

# *What does a new, better nTOF buy us?*

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# Improving the nTOFs will drive innovation at NIF

- What does an additional nTOF get us?
  - Capturing Tion variation
    - Choose your line of sight carefully – antipodal? Fill a gap?
    - What does it take to see expected Tion anisotropy? 100 eV error bar.
  - North pole
    - Removes systematics in Tion and bulk velocity
    - Gets odd modes
    - Chance at thermal ion temperature
- What do improvements to nTOF measurements or analysis get us?
  - Moments of the peak reveal thermal and fluid state
  - 1<sup>st</sup> (shift) – bulk velocity
  - 2<sup>nd</sup> (width) – Tion and flow variance
  - 3<sup>rd</sup> (skew) – correlation of temperature and flow. Is the hot stuff moving fast?
  - 4<sup>th</sup> (kurtosis) – correlation of temperature variance and flow. Is burn happening over a broad temperature range?

# Apparent Tion varies with line of sight AND antipodal Tions are the same

- Fluid velocity variance increases the apparent temperature

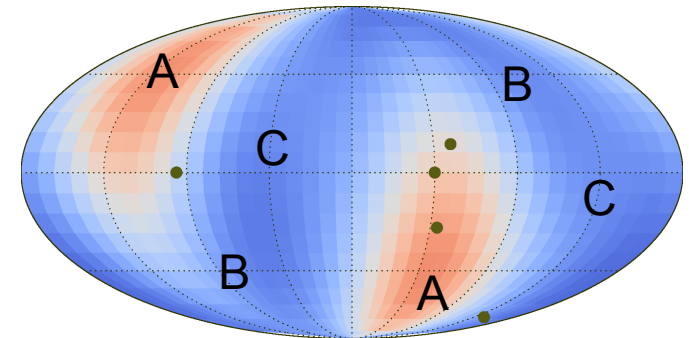
$$T_{Brysk} = \left( \frac{m_D + m_T}{k} \right) \sigma_v^2 + T_{thermal}$$

Murphy PoP

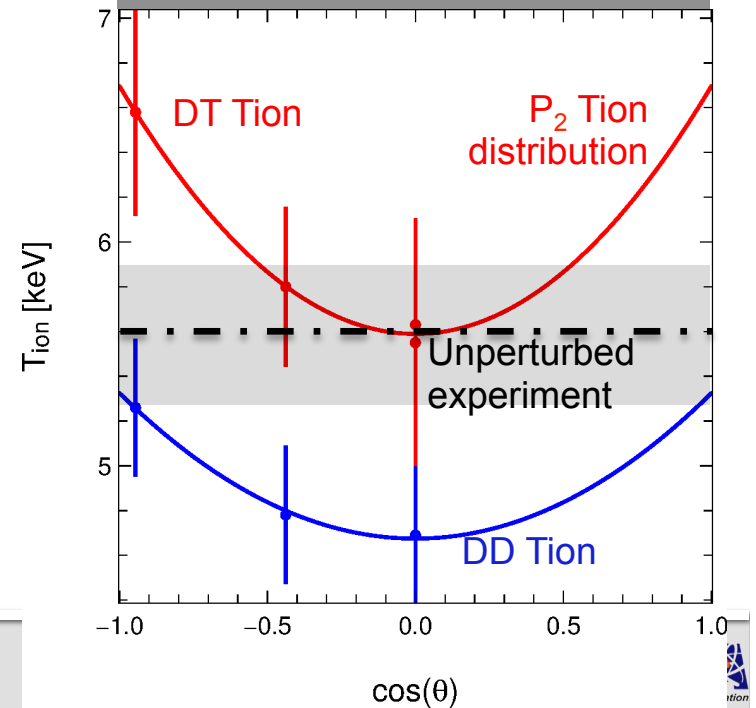
- Apparent temperature has an L=2, ellipsoidal distribution
  - Varies with line of sight
  - Equal on antipodal (opposite) lines of sight (LOS)

Some lines of sight are more valuable than others.

Simulation and theory show antipodal temps are identical



P<sub>1</sub> experiment suggests antipodal temps are the same

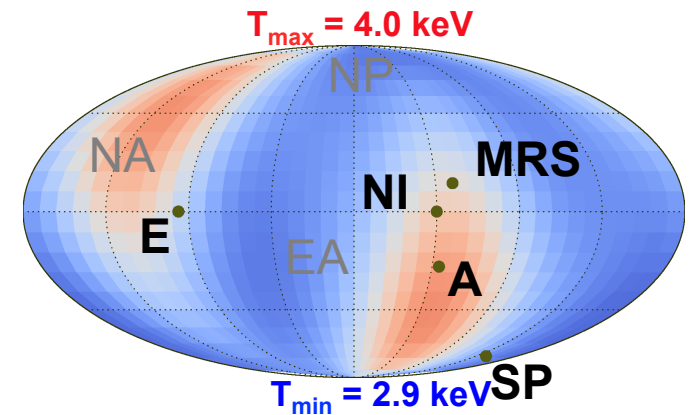


# An additional nTOF increases the ability to capture differences with line of sight

Detector	$T_{\text{Brysk}}$
SpecE	3.49
SpecA	3.56
SpecSP	2.96
NITOF	3.50
MRS	3.39

**Simulated detectors caught 55% of PTV**

**Antipodal temps are identical**



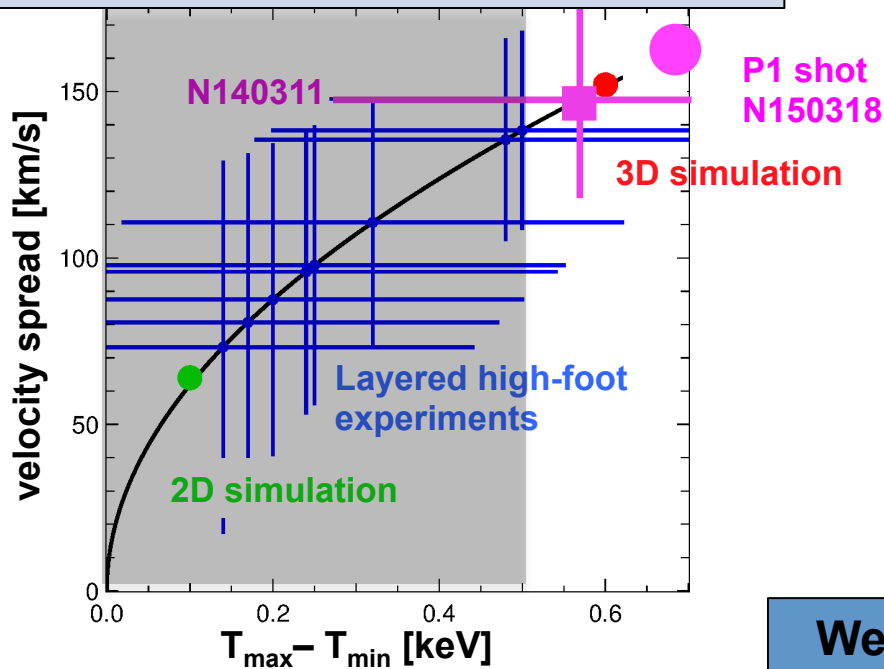
- With 4 nTOFs: capture 50 - 65% of Tion peak-to-valley (PTV)
- Add Spec NP: capture 55 - 70% (not exactly opposite SP)
- Add still another
  - antipodal specA: stays 55-70% of PTV
  - Collinear E and A: get 75 - 80 %
    - Larger percentage
    - Less variation (dependence on ellipsoid shape)

**Some lines of sight are more valuable than others ... for capturing PTV.**



# So, is the high foot apparent $T_{ion}$ usually isotropic or not?

The NIF data cannot (currently) distinguish between isotropy and the expected level of anisotropy



$$T_{Brysk} = \left( \frac{m_D + m_T}{k} \right) \sigma_v^2 + T_{thermal}$$

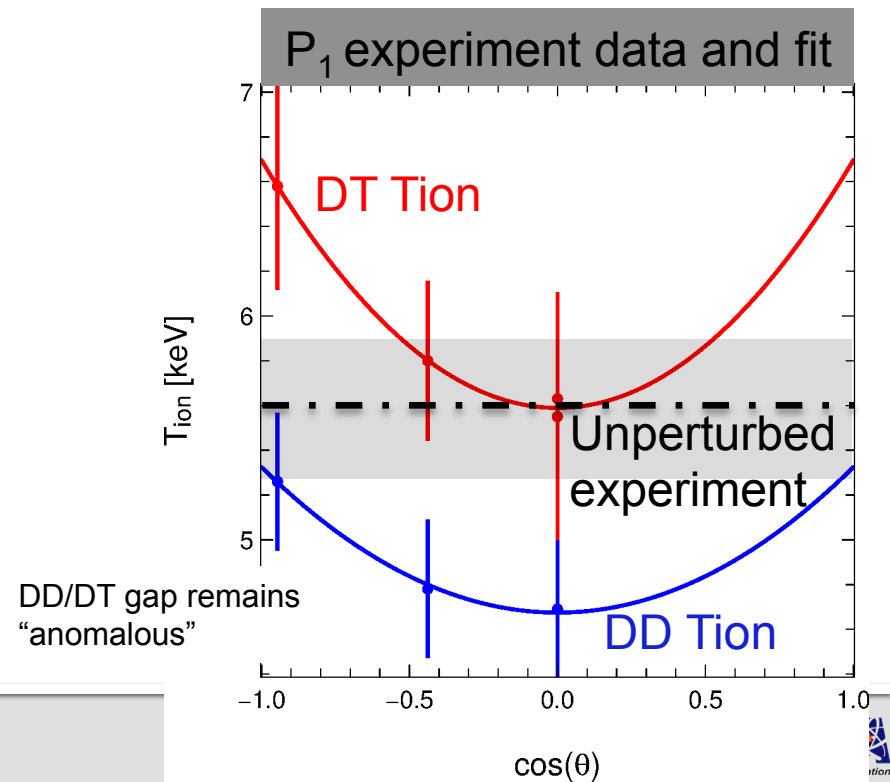
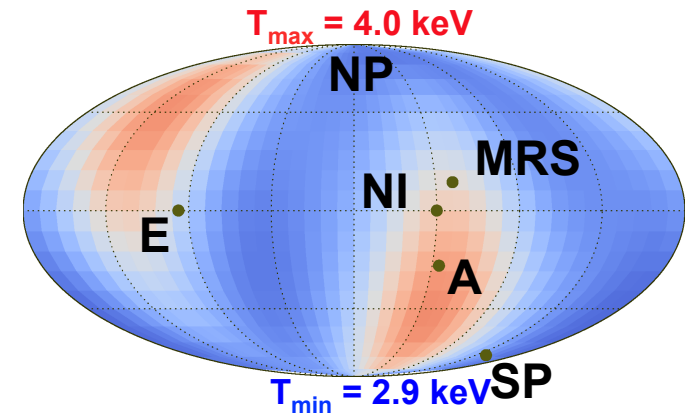
- Post shot simulations suggest Tion anisotropy of ~ 300 - 400 eV
- Detectors would typically sample ~ 150-200 eV
- Detectors can measure down to 500 eV anisotropy (PTV)

See M. Gatu Johnson paper

We need neutron spectrometers that can measure 300 eV anisotropy – that’s about a 100 eV error

# Antipodal nTOF removes systematic errors in analysis

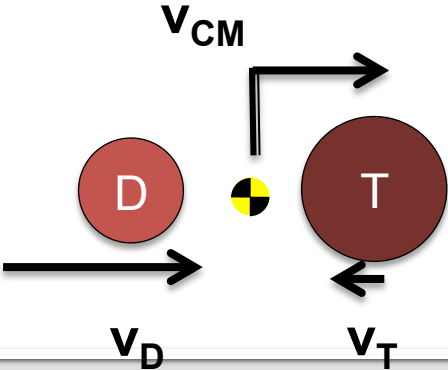
- Apparent temperature should be the same on opposing sides
  - Sources of differences
    - underappreciated physics influences (scattering)
    - instrumental or analysis systematics
- Bulk velocity
  - Equal and opposite on opposing sides
  - Similar sources of differences
- Odd mode DSR
- Can shine light on
  - Tion uncertainty
  - DD and DT Tion differences
  - DD and DT bulk velocity differences
  - Polar areal density ice caps



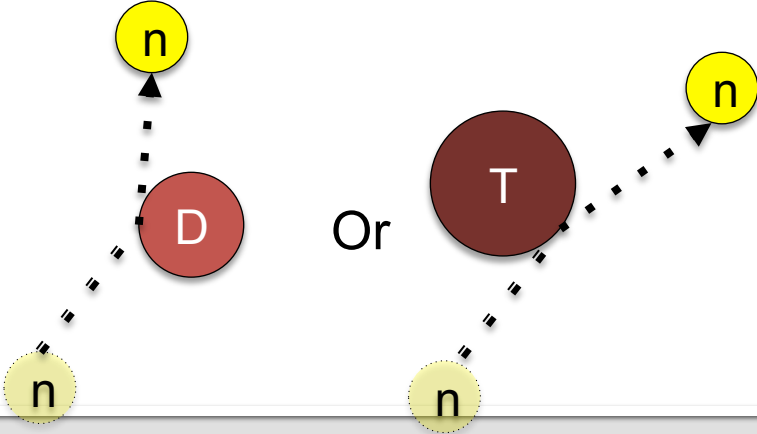
# At least three effects change the neutron spectral peak location

- Center of mass drift velocity of the burning volume
- Relative kinetic energy of the fusing ions
- Scattering of the neutrons by compressed fuel
  - scattering into the peak region by fuel
  - scattering out of the peak region by fuel

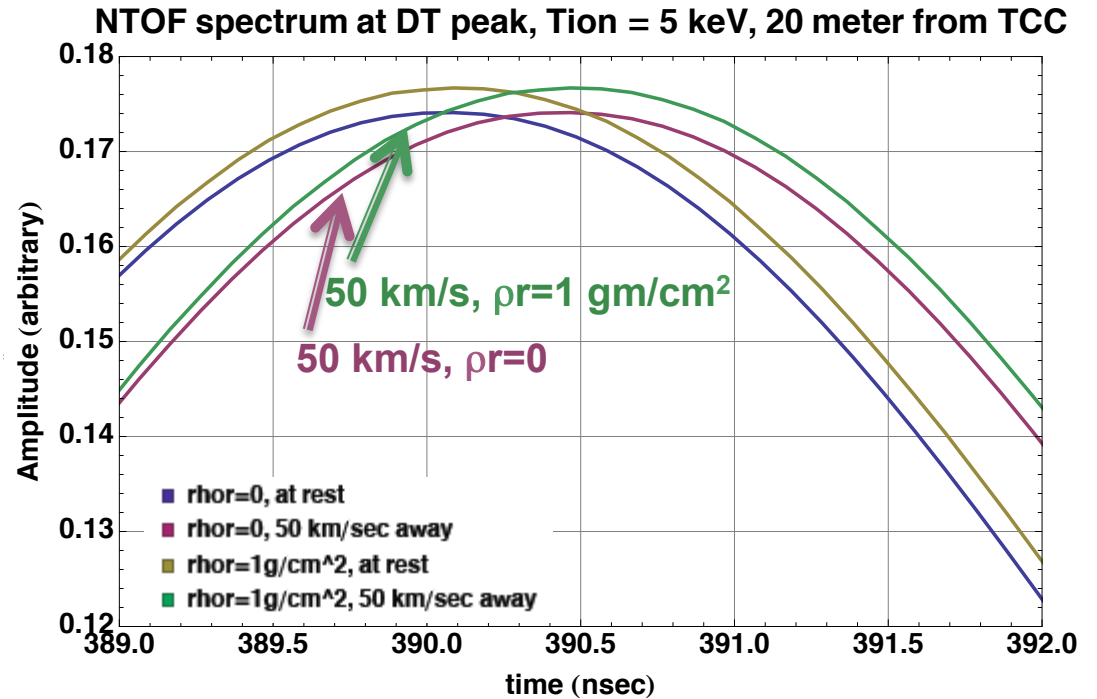
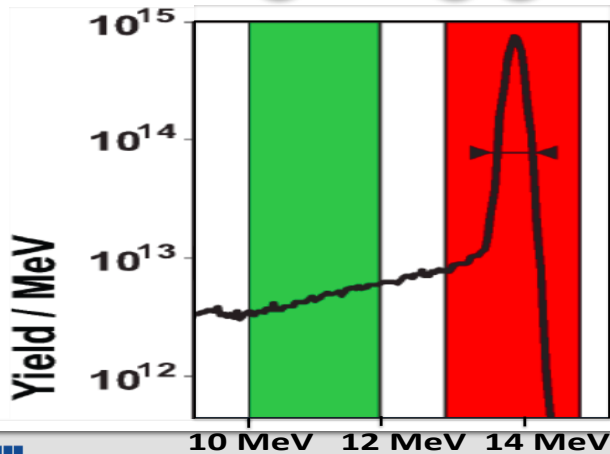
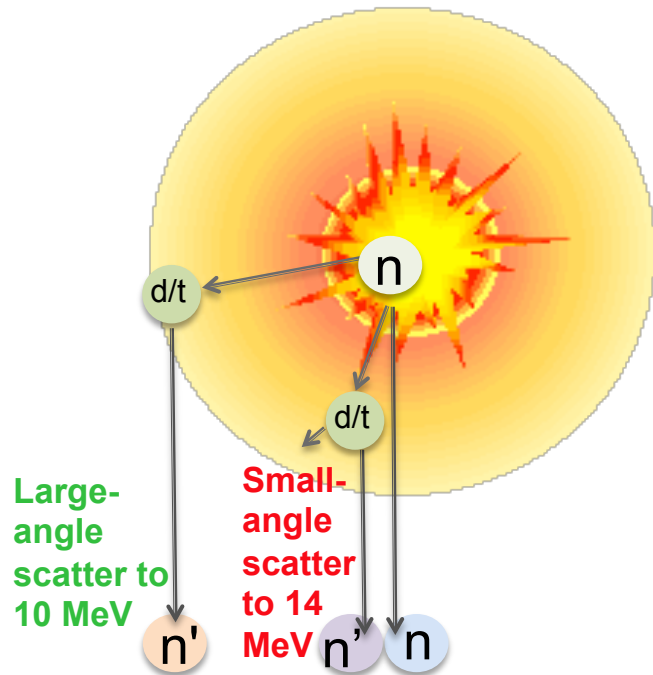
Reaction kinematics



Scattering kinematics



# Neutron scattering by dense DT shifts the central peak



Shift in central peak due to scattering is small

# Antipodal detectors provide a better measure of drift velocity

- The centroid of the DT peak shifts
  - Mainly due to neutron-weighted bulk flow (rigid translation)
  - Smaller correction due to “Ballabio” effect – neutron boost from reactant KE

$$v_{p,i} = \underbrace{v_0 + v_{Ballabio}}_{\text{Constant over LOS}} + v_{drift,i}$$

	$v_0$ [km/s]	$v_{drift}$ [km/s]	$v_{Ballabio}$ [km/s]
DT	51233.6	100	35
DD	21601.9	100	70

- Differencing opposing lines of sight leaves drift term

$$\frac{v_{p,1} - v_{p,2}}{2} = v_{drift}$$

Need 15-30 km/s precision to measure drift velocity

# Antipodal detectors provide a measure of thermal ion temperature – in theory!

- The centroid of the DT peak shifts
  - Mainly due to neutron-weighted bulk flow (rigid translation)
  - Smaller correction due to “Ballabio” effect – neutron boost from reactant KE

$$v_{p,i} = \underbrace{v_0 + v_{Ballabio}}_{\text{Constant over LOS}} + v_{drift,i}$$

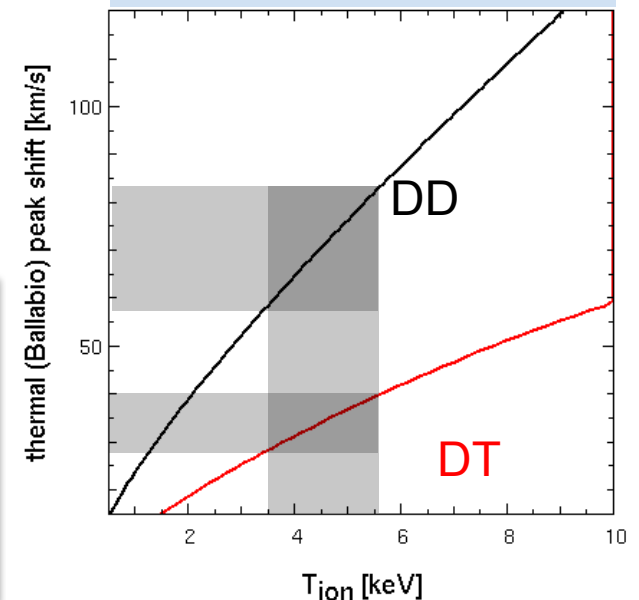
	$v_0$ [km/s]	$v_{bulk}$ [km/s]	$v_{Ballabio}$ [km/s]
DT	51233.6	100	35
DD	21601.9	100	70

- Averaging opposing lines of sight leaves Ballabio term

$$\frac{v_{p,1} + v_{p,2}}{2} - v_0 = v_{Ballabio}$$

DT needs 5 km/s precision  
DD needs 15 km/s precision

Ballabio shift depends only on **thermal**  $T_{ion}$



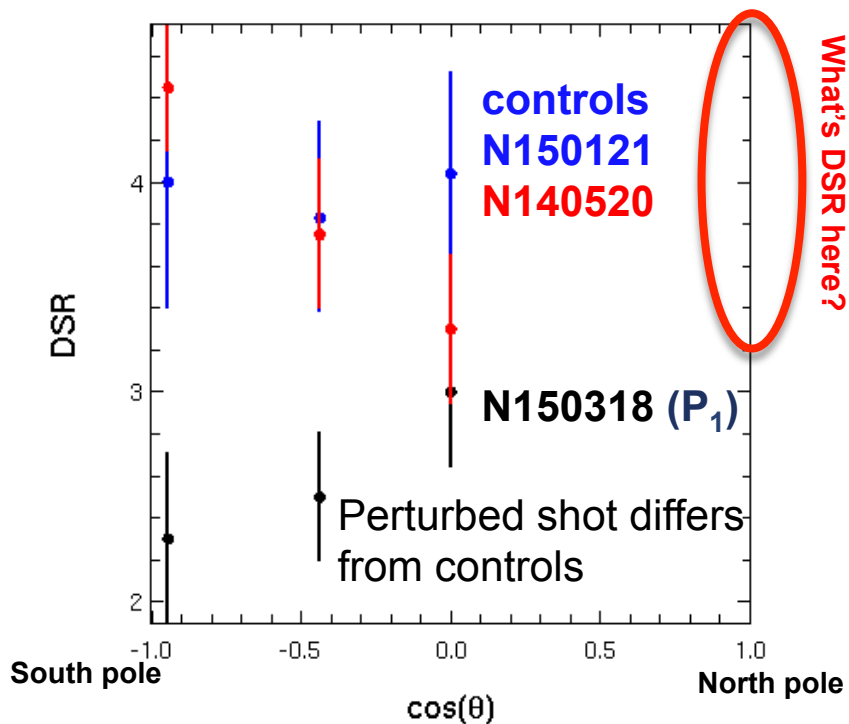
Alas, some things are harder to do in experiment than in theory



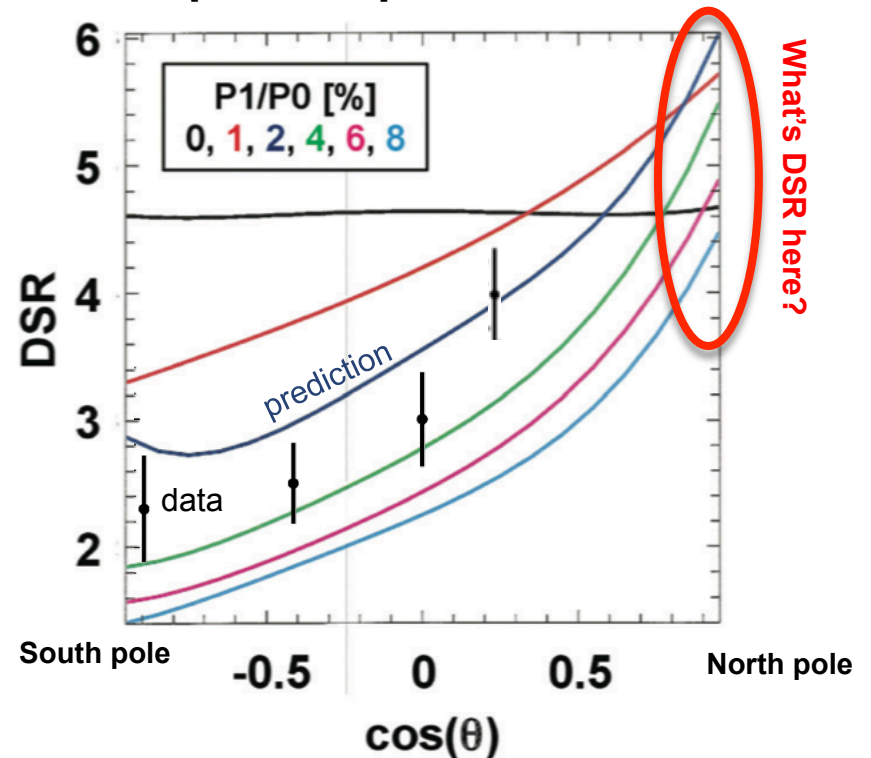
# North pole nTOF measures odd modes in cold shell

- DSR in P1 shots has strong odd mode (mode 1)
- Is the north pole as predicted? Is it different from control shots?

**Perturbed shot is different from control shots**



**Experiments compare nicely with preshot predictions**

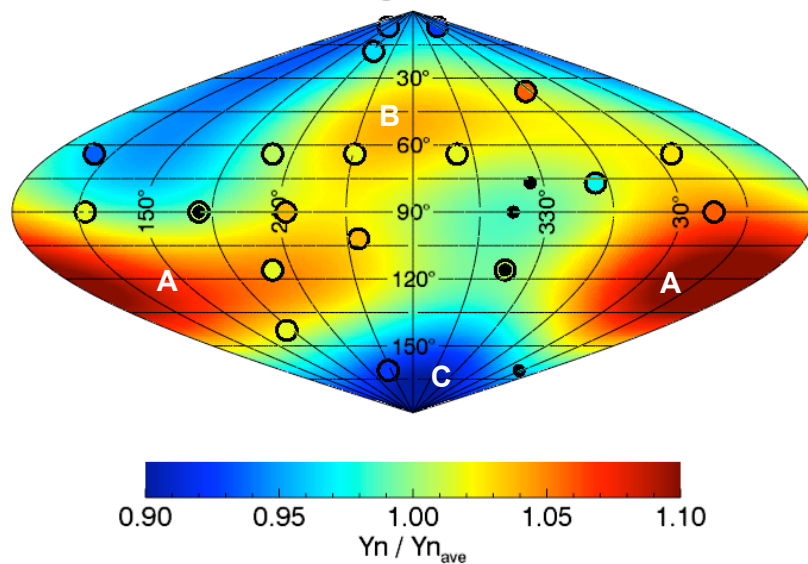


North pole nTOF provides a missing diagnosis of DSR asymmetry

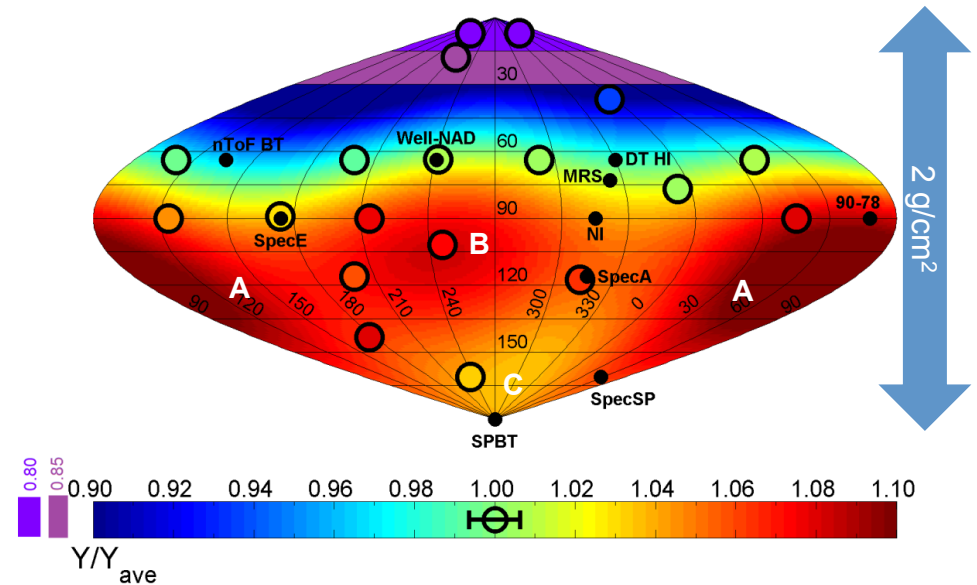
# North pole nTOF could help resolve polar ice cap issues

- DSR in P1 shots has strong odd mode (mode 1)
- Is the north pole as predicted? Is it different from control shots
- Does DSR variation compare well with fNADS, especially pole to pole?

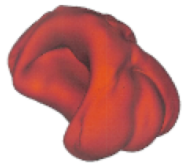
N140520 control shot



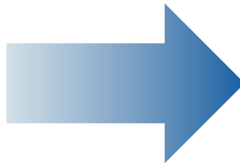
N150318 P<sub>1</sub> shot



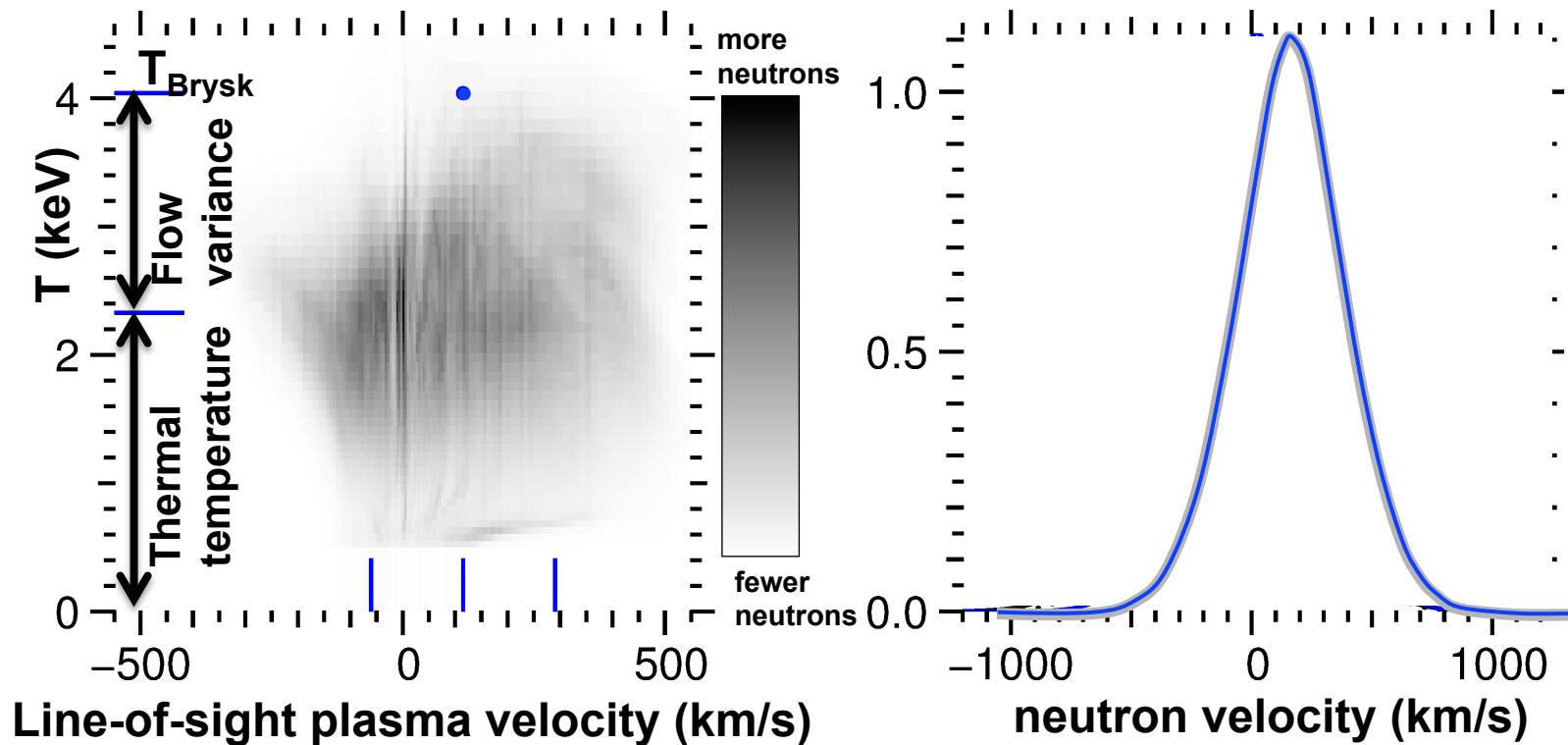
# Spectral peak depends on the distribution of neutron production in temperature and velocity space



Neutrons produced over a range of temperatures and velocities



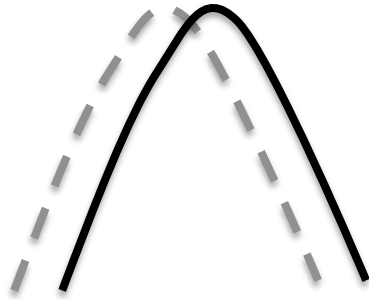
Peak width records neutron-weighted thermal temperature and the flow variance



Simulations have to get a lot right to capture the temperature variation

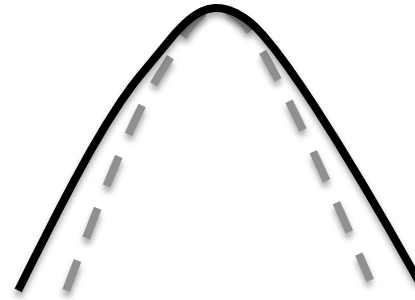
# Stagnation measurements can be much more informative

First moment:  
peak shift  $\sim f(\text{bulk velocity}, T_{\text{thermal}})$



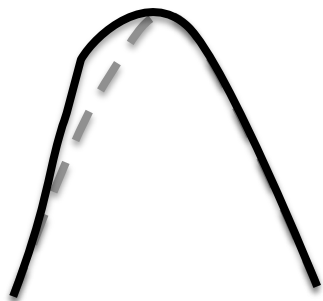
What's the  
bulk  
velocity?

Second moment:  
Width  $\sim f(T_{\text{thermal}}, \text{flow variance})$



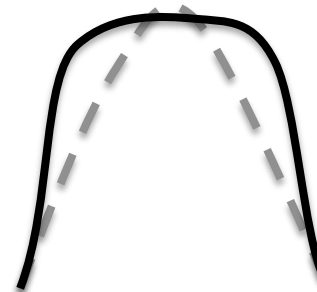
What's the  
apparent  
temp, thermal  
temp, residual  
flow?

Third moment:  
Skew  $\sim \text{cov}(T_{\text{thermal}}, \text{flow})$



Is the hot stuff  
moving fast?

Fourth moment:  
Kurtosis  $\sim \text{variance of } T_{\text{ion}}$



How broad is  
the distribution  
of thermal  
temperatures?

# High convergence NIC capsule view at 10 KeV

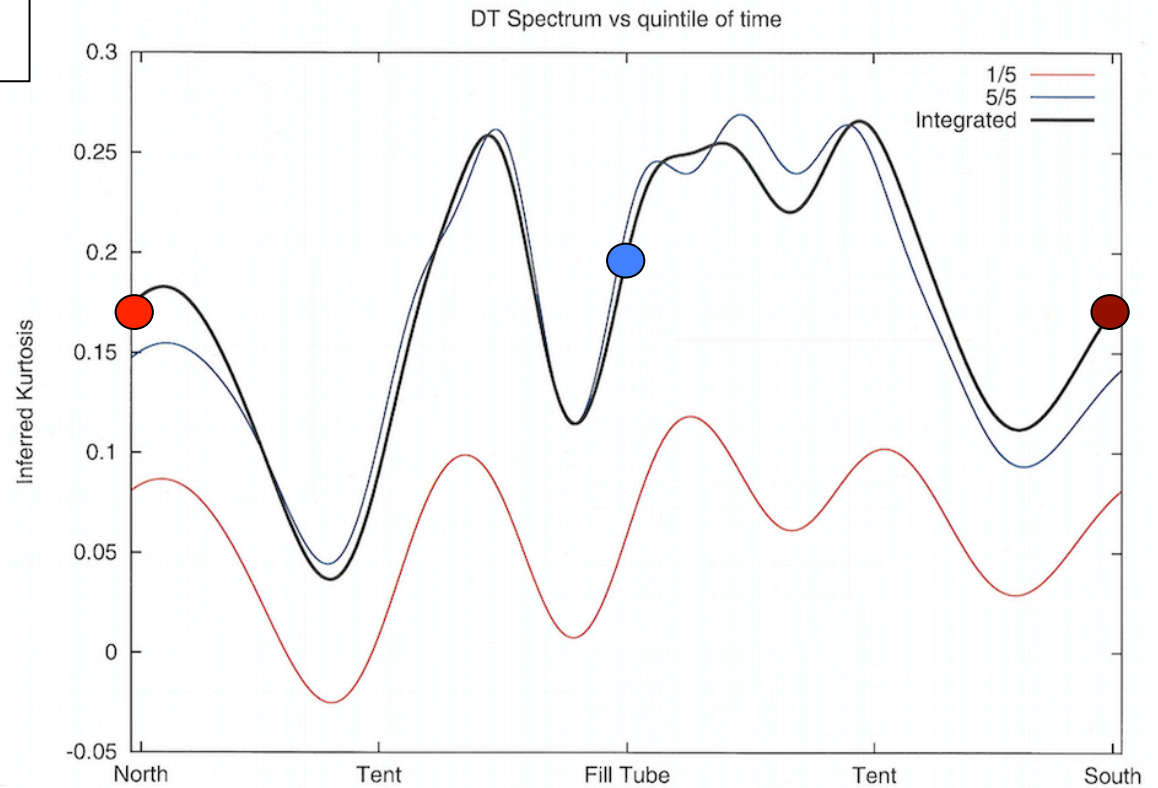
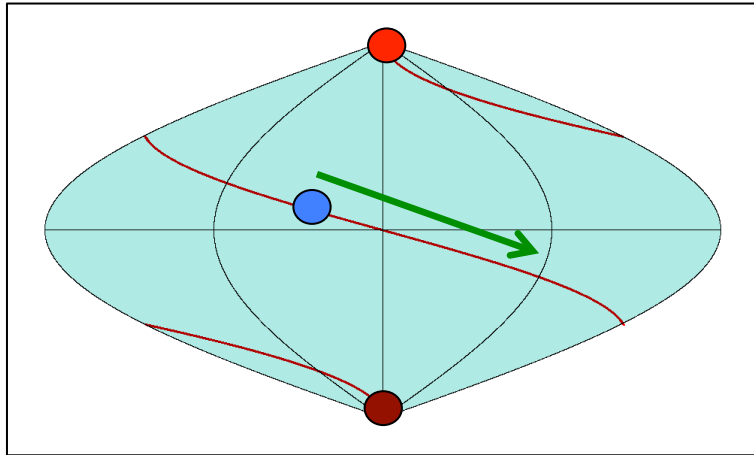


(different scales)

N120321 HYDRA hi-resolution simulation with 470M zones  
by D. Clark and C. Weber

# The kurtosis shows hot spot cooling and flow effects.

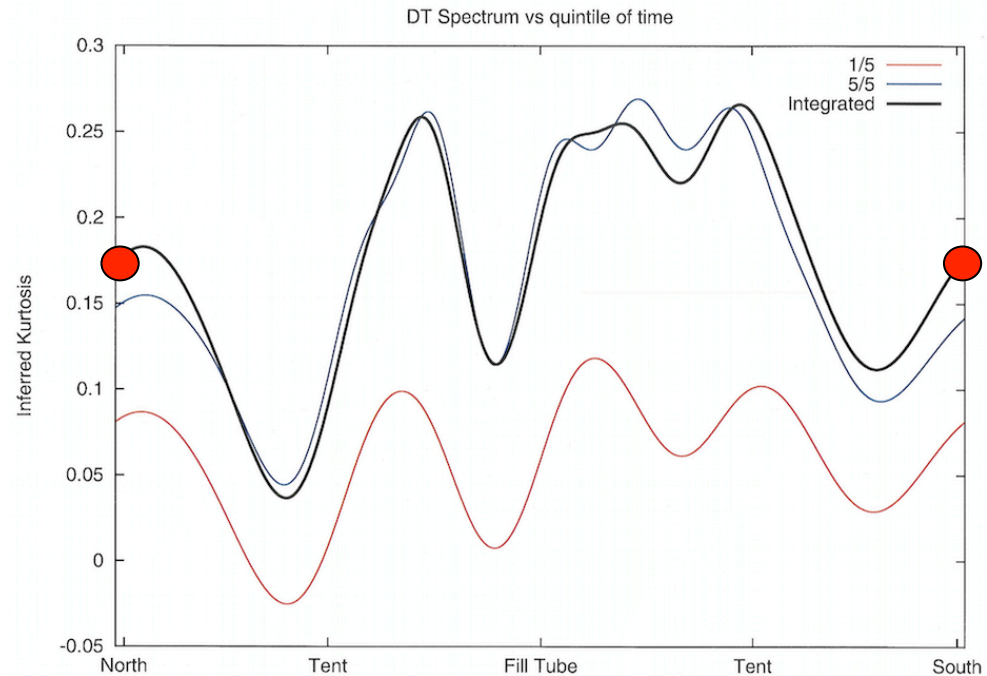
VIEW PATH





# The kurtosis shows hot spot cooling and flow effects

1. Positive kurtosis suggests temperature variation during burn
2. Negative kurtosis implies velocity variation.
3. Variation with angle is due to velocity.
4. Kurtosis would be constant with LOS in a spherical or stagnant implosion



scalar

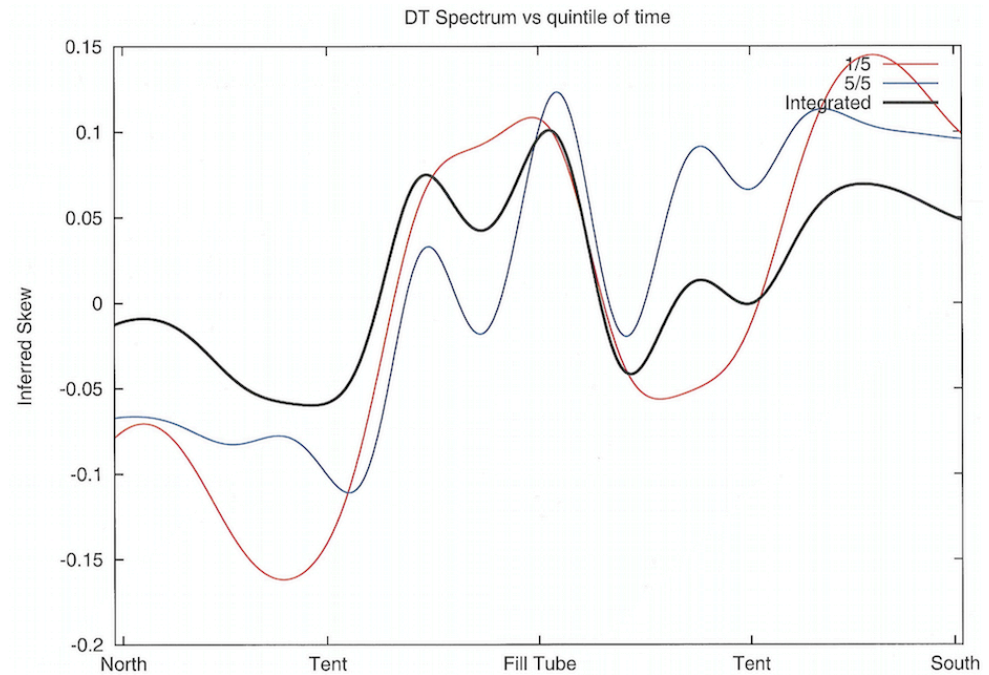
Vary with line of sight (tensors)

$$\text{Kurt}(\omega) = \frac{3 \text{Var}(\tau) + 6 \text{Cov}(\tau, u_{\Omega}, u_{\Omega}) + \text{Cov}(u_{\Omega}, u_{\Omega}, u_{\Omega}, u_{\Omega}) - 3 \text{Var}(u_{\Omega})^2 + \dots}{\text{Var}(\omega)^2}$$

**L=0, 2, 4 in direction → antipodes are identical**

Kurtosis variation with line of sight is another direct measure of stagnation and stagnation asymmetry – need it to ~ 5% precision

1. Skew gives correlation of temperature and velocity
2. Is the hottest material moving fast? Slow?



Vary with line of sight (tensor)

$$\text{Skew}(\omega) = \frac{3\text{Cov}(\tau, u_{\Omega}) + \text{Cov}(u_{\Omega}, u_{\Omega}, u_{\Omega}) + \dots}{\text{Var}(\omega)^{3/2}}$$

**L=1, L=3 in direction → antipodes measure odd modes**

Skewness gives us a picture of the partition of mechanical and thermal energy – need it to ~ 3-5% precision

# Capturing the spatial variation of spectral signatures requires tradeoffs in instrument number and precision

- Consider a test case
  - 400eV PTV Tion variation
  - 15% skew variation
  - 20% kurtosis variation
- Current suite + SpecNP (50% sampling efficiency)
  - 1<sup>st</sup> moment – peak location to 15-30 km/s, needed on at least 3 LOS
  - 2<sup>nd</sup> moment – sample 200 eV PTV → 100 eV precision
  - 3<sup>rd</sup> moment – sample 7% skew PTV → 3% precision
  - 4<sup>th</sup> moment – sample 10% kurtosis PTV → 5% precision
- Current suite + SpecNP + collinear EA (75% sampling efficiency)
  - 1<sup>st</sup> moment – peak location to 15-30 km/s, needed on at least 3 LOS
  - 2<sup>nd</sup> moment – sample 300 eV PTV → 150 eV precision
  - 3<sup>rd</sup> moment – sample 12% skew PTV → 6% precision
  - 4<sup>th</sup> moment – sample 15% kurtosis PTV → 8% precision

# Improving the nTOFs will drive innovation at NIF

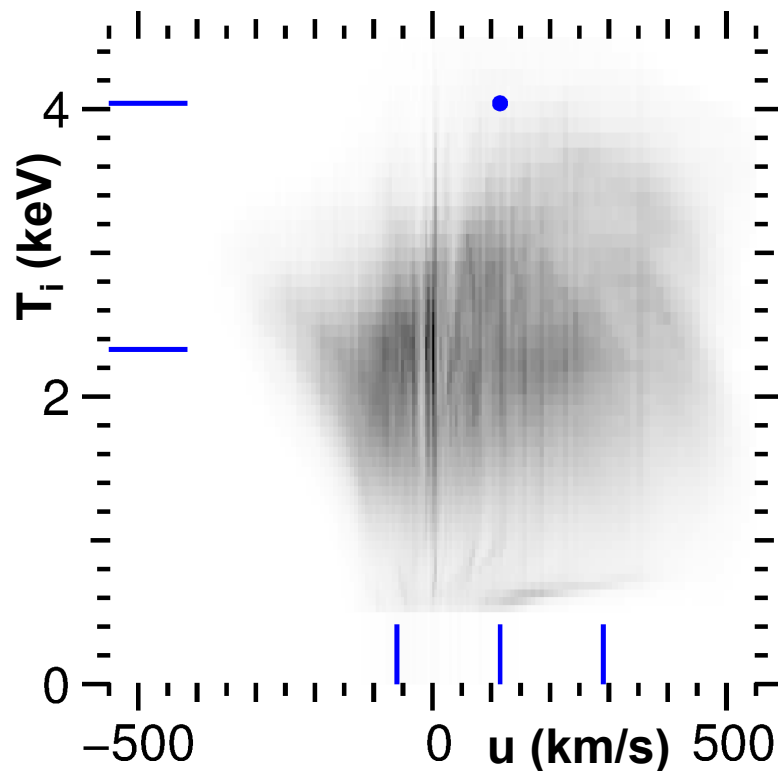
- What does an additional nTOF get us?
  - Captures peak-to-valley Tion variation
- OR
- Removes systematics in Tion and bulk velocity
- Gets odd modes
- Gives a chance at thermal ion temperature
- What do improvements to nTOF measurements or analysis get us?
  - Moments of the peak reveal thermal and fluid state
  - 1<sup>st</sup> (shift) – bulk velocity
  - 2<sup>nd</sup> (width) – Tion and flow variance
  - 3<sup>rd</sup> (skew) – Is the hot stuff moving fast?
  - 4<sup>th</sup> (kurtosis) – Is hot spot burning over a broad range of temperatures?
  - Do the nuclear signatures provide a signature of asymmetry?



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# Neutron spectral moments and LOS dependence are important clues

burn T-u distribution (3D simulation)



burning plasma exceedingly non-uniform, neutrons produced in wide range of  $T_i$  and fluid  $u$

shift of spectral peak only tells us mean  $\langle u \rangle + \text{shift}(\langle T_i \rangle)$

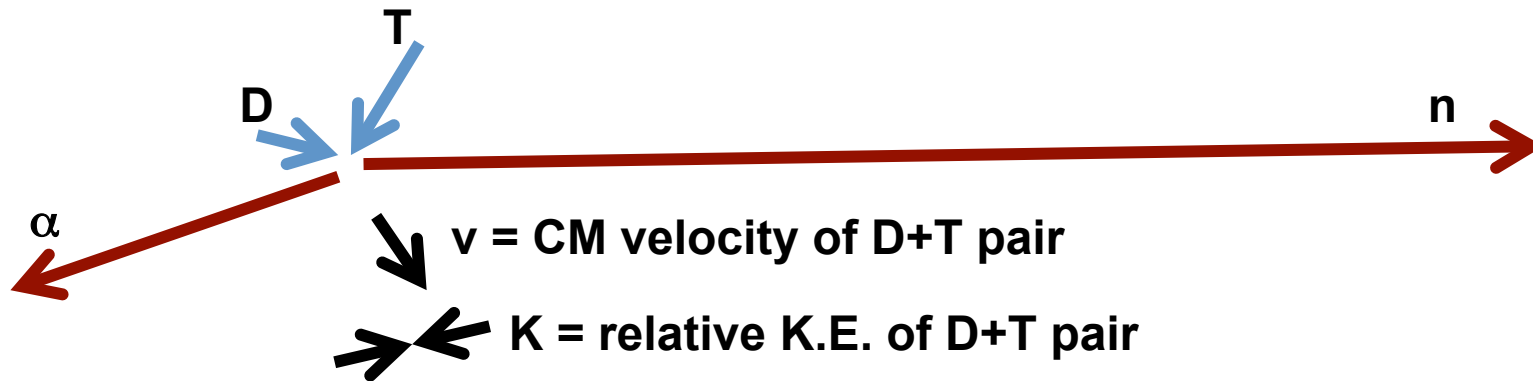
variance of spectral peak only captures  $\langle T_i \rangle + \text{Var}(u)$

skew and kurtosis of spectral peak tell us about T-u correlations and  $\text{Var}(T)$

$u$  = fluid velocity component along LOS



## Each D+T (or D+D) reaction makes n with slightly different momentum



$$m^2 = E^2 - p^2 = E'^2 - p'^2$$

Lorentz invariants for neutron boost

$$E' = \gamma(E + \mathbf{v} \cdot \mathbf{p}) = \gamma(E + v_{\Omega} p)$$

Boost CM 4-momentum by CM  $\mathbf{v}$

$$p' = p + E_0 v_{\Omega} - \frac{v_{\perp}^2}{2v_0^2} v_0 + \frac{v^2 + v_{\perp}^2}{2} v_0 + O(v^3) \quad \begin{array}{l} E_0 = m_n + K_0, K_0 \sim 14 \text{ MeV} \\ v_0 \sim 51000 \text{ km/s (14 MeV)} \end{array}$$

$$\omega \equiv \frac{p'}{E_0} - v_0 = v_{\Omega} + \frac{p}{E_0} - v_0 - \frac{v_{\perp}^2}{2v_0^2} v_0 + \frac{v^2 + v_{\perp}^2}{2} v_0 + O(v^3)$$

# Shifted, scaled neutron momentum is best variable for spectrum

$$\omega \equiv \frac{p'}{E_0} - v_0 = v_\Omega + \frac{p}{E_0} - v_0 - \frac{v_\perp^2}{2v_0^2} v_0 + \frac{v^2 + v_\perp^2}{2} v_0 + O(v^3)$$

CM velocity component  
thermal motion T,  
fluid motion u

$p=p(K)$  relative K.E.  
thermal motion  $K \sim 5T$

$$M = m_D + m_T \quad \langle v_\Omega^2 \rangle_{\text{thermal}} = \frac{T}{M} \equiv \tau \quad T \text{ in units of velocity}^2$$

1 keV  $\rightarrow$  (139 km/s)<sup>2</sup> DT  
(155 km/s)<sup>2</sup> DD

$$\frac{p}{E_0} - v_0 \approx \frac{1}{v_0} \left( \frac{M}{E_0} - 1 \right) \frac{K}{M} \equiv K \quad K \text{ in units of velocity}$$

10 keV  $\rightarrow$  14.7 km/s DT, 33.1 km/s DD

## For given T, u, and K, can integrate over directions, Maxwellian exactly

fixed K = relative K.E. defers needing to know reaction cross section

$$\frac{dN}{d\Omega dp_n''} \sim \frac{p_n''^2}{E_n'' p_n' p_n} \exp\left(-(\gamma - 1) \frac{M + K}{T}\right)$$

unprimed is CM  
' is fluid frame  
" is lab frame

This spectrum exact Maxwell-Juttner averaged relativistic kinetics  
Can also integrate momentum moments analytically

Averages over the distribution of K for given T done by expanding  
in  $K/K_0$  and  $K/M$  – this averaging requires reaction cross section

Finally, average over T, u distribution

## Use neutron momentum spectrum, scaled to units of velocity

$$\omega = p_n / (m_n + K_0) - v_0 \quad \text{scaled and shifted neutron momentum}$$

- very nearly CM velocity of reacting DT pair

$$4\pi \frac{dN}{d\omega d\Omega} \quad \text{momentum spectrum = number of neutrons per sphere}$$

within  $d\omega$  of “velocity”  $\omega$  and within  $d\Omega$  of direction  $\Omega$

$$\langle \omega^n \rangle = \frac{\int d\omega \omega^n \frac{dN}{d\omega d\Omega}}{\int d\omega \frac{dN}{d\omega d\Omega}} \quad \text{n}^{\text{th}} \text{ moment of scaled momentum spectrum}$$

$$\tau = T / (m_D + m_T) \quad \text{fluid temperature T as a velocity variance}$$

$$u_\Omega = \mathbf{u} \cdot \boldsymbol{\Omega} \quad \text{fluid velocity component along LOS}$$

$$\bar{K} = \frac{1}{v_0} \left( \frac{m_D + m_T}{m_n + K_0} - 1 \right) \frac{\bar{K}(T)}{m_D + m_T} \approx \bar{\omega}(T) \quad \text{“velocity” for mean DT K.E.(T)}$$

(“Ballabio shift”)

## Each spectral moment constrains moments of (T,u) burn distribution

$f(T, \mathbf{u}) dT d^3 \mathbf{u}$       fraction of neutrons produced in plasma at temperature T within dT, velocity u within du

$\langle XY \rangle = \int XY f(T, \mathbf{u}) dT d^3 \mathbf{u}$       burn average of quantity XY

$\int d\omega 4\pi \frac{dN}{d\omega d\Omega} = 1 + \frac{2}{v_0} \langle u_\Omega \rangle + \frac{1+v_0^2}{2v_0^2} (3\langle u_\Omega^2 \rangle - \langle u^2 \rangle) + \dots$       LOS dependence of yield

$\langle \omega^1 \rangle = \langle u_\Omega \rangle + \langle \kappa \rangle + (1 + \frac{1}{2} v_0^2) \langle \tau \rangle / v_0 + \dots$       centroid of spectrum

$\langle \omega^2 \rangle = \langle \tau \rangle + \langle u_\Omega^2 \rangle + 2\langle \kappa u_\Omega \rangle + \dots$

$\langle \omega^3 \rangle = 3\langle \tau u_\Omega \rangle + \langle u_\Omega^3 \rangle + \dots$       (showing only largest contributions)

$\langle \omega^4 \rangle = 3\langle \tau^2 \rangle + 6\langle \tau u_\Omega^2 \rangle + \langle u_\Omega^4 \rangle + \dots$

## Compute cumulants to see deviation from Gaussian spectrum

$$\text{Cov}(X, Y) = \langle (X - \langle X \rangle)(Y - \langle Y \rangle) \rangle = \langle XY \rangle - \langle X \rangle \langle Y \rangle$$

$$\text{Var}(X) = \text{Cov}(X, X) = \langle X^2 \rangle - \langle X \rangle^2$$

$$\text{Cov}(X, Y, Z, \dots) = \langle (X - \langle X \rangle)(Y - \langle Y \rangle)(Z - \langle Z \rangle) \dots \rangle$$

$$\text{Skew}(X) = \text{Cov}(X, X, X) / \text{Var}(X)^{3/2}$$

$$\text{Kurt}(X) = \text{Cov}(X, X, X, X) / \text{Var}(X)^2 - 3$$

**skew, kurtosis  
zero for Gaussian  
distribution**

$$\text{Var}(\omega) = \langle \tau \rangle + \text{Var}(u_\Omega) + 2 \text{Cov}(\kappa, u_\Omega) + \dots$$

**L=0, L=2, L=1 in direction**

$$\text{Skew}(\omega) = \frac{3 \text{Cov}(\tau, u_\Omega) + \text{Cov}(u_\Omega, u_\Omega, u_\Omega) + \dots}{\text{Var}(\omega)^{3/2}}$$

**L=1, L=3 in direction**

$$\text{Kurt}(\omega) = \frac{3 \text{Var}(\tau) + 6 \text{Cov}(\tau, u_\Omega, u_\Omega) + \text{Cov}(u_\Omega, u_\Omega, u_\Omega, u_\Omega) - 3 \text{Var}(u_\Omega)^2 + \dots}{\text{Var}(\omega)^2}$$

**L=0, 2, 4**





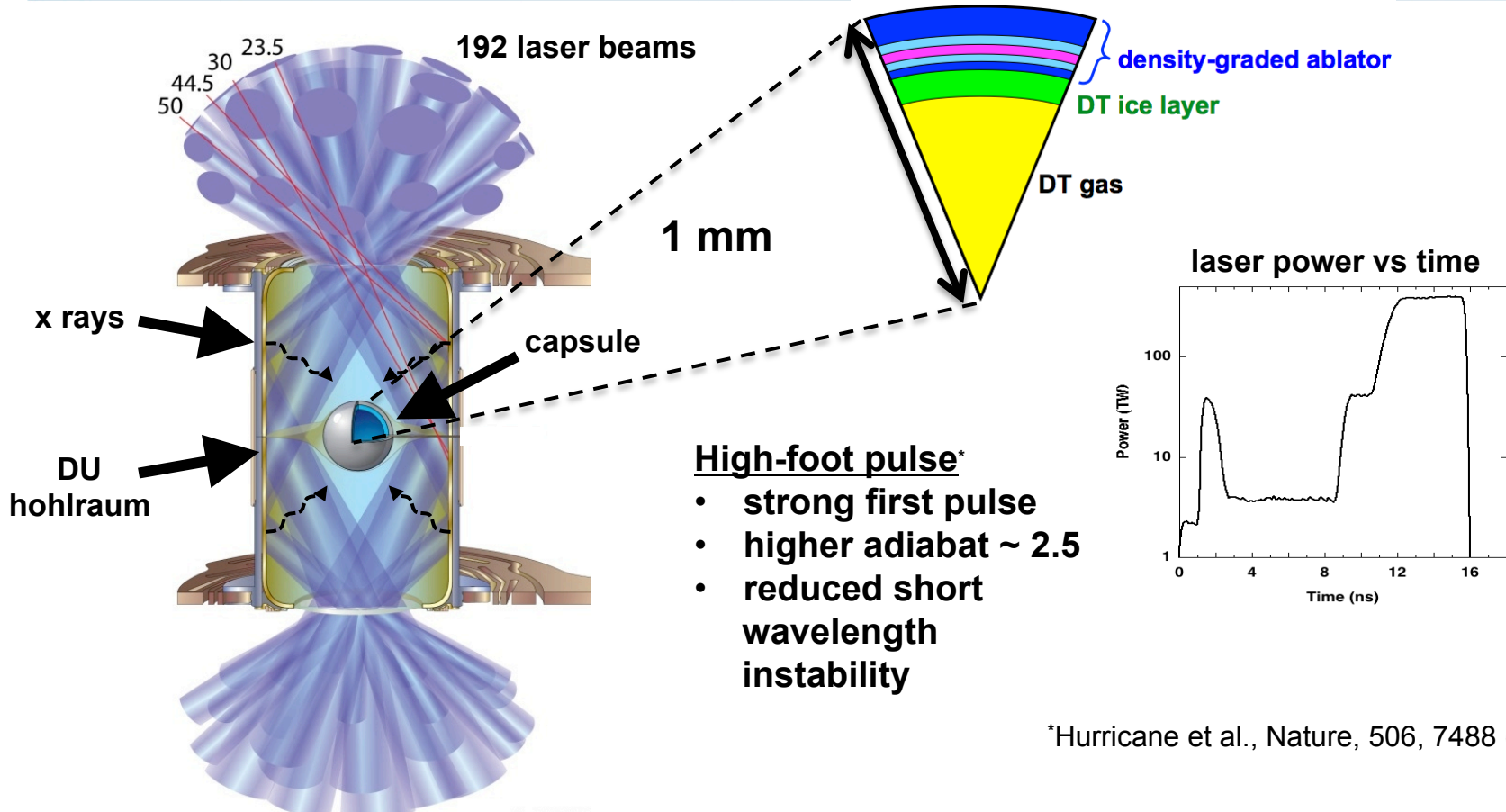
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# Nuclear diagnosis at NIF provides an unprecedented picture of stagnated ICF implosions

- Hohlraum and capsule symmetry respond to large drive perturbations ( $P_1$ ) as predicted
- Nuclear diagnostics capture the thermodynamics and flow of the hot spot and cold shell
- Simulated hot spot and cold shell diagnostics match experimental observables
- The repeatability of the high foot implosion platform supports perturbed stagnation experiments

Our codes and diagnostics have captured the detailed effects of intentional perturbations

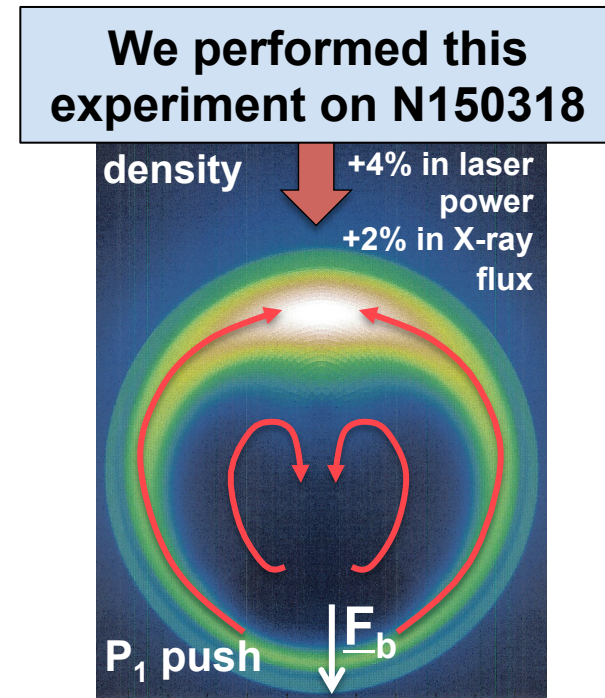
# We used high-adiabat implosions with reduced high-mode instability



**High-adiabat implosions allow investigation of asymmetry and stagnation processes**

# Top-to-bottom drive imbalance (mode 1) is an ideal symmetry perturbation

- Implosions are sensitive to mode 1
  - Buoyancy force on hot spot due to  $P_1$  acceleration
  - Hot spot flows
  - Shell asymmetry
  - Similar flows result from ice layer asymmetry



Spears, PoP 2014  
Chittenden et al

- Mode 1 effects are observable by nuclear diagnosis
- Signatures of mode 1 are present in many high foot implosions

# Asymmetrically driven implosions are relevant to the stockpile stewardship mission on NIF

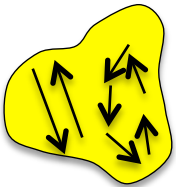
- Provide an experimental platform with asymmetric radiation flow
- Detailed measurements of the stagnating plasma
- Detailed code predictions of observable signatures (neutron spectra)

**Perturbed implosions provide an integrated test of our code capabilities**

# We measure multiple stagnation quantities by neutron spectrometry

## Ion temperature

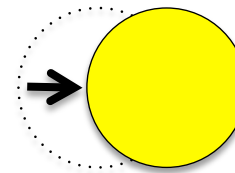
Neutron spectral peak width.  
Temperature and hot spot flow.



Shear  
Swirling  
Velocity field variance

## Bulk velocity

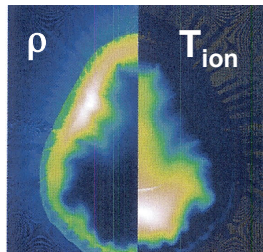
Neutron spectral peak shift.  
One-sided imbalance drives this.



Rigid-body  
translation

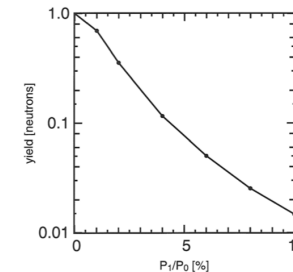
## Shell uniformity

Neutron scattering.  
Asymmetries perturb the shell.



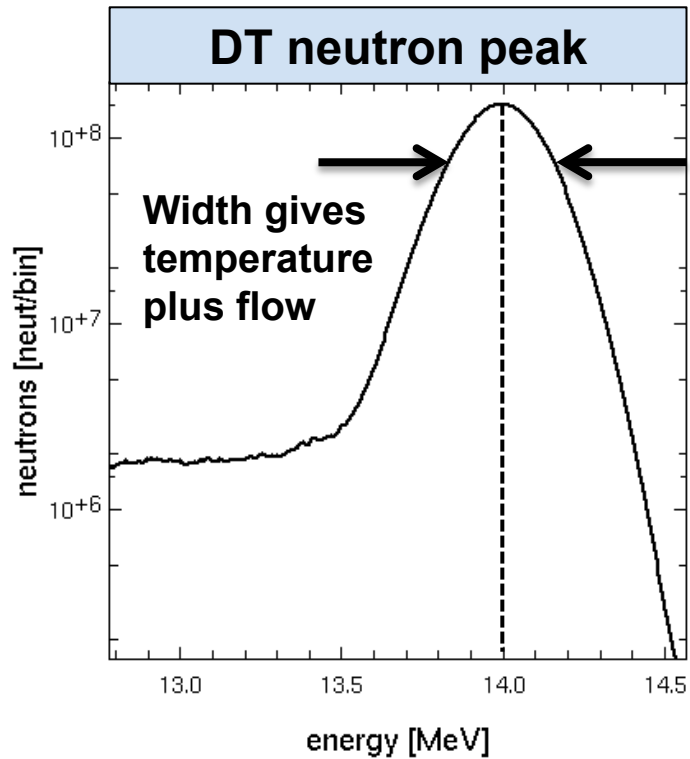
## Neutron yield

Integrated performance metric.  
Incomplete stagnation reduces yield.



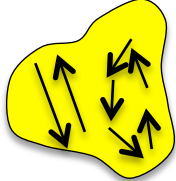
Implosion asymmetry alters stagnation phase properties

# Neutron spectrometers measure *apparent* ion temperature from spectral peak width



- Peak is broadened by:
1. thermal temperature
  2. fluid flow

Flowing hot spot:  
 Shear  
 Swirling  
 Velocity field variance



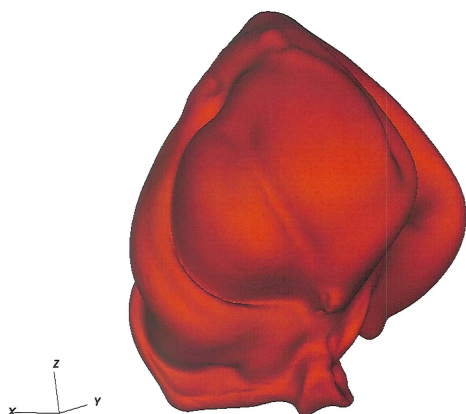
spread in fluid velocity

$$T_{Brysk} = \left( \frac{m_D + m_T}{k} \right) \sigma_v^2 + T_{thermal}$$

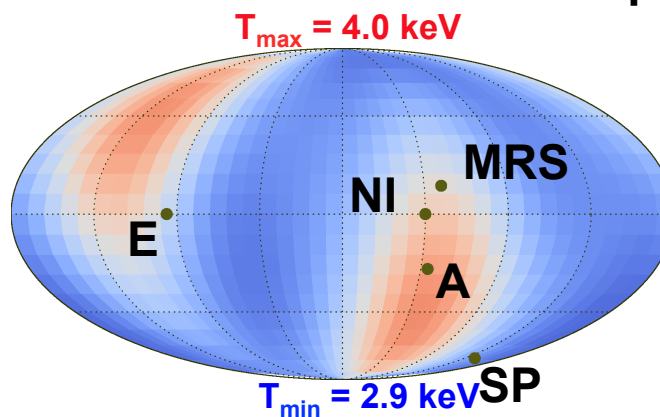
**Hot spot flows increase the apparent (Brysk) temperature**

# Asymmetric 3D simulations show angular temperature variations due to flow

Asymmetric flow in distorted hot spot



Apparent temperature distribution from simulated peak widths



Detector	$T_{\text{Brysk}}$
SpecE	3.49
SpecA	3.56
SpecSP	2.96
NITOF	3.50
MRS	3.39

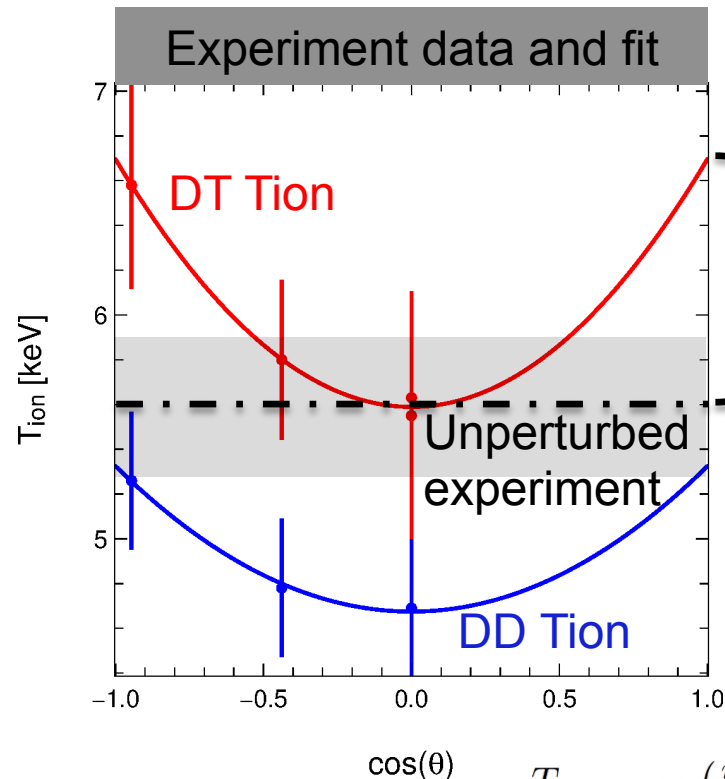
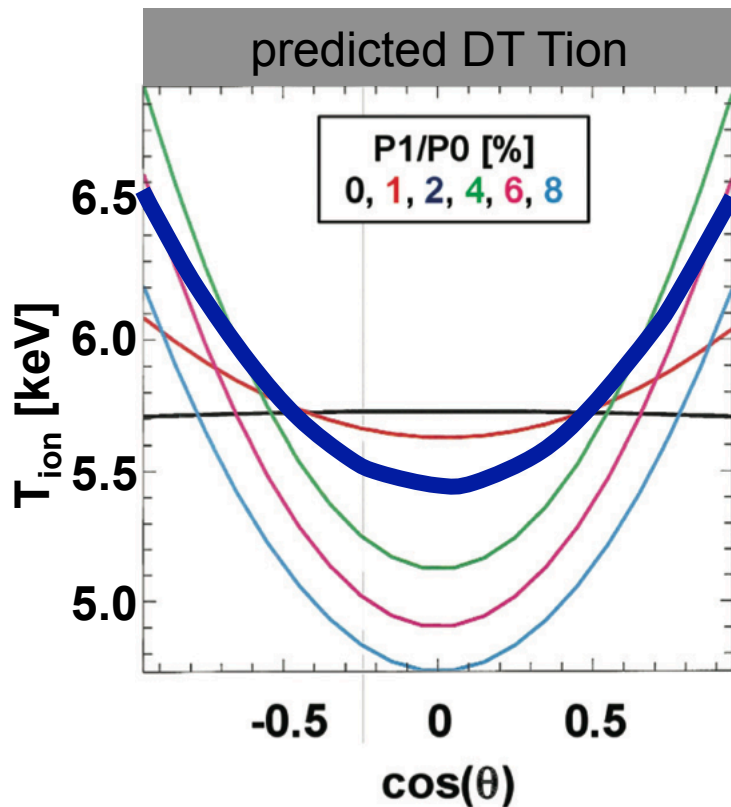
- Thermal temperature is 2.3 keV
- Apparent temperatures span 2.9 to 4.0 keV – depending on direction
- Detector array typically samples 50% of full PTV

Hot spot flow can be estimated from temperature differences



# P<sub>1</sub> perturbed experiments confirm our ability to measure flow-induced temperature variation

- Preshot simulations predict 1 keV temperature variation due to flow
- Experiments show very similar variation, amplitude and shape

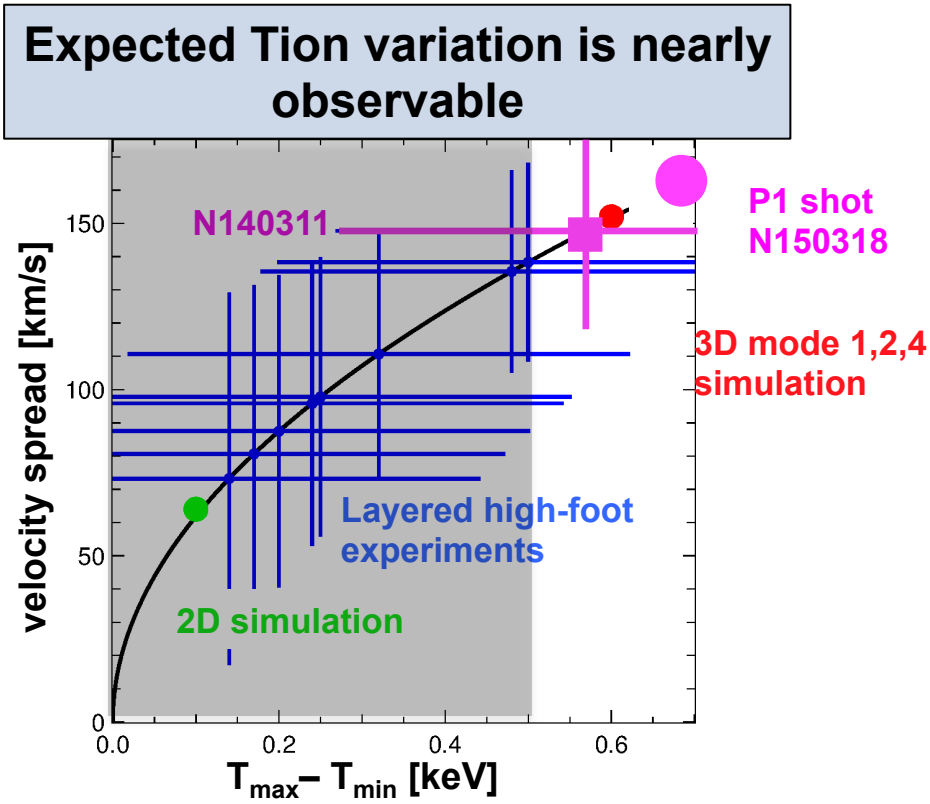


DD/DT gap remains "anomalous"

$$T_{Brysk} = \left( \frac{m_D + m_T}{k} \right) \sigma_v^2 + T_{thermal}$$

**We can measure 1 keV apparent T<sub>ion</sub> anisotropy**

# So, is the high foot apparent $T_{ion}$ usually isotropic or not?



The NIF data cannot (currently) distinguish between isotropy and the expected level of anisotrop

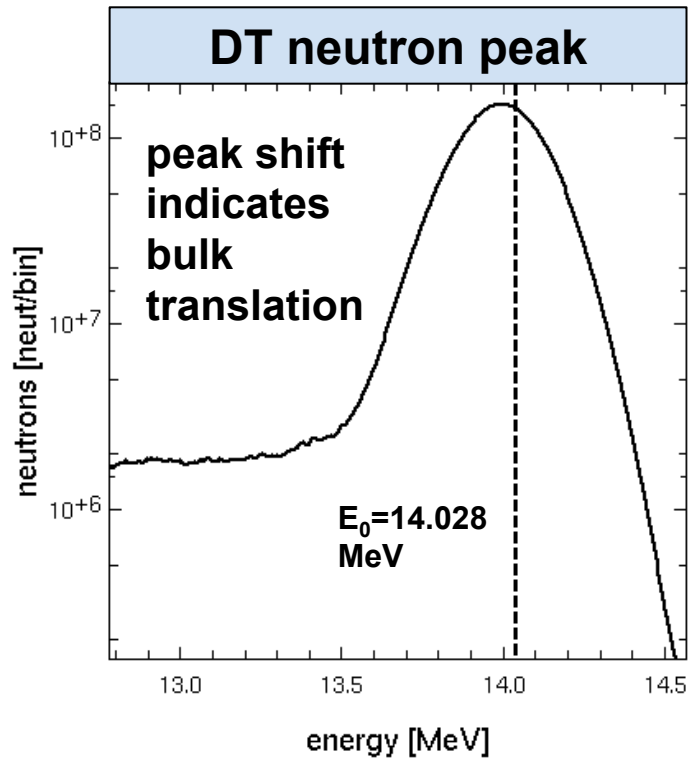
- Post shot simulations suggest Tion anisotropy of ~ 300 - 400 eV
- Detectors would typically sample ~ 150-200 eV
- Detectors can measure down to 500 eV anisotropy

$$T_{Brysk} = \left( \frac{m_D + m_T}{k} \right) \sigma_v^2 + T_{thermal}$$

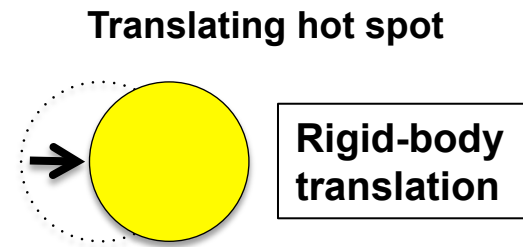
See M. Gatu Johnson paper

**We need neutron spectrometers that can measure 300 eV anisotropy**

# Neutron spectrometers measure bulk velocity from spectral peak shift



Primary neutron peak location gives translational or bulk velocity



$$t = \frac{d}{v_n + \bar{v}_{fluid}}$$

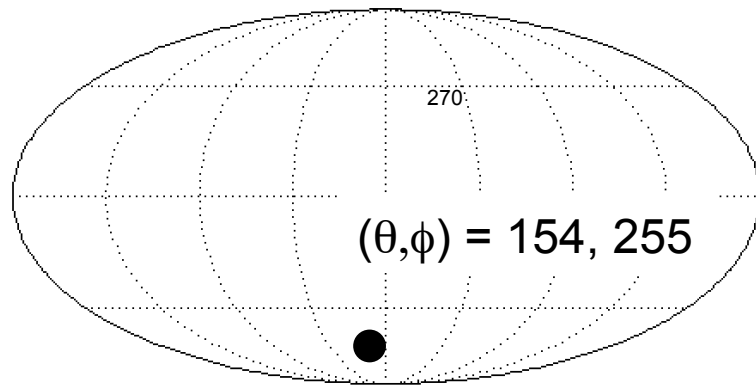
Velocity components measured on 3 nearly orthogonal lines of sight

## Measure speed and direction of hot spot translation

# Mode 1 perturbed experiments confirm our ability to measure bulk flow velocity

## Experimental measurement

85 +/- 15 km/s resultant  
26 degrees off vertical

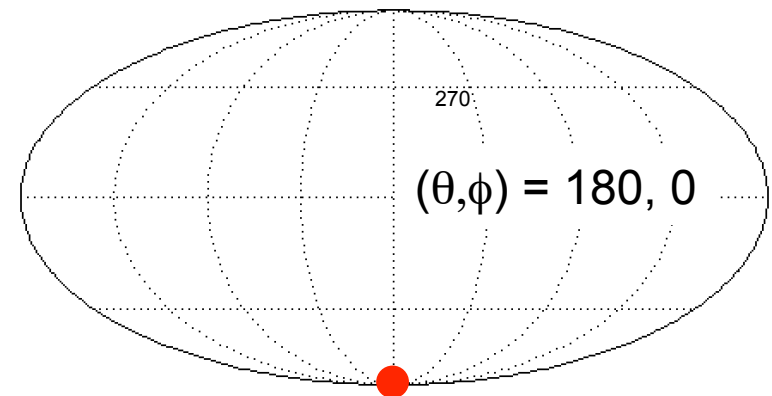


85 +/- 15 km/s

3D effects drive the flow off axis

## Preshot prediction

90 km/s resultant  
directly downward

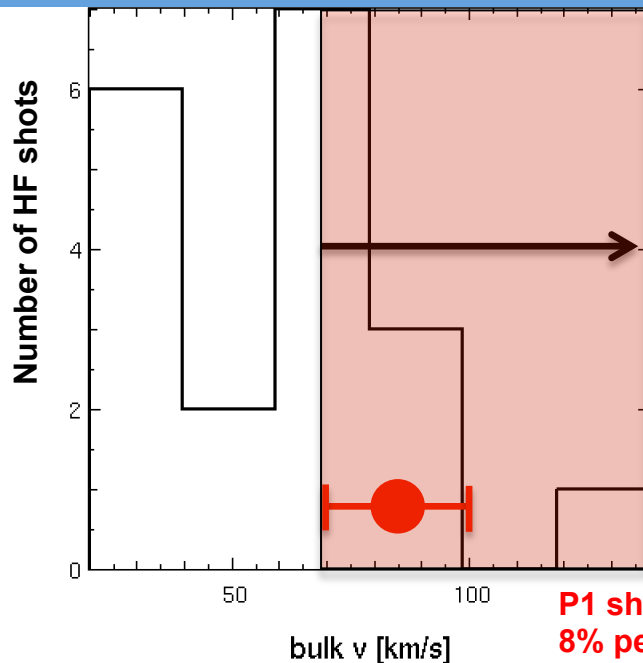


90 km/s

Composition of multiphysics effects (laser propagation, LPI, radiation transport, implosion hydrodynamics) is mainly captured by HYDRA

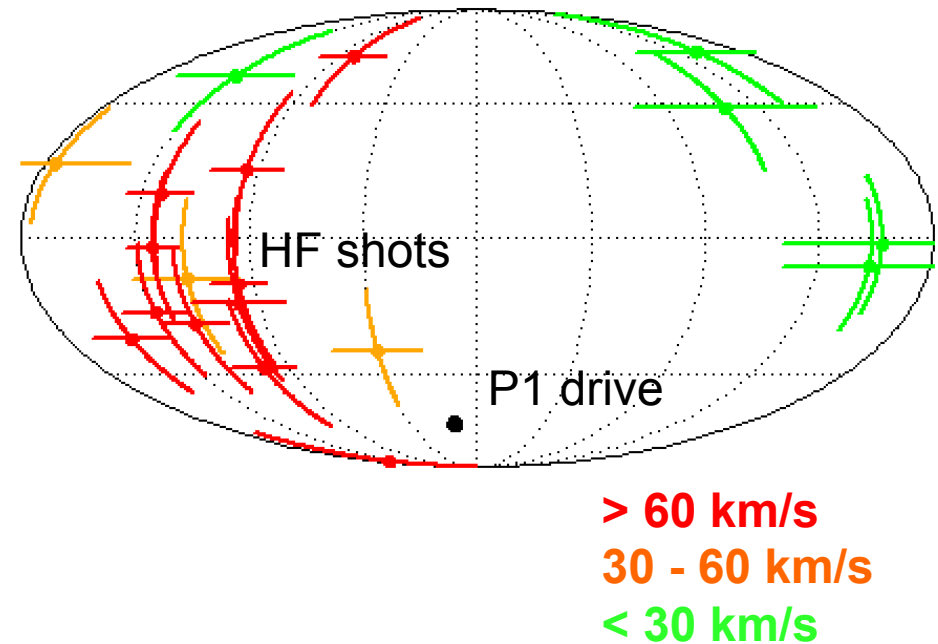
# The average high foot shot bulk velocity is 70% of the intentional $P_1$

Average HF bulk velocity is 60 km/s;  $P_1$  was 85 km/s



**P1 shot N150318**  
8% peak to valley  
power imbalance

Large bulk velocities tend to cluster

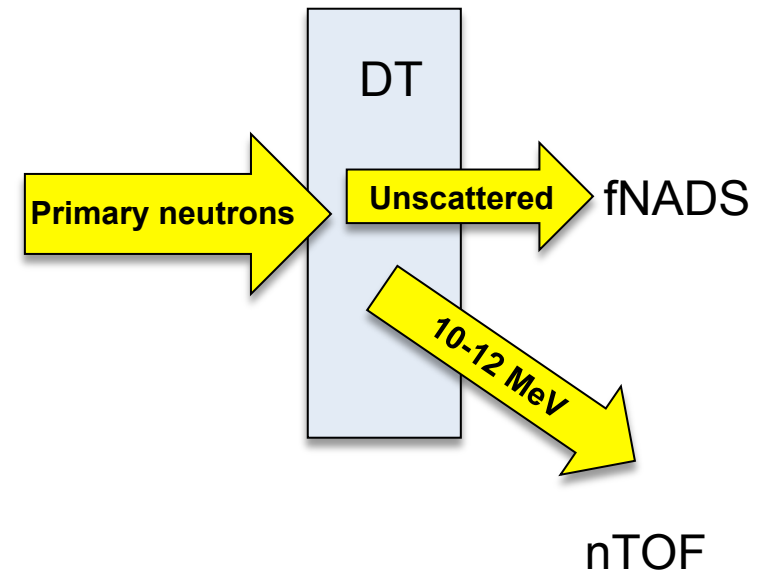


8 of 19 HF shots have velocities larger than the  $P_1$  shot

We haven't yet identified what is producing these perturbations

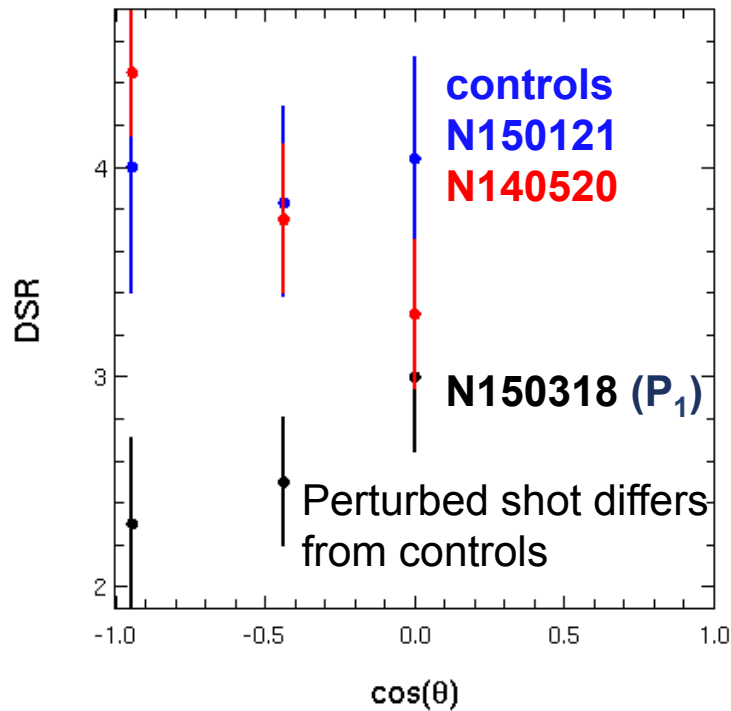
# The cold shell conformation is probed by exiting neutrons

- Neutron spectrometers (nTOF) measure downscattered neutrons
  - High areal density DT scatters into 10 – 12 MeV band
  - Multiple lines of sight measure the asymmetry
- Flange Neutron Activation Diagnostics (fNADS) measure unscattered primary neutrons
  - Zr activated by neutrons above 1X.XX MeV threshold
  - 19 locations on chamber
  - Complementary to DSR

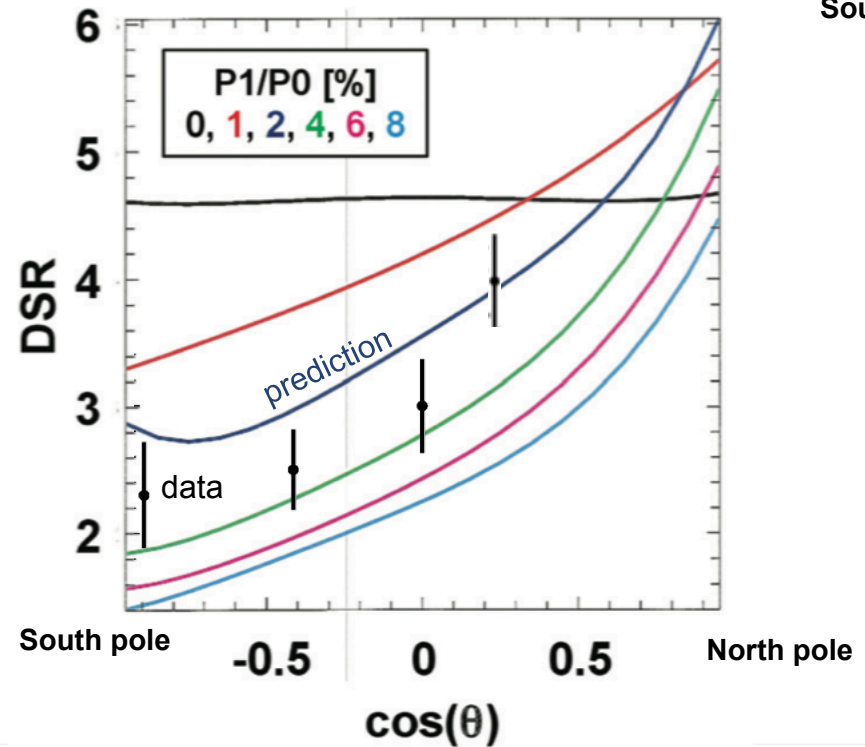
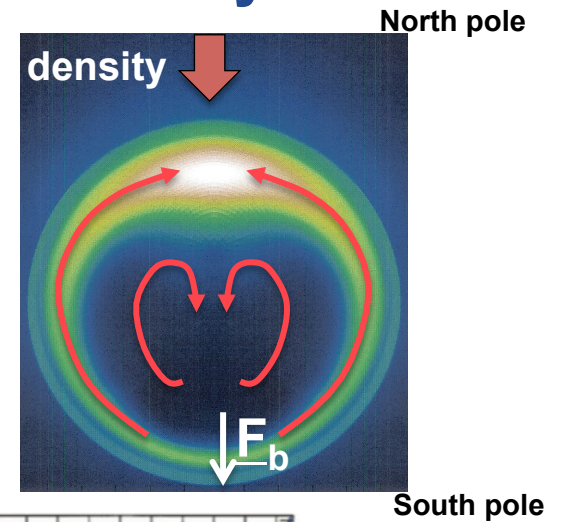


# Mode 1 perturbed experiments confirm our ability to measure angular variation in DSR

Perturbed shot is different from control shots

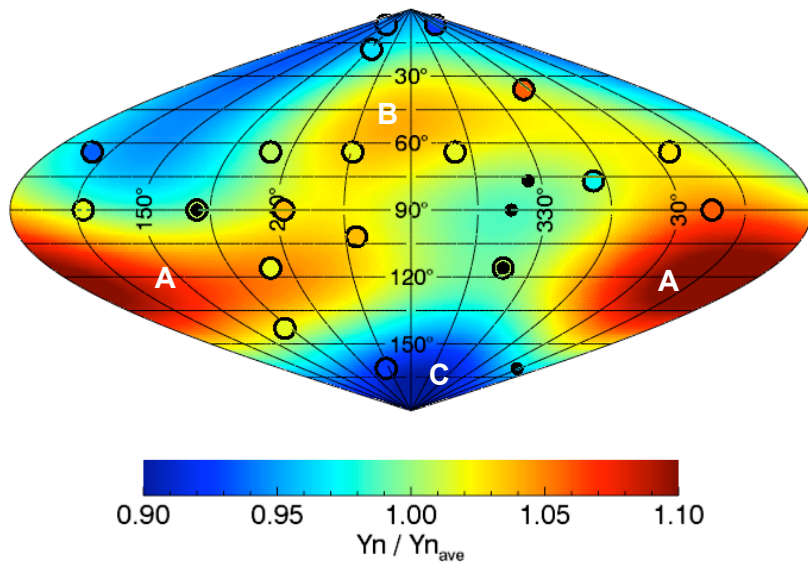


Experiments compare nicely with preshot predictions

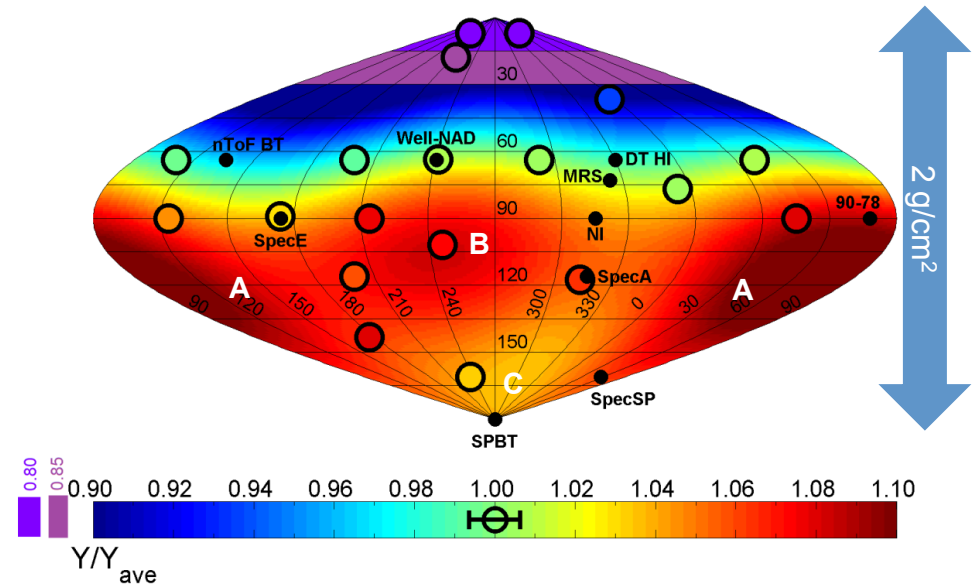


# fNADS measured the predicted angular distribution of escaping primary neutrons

N140520 control shot



N150318 P<sub>1</sub> shot



- Predicted fNADS variation of ~ 25% peak to valley → measured 30%
- Expected P<sub>1</sub> asymmetry → observed P1 plus 3D similar to control shot

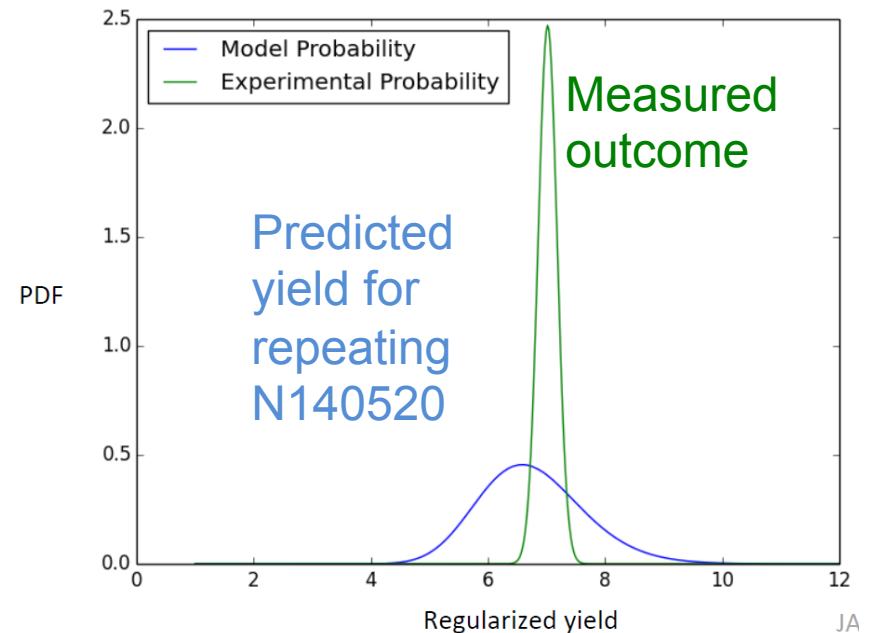
We can predict aspects of the cold shell areal density distribution when the perturbation is large enough



# The repeatability of the unperturbed implosion supports the perturbed results

- We have 3 nominal repeats
  - Yield:  $\mu=7.0e15$ ,  $\sigma=0.5e15$
  - $T_{ion}$ :  $\mu=5.44$ ,  $\sigma=0.087$
- We developed a statistical model of variability using the growing database and Callahan scaling
  - Uses both repeats and other high foot shots
  - Predicted variability compared favorably with a blind test on a repeat shot
- Stagnation properties are repeatable, even if not perfected

Calibration { N140520 = 7.6e15  
N150121 = 6.3e15  
Prediction 6.5e15 +/- 1e15  
Outcome N150409 = 6.9e15



Jim Gaffney, Tammy Ma, Dan Casey, Niko Izumi, Debbie Callahan, Brian Spears

The repeatability of the platform is sufficient for testing perturbation effects

# Reduction in yield was smaller than predicted by single failure mode simulations

- Control shots:  $7.0 \times 10^{15} \pm 0.5 \times 10^{15}$
- P1 shot gave  $4.8 \times 10^{15}$ 
  - Experiment degradation was 30%, observed  $3\sigma$  reduction from control
  - Expected degradation was 60%, observed  $3\sigma$  above expectations

The yield is different from the controls

The yield is different from the prediction

## Control shots

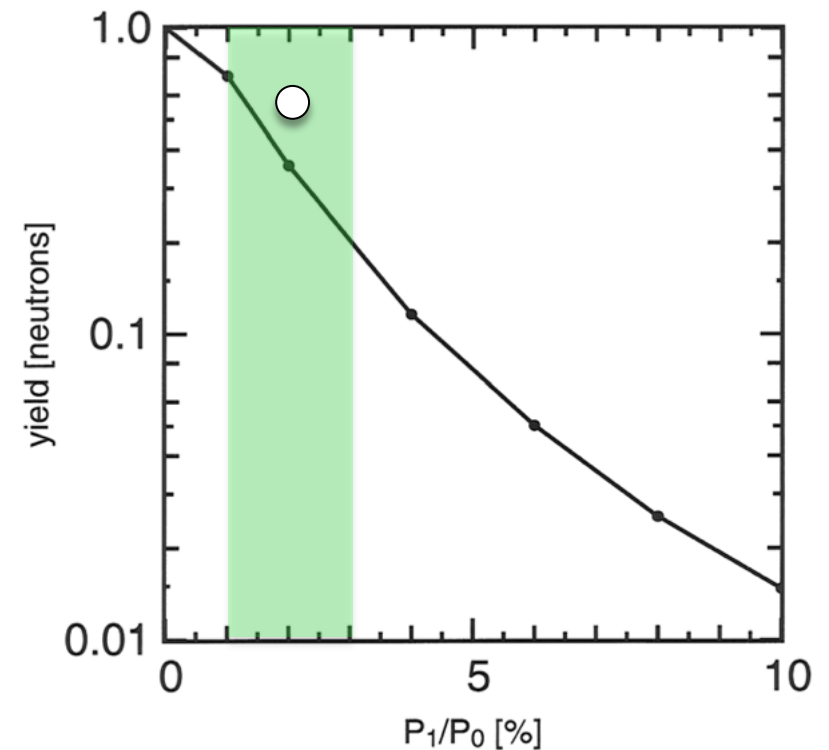
N140520 =  $7.6 \times 10^{15}$

N150121 =  $6.3 \times 10^{15}$

N150409 =  $6.9 \times 10^{15}$

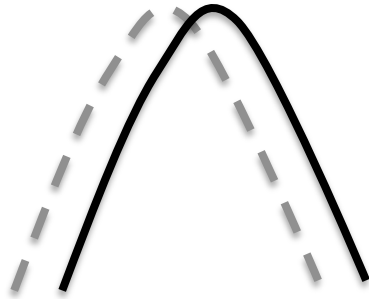
## P<sub>1</sub> shot

N150318 =  $4.8 \times 10^{15}$



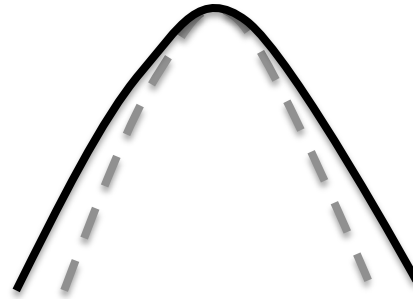
# Stagnation measurements can be much more informative

First moment:  
peak shift  $\sim f(\text{bulk velocity}, T_{\text{thermal}})$



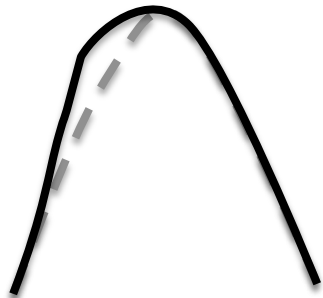
What's the  
bulk  
velocity?

Second moment:  
Width  $\sim f(T_{\text{thermal}}, \text{flow variance})$



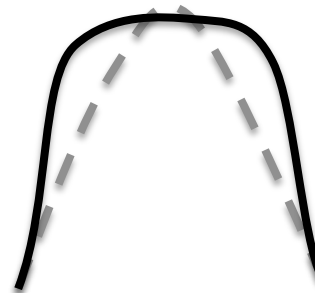
What's the  
apparent  
temp, thermal  
temp, residual  
flow?

Third moment:  
Skew  $\sim \text{cov}(T_{\text{thermal}}, \text{flow})$



Is the hot stuff  
moving fast?

Fourth moment:  
Kurtosis  $\sim \text{variance of } T_{\text{ion}}$



How broad is  
the distribution  
of thermal  
temperatures?

New measurements provide increasingly detailed picture for code validation

# Nuclear diagnosis at NIF provides an unprecedented picture of stagnated ICF implosions

- Hohlraum and capsule symmetry respond to large drive perturbations ( $P_1$ ) as predicted
- Nuclear diagnostics capture the thermodynamics and flow of the hot spot and cold shell
- Simulated hot spot and cold shell diagnostics match experimental observables
- The repeatability of the high foot implosion platform supports perturbed stagnation experiments

Precision diagnostics, platforms, and codes are advancing our validation efforts



**Lawrence Livermore  
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