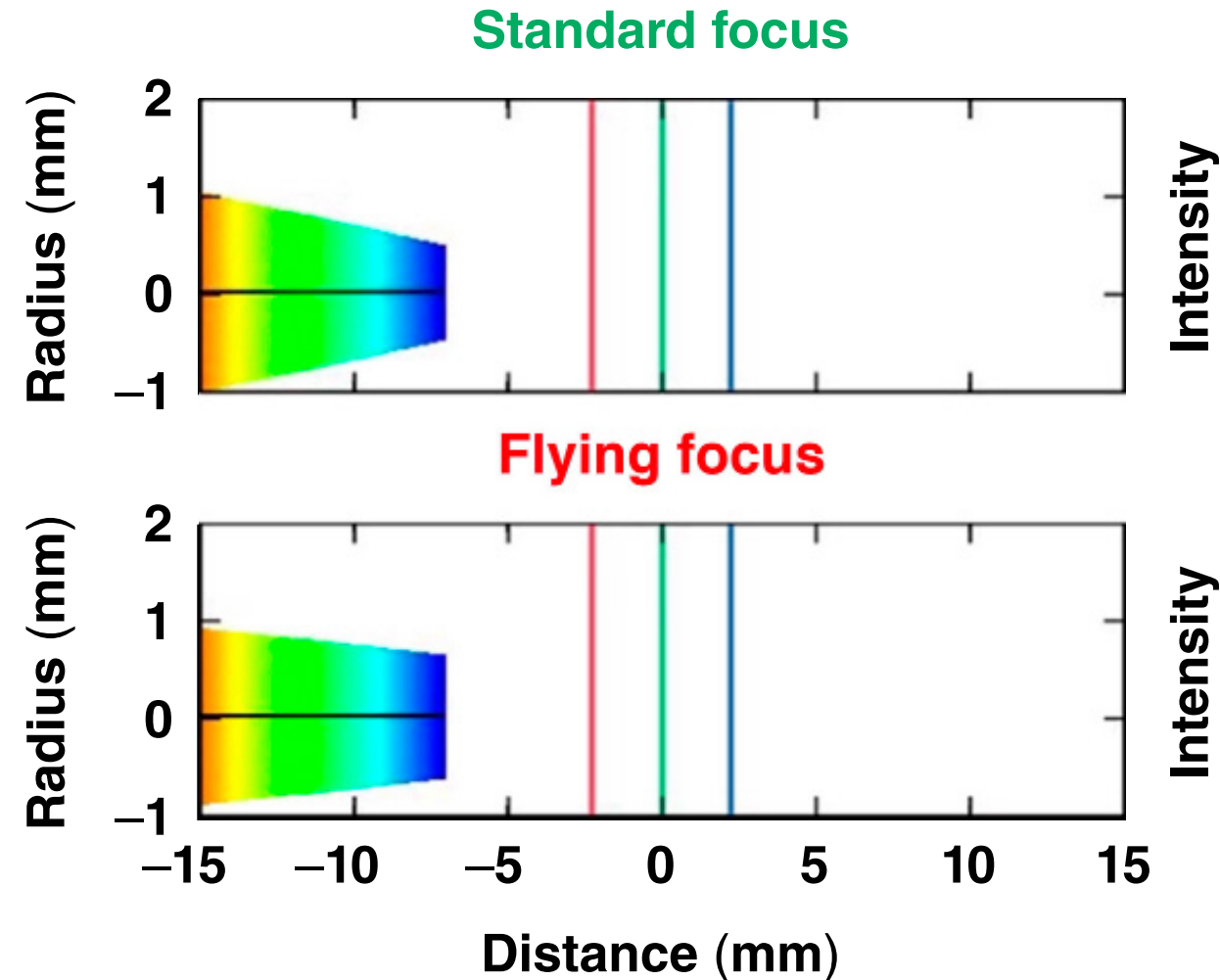
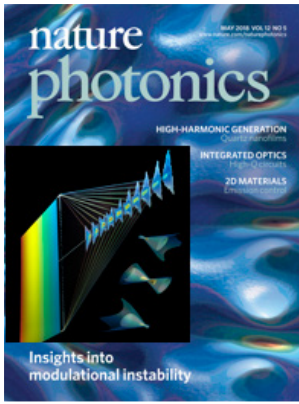


Flying Focus: Spatiotemporal Control of Intensity for Laser-Based Applications



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Plasma Physics Group Leader, Laboratory for Laser Energetics
University of Rochester

60th Annual Meeting of the
American Physical Society
Division of Plasma Physics
Portland, OR
5–9 November 2018



Summary

A chirped laser pulse focused by a chromatic lens exhibits a dynamic, or “flying,” focus*



- The flying focus provides unprecedented spatiotemporal control over laser–plasma interactions by decoupling
 - the spot size of the pulse from the focal range
 - the velocity of the peak intensity from the group velocity
- Experiments have demonstrated the flying focus and the ability to generate ionization waves at any velocity (IWAV)
- Flying focus was applied to several applications
 - photon accelerator: IWAV’s can shift visible laser light to the XUV
 - Raman amplification: flying focus could overcome several challenges of laser-plasma amplifiers
 - Cherenkov radiation: flying focus allows new radiation sources
 - vacuum electron acceleration: flying focus enables vacuum acceleration

*D. H. Froula *et al.*, Nat. Photonics **12**, 262 (2018);
A. Sainte-Marie, O. Gobert, and F. Quéré, Optica **4**, 1298 (2017).
XUV: extreme ultraviolet

We have an outstanding research team working on the flying focus and its applications



J. P. Palastro, D. Turnbull, T. J. Kessler, A. Davies, P. Franke, A. Howard, L. Nguyen, D. Ramsey, G. W. Jenkins, S.-W. Bahk, I. A. Begishev, R. Boni, J. Bromage, S. Bucht, R. K. Follett, D. Haberberger, J. Katz, and J. L. Shaw

**University of Rochester
Laboratory for Laser Energetics**

**F. A. Hegmann and D. Purschke
University of Alberta**

**N. Vafaei-Najafabadi
Stoney Brook University**

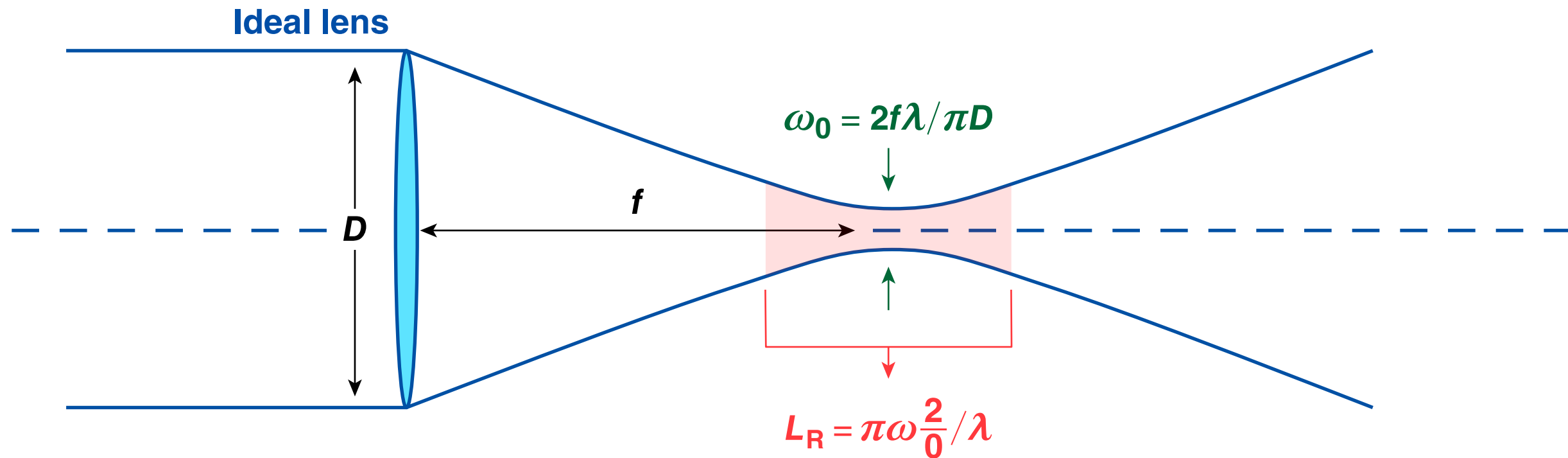
**J. Vieira
IST Lisbon**

**F. Quéré
CEA**

This material is based upon work supported by the U.S. Department of Energy under Award Number DE-SC0017950, and by the Department of Energy National Nuclear Security Administration under Award Number DE-NA0003856.



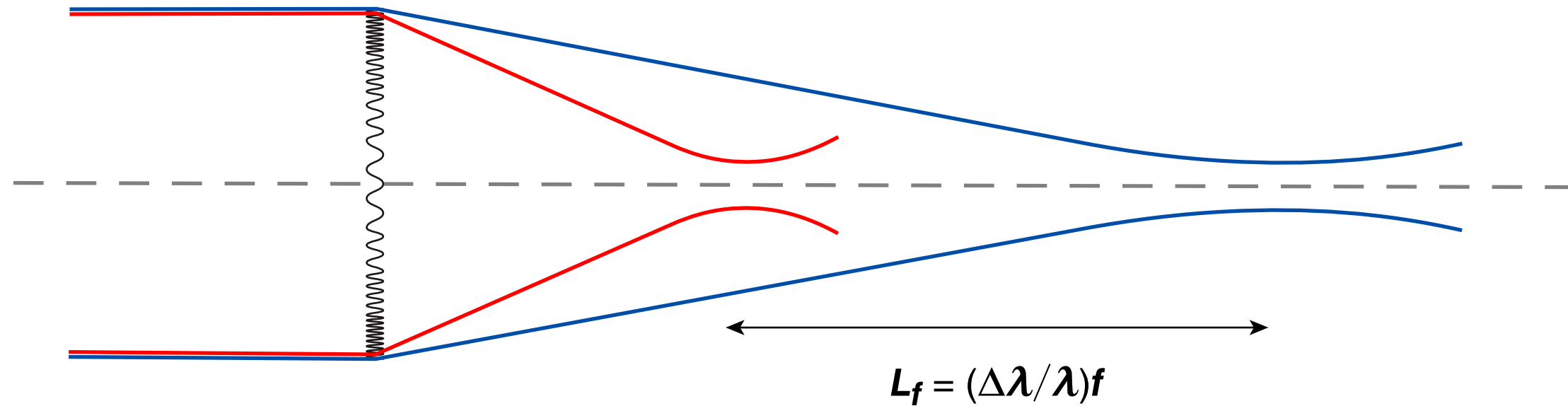
Ideal lenses limit the region of high intensity to the Rayleigh range



For fixed pulse power (P), increasing the Rayleigh range necessarily *decreases* the intensity, $I \sim P/\omega_0^2 \sim P/L_R$.

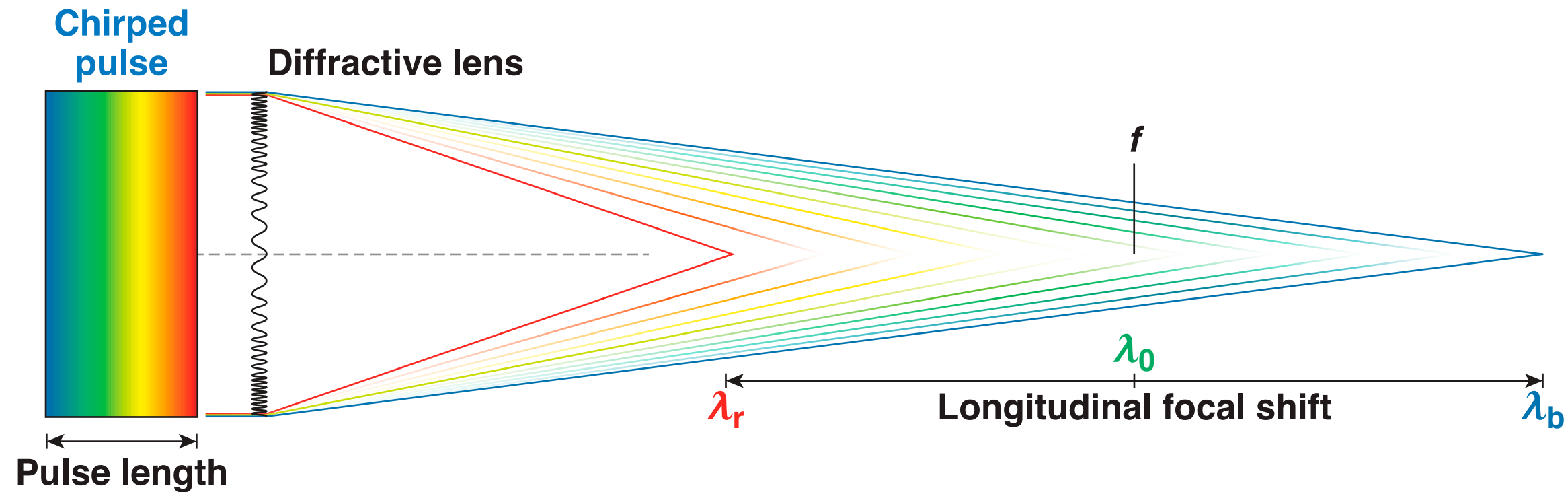
A diffractive lens has a different focal length for each color

Diffractive lens



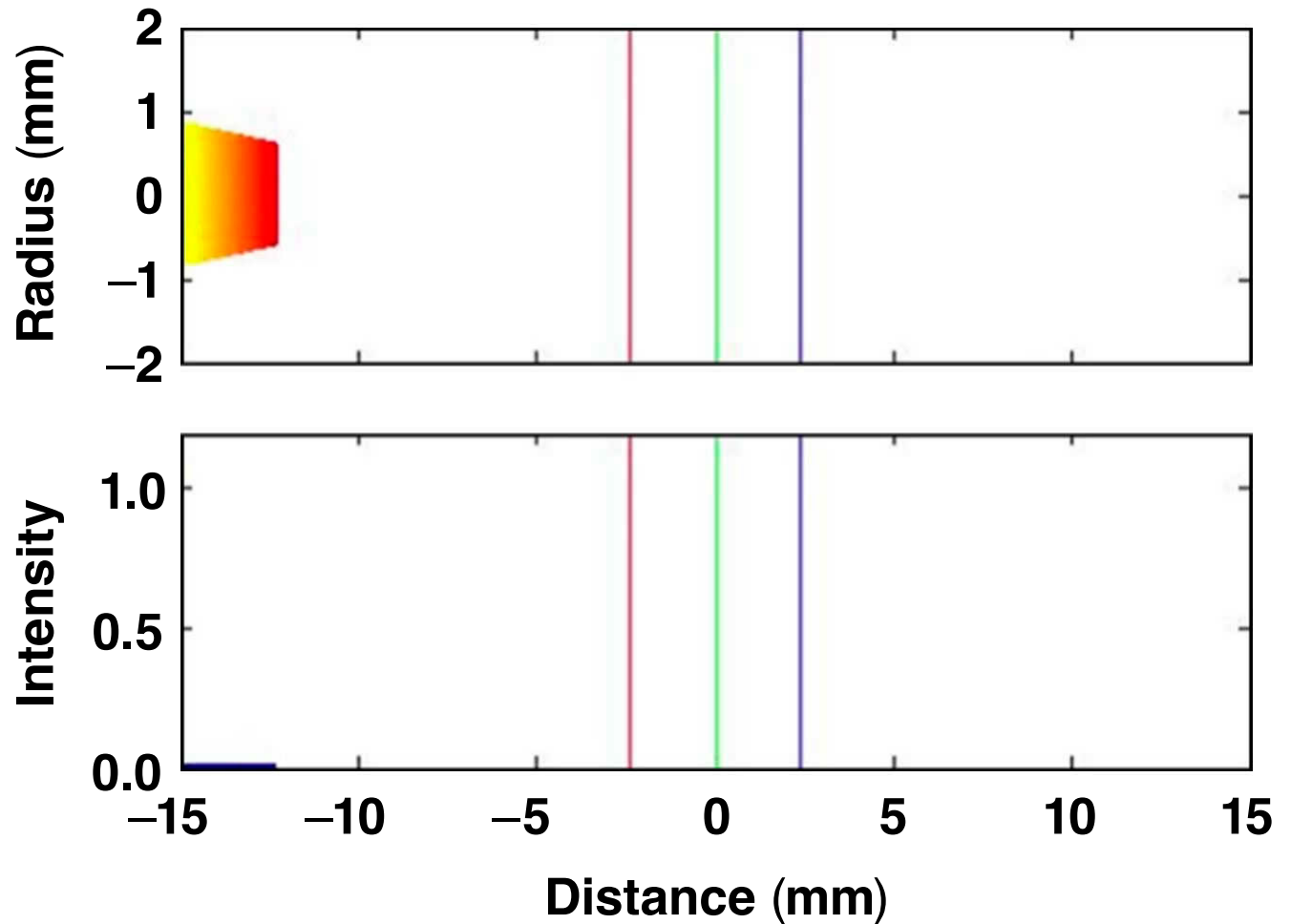
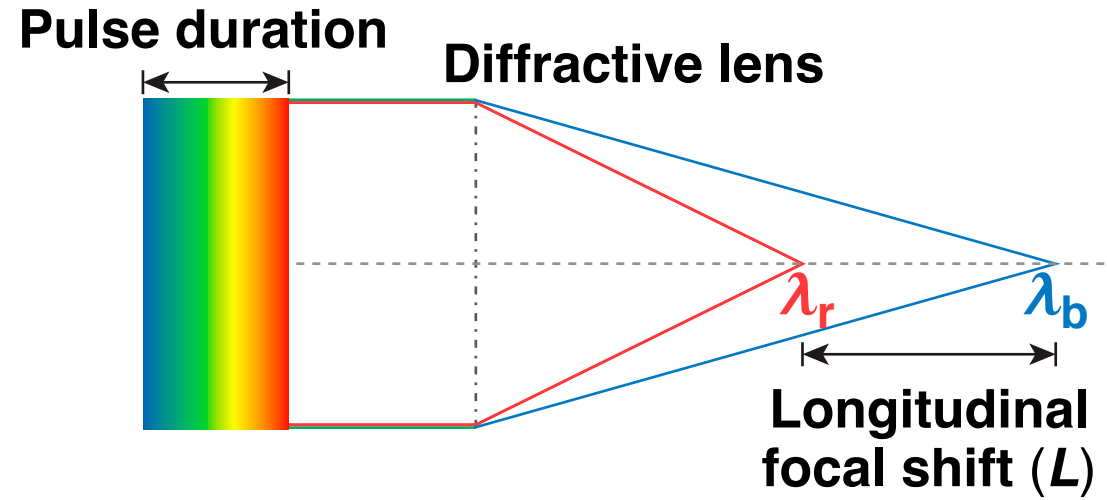
With only 10 nm of bandwidth, the distance separating focused colors can be $\sim 100\times$ greater than the Rayleigh length, *extending the range of high intensity.*

Combining a diffractive lens with a chirped laser pulse provides spatiotemporal control over the focus



The spectral phase of the pulse determines the time at which color reaches focus, resulting in a peak intensity with a dynamic trajectory.

The dynamic focus can propagate over 100× the Rayleigh length of the system



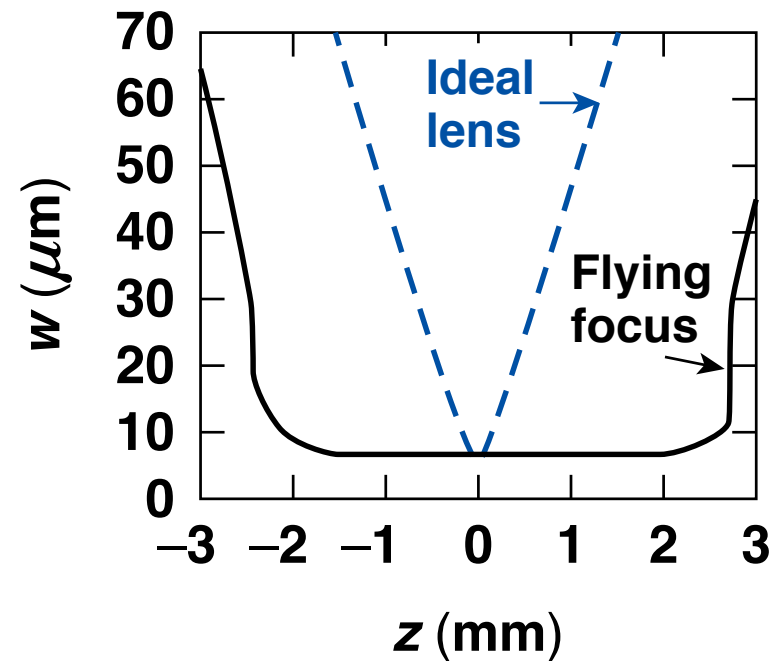
The time between red and blue colors (i.e., pulse duration) determines the time to move from the first focus to the last (velocity).

$$L = \frac{\Delta\lambda}{\lambda} f_0 \cong 4.5 \text{ mm}$$

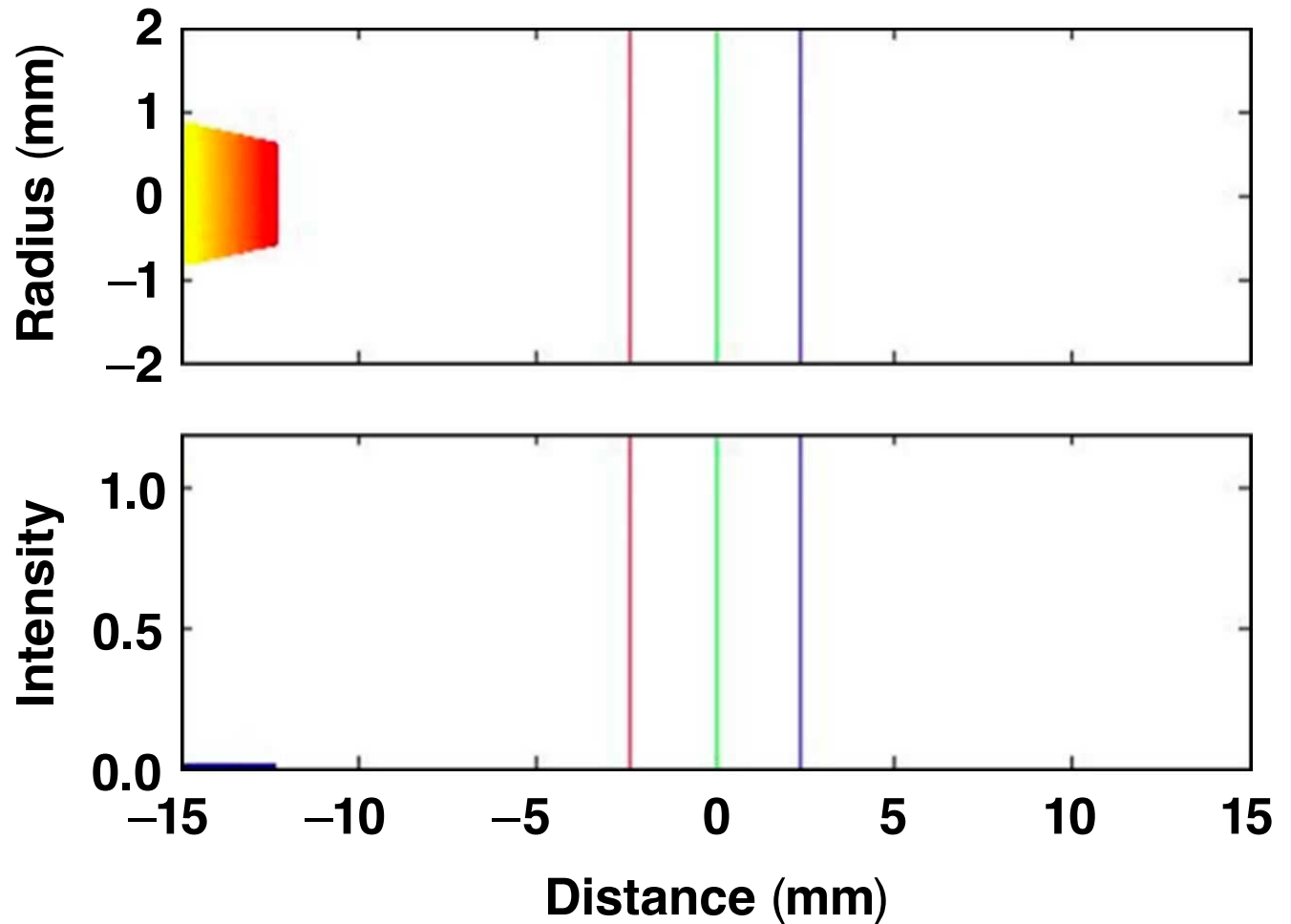
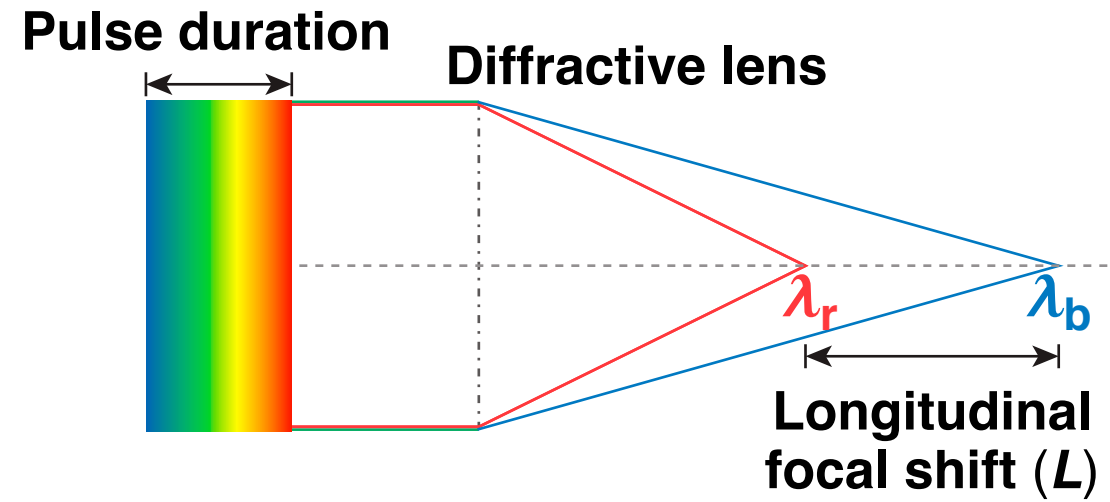
$$\lambda_0 = 1054 \text{ nm}$$

$$\Delta\lambda = 9.2 \text{ nm}$$

$$f_0 = 511 \text{ mm}$$



The dynamic focus can propagate over 100× the Rayleigh length of the system



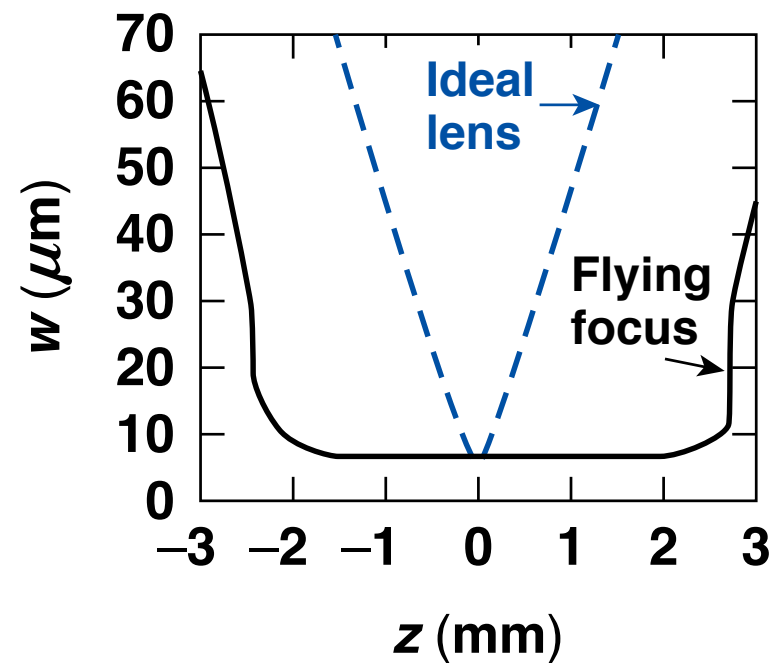
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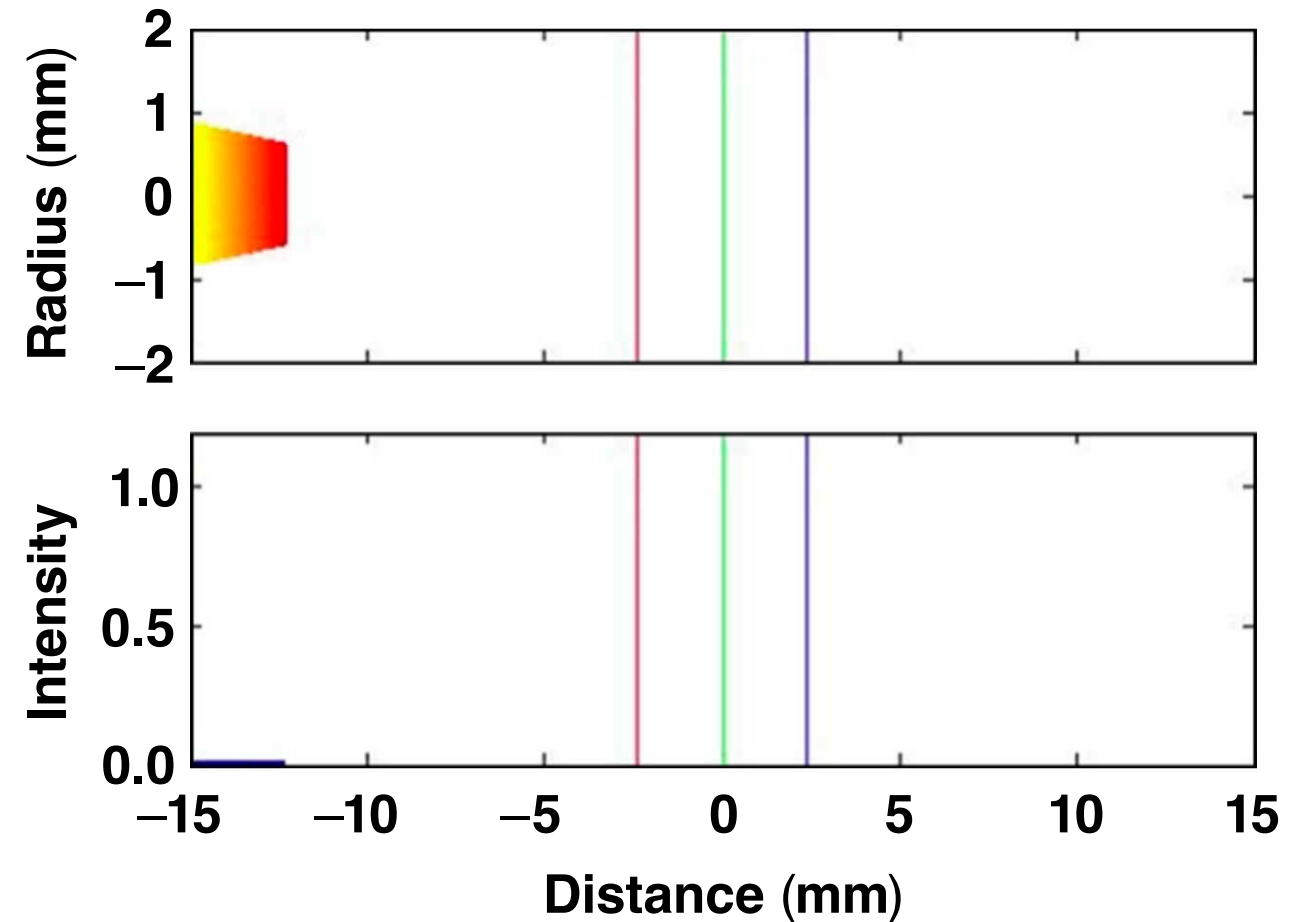
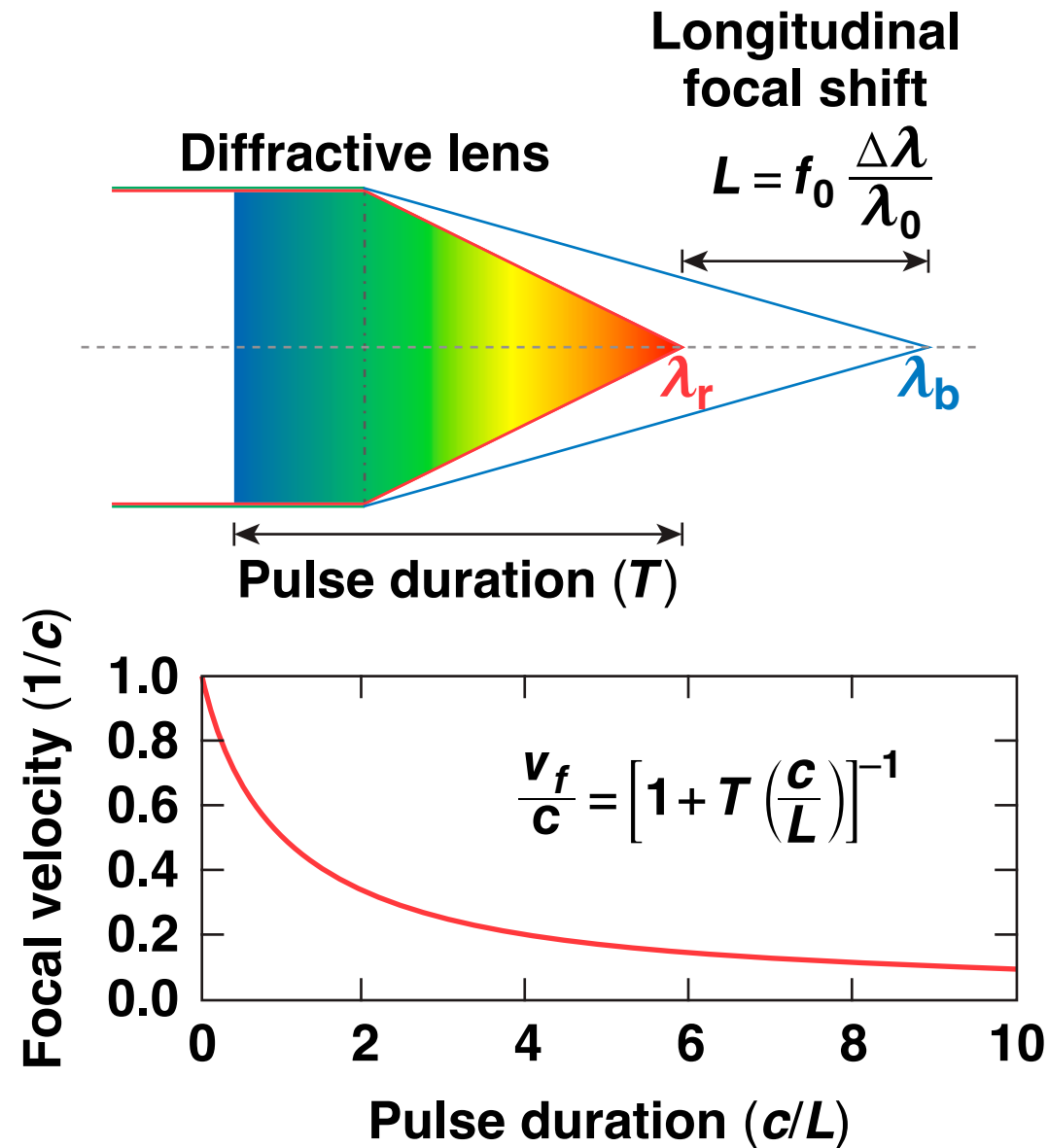
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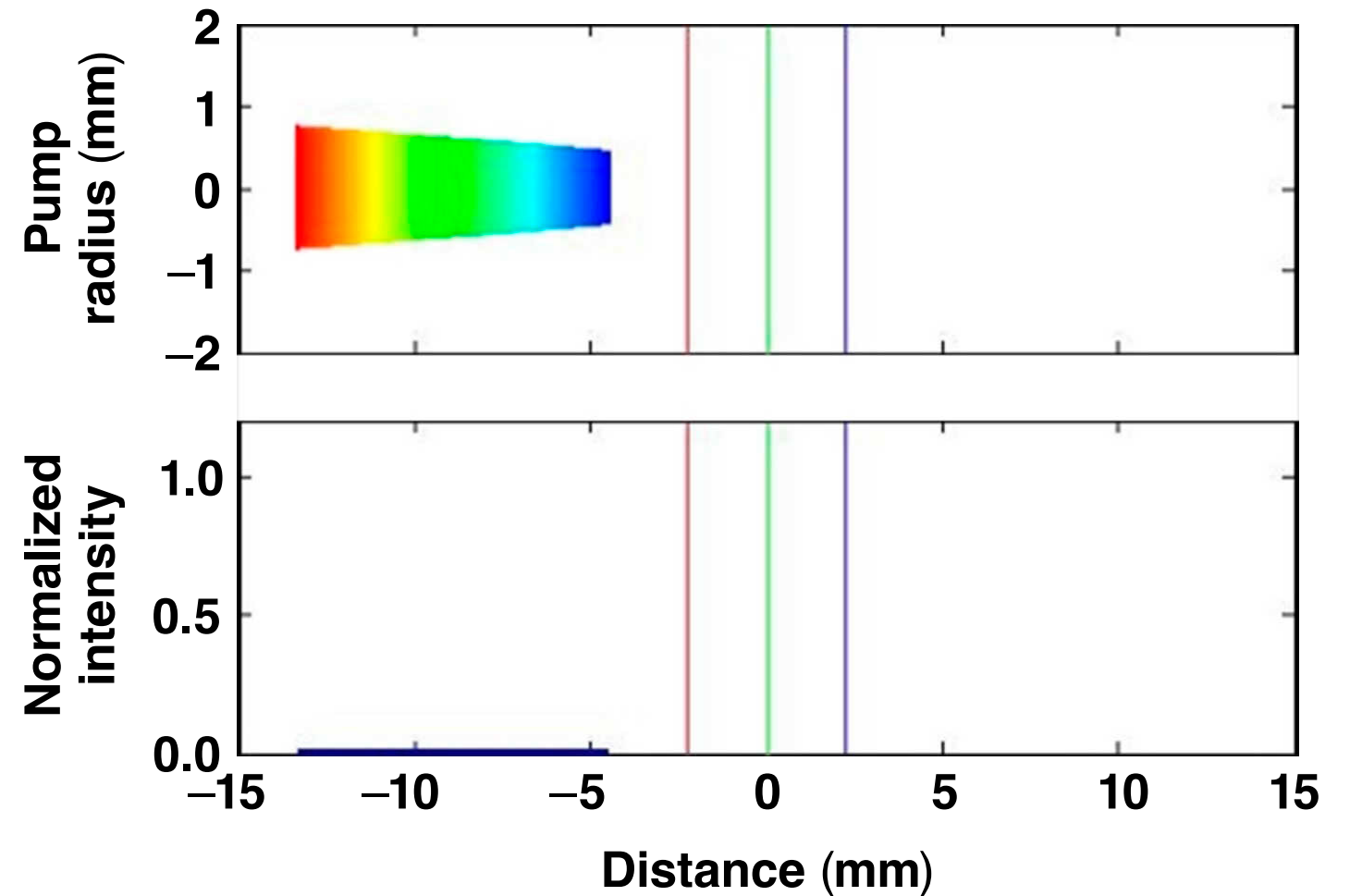
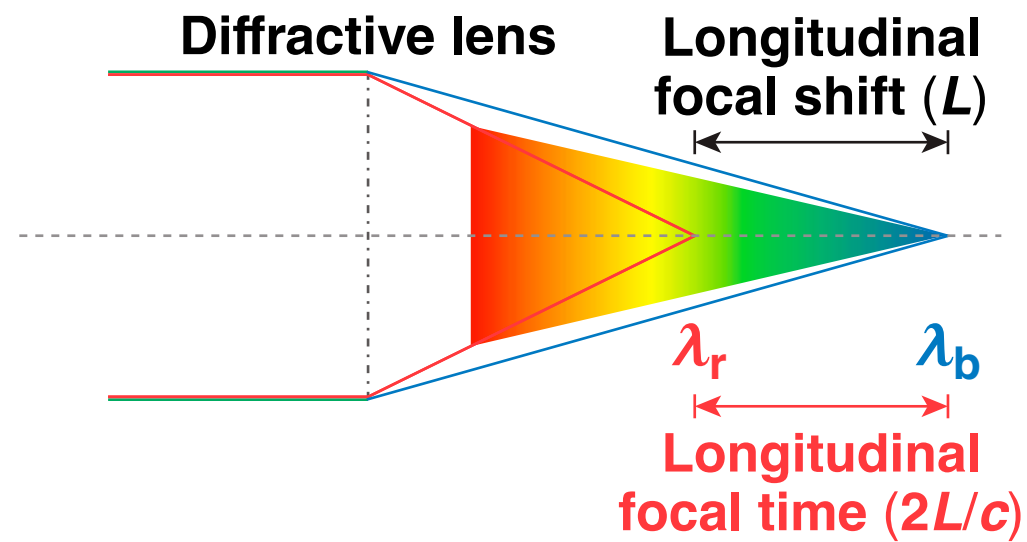


By varying the pulse duration (chirp) of the laser (T), the velocity of the focus can be controlled

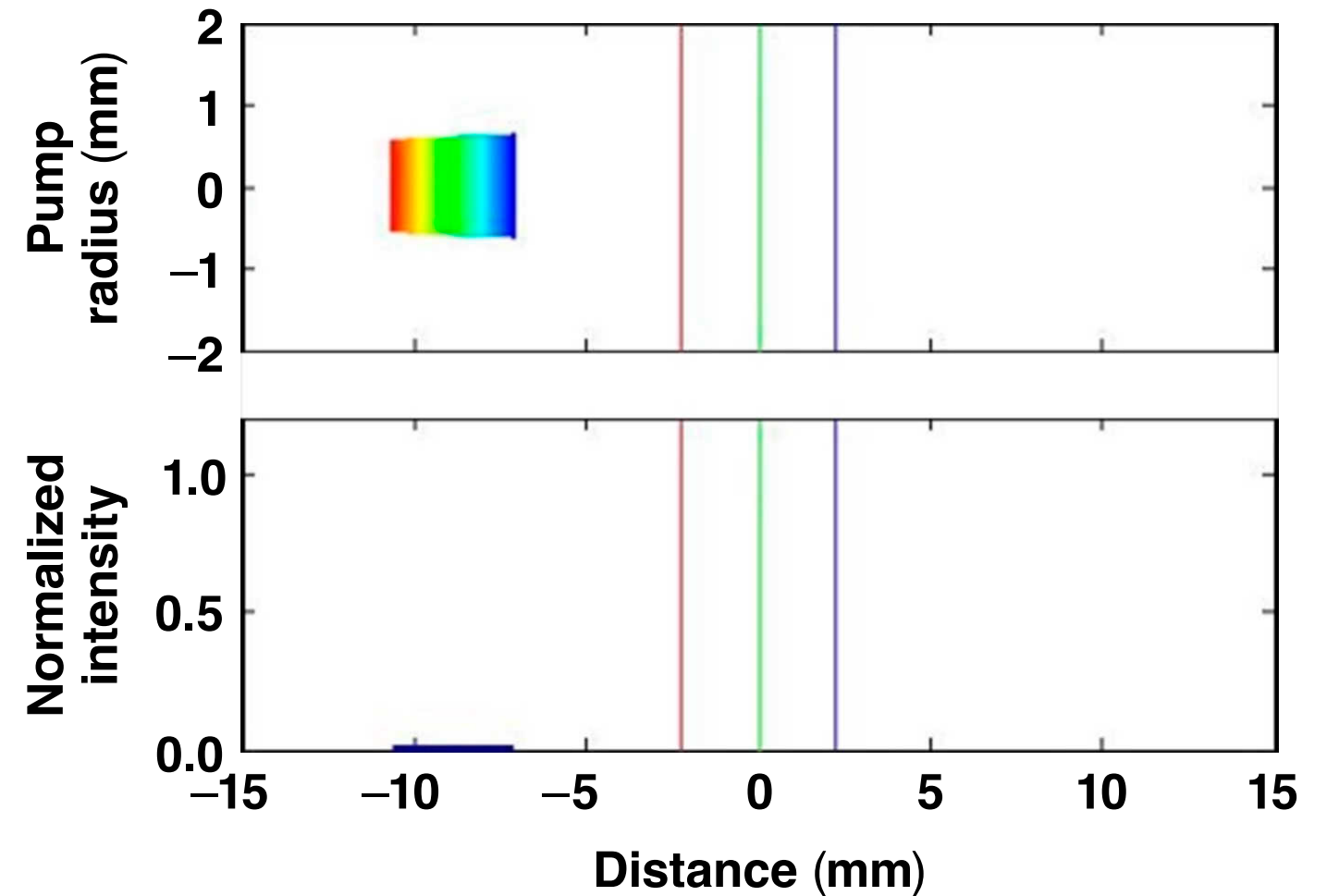
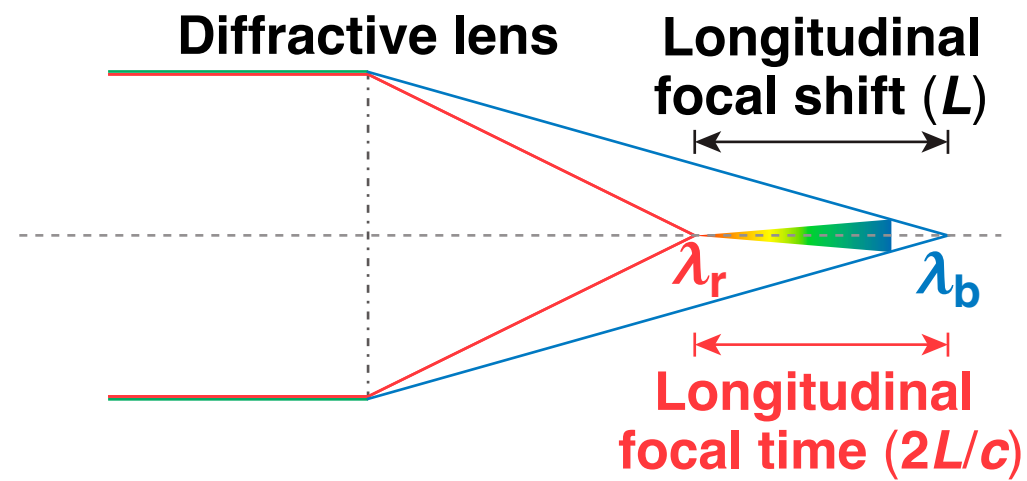


The time between red and blue colors (i.e., pulse duration) determines the time to move from the first focus to the last (velocity).

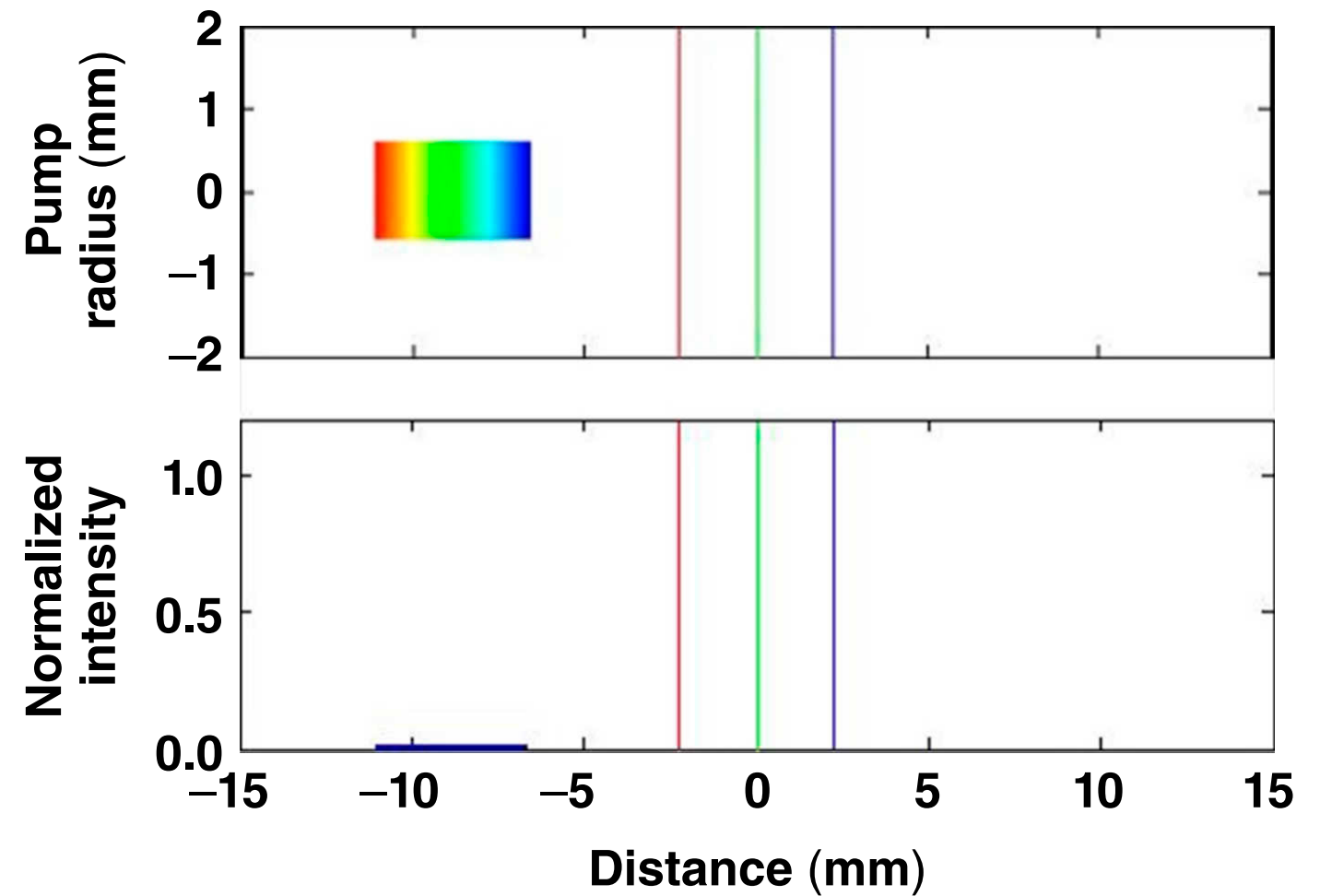
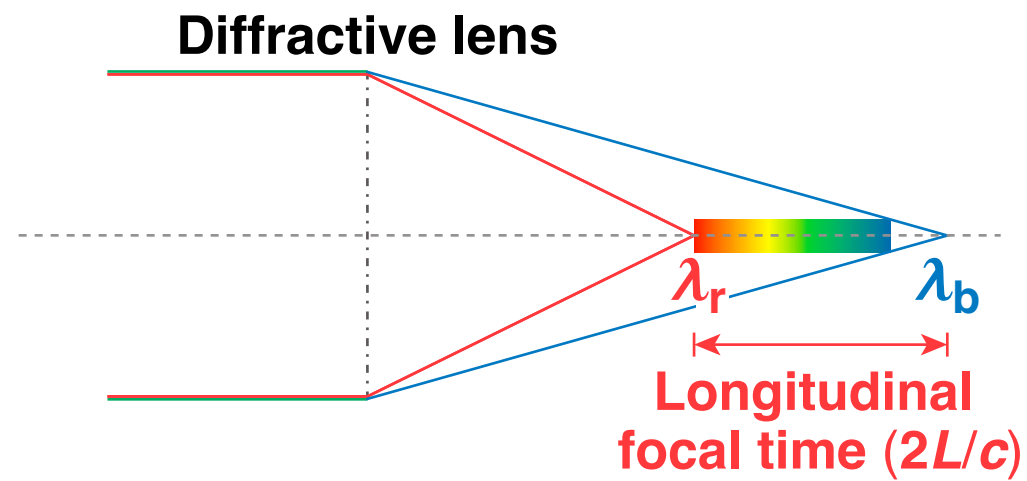
By changing the direction of the chirp (blue to red)
the focus can be made to counter-propagate



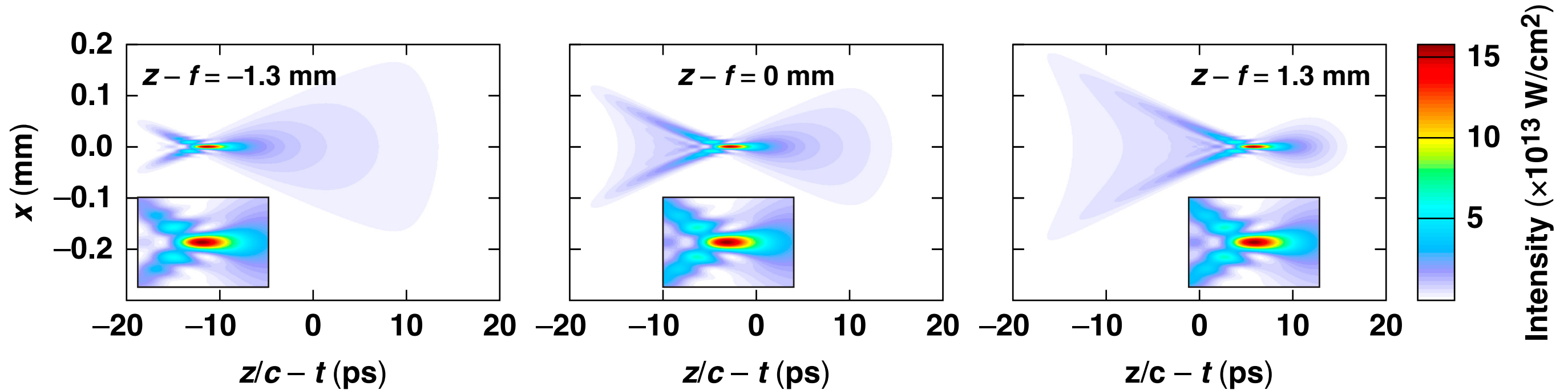
Reducing the pulse duration of the negatively chirped beam (blue to red) produces a focal velocity faster than the speed of light



Setting the pulse duration equal to the focal range (L/c) results in an “infinitely” fast focal velocity (line focus)



A simulation of the focal region shows that the peak intensity of the flying focus propagates with a self-similar form



The longitudinal profile of the intensity peak depends only on space and time in the combination

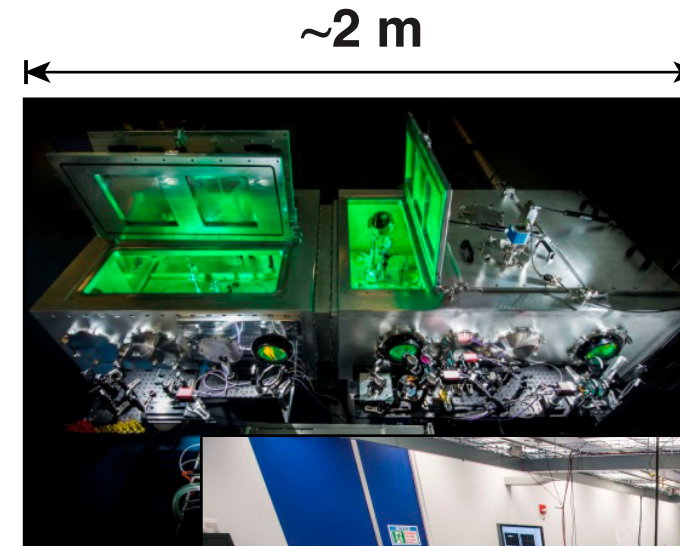
$$\left[1 - \frac{\Delta\lambda(z - ct)}{\lambda c T} \right]^{-1} z$$

Outline

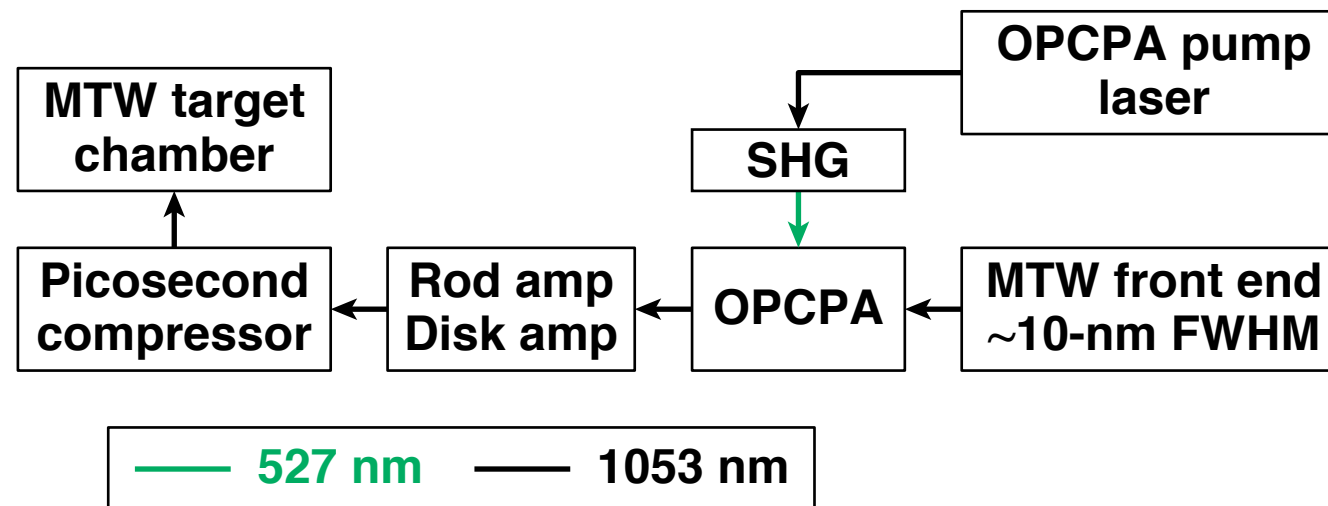
- Description of the flying focus
- **Experimental demonstration of the flying focus**
- Ionization waves of arbitrary velocity (IWAV's)
- Applications of the flying focus

The Multi-Terawatt (MTW) laser at the University of Rochester's Laboratory for Laser Energetics (LLE) was used to demonstrate the flying focus

- OPCPA front end and Nd:glass amplifiers (1053-nm, 10-nm bandwidth)
- 0.7 ps to 300 ps up to 50 J

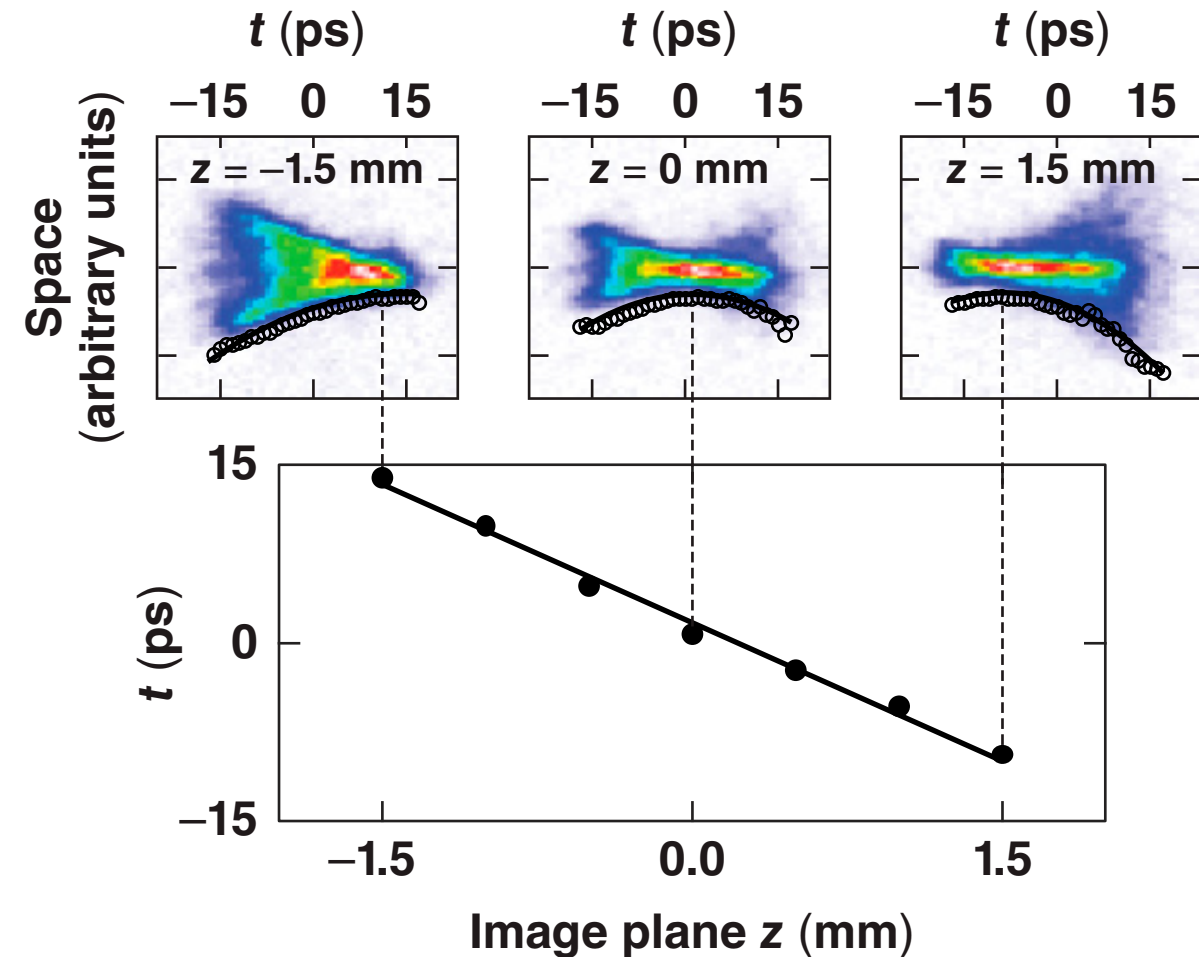
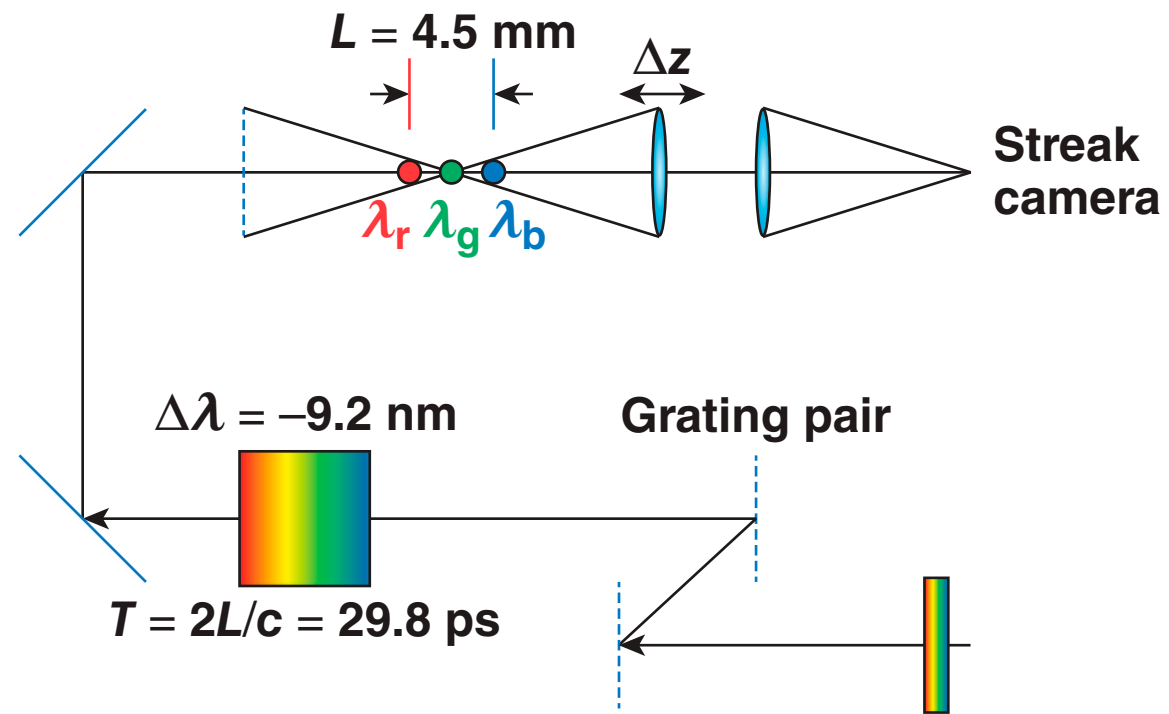


MTW target area



OPCPA: optical parametric chirped-pulse amplifier
 SHG: second-harmonic generation

A picosecond optical streak camera imaged the intensity profile of the flying focus



Several locations in the focal region were imaged onto a streak camera, providing the spatiotemporal profile of the flying focus pulse.

The measurements show excellent agreement with the analytic calculations

Focal velocity:

$$\frac{v_f}{c} = \left(1 \pm \frac{cT}{L}\right)^{-1}$$

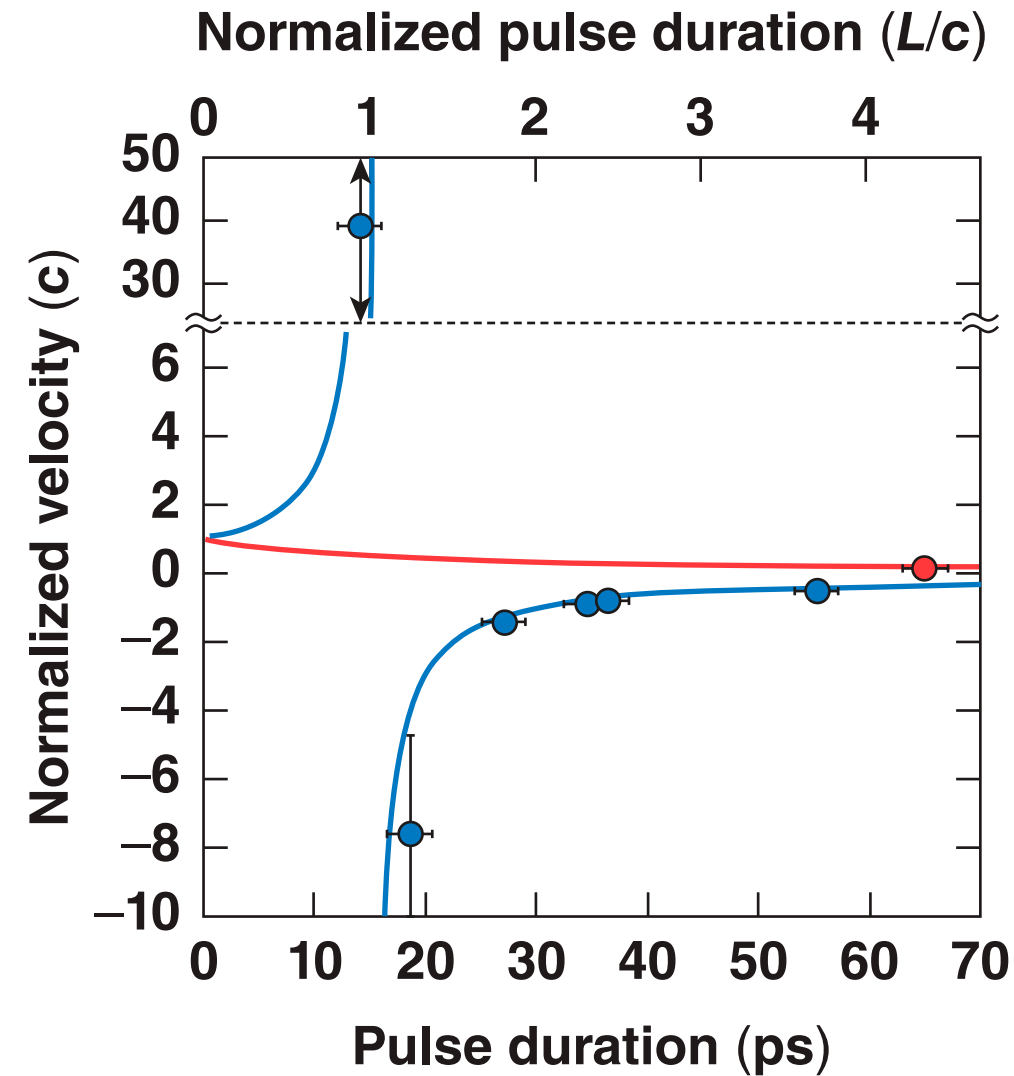
T : stretched pulse duration

$$L = f_0 \frac{\Delta\lambda}{\lambda_0} \cong 4.5 \text{ mm}$$

$$\lambda_0 = 1054 \text{ nm}$$

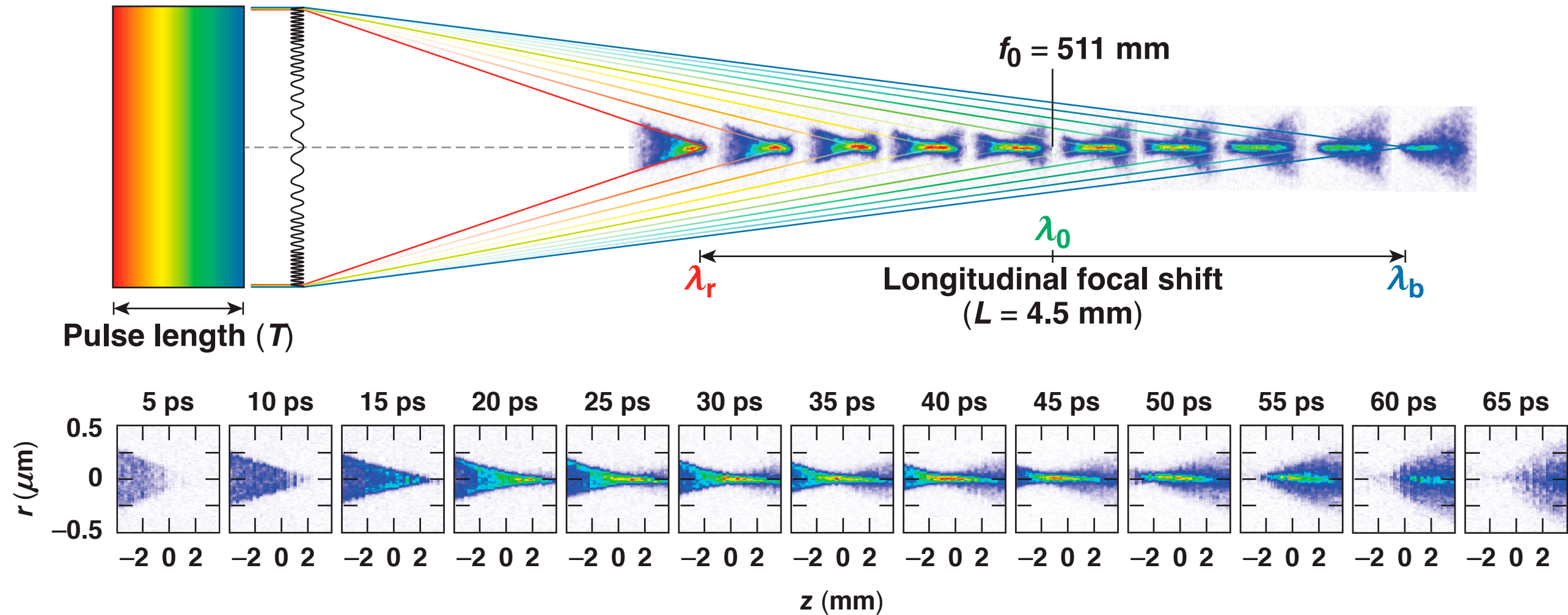
$$\Delta\lambda = 9.2 \text{ nm}$$

$$f_0 = 511 \text{ mm}$$



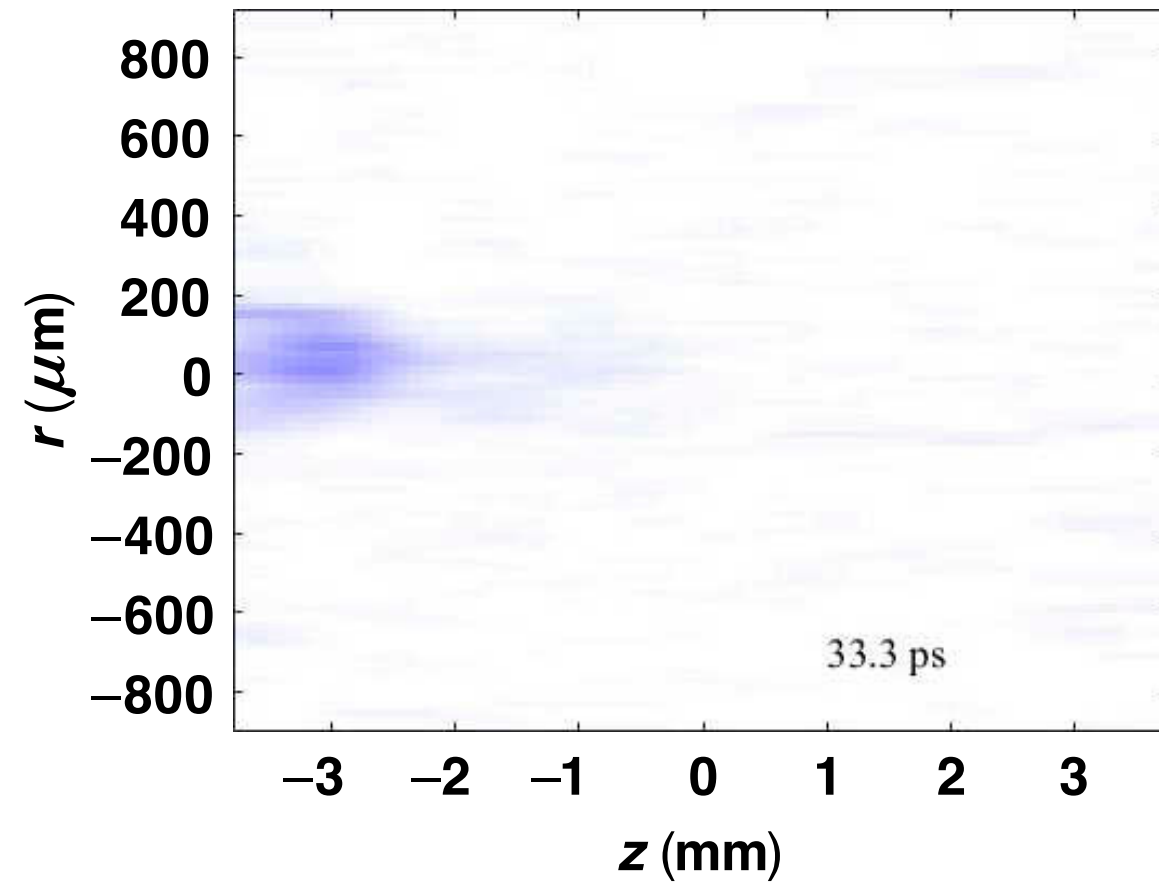
The measured images provided space and time information that were reconstructed to generate the focal intensity in space

Negative chirp
(blue to red)
 $\Delta\lambda = -9.2 \text{ nm}$ Diffractive lens

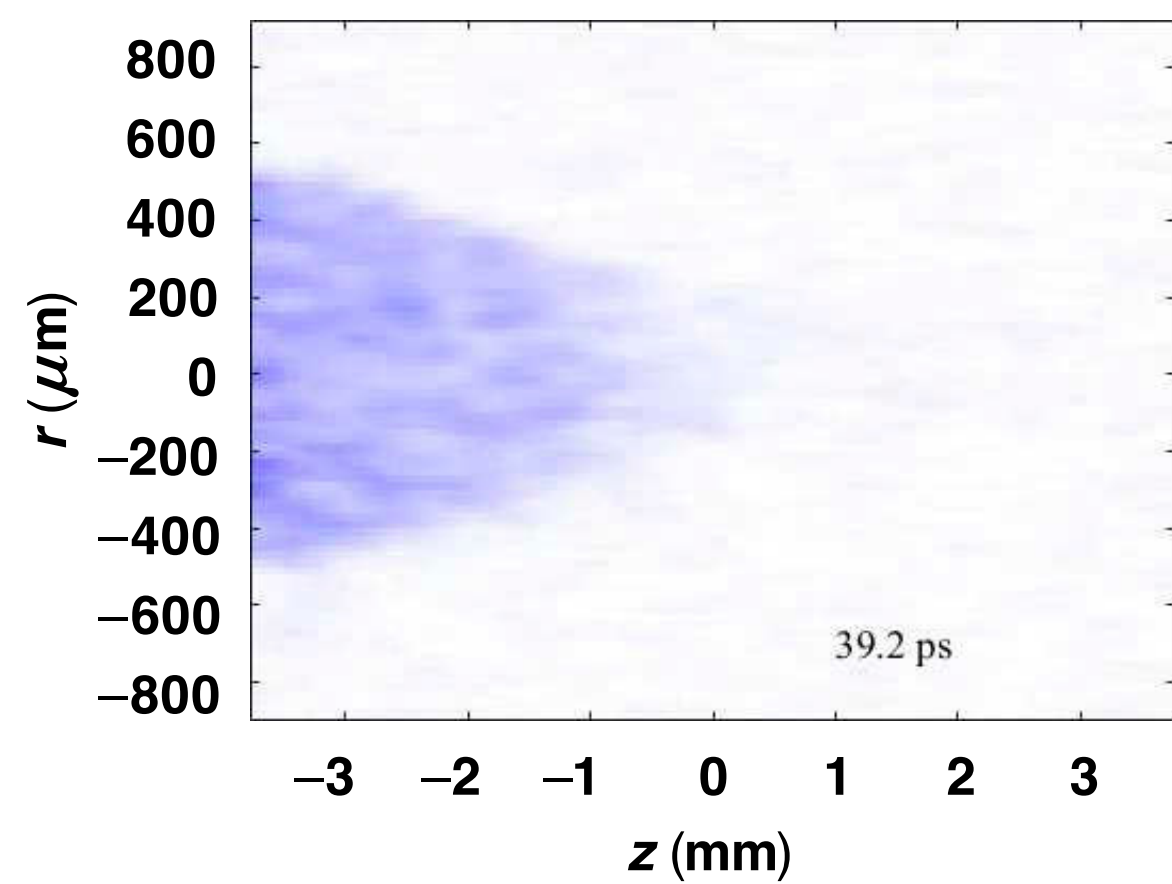


The images form a movie of the flying focus (~ 3 ps/frame)

Positive chirp (65 ps)

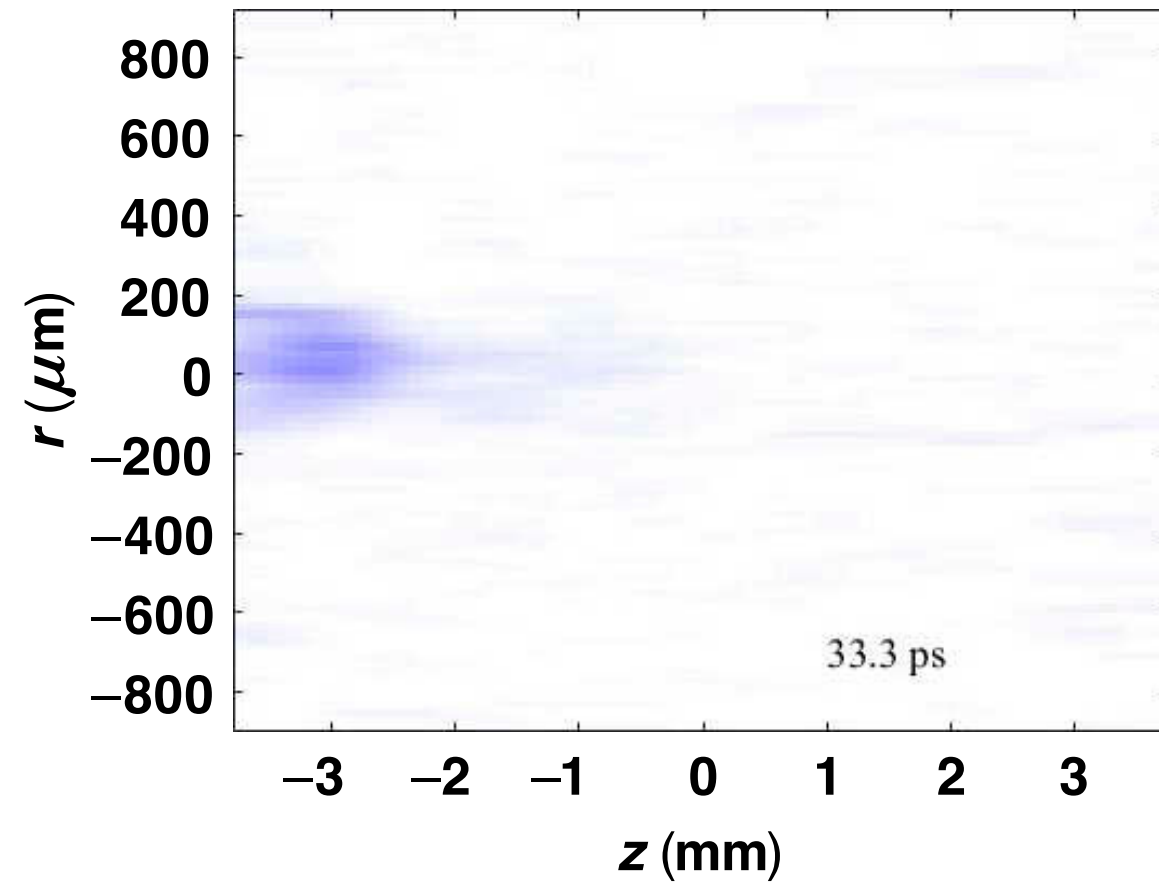


Negative chirp (55 ps)

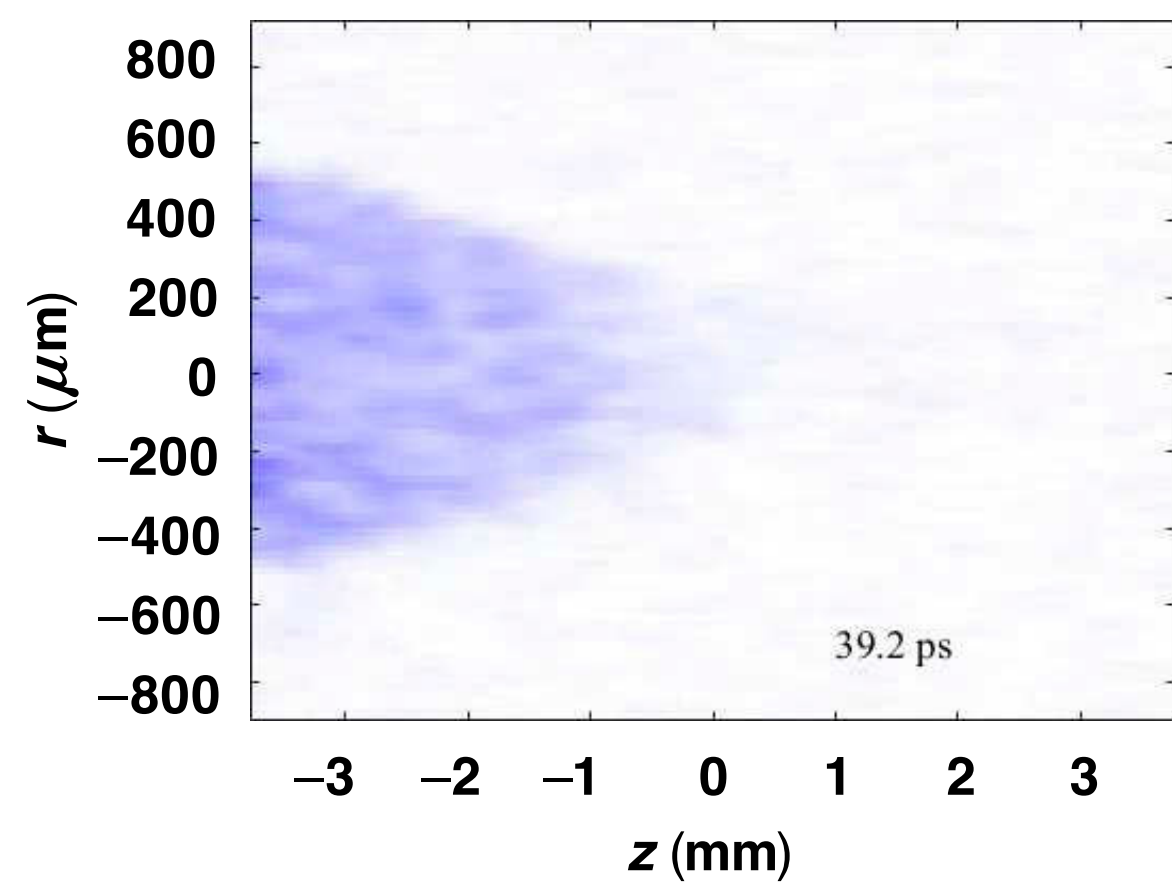


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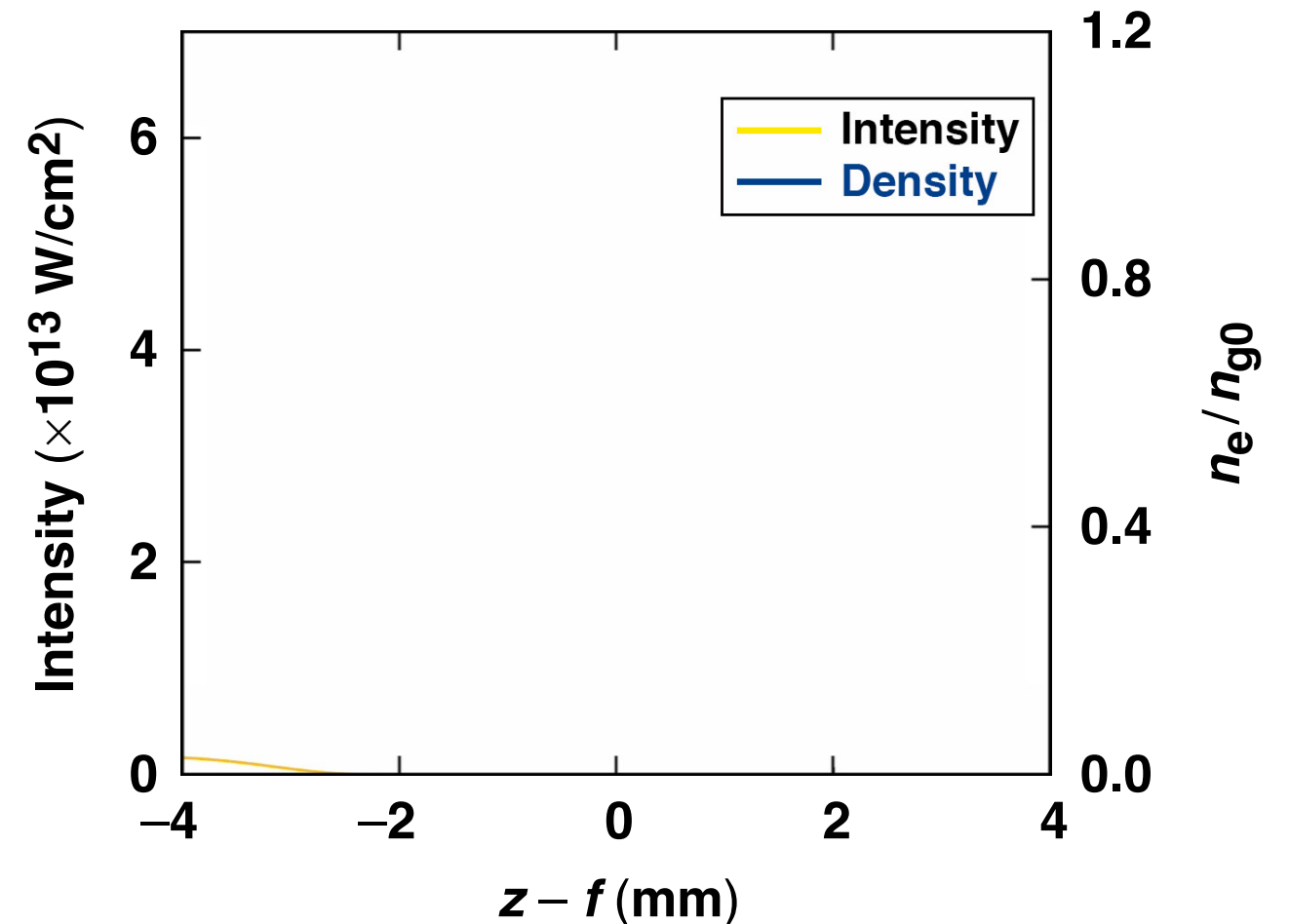
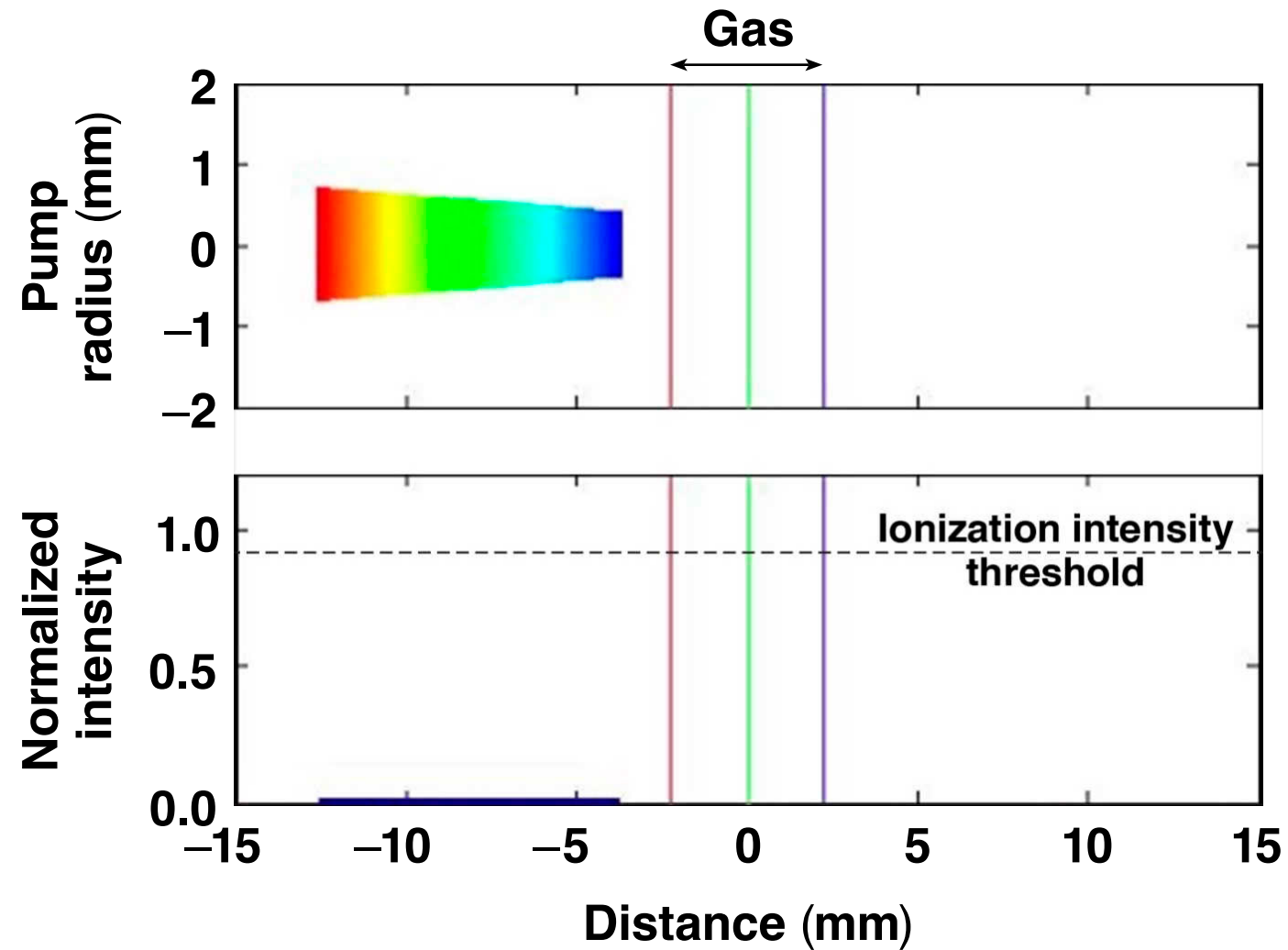
Negative chirp (55 ps)



Outline

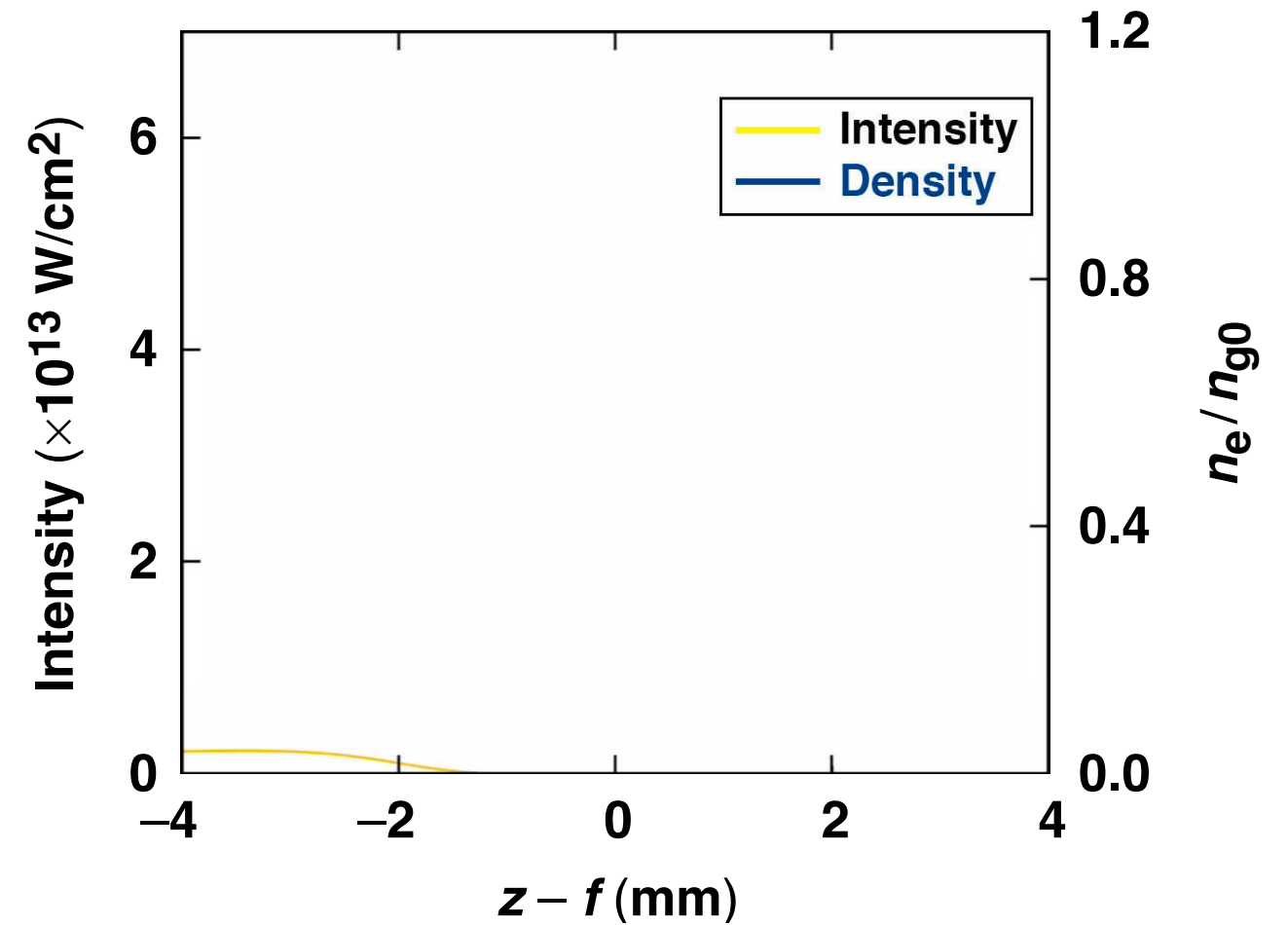
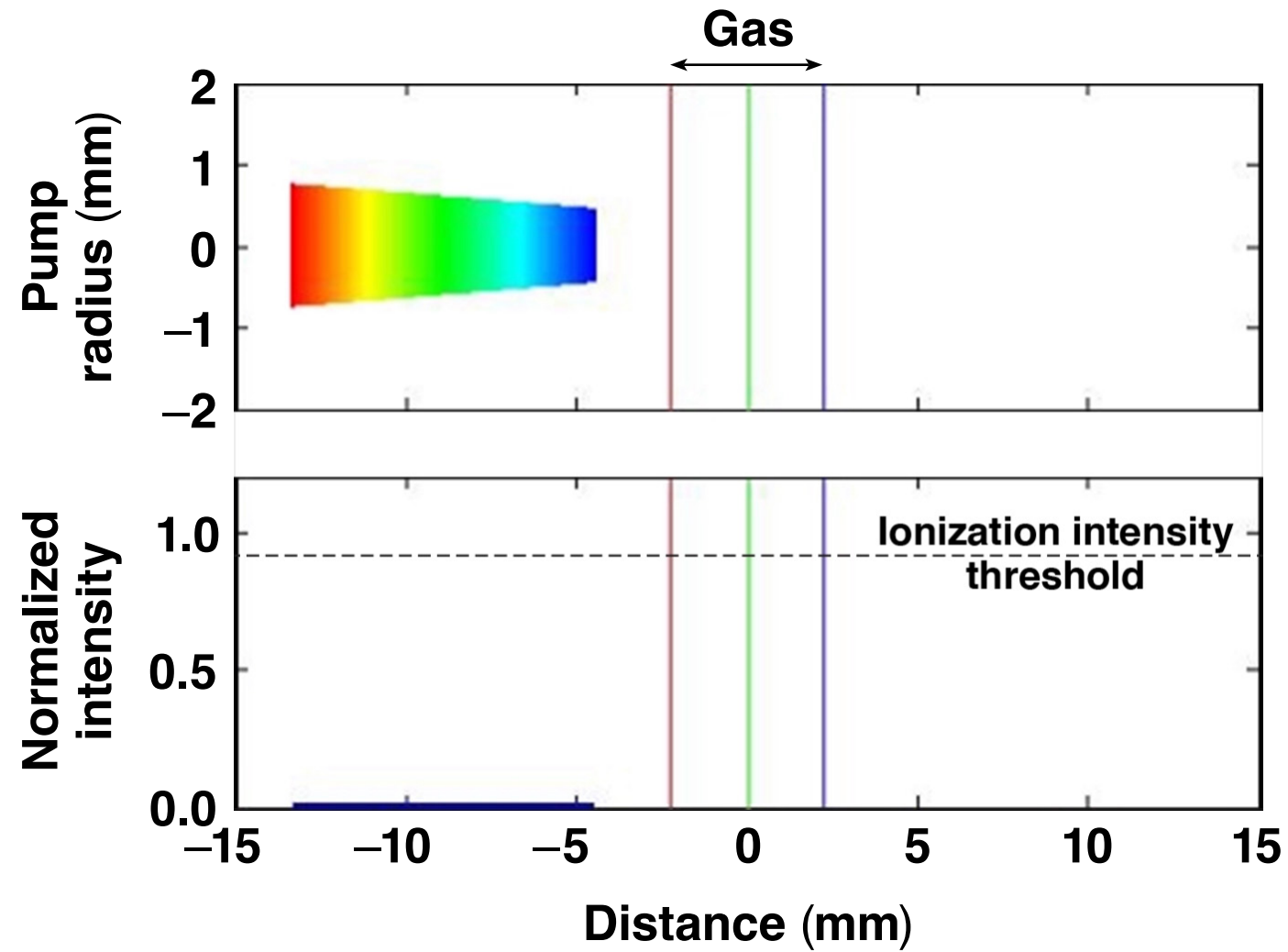
- Description of the flying focus
- Experimental demonstration of the flying focus
- **Ionization waves of arbitrary velocity (IWAV's)**
- Applications of the flying focus

The flying focus pulse can be used to generate an ionization wave of arbitrary velocity (IWAV)

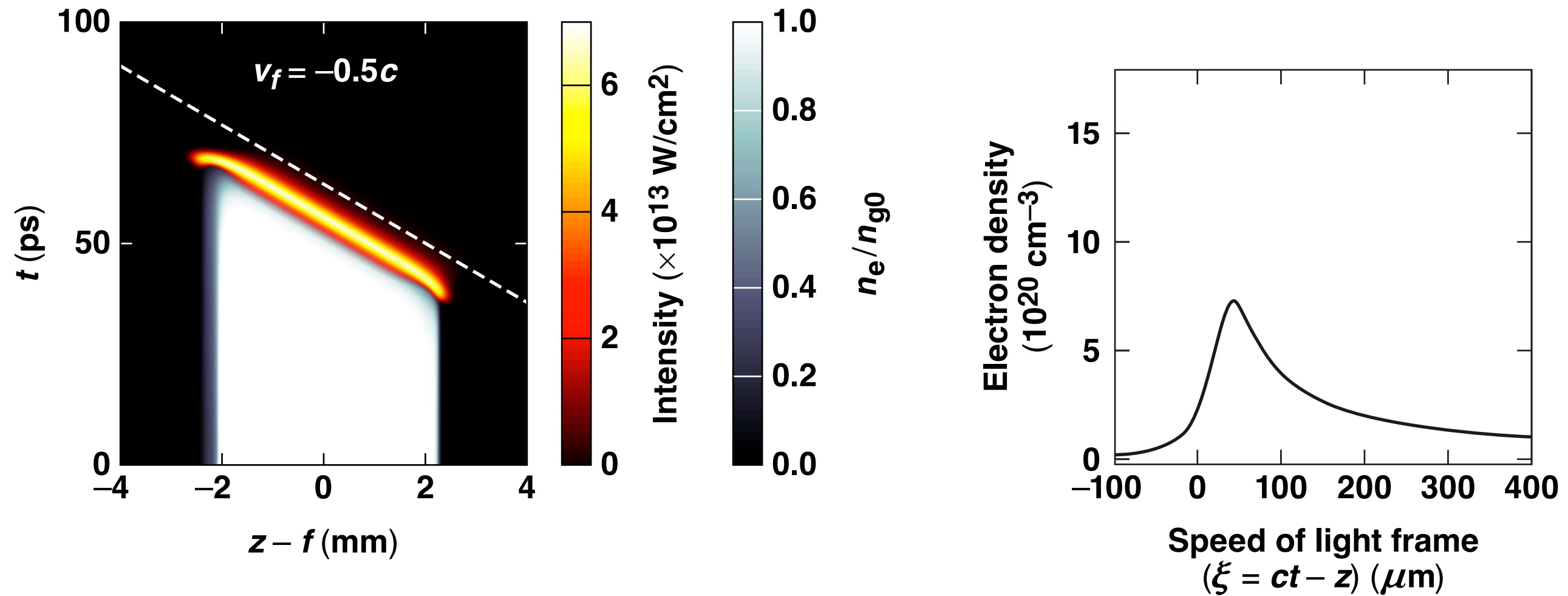


Counter-propagating flying focus mitigates ionization-induced refraction.

The flying focus pulse can be used to generate an IWAV

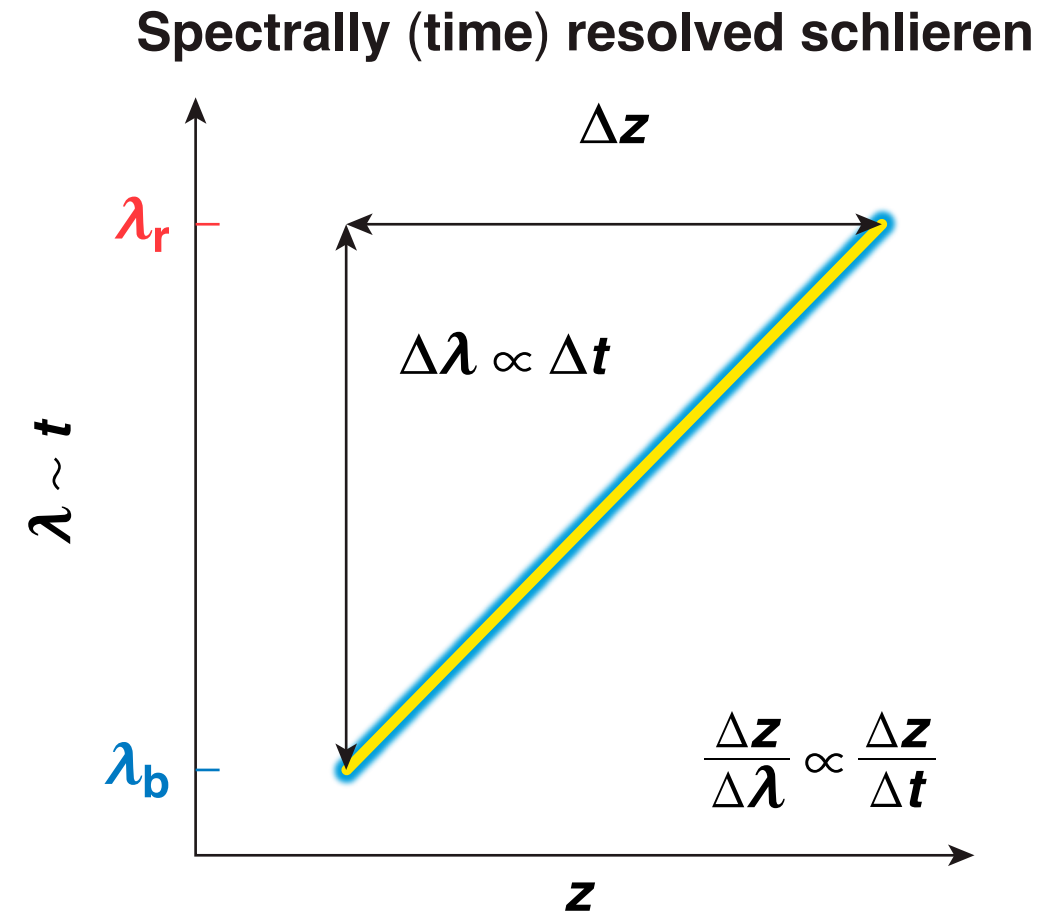
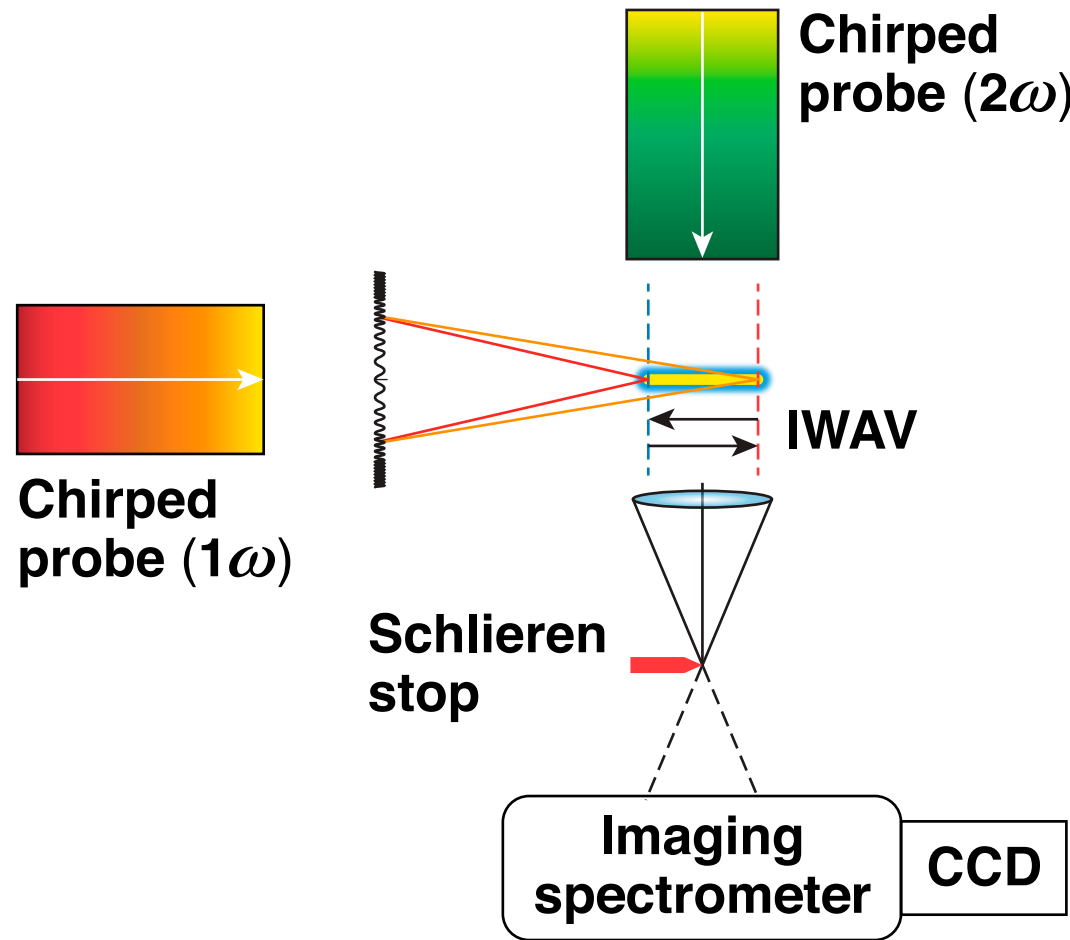


A counter-propagating flying focus mitigates plasma refraction and produces a sharp ionization front



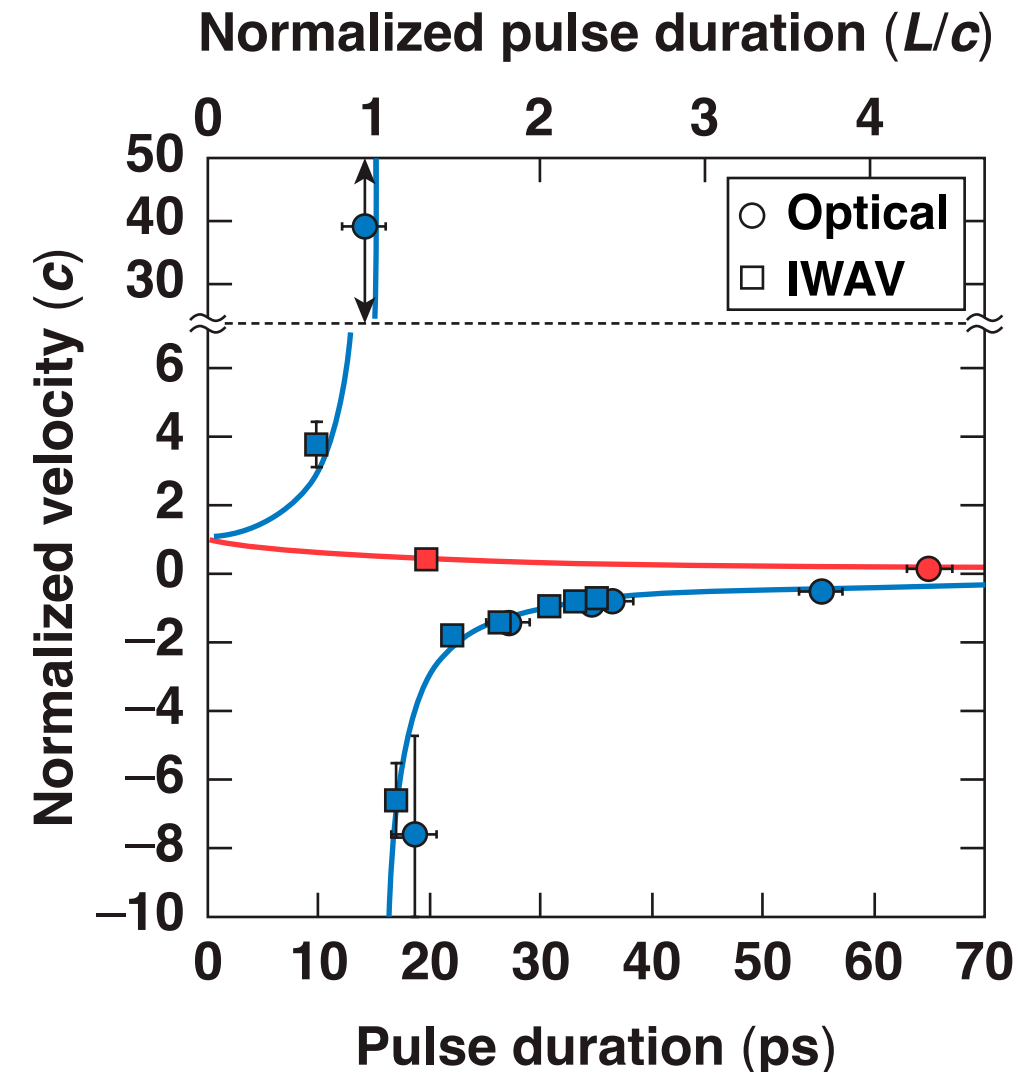
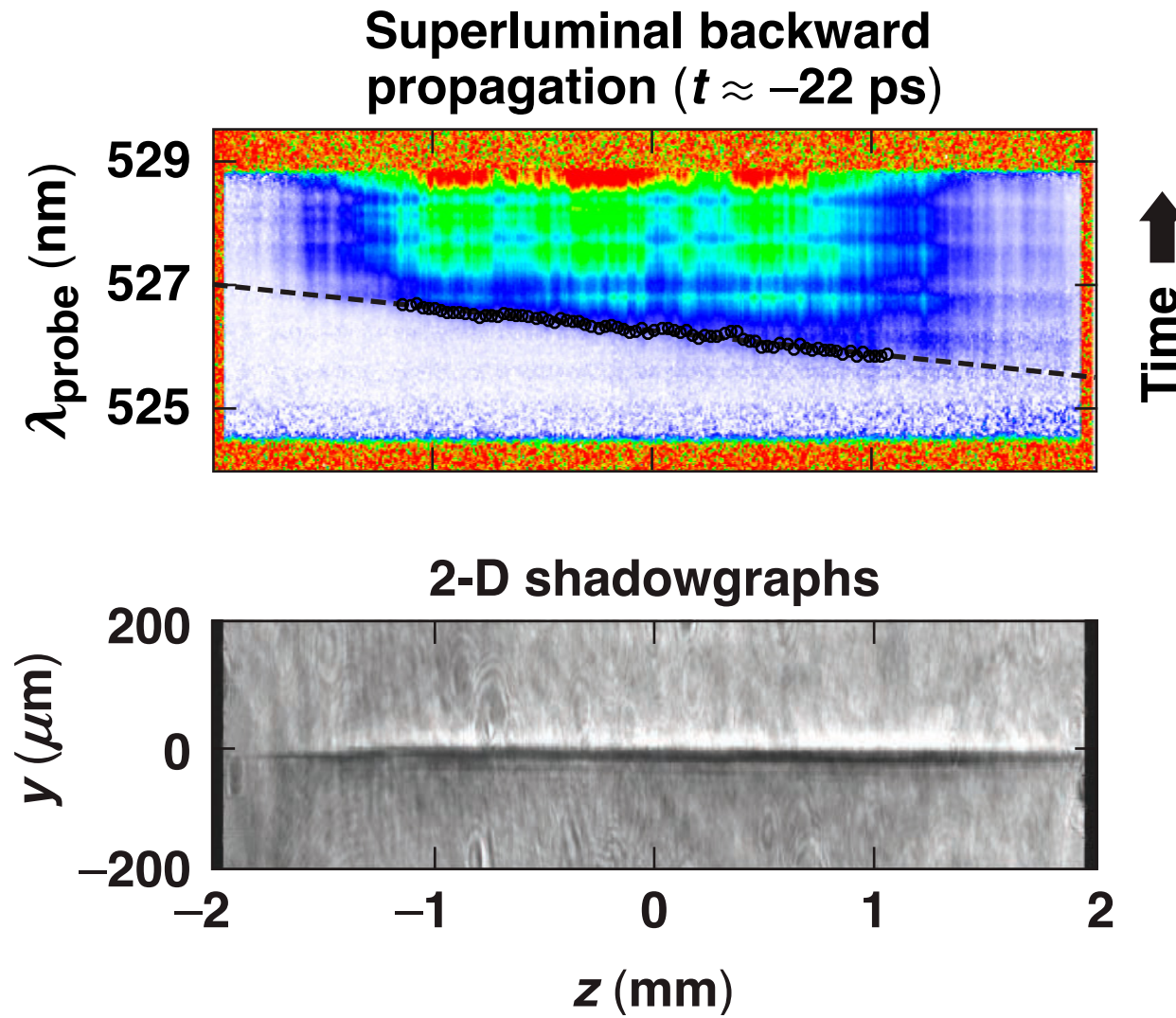


To study IWAV's in the laboratory, a spectrally resolved schlieren diagnostic was used



P. Franke, UP11.00093 this conference.
CCD: charge-coupled device

The analytic calculations and simulations are in excellent agreement with the measurements



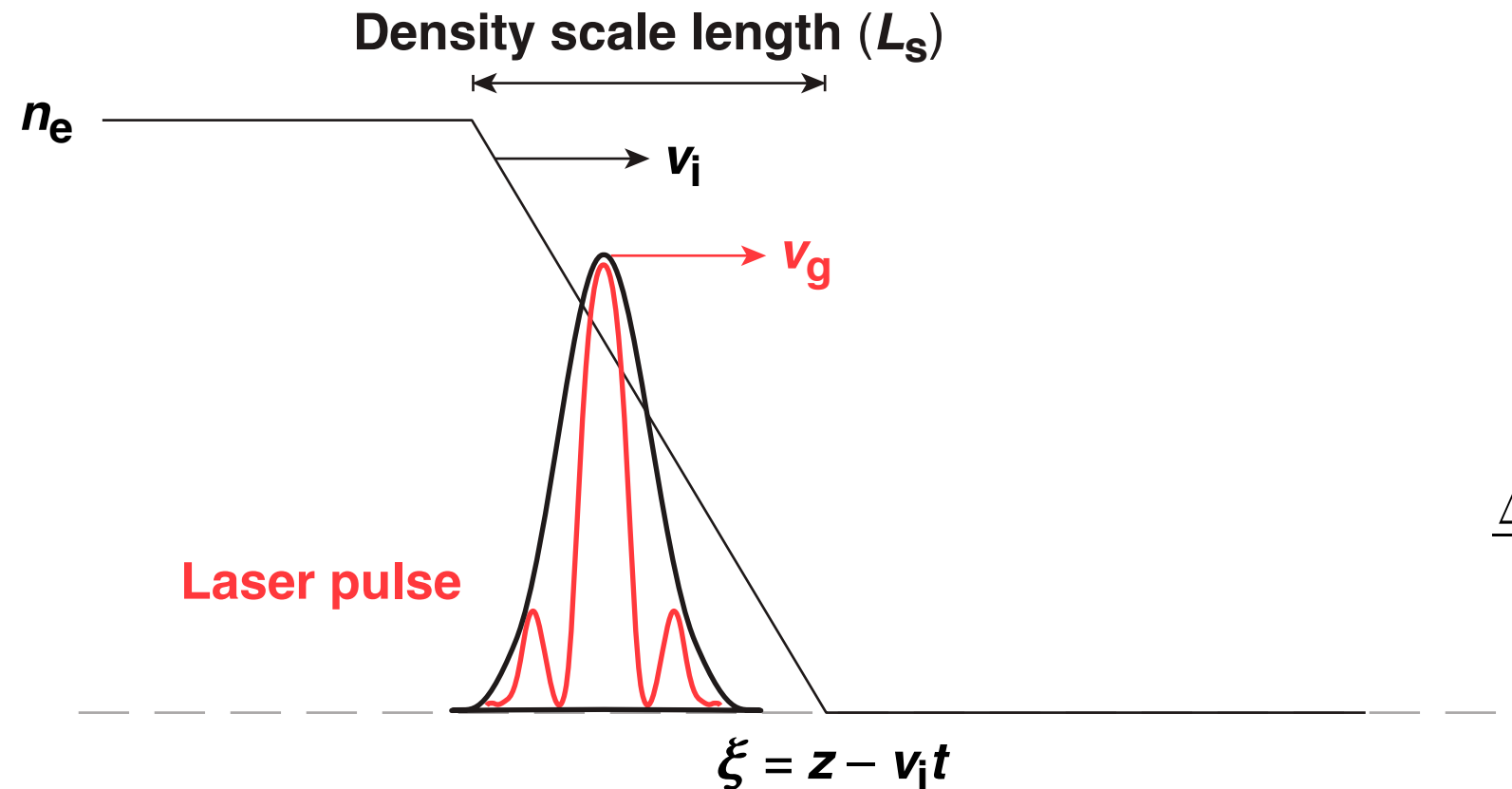
D. Turnbull *et al.*, Phys. Rev. Lett. **120**, 024801 (2018);
 D. Turnbull *et al.*, CO8.00012 this conference.

Outline

- Description of the flying focus
- Experimental demonstration of the flying focus
- Ionization waves of arbitrary velocity (IWAV's)
- **Applications of the flying focus**

The flying focus could be used extend the interaction length in a photon accelerator

A photon accelerator uses a time-varying refractive index to increase the group velocity of light



Index of refraction

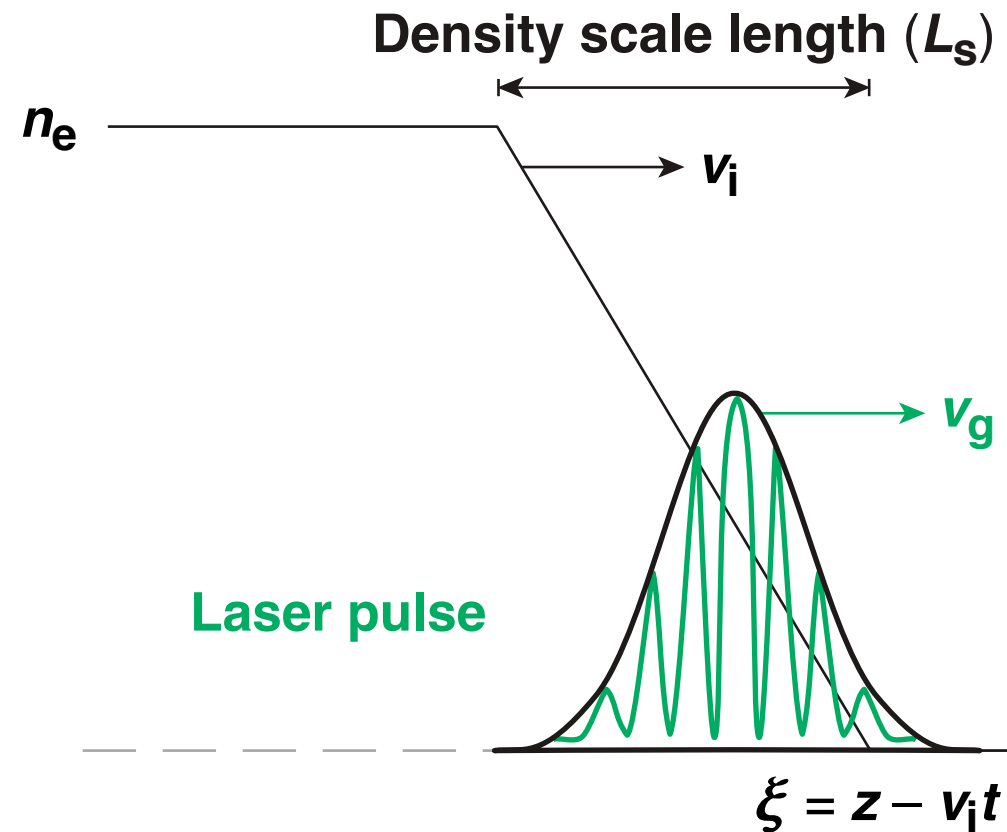
$$n = \sqrt{1 - \left(\frac{\omega_p}{\omega}\right)^2} \quad \omega_p \propto \sqrt{n_e}$$

Frequency shift

$$\frac{\Delta\omega}{\omega} = \sqrt{1 - \left(\frac{\omega_p}{\omega}\right)^2} \frac{z_{\text{effective}}}{L_s}$$

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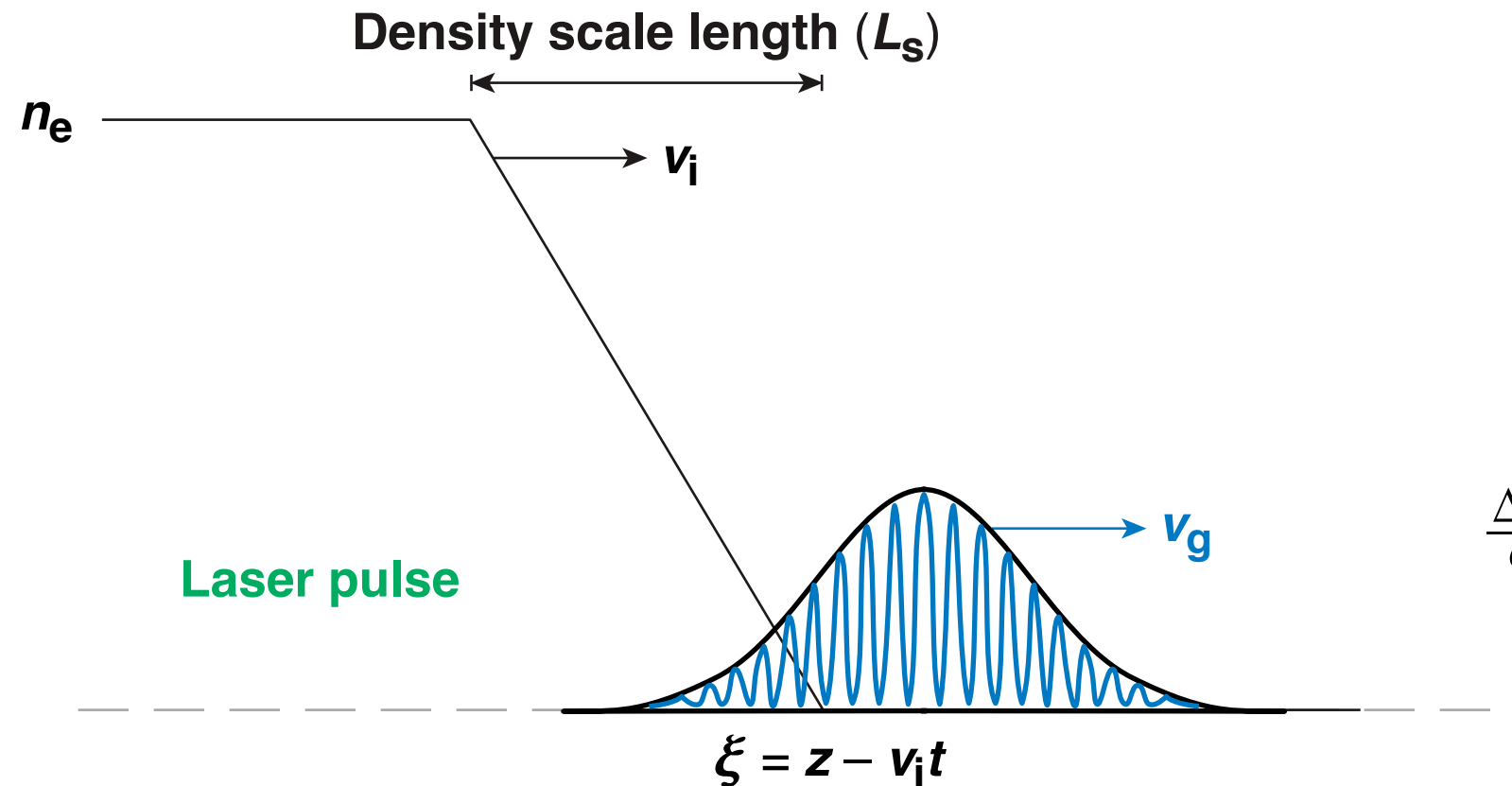
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Index of refraction

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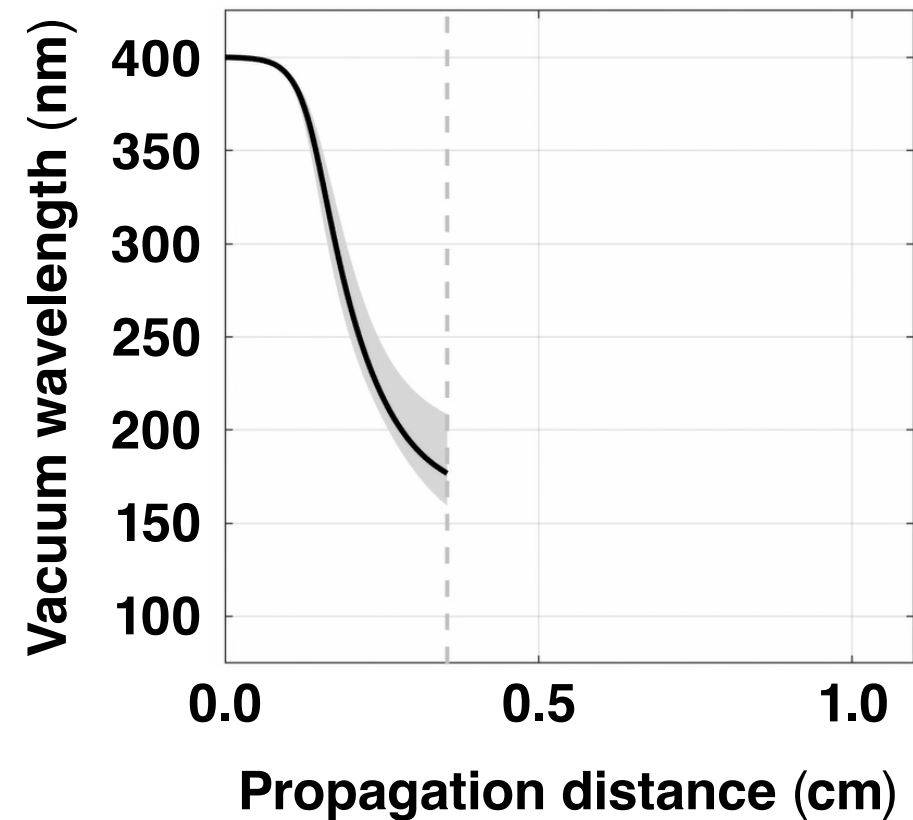
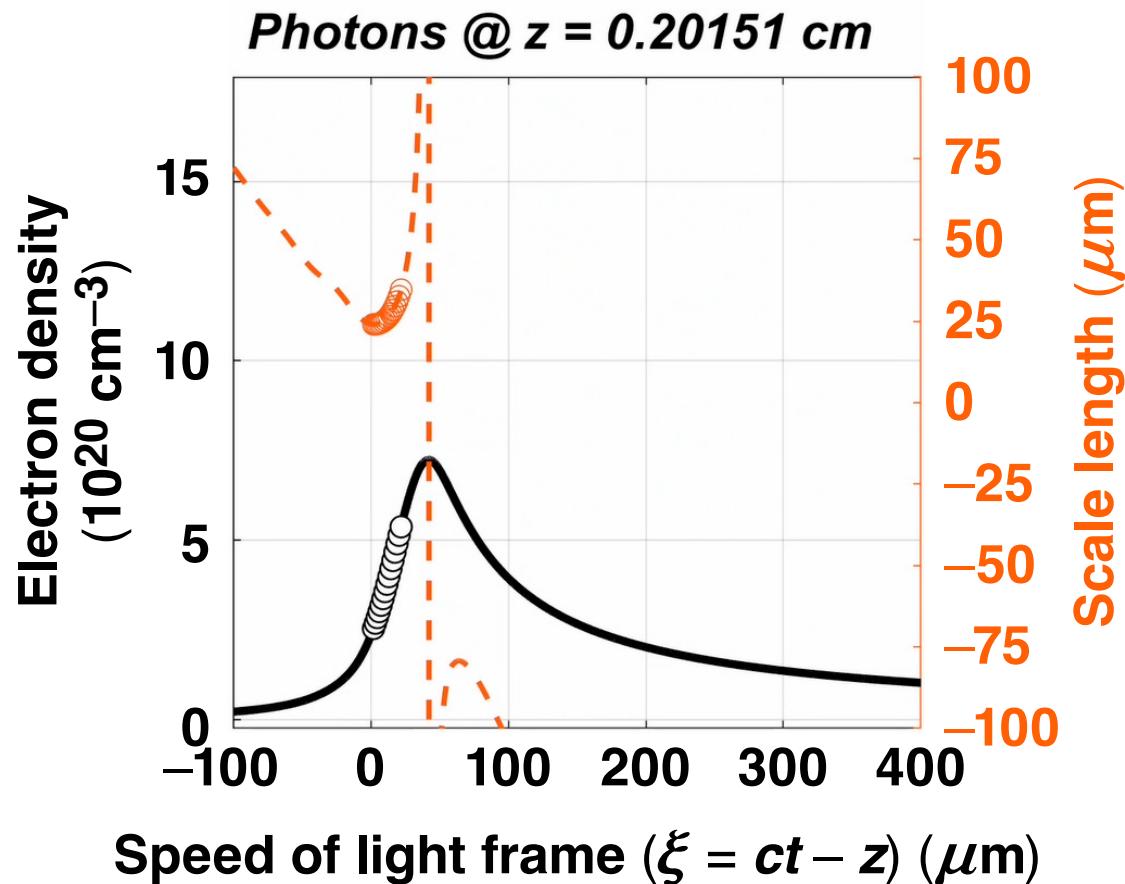
$$\frac{\Delta\omega}{\omega} = \sqrt{1 - \left(\frac{\omega_p}{\omega}\right)^2} \frac{z_{\text{effective}}}{L_s}$$

As the photons accelerate they eventually outrun the ionization front in a “conventional” photon accelerator.



PAS Undergraduate

By using the flying focus to generate ionization waves propagating at the speed of light, optical light can be frequency converted to the extreme ultraviolet

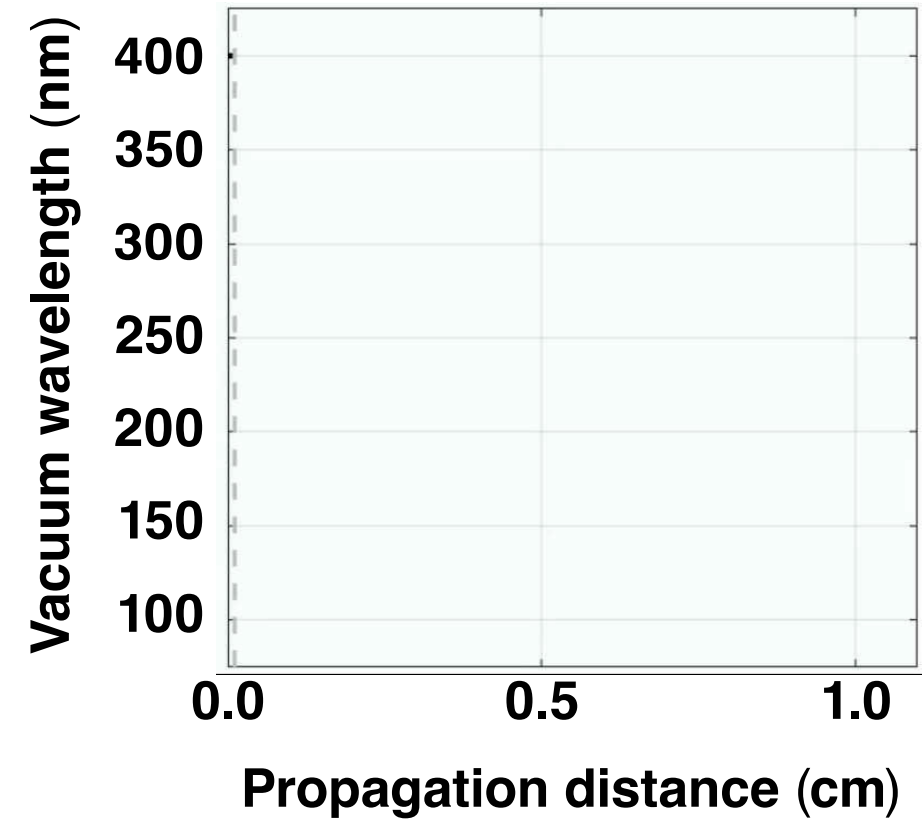
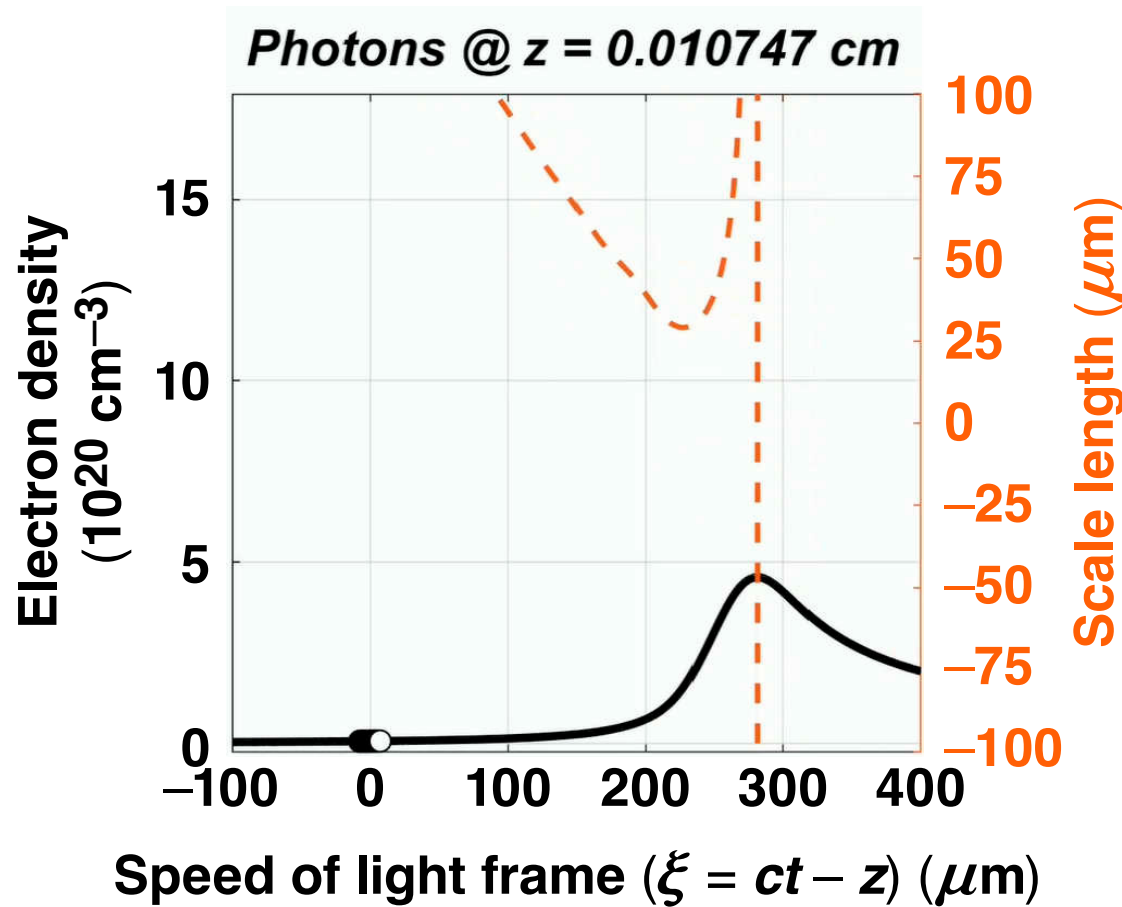


The photon accelerator driven by a flying focus could convert optical light to the extreme ultraviolet (100 nm).



PAS Undergraduate

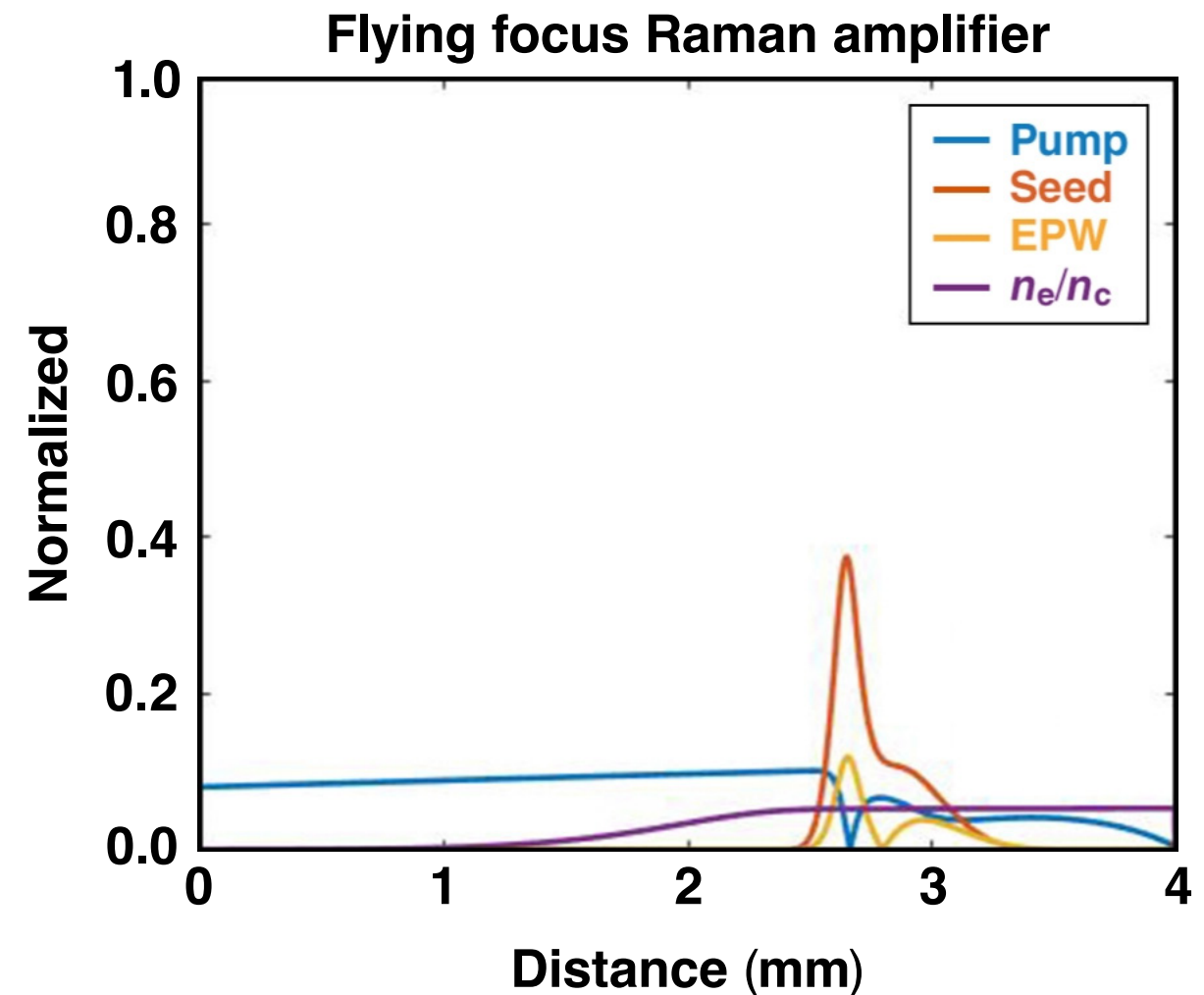
By using the flying focus to generate ionization waves propagating at the speed of light, optical light can be frequency converted to the extreme ultraviolet



The photon accelerator driven by a flying focus could convert optical light to the extreme ultraviolet (100 nm).

A counter-propagating ionization wave could overcome several challenges of Raman amplification

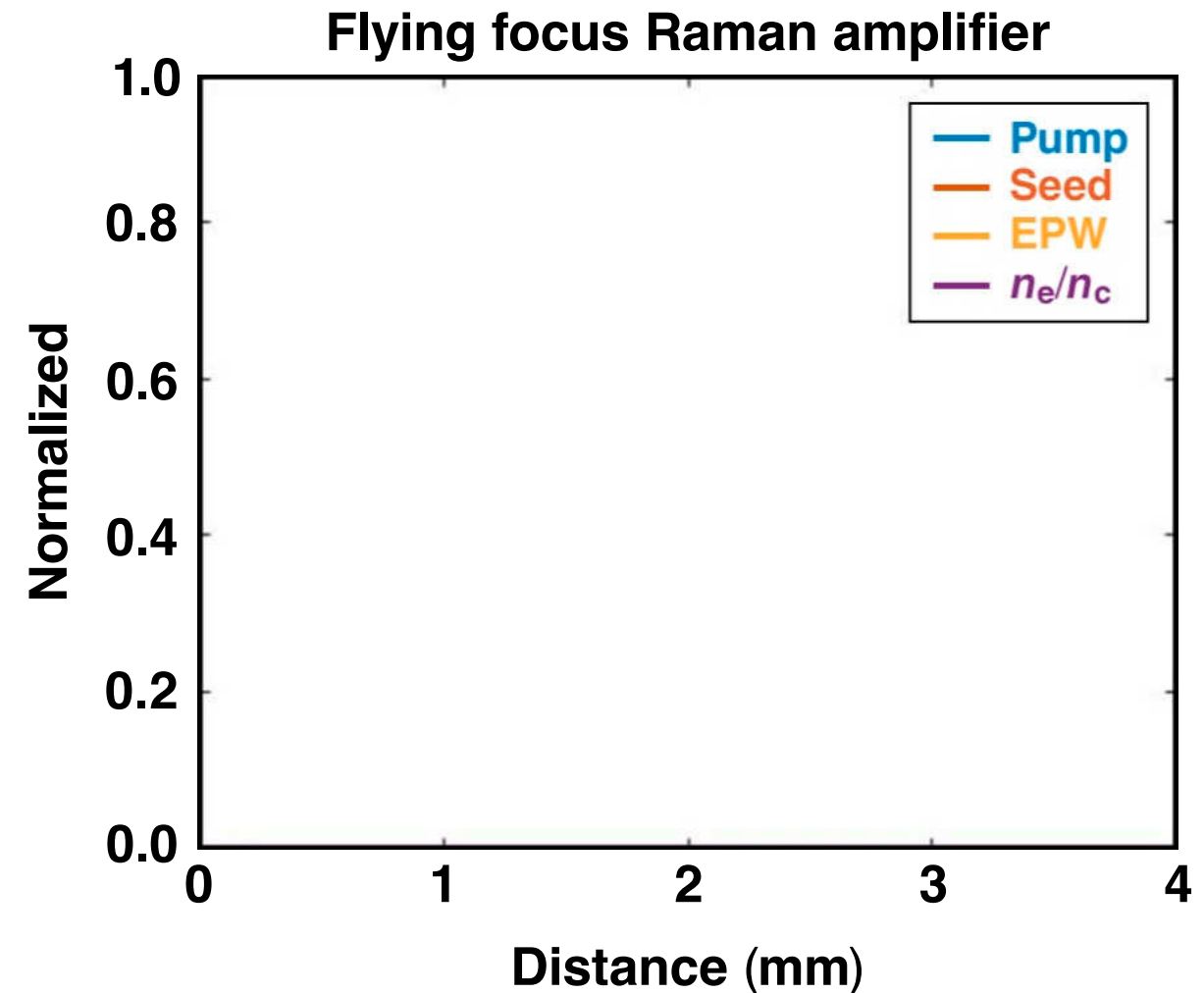
- **Constant longitudinal intensity:** the seed pulse experiences a constant pump intensity over the entire amplifier length
- **Counter-propagating ionization wave:** the pump will propagate through gas-eliminating parasitic instabilities
- **Plasma conditions:** the plasma conditions observed by the seed will be constant and controllable



D. Turnbull *et al.*, Phys. Rev. Lett. **120**, 024801 (2018).
EPW: electron plasma wave

A counter-propagating ionization wave could overcome several challenges of Raman amplification

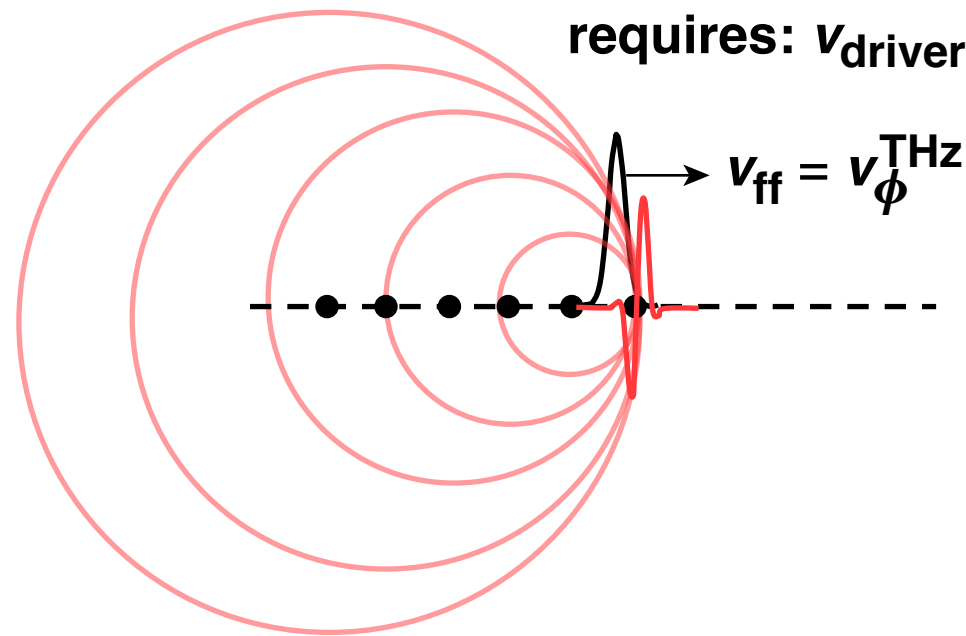
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The flying focus could enable new sources by decoupling the group velocity of light from the driver velocity

Cherenkov radiation
(THz source*)

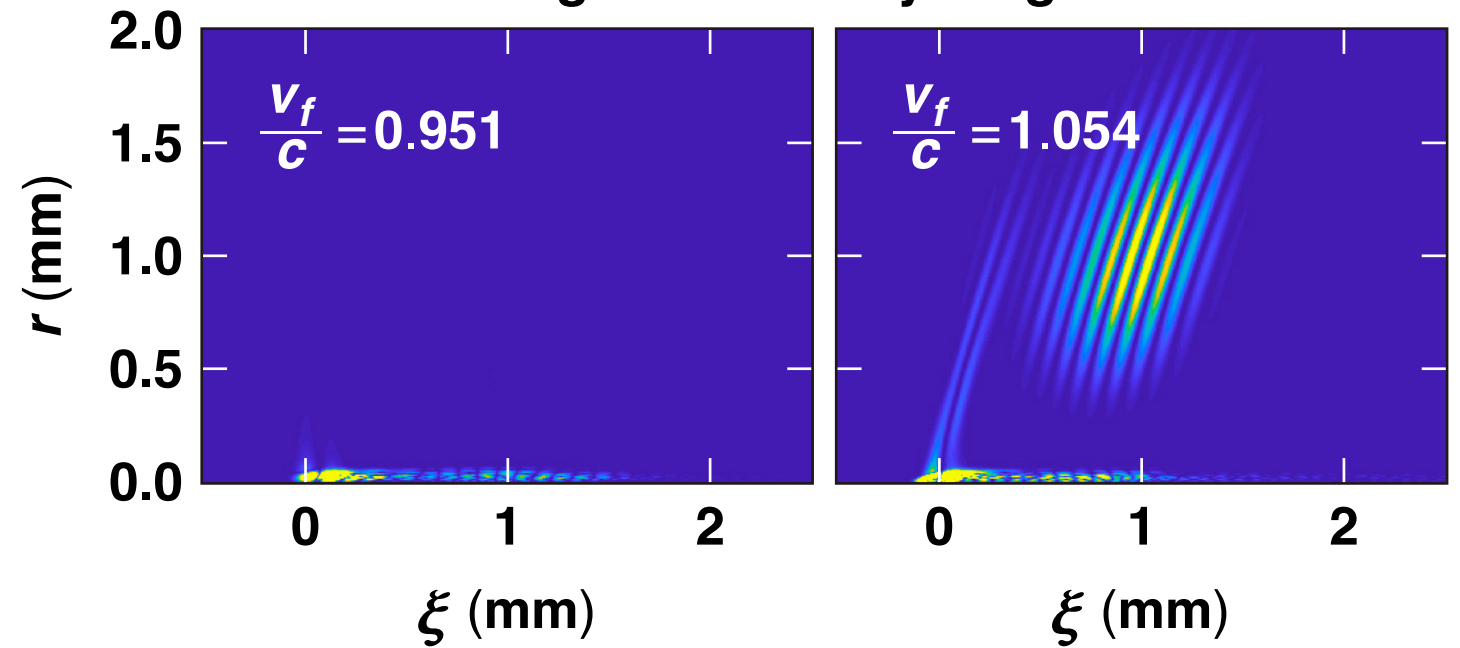
Cherenkov radiation
requires: $v_{\text{driver}} > v_{\phi}$



Cherenkov radiation is typically prohibited in a non-magnitized plasma ($v_{\text{gr}} < v_{\phi}$).

The flying focus can generate Cherenkov radiation in a plasma

Magnitude of Poynting flux

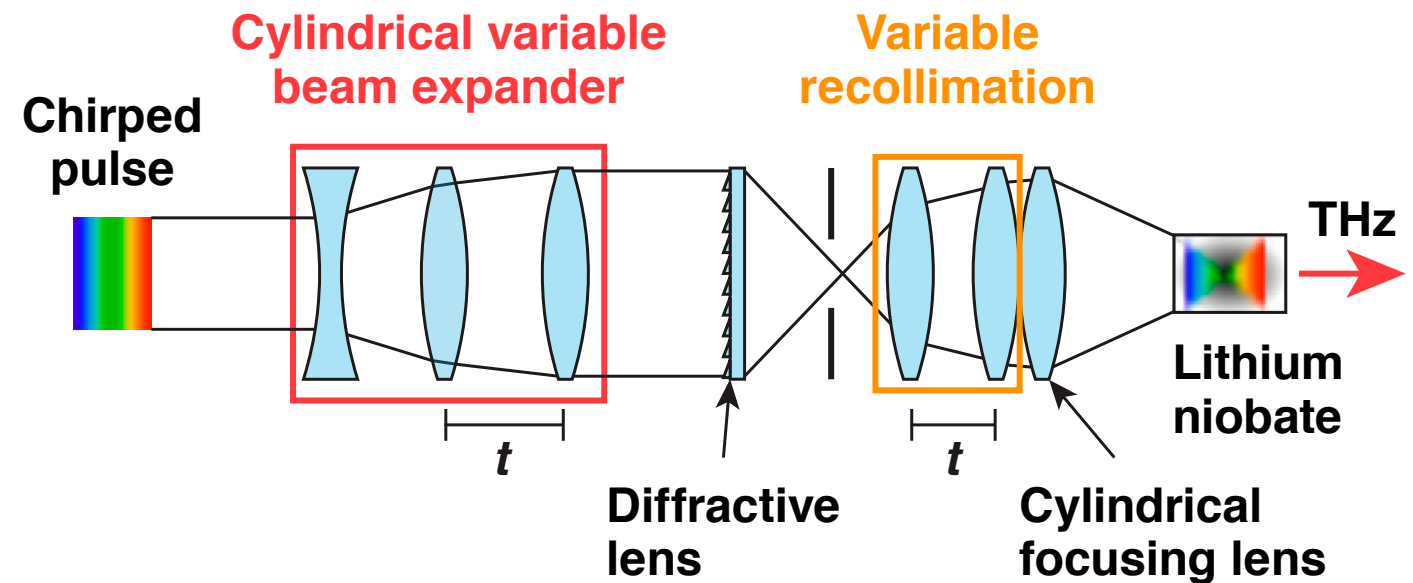
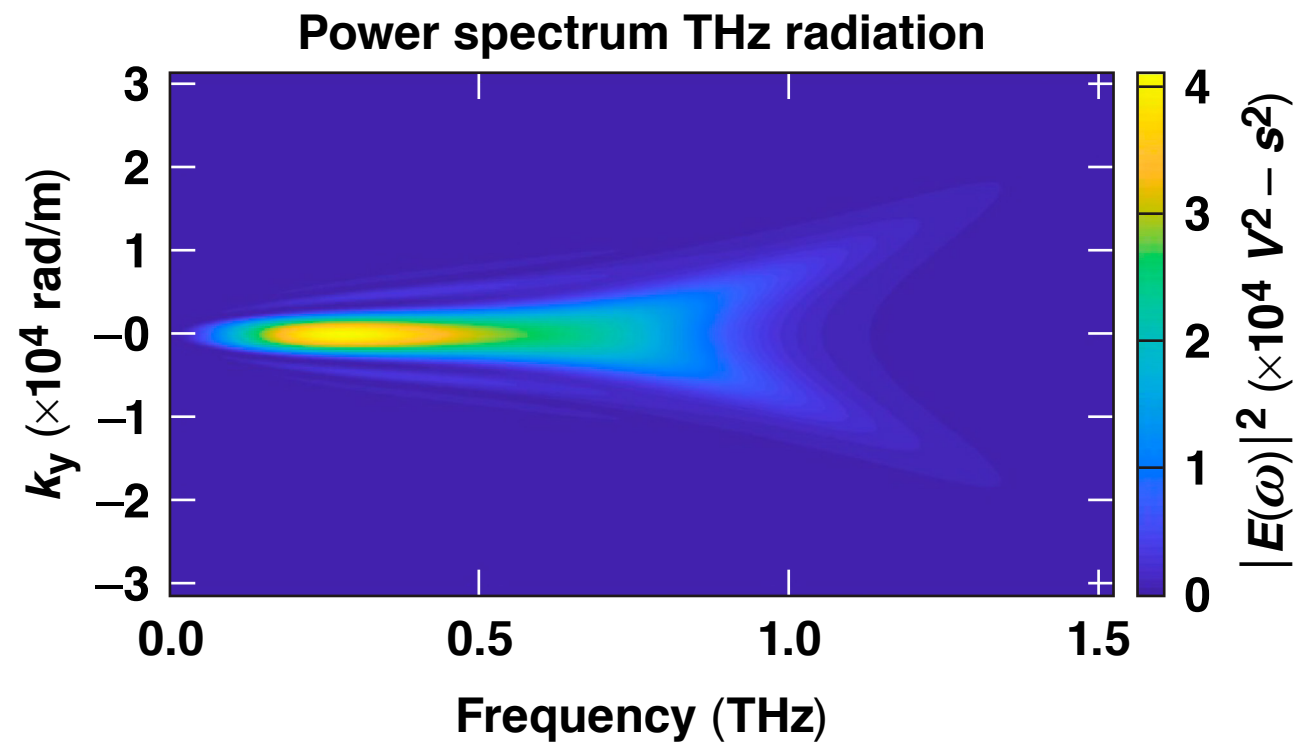


The flying focus enables Cherenkov radiation in a plasma by propagating the driver faster than the phase velocity of the radiation.

*J. Palastro *et al.*, in preparation Phys. Rev. Lett.;
J. Palastro *et al.*, JO8.00008, this conference.

PAS
Grad student

The flying focus can be used to phase match THz radiation in a crystal and extend the frequency-conversion process by orders of magnitude

Inst. Optics
Grad student

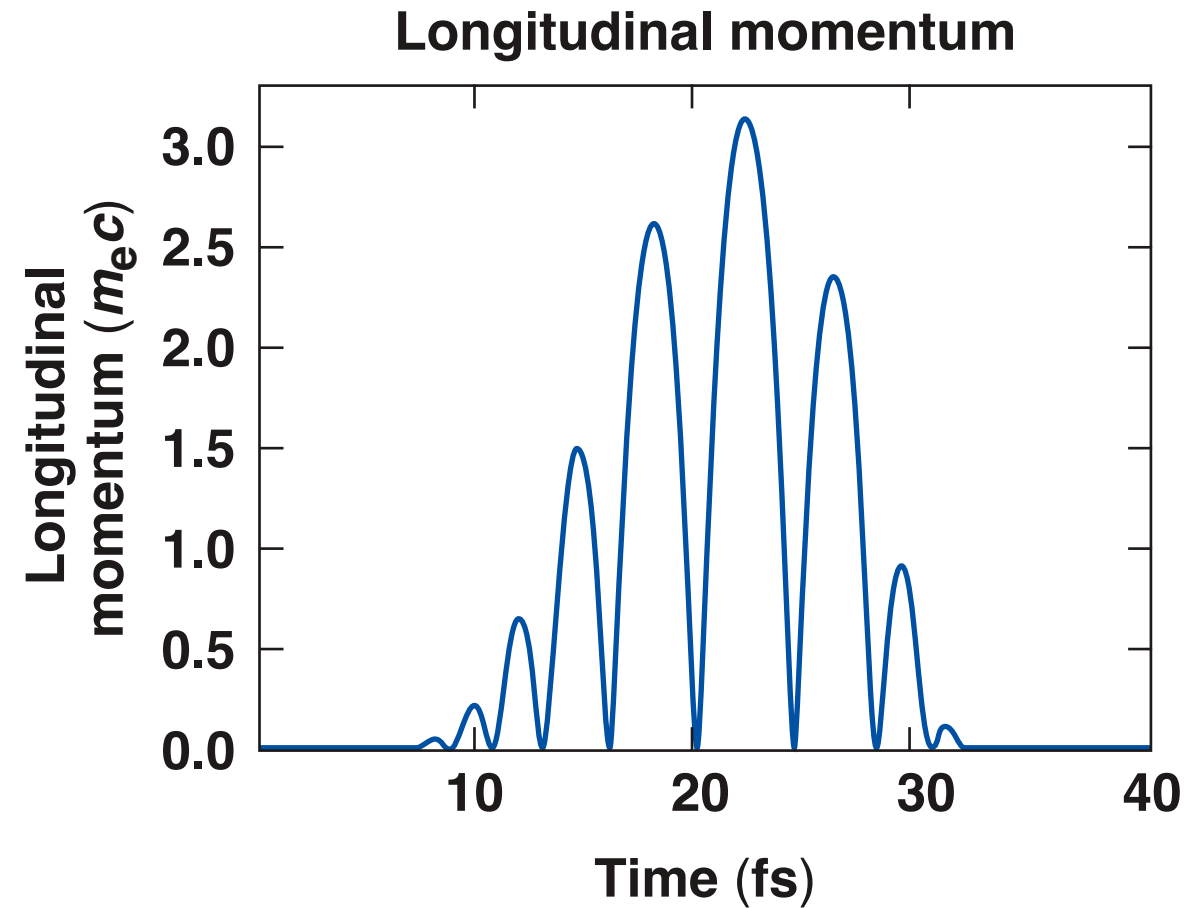
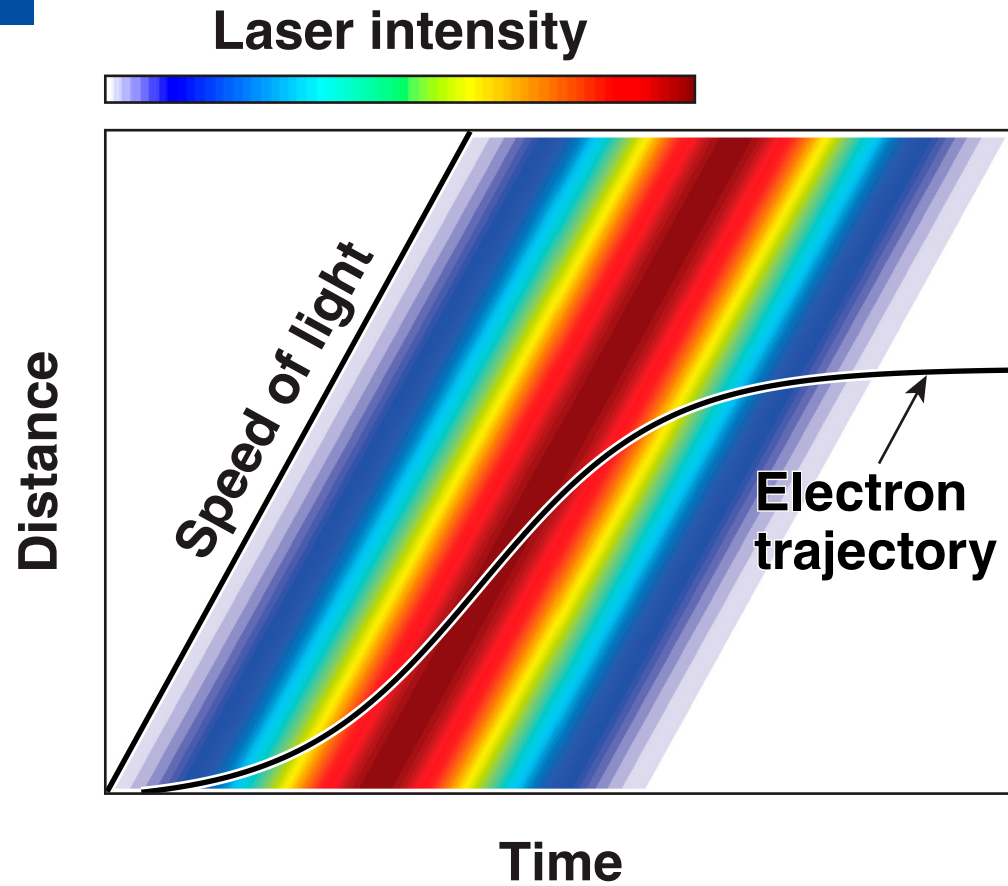
Simulations show on-axis phase matching in crystals over 1 cm,* otherwise constrained to off-axis phase matching.

A zoom lens system enables the flying focus velocity to be tuned for a given chirp.



PAS
Grad student

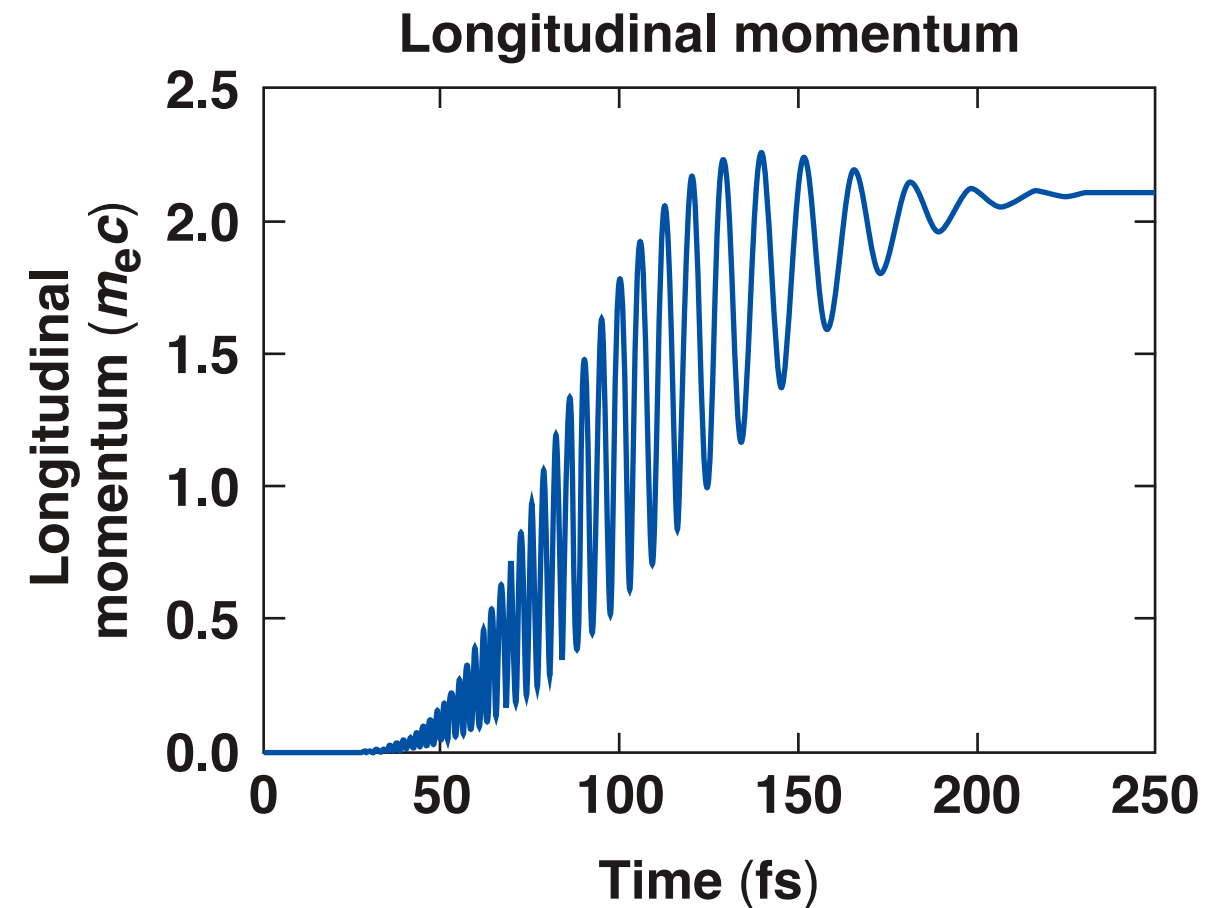
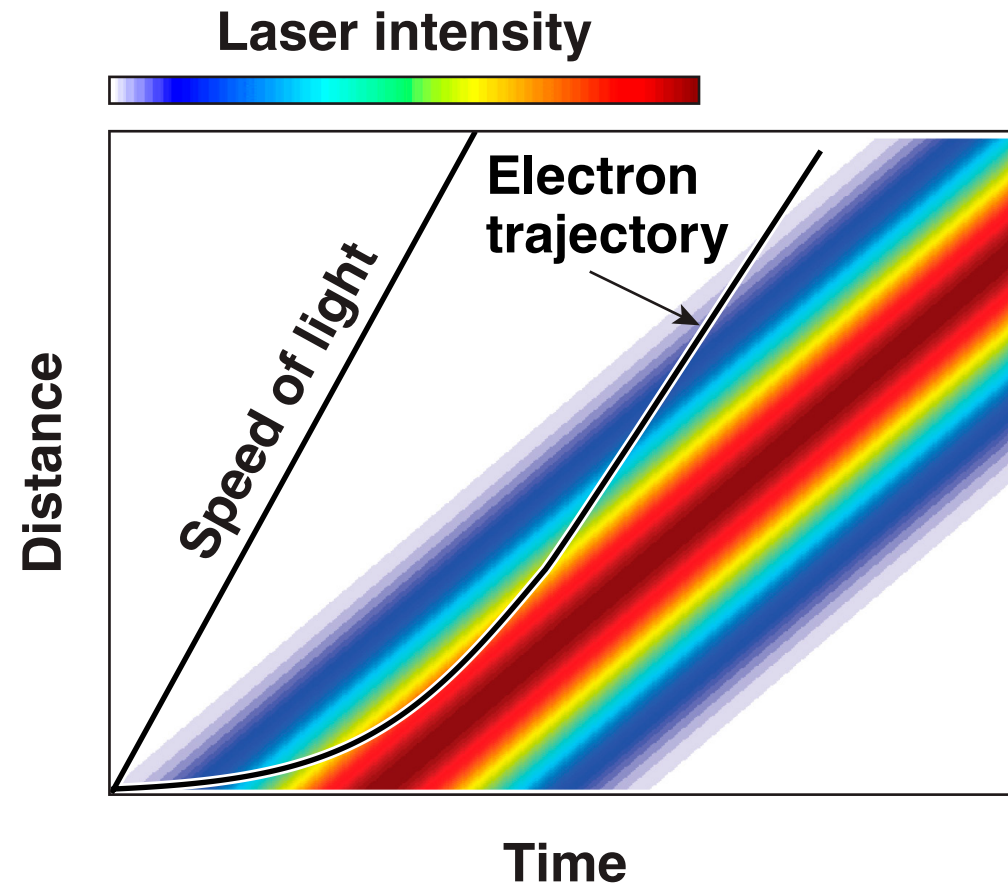
The flying focus enables a novel mechanism for direct vacuum electron acceleration*



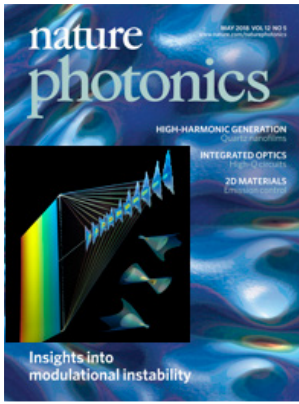
Lawson–Woodward Theorem: the net energy gain for an electron in a laser pulse is zero.

*D. Ramsey *et al.*, in preparation Phys. Rev. Lett.

In a flying focus pulse, electrons that overtake the laser pulse will exit with a longitudinal momentum



Flying focus overcomes the Lawson–Woodward Theorem: electrons can extract energy from the laser pulse.



Summary/Conclusions

A chirped laser pulse focused by a chromatic lens exhibits a dynamic, or “flying,” focus*



- The flying focus provides unprecedented spatiotemporal control over laser–plasma interactions by decoupling
 - the spot size of the pulse from the focal range
 - the velocity of the peak intensity from the group velocity
- Experiments have demonstrated the flying focus and the ability to generate ionization waves at any velocity (IWAV)
- Flying focus was applied to several applications
 - photon accelerator: IWAV’s can shift visible laser light to the XUV
 - Raman amplification: flying focus could overcome several challenges of laser-plasma amplifiers
 - Cherenkov radiation: flying focus allows new radiation sources
 - vacuum electron acceleration: flying focus enables vacuum acceleration

This decoupling has the potential to enable or improve several laser-based applications.

*D. H. Froula *et al.*, *Nat. Photonics* **12**, 262 (2018).
A. Sainte-Marie, O. Gobert, and F. Quéré, *Optica* **4**, 1298 (2017).

Thank you for your attention



3-D calculations (counter-propagating flying focus)

D. Turnbull **Ionization Waves of Arbitrary Velocity**
CO8.00012 4:12 pm, Monday, Nov. 5

J. Palastro **Cherenkov Radiation from a Plasma**
JO8.00008 3:24 pm, Tuesday, Nov. 6

A. Howard **Photon Acceleration in the**
Ionization Front of a Flying Focus
JO8.00015 4:48 pm, Tuesday, Nov. 6

P. Franke **Ionization Waves of Arbitrary Velocity**
UP11.00093 2:00 pm, Thursday, Nov. 8

