \frac{d\mathbf{x}'}{\lambda \cos \theta}, \quad \xi = \int_{\mathbf{x}'}^{\mathbf{x}} \frac{d\mathbf{x}''}{\lambda \cos \theta}$$

$$\bullet \quad \mathbf{j}_x = e \int d^3v \mathbf{v} v_x \mathbf{f},$$

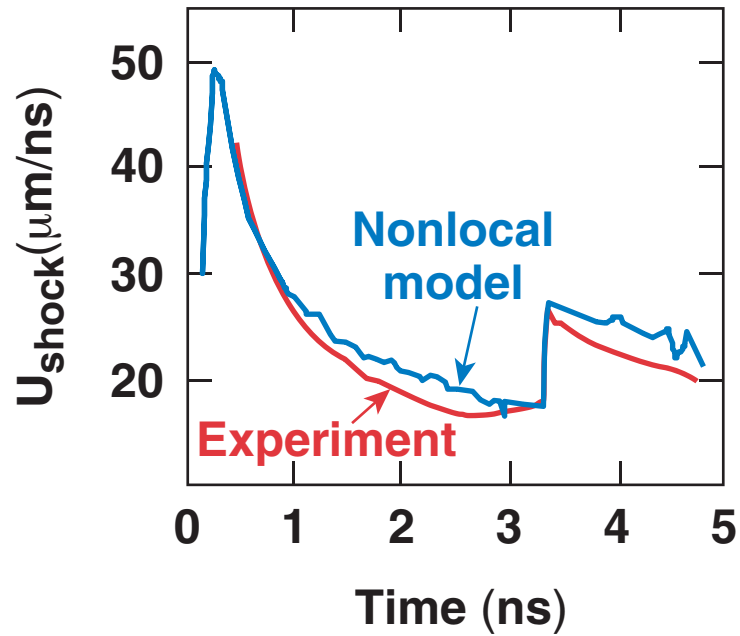
$$q_{NL} = \frac{m}{2} \int d^3v v v^2 v_x \mathbf{f}$$

$$\mathbf{j}_x = 0 \Rightarrow \mathbf{E}$$

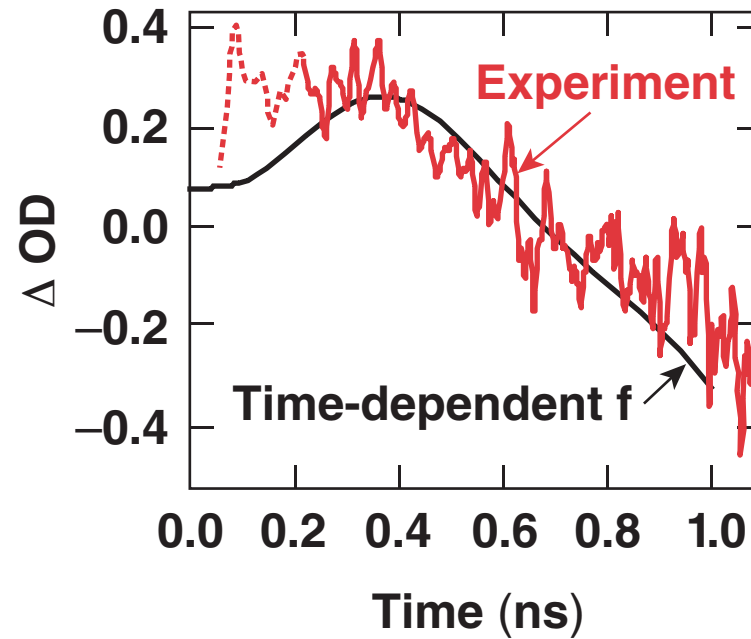


The nonlocal transport model is consistent with shock timing and RM growth measurements

Shock timing

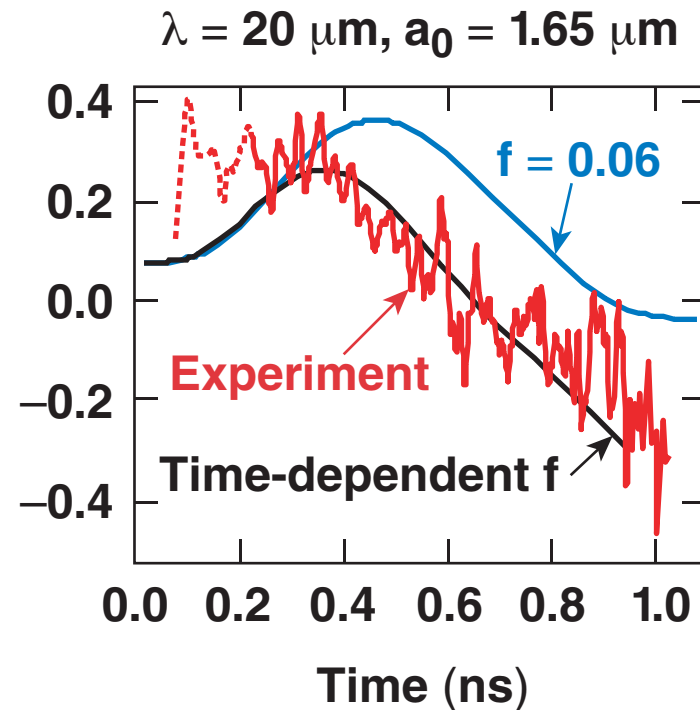
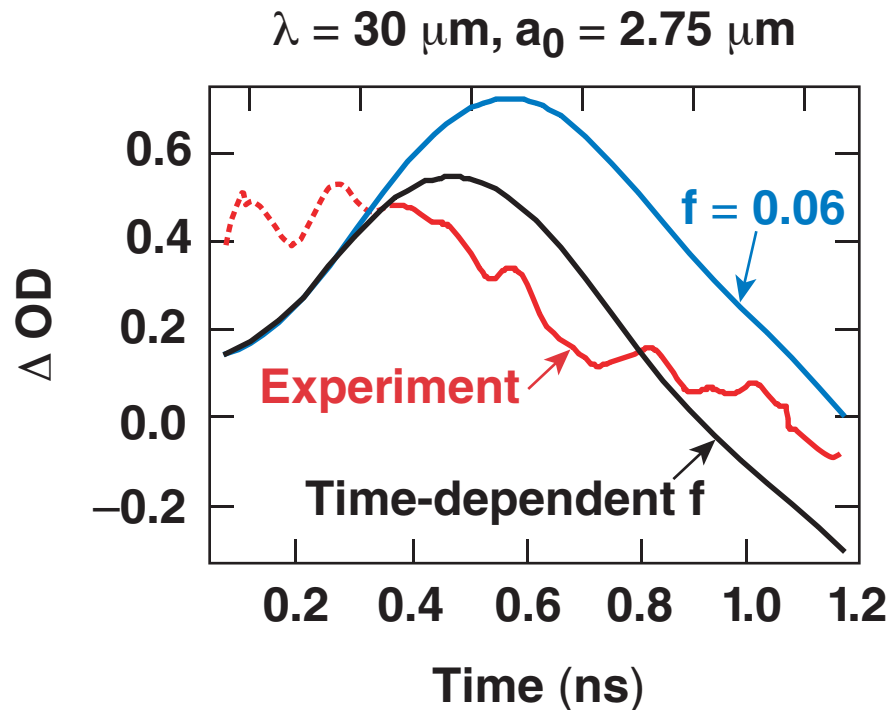


RM evolution



Results of 2-D calculations with the time-dependent flux limiter are consistent with the RM growth measurements for different wavelengths

- 2-ns square pulse with peak intensity $I = 420 \text{ TW/cm}^2$, $d_{\text{CH}} = 60 \mu\text{m}$
- Simulations use SESAME EOS and full 2-D raytracing



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