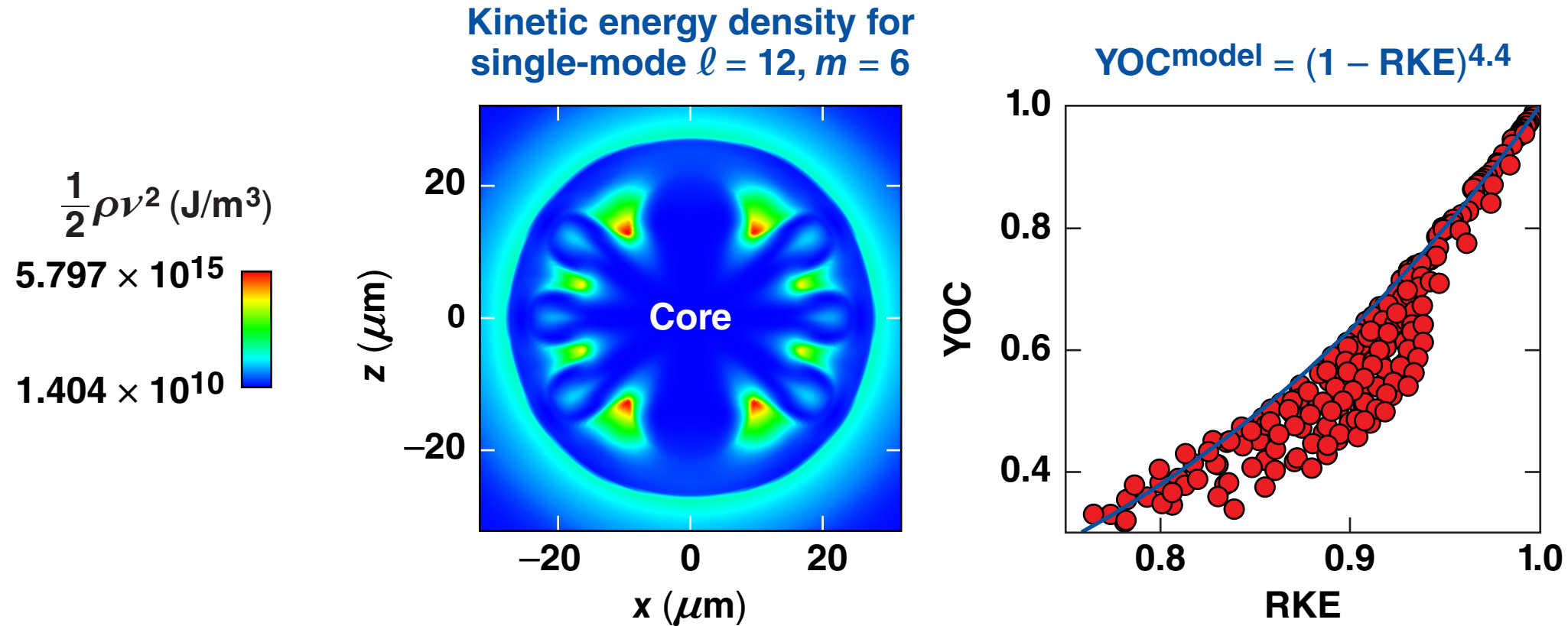


Three-Dimensional Studies of the Effect of Residual Kinetic Energy on Yield Degradation



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Summary

A hot-spot model indicates that the yield degradation caused by low- and mid-mode nonuniformities is a strong function of the residual kinetic energy



- A synthetic single-mode database ranging from low mode ($\ell = 1$) to mid mode ($\ell = 12$) was built using the 3-D hydrocode *DEC3D** applied to the deceleration phase of inertial confinement fusion (ICF) implosions
- It is shown that the yield-over-clean (YOC) is strongly correlated to residual kinetic energy (RKE) at bang time
- The simulation results are also confirmed by a simple analytical hot-spot model

*K. M. Woo *et al.*, “Three-Dimensional Studies of the Effect of Residual Kinetic Energy on Yield Degradation,” to be submitted to *Physics of Plasmas*.

Collaborators



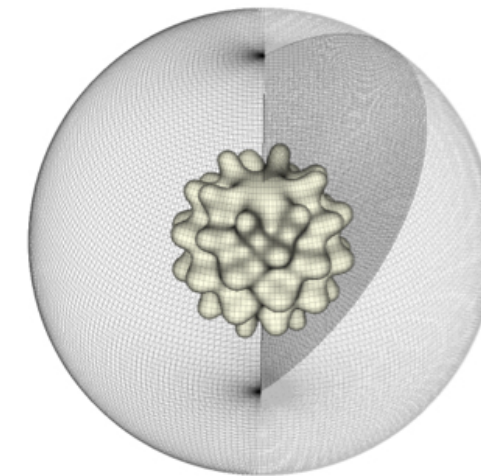
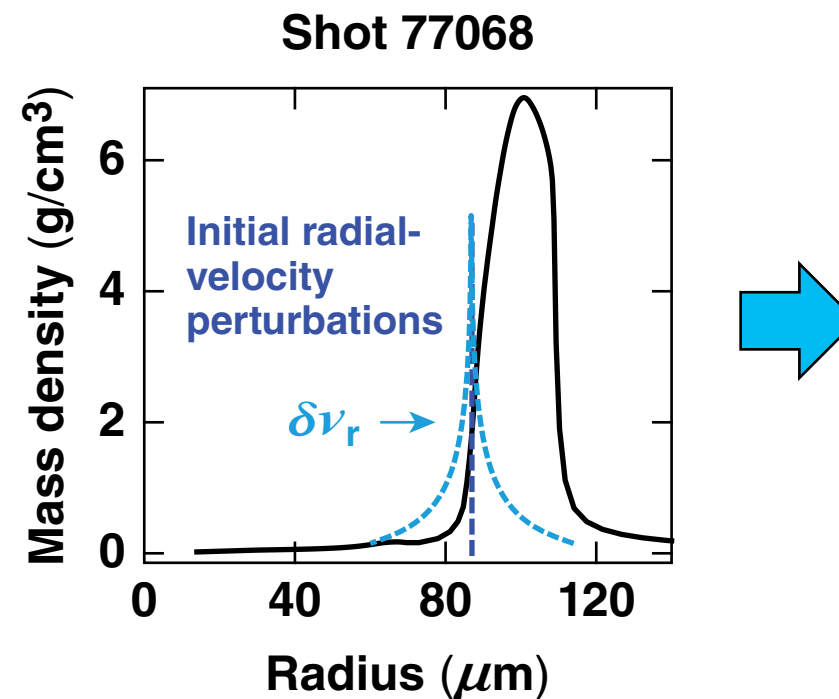
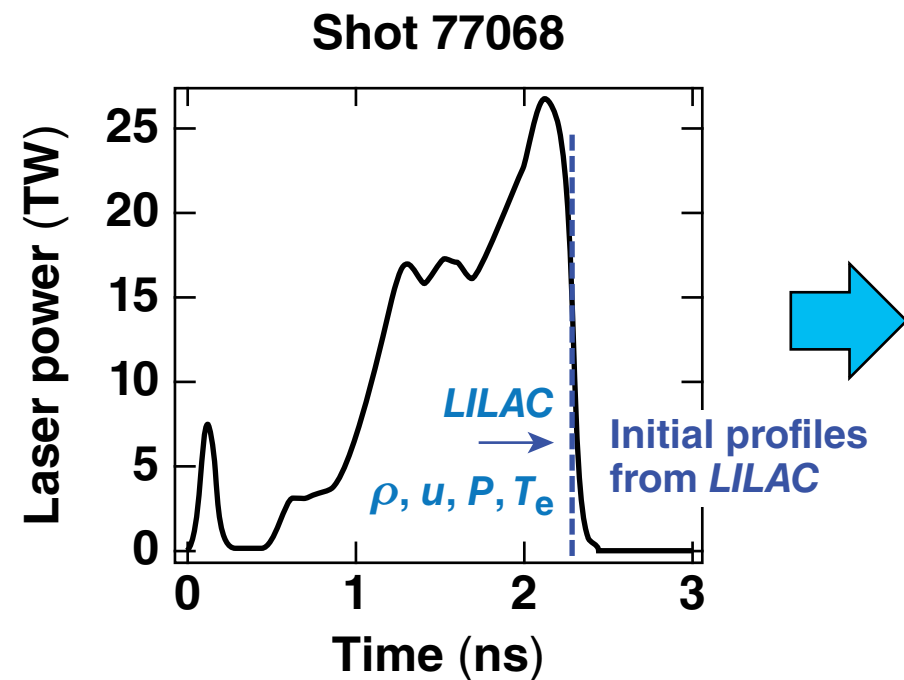
R. Betti, A Bose, D. Patel, and V. Gopaldaswamy

**University of Rochester
Laboratory for Laser Energetics**

A synthetic single-mode database was built using *DEC3D* to study yield degradation caused by Rayleigh–Taylor instabilities (RTI) in the deceleration phase

Simulation method

3-D radiation–hydrodynamic deceleration-phase code *DEC3D*



$\delta v_r = \Delta v/v_0$ is chosen appropriately to degrade the yield

1-keV T_e contour surface for single-mode $\ell = 12, m = 6$ at stagnation

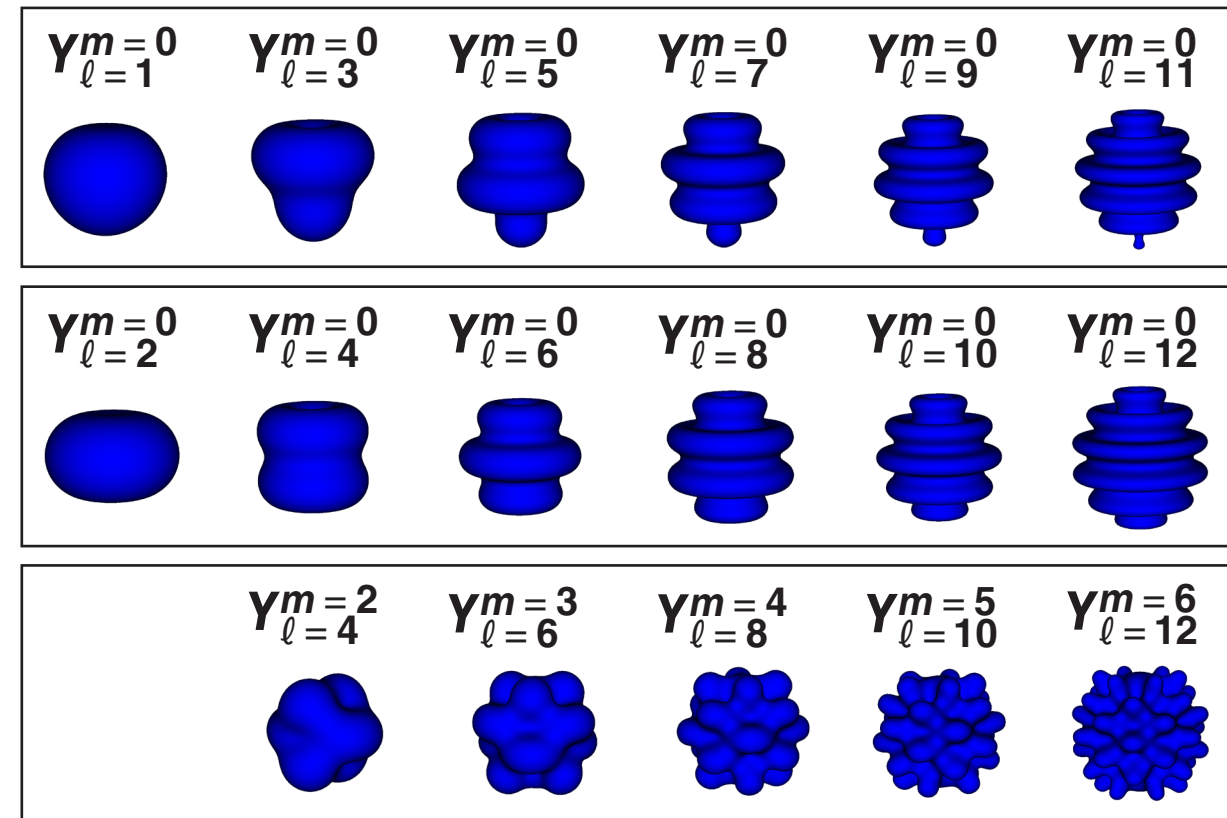
The synthetic database includes 3-D simulations using different velocity perturbations with spherical-harmonic single modes from $\ell = 1$ to $\ell = 12$

Initial velocity perturbation

A shape function to spread out the perturbation over space

$$v_r^{3-D}(r, \theta, \varphi, t_0) = v_r^{1-D}(r, t_0) + \underbrace{\frac{\Delta v}{v_0} f(r) Y_\ell^m(\theta, \varphi)}_{\delta v_r}$$

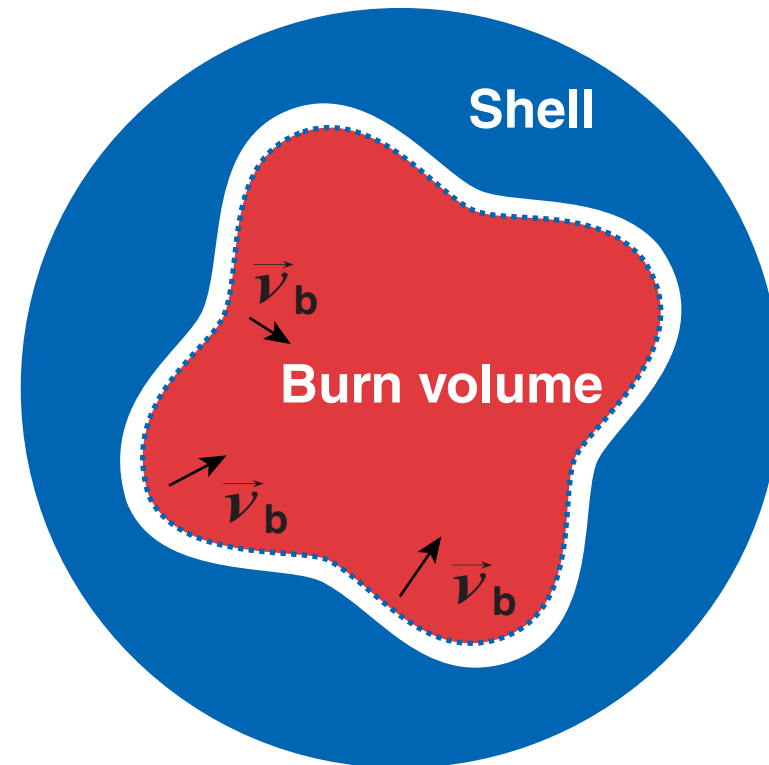
1-keV T_e contour surface at stagnation



- PPM* Riemann solver
- *HYPRE* thermal diffusion
- Resolution = $128 \times 128 \times 256$ ($r \times \theta \times \phi$ zones)

A simple 3-D hot-spot model is derived using energy conservation and an adiabatic condition, and neglecting the heat flux flowing into the cold bubbles

Three-dimensional hot-spot model



Neutron yield: $Y \simeq n^2 \langle \sigma v \rangle V \tau$

Yield-over-clean: $YOC \simeq \left[\frac{n_{3-D}}{n_{1-D}} \right]^2 \left[\frac{T_{3-D}}{T_{1-D}} \right]^4 \left[\frac{V_{3-D}}{V_{1-D}} \right] \left[\frac{\tau_{3-D}}{\tau_{1-D}} \right]$

Adiabatic implosion:

$$P_{3-D} V_{3-D}^{5/3} = P_{1-D} V_{1-D}^{5/3}$$

$$\longrightarrow \frac{P_{3-D}}{P_{1-D}} = \left(\frac{IE_{3-D}}{IE_{1-D}} \right)^{5/2} \quad \text{and} \quad \frac{V_{3-D}}{V_{1-D}} = \left(\frac{IE_{3-D}}{IE_{1-D}} \right)^{-3/2}$$

Energy conservation: $IE_{HS}^{stag} = KE_{tot}^{max} - KE_{tot}^{stag} - IE_{SH}^{stag}$

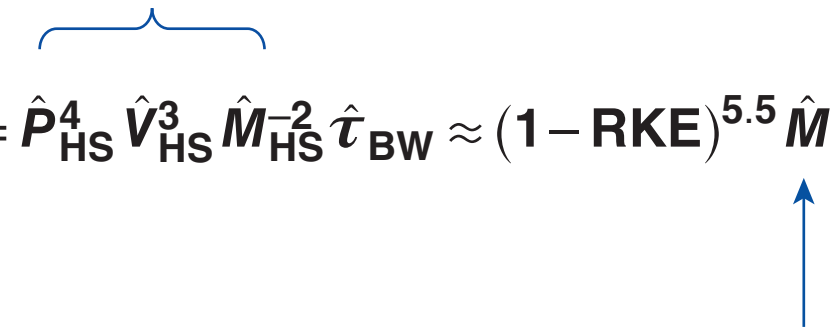
Energy conservation and adiabatic condition are used to derive the YOC dependence in the residual kinetic energy

Definitions

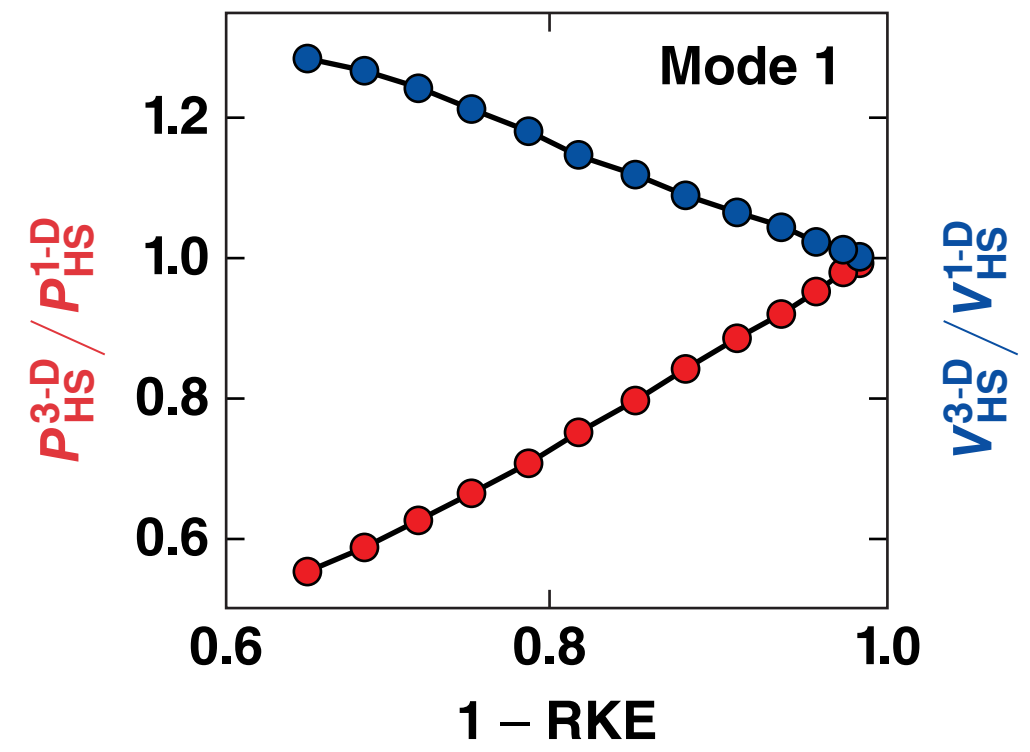
$$\hat{Q} = Q_{3-D} / Q_{1-D} \quad \text{and} \quad \text{RKE} = \frac{\text{KE}_{\text{shell}}^{3-D}(t_{\text{stag}}^{3-D}) - \text{KE}_{\text{shell}}^{1-D}(t_{\text{stag}}^{1-D})}{\text{KE}_{\text{tot}}^{\text{max}}} \rightarrow \frac{\text{IE}_{\text{HS}}^{3-D}}{\text{IE}_{\text{HS}}^{1-D}} \simeq 1 - \text{RKE}$$

$\hat{P} = (1 - \text{RKE})^{5/2}, \hat{V} = (1 - \text{RKE})^{-3/2}$

$$\text{YOC}^{\text{model}} = \hat{P}_{\text{HS}}^4 \hat{V}_{\text{HS}}^3 \hat{M}_{\text{HS}}^{-2} \hat{\tau}_{\text{BW}} \approx (1 - \text{RKE})^{5.5} \hat{M}_{\text{HS}}^{-2} \hat{\tau}_{\text{BW}}$$



Using 1-D scaling for the mass ablation rate (next slide)



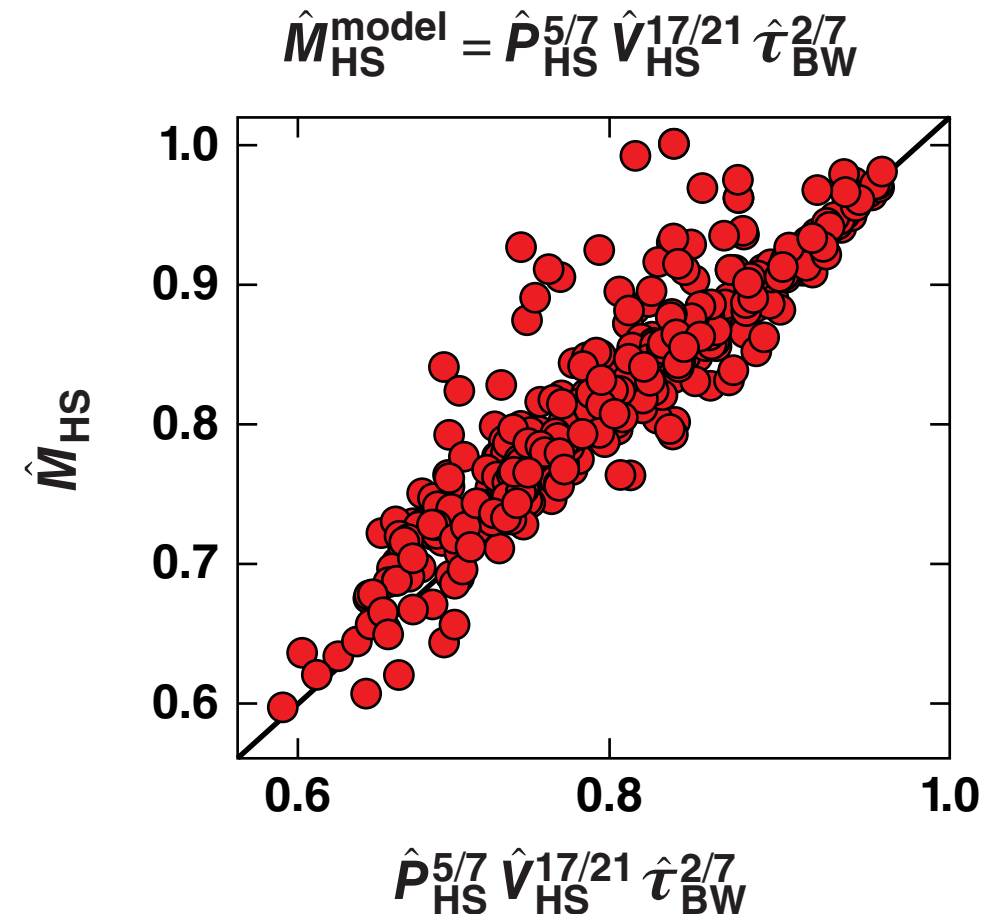
For low modes, a 1-D scaling of the mass ablation rate is used to derive the 3-D hot-spot mass

One-dimensional approximations for mass ablation rate* and hot-spot surface area

$$\frac{M_{\text{HS}}}{\tau} \simeq \underbrace{\dot{m}_{\text{abl}}}_{T_{\text{HS}}^{5/2}/R} \underbrace{S_{\text{HS}}}_{R^2} \sim T_{\text{HS}}^{5/2} R \sim (P_{\text{HS}} V_{\text{HS}} / M_{\text{HS}})^{5/2} V_{\text{HS}}^{1/3}$$

Scaling for 3-D hot-spot mass

$$\hat{M}_{\text{HS}} = \hat{P}_{\text{HS}}^{5/7} \hat{V}_{\text{HS}}^{17/21} \hat{\tau}_{\text{BW}}^{2/7} = (1 - \text{RKE})^{4/7} \hat{\tau}_{\text{BW}}^{2/7}$$



The YOC is a strong function of the residual kinetic energy

Yield degradation driven by pure hydrodynamics



$$\text{YOC} \approx (1 - \text{RKE})^{5.5} \hat{M}_{\text{HS}}^{-2} \hat{\tau}_{\text{BW}}$$

Yield degradation includes the mass ablation effect

$$\hat{M}_{\text{HS}} \approx (1 - \text{RKE})^{4/7} \hat{\tau}_{\text{BW}}^{2/7}$$

$$\text{YOC} \approx (1 - \text{RKE})^{4.4} \hat{\tau}_{\text{BW}}^{3/7} \approx (1 - \text{RKE})^{4.4}$$




Neglect/slowly varying

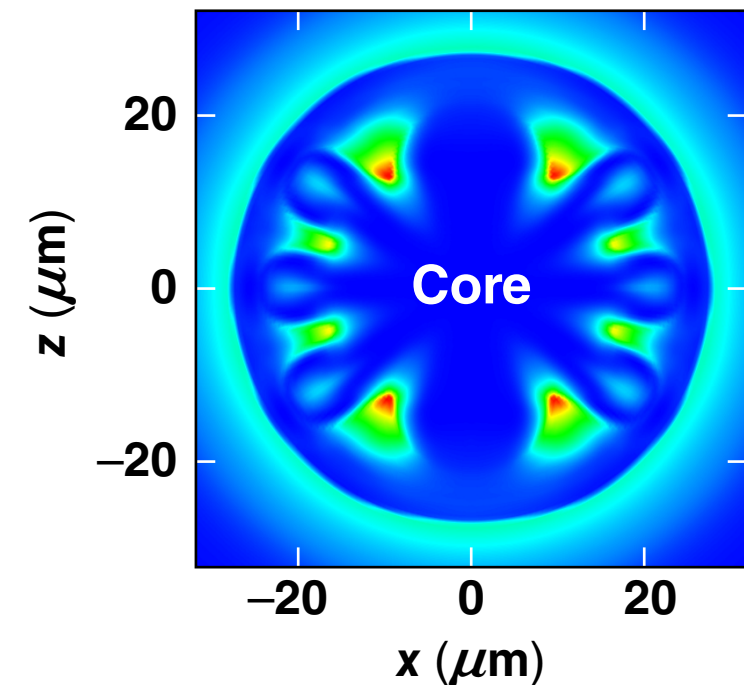
$$\frac{1}{2} \rho v^2 \text{ (J/m}^3\text{)}$$

5.797 × 10¹⁵

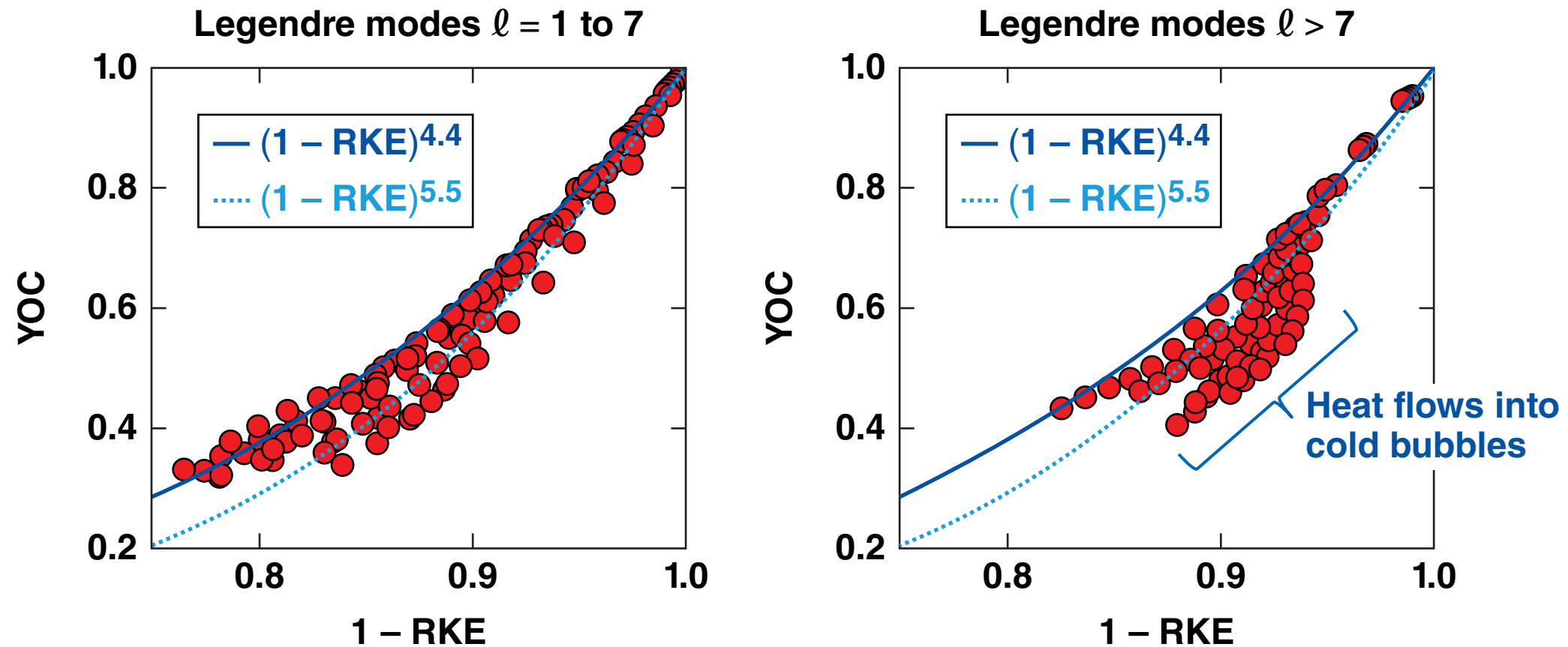
1.404 × 10¹⁰



Kinetic energy density for single-mode $\ell = 12, m = 6$



The RKE model provides a reasonable approximation for the YOC for low to mid modes and provides an upper bound for the YOC for mid modes



The results are consistent with 2-D *HYDRA* simulations by Kritcher *et al.**

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