Mitigating Imprint in Direct-Drive Implosions Using Rarefaction Flows

3-D ASTER simulations of OMEGA cryogenic implosion designs with laser imprint

150

50

-50

t=2.55ns

YOU=0.457

-50

X (um)

150 - 100



Design using unsupported shock

100

50

150

Targets at CR ~ 4

UR

I. V. Igumenshchev University of Rochester Laboratory for Laser Energetics

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Summary

Imprint in direct-drive implosions can be mitigated employing optimized picket

pulses, which produce rarefaction post-shock flows

- Effects of imprint in OMEGA implosions were studied using the 3-D hydrodynamic code ASTER¹
- Simulations show that the development of imprint modulations mainly depends on the type of flows developing inside compressed implosion shells
- Laser pulses with continuous drive launch supported shocks with almost uniform post-shock flows, which lead to enhanced short-scale (*l* ~ 100 to 200) imprint modulations
- Laser pulses with picket(s) launch unsupported shocks with rarefaction post-shock flows, in which sound waves convect imprint modulations outward,² consequently reducing effects of these modulations
- Optimized laser pulses with picket(s) are required to successfully mitigate imprint

¹ I. V. Igumenshchev *et al.*, Phys. Plasmas **23**, 052702 (2016).

² A. L. Velikovich *et al.*, Phys. Plasmas **8**, 592 (2001).





R.C. Shah, R. Betti, E.M. Campbell, V.N. Goncharov, J.P. Knauer, S.P. Regan, and A. Shvydky

University of Rochester Laboratory for Laser Energetics

A.L. Velikovich and A.J. Schmitt

Naval Research Laboratory



Reduction of imprint modulations in implosions driven by laser pulses with

picket(s) has been observed in 2-D simulations for a long time¹

- Suppression of the short-time RT growth of mass modulations because of presence of the fuel-ablator interface¹
- Decay (or relaxation) shock produced by a picket pulse decreases $\gamma_{\rm RT}$ because of increasing the ablative stabilization ("adiabat shaping")²⁻⁴
- Modern numerical simulations, with advances of including the third dimension and using higher spatial resolution, make it possible to reveal another important mechanism of imprint mitigation
 - ¹ T.J.B. Collins and S. Skupsky, Phys. Plasmas 9, 275 (2002).
 - ² V. N. Goncharov *et al.*, Phys. Plasmas **10**, 1906 (2003).
 - ³ K. Anderson and R. Betti, Phys. Plasmas **11**, 5 (2004).
 - ⁴ T.J.B. Collins *et al.*, Phys. Plasmas **11**, 1569 (2004).



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Picketless and picket pulses produce different type flows inside compressed





• The evolution of imprint modulations depends on the type of post-shock flows



Supported shocks lead to enhanced imprint modulations



 High l modes grow as predicted by Goncharov's imprint model*

- Modes with *l* ~100 to 200 grow most efficiently and become dominant by the end of imprinting phase
- These modes are efficiently amplified by the RT growth

$$A_{\ell}(t) = A_{\ell}^0 \cdot \exp(\gamma_{\rm RT} t)$$

* V. N. Goncharov *et al.*, Phys. Plasmas **7**, 2062 (2000).



Unsupported shocks and following rarefaction flows result in decay of short-

wavelength imprint modulations



- Imprint modes experience phase inversion, which are caused by decaying areal mass oscillations*
- This phase inversion results in suppression of high *l* modes
- The spectrum at the end of imprinting phase is dominated by modes with *l* ~ 30
- Modes with l < 30 are hydrodynamically decoupled and not affected by the RT growth during the acceleration phase

* A. L. Velikovich et al., Phys. Rev. E **72**, 046306 (2005).

The shell compressed by the unsupported shock shows much less damages because of imprinting



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Single-picket pulse design (unsupported shock)

* YOU is the neutron yield over the yield in uniform (1-D) simulations

Efficient imprint mitigation strategy requires an optimization of picket pulses



The difference in performance of shots 69236 and 77066 is reproduced by ASTER

simulations



Summary/Conclusions

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