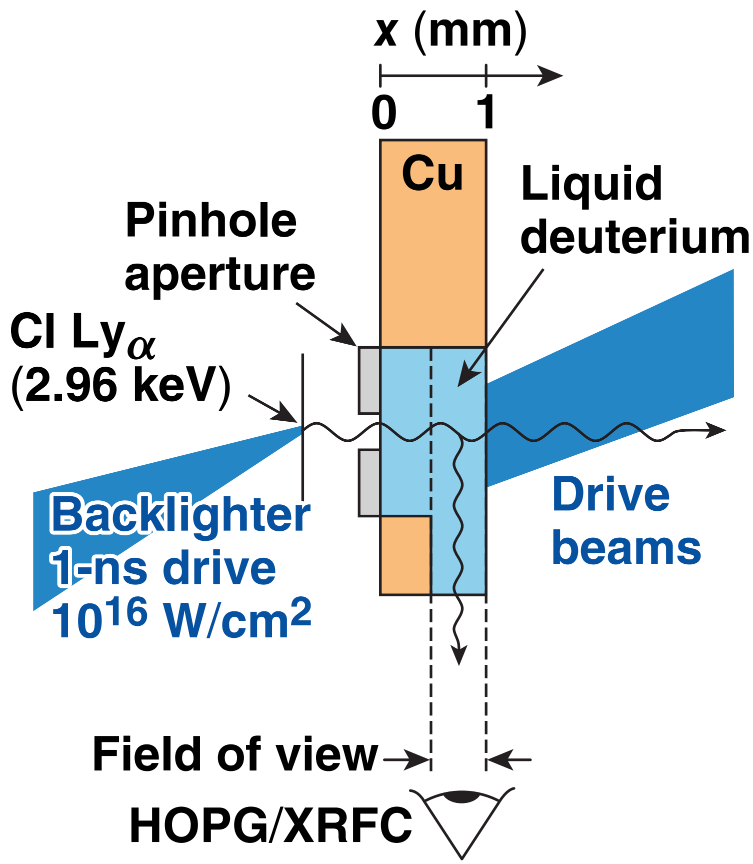
X-Ray Thomson-Scattering: Incisive Probe for Warm, Dense Matter



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Division of Plasma Physics
Providence, RI
29 October–2 November 2012

Summary

X-ray Thomson scattering* (XRTS) is a primary diagnostic for high-energy-density-physics (HEDP) experiments

- The conditions of dense plasmas are probed with XRTS, especially warm, dense matter with $T_e \sim T_F$ and the ratio of potential energy to kinetic energy of the ions greater than unity
- The elastic and inelastic x-ray scattering features are spectrally resolved
  - scattering from electrons (noncollective) → T_e, Z
  - scattering from plasmons (collective) → n_e
  - elastic scattering → structure of matter
- A review of XRTS experiments is presented
  - inertial confinement fusion (ICF)
  - radiation heated
  - shock heated and compressed
  - proton heated
  - laboratory astrophysics (radiative shocks and planetary interiors)

Many XRTS experiments need spatially resolved spectral measurements.

*S. H. Glenzer and R. Redmer, Rev. Mod. Phys. 81, 1625 (2008).
O. L. Landen *et al*., J. Quant. Spectrosc. Radiat. Transf. 71, 465 (2001).
G. Gregori *et al*., Phys. Rev. E 67, 026412(2003).

XRTS research involves many international collaborations

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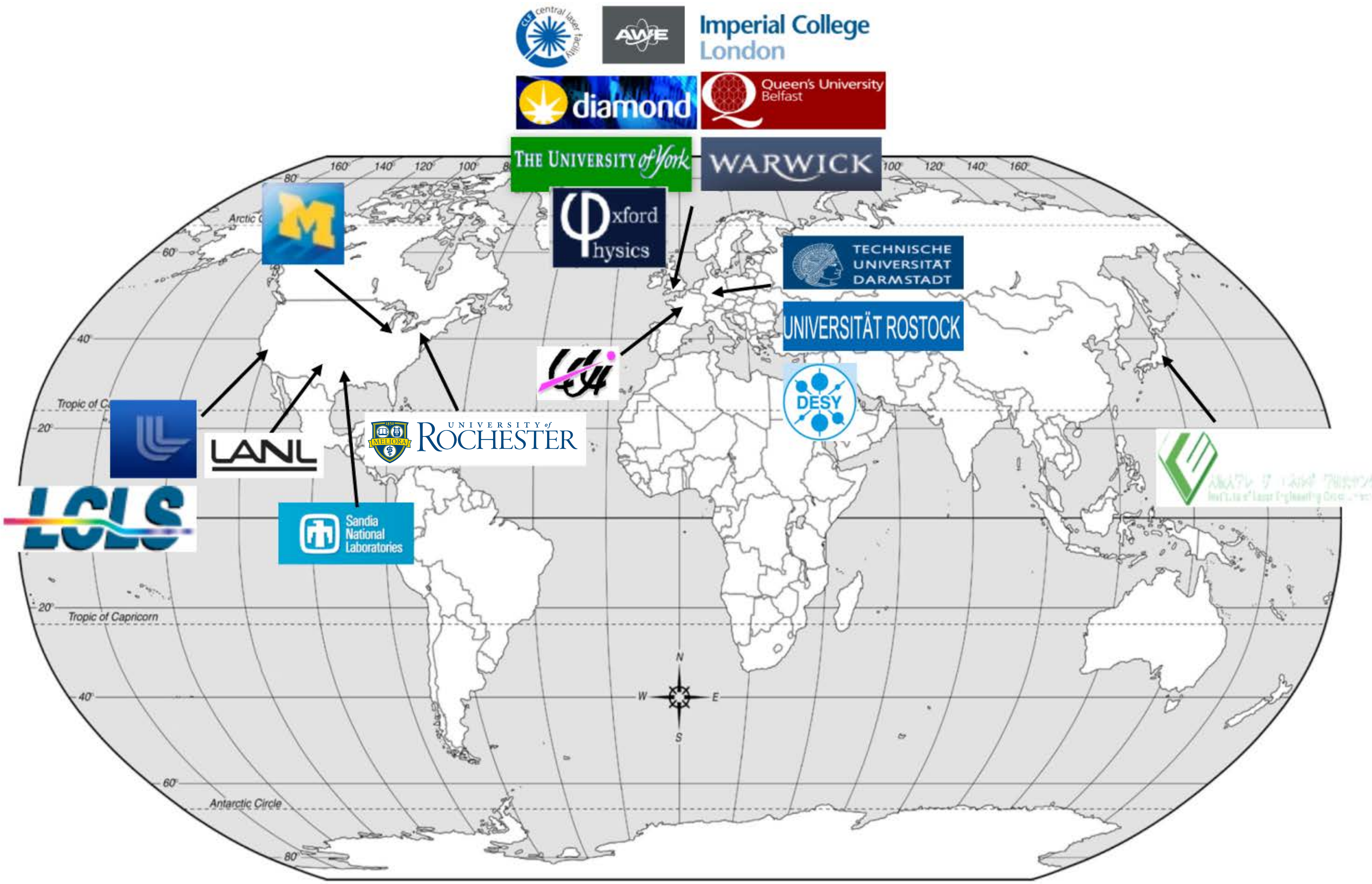

Collaborators for XRTS experiment of shocked liquid deuterium

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B. Crowley,[2,3] S. H. Glenzer,[4] O. L. Landen,[4] D. O. Gericke,[5]
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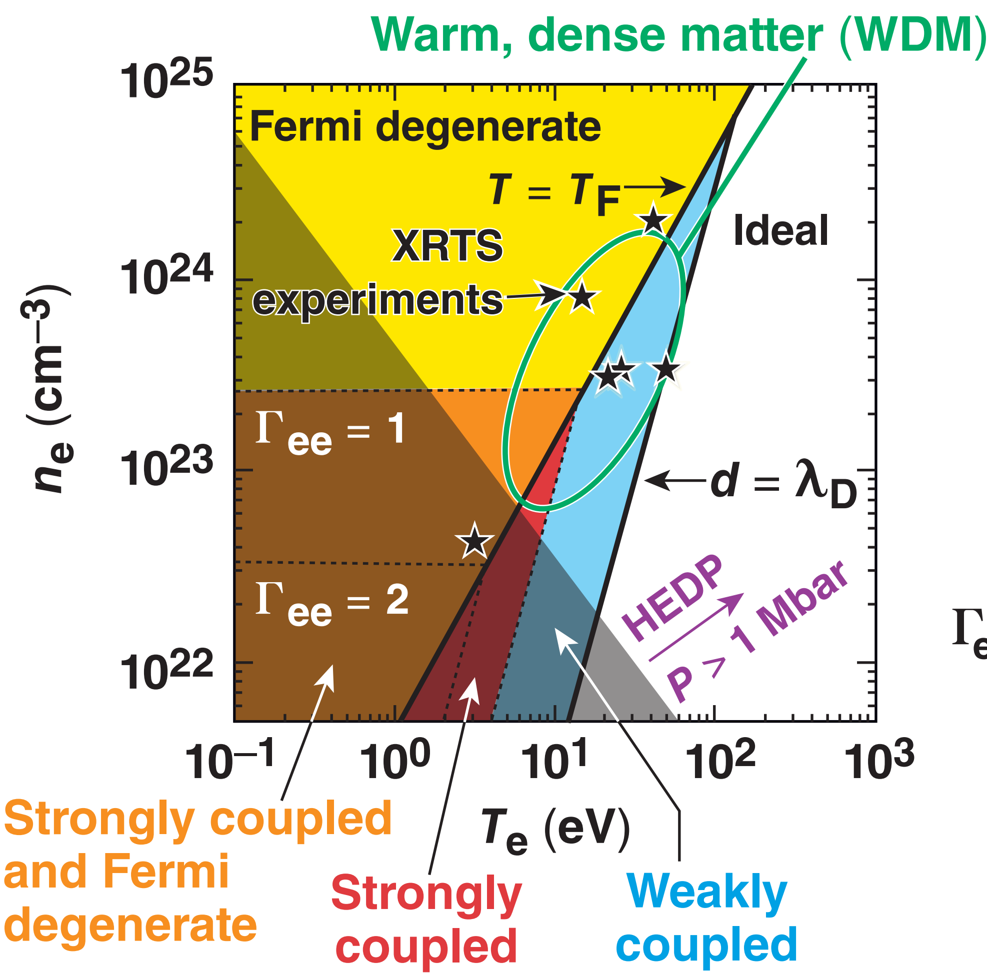
Outline

- **Motivation**
- **XRTS**
- **Experiments**

X-ray Thomson scattering* (XRTS) is a primary diagnostic for HEDP experiments



Degeneracy parameter

$$\theta = \frac{k_B T_e}{E_F}, \quad E_F = \frac{\hbar^2}{2m_e}\left(3\pi^2 n_e\right)^{2/3}$$

Coupling parameter

$$\Gamma_{ee} = \frac{e^2}{4\pi\varepsilon_0 \bar{d} k_B T_e}, \quad \bar{d} = \left(\frac{4\pi n_e}{3}\right)^{-1/3}$$

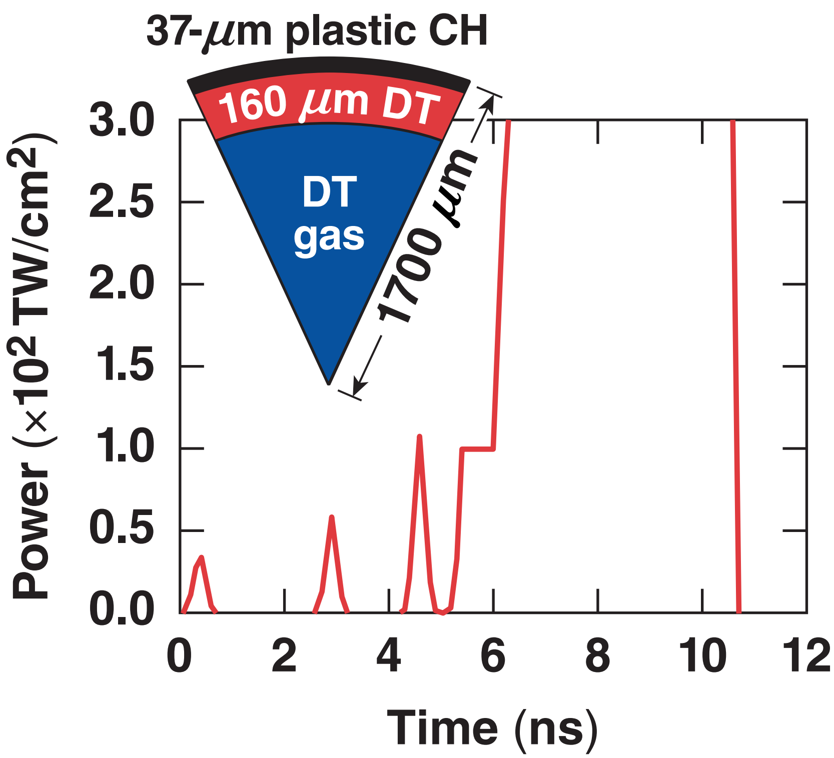
Penetrating x rays are needed to probe dense plasmas.

E16639b

*S. H. Glenzer and R. Redmer, Rev. Mod. Phys. 81, 1625 (2008).
O. L. Landen *et al.*, J. Quant. Spectrosc. Radiat. Transf. 71, 465 (2001).
G. Gregori *et al.*, Phys. Rev. E 67, 026412(2003).

Diagnosing the shock heated and compressed-shell conditions in an ICF target is an ideal XRTS application

Triple-picket, direct-drive design for the NIF[1]



$$E_{\text{min}} \sim \frac{\alpha^{1.8}}{V_{\text{imp}}^{5.8}}$$

(minimum laser energy for ignition)[2,3]

$$\text{Minimize} \rightarrow \alpha = \frac{P_{\text{fuel}}}{P_{\text{Fermi}}}$$

The minimum energy needed for ignition depends on the plasma conditions in the fuel layer (P_{fuel}), which can be diagnosed with XRTS.

[1]V. N. Goncharov *et al.*, Phys. Rev. Lett. 104, 165001 (2010).
[2]M. C. Hermann, M. Tabak, and J. Lindl, Nuc. Fusion 41, 99 (2001).
[3]R. Betti *et al.*, Phys. Plasmas 9, 2277 (2002).

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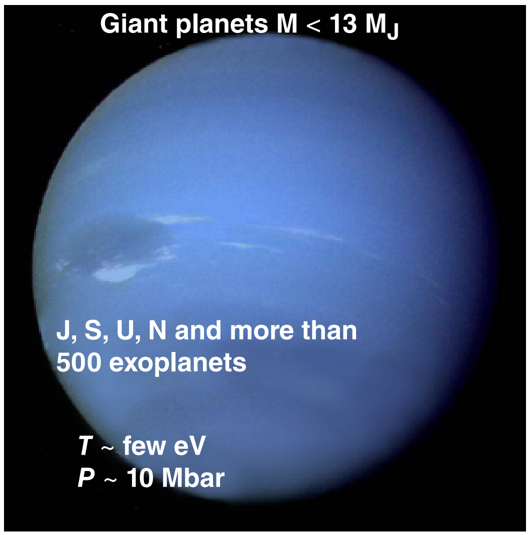
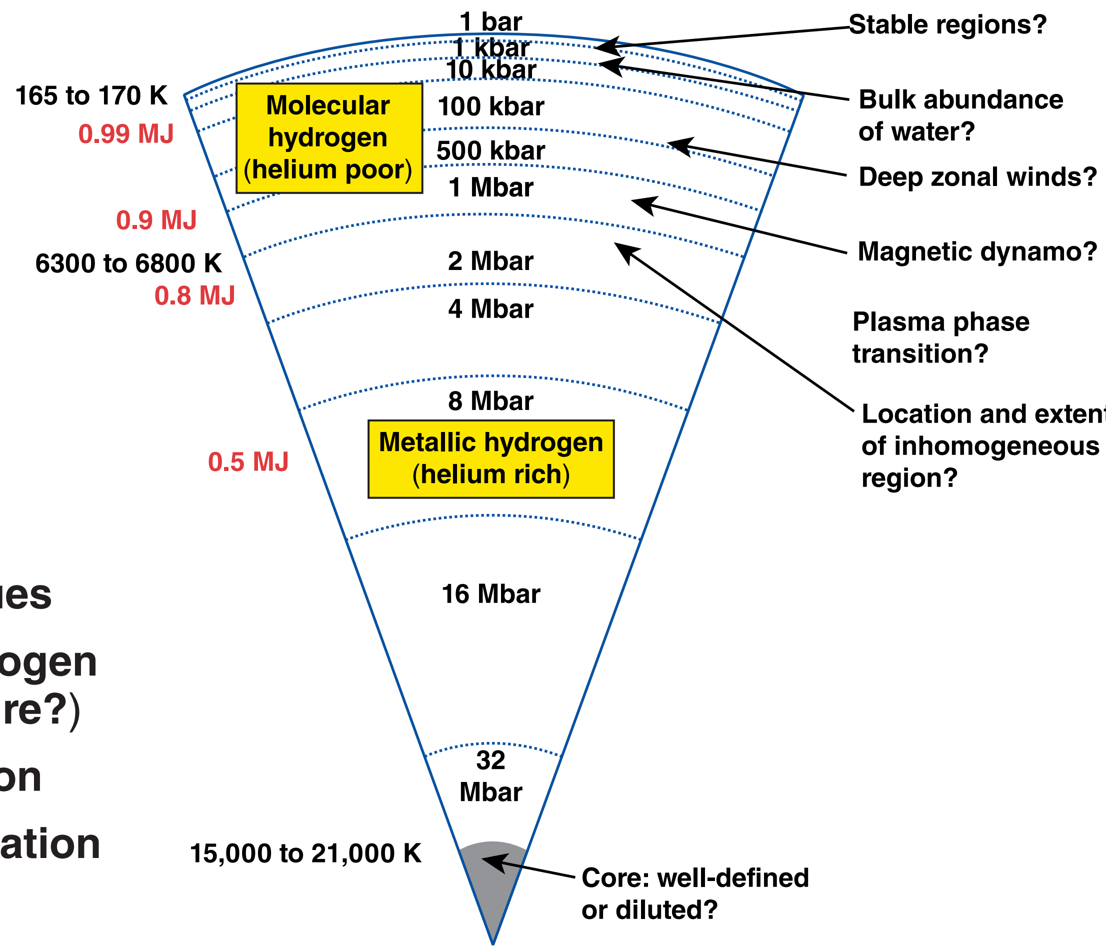
XRTS can provide accurate measurements of conditions in HEDP plasmas designed for planetary science



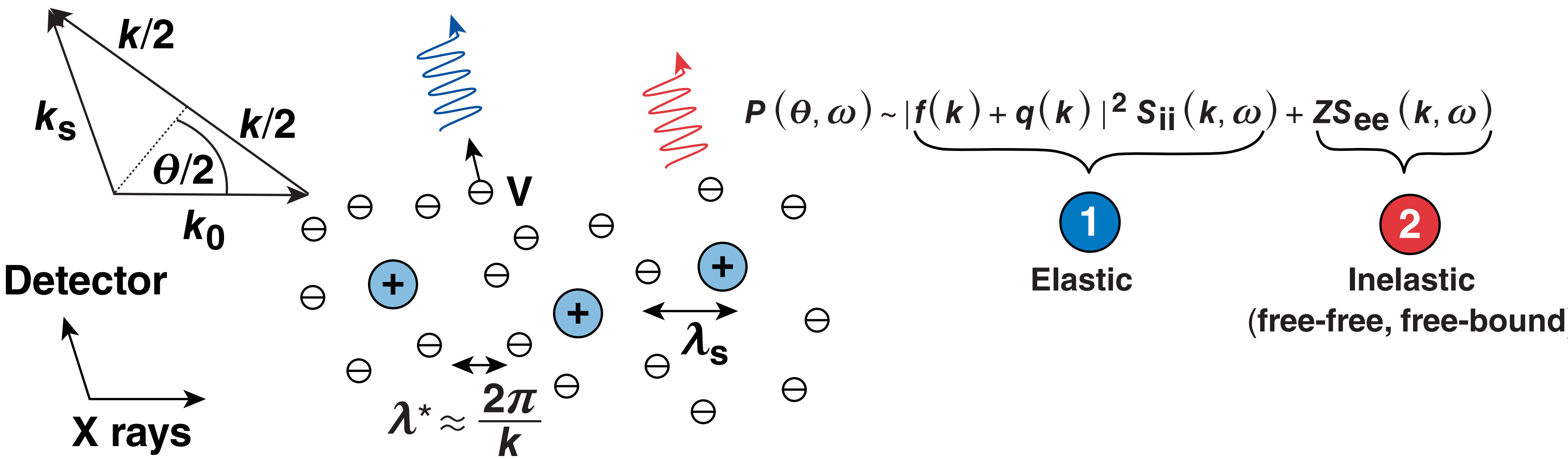


- Outstanding physics issues
  - metallization of hydrogen (pressure, temperature?)
  - H/He phase separation
  - magnetic-field generation

Outline

- Motivation
- **XRTS**
- Experiments

The properties of dense plasmas can be diagnosed with XRTS

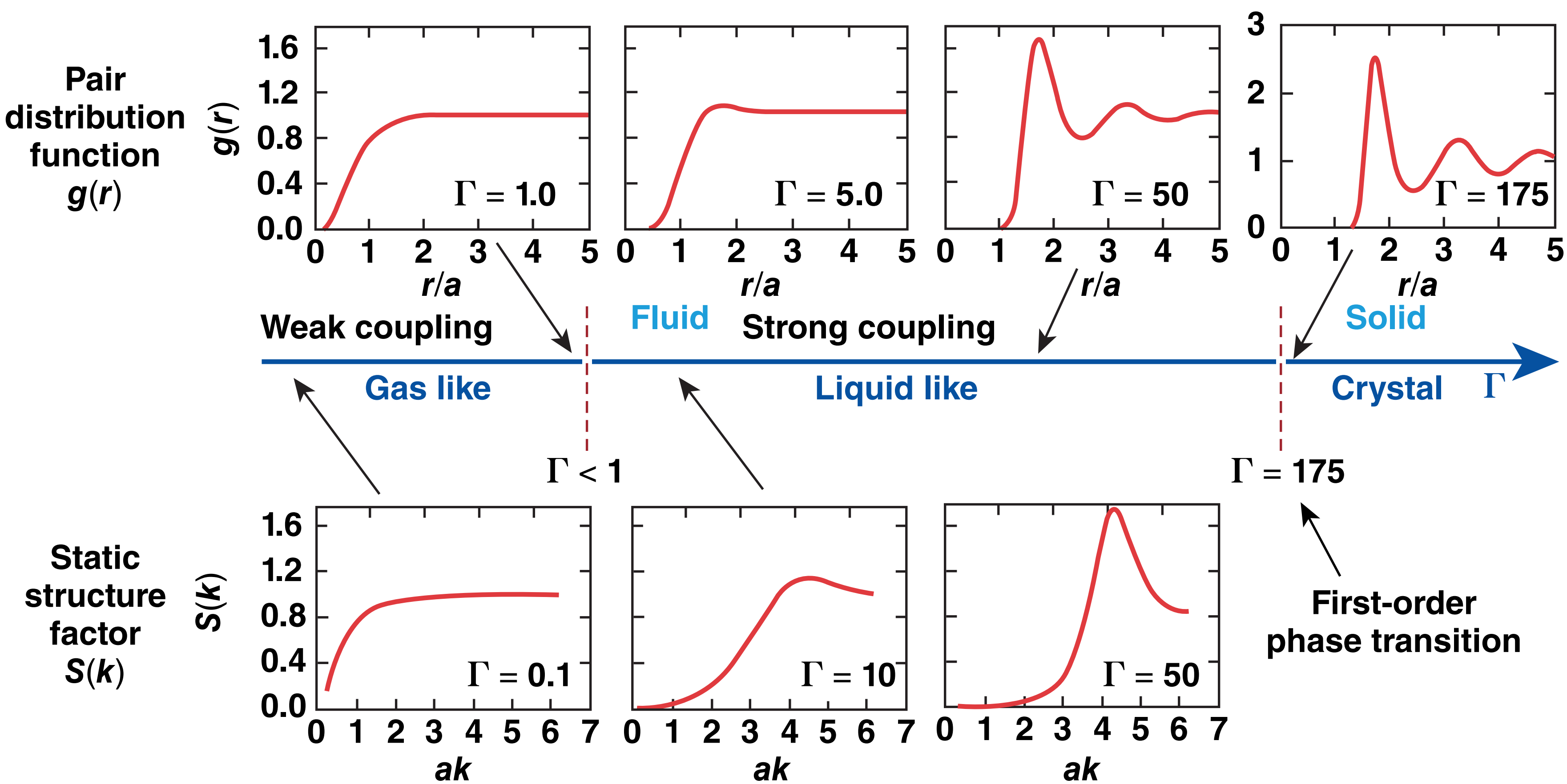


1. **Ion–ion correlations**
   – scattering as a function of the scattering angle exhibit peaks that are representative of the structure (diffraction)

2. **Electron–electron correlations**
   – scattering as a function of energy describes either single-particle dynamics or correlated plasmons

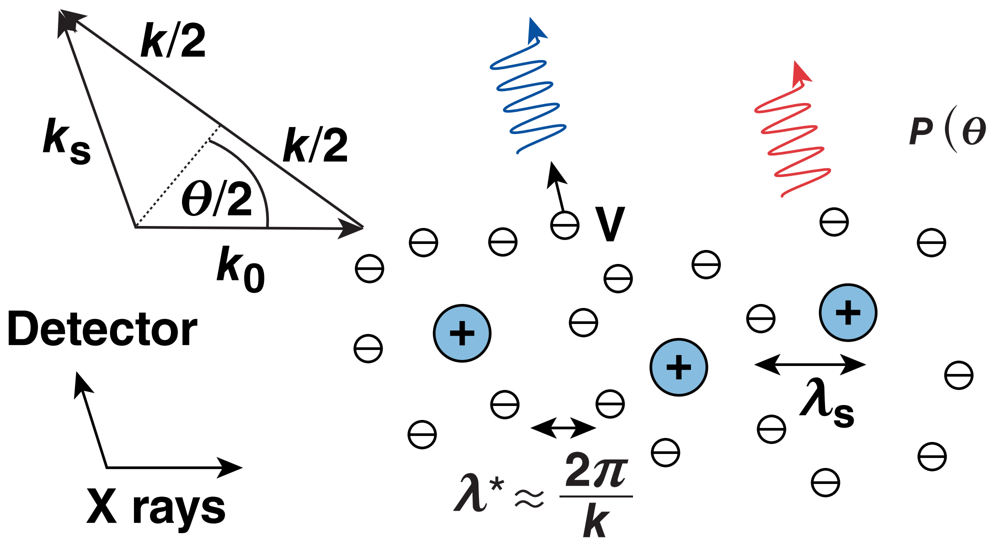
S. H. Glenzer and R. Redmer, Rev. Mod. Phys. 81, 1625 (2008).
G. Gregori *et al*., Phys. Rev. E 67, 26412 (2003).

Elastic scattering (diffraction) is used to test the degree of correlation between ions



- As the plasma gets denser (or cooler) correlations among ions emerge with the formation of crystalline structure

The microscopic characterization of warm, dense matter is inferred from inelastic scattering



$$P(\theta,\omega) \sim \underbrace{|f(k)+q(k)|^2 S_{ii}(k,\omega)}_{1} + \underbrace{ZS_{ee}(k,\omega)}_{2}$$

1 Elastic

2 Inelastic (free-free, free-bound)

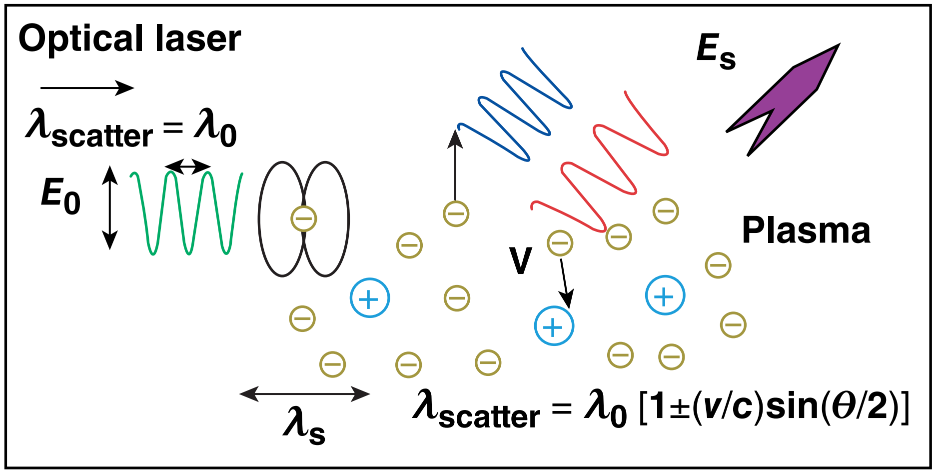
1. Ion–ion correlations
   - scattering as a function of the scattering angle exhibit peaks that are representative of the structure (diffraction)

2. Electron–electron correlations
   - scattering as a function of energy describes either single-particle dynamics or correlated plasmons

S. H. Glenzer and R. Redmer, Rev. Mod. Phys. 81, 1625 (2008).
G. Gregori *et al*., Phys. Rev. E 67, 26412 (2003).

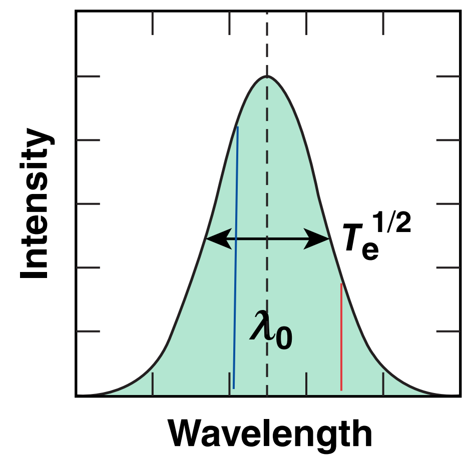
Inelastic x-ray scattering probes the electron-velocity distribution function

Noncollective Thomson scattering ($\lambda_{scatter} < \lambda_s$)

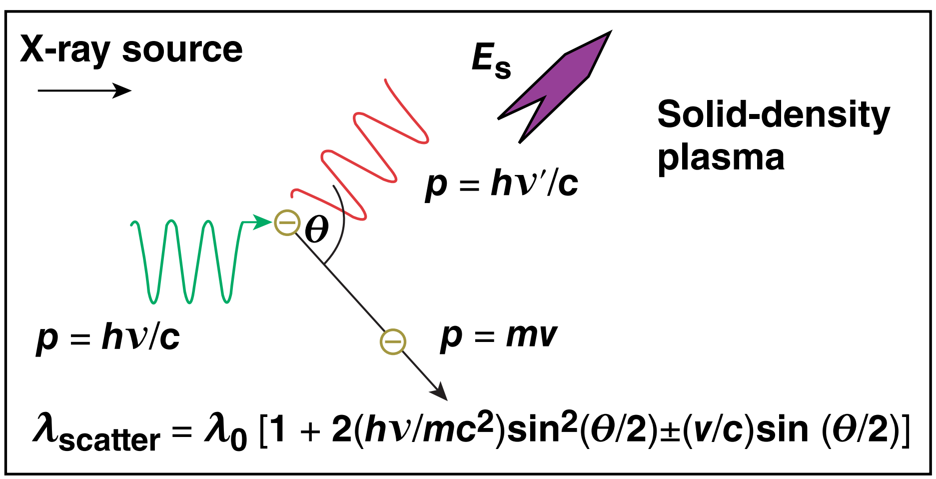


Scattering on free electrons

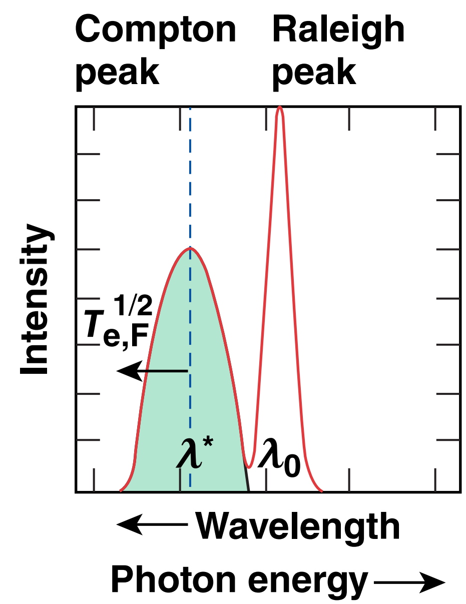
Boltzman distribution



X-ray Compton scattering



Scattering on free and weakly bound electrons



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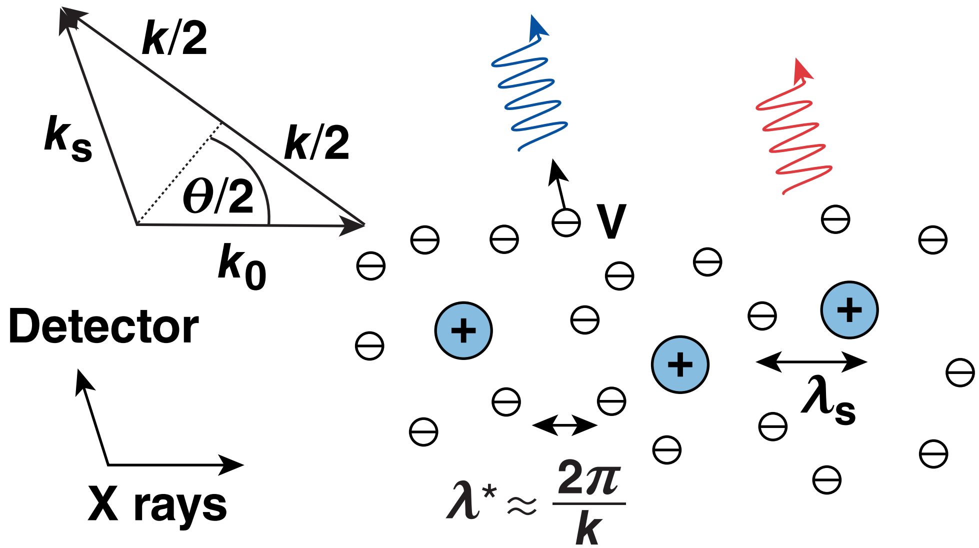
Collective inelastic x-ray scattering is used to diagnose the electron density



Scattering parameter

$$\alpha = \frac{\lambda^*}{2\pi\,\lambda_s} = \frac{1}{k\lambda_s}$$

$$k = |\mathbf{k}| = 4\pi \frac{E_0}{hc} \sin(\theta/2)$$

- $\alpha < 1$: scattering from electrons (noncollective) → T_e, Z
- $\alpha > 1$: scattering from plasmons (collective) → n_e

$$\omega_{plasmon} \sim \omega_{plasma} = \sqrt{n_e e^2/\varepsilon_0 m_e}$$

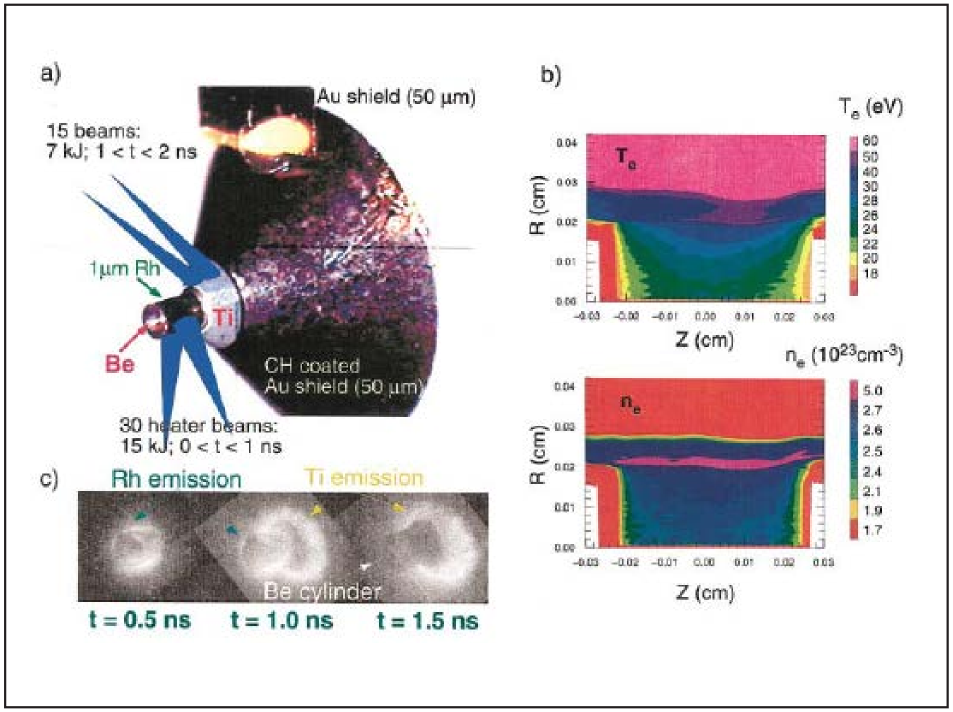
S. H. Glenzer and R. Redmer, Rev. Mod. Phys. 81, 1625 (2008).
G. Gregori *et al.*, Phys. Rev. E 67, 26412 (2003).

Outline

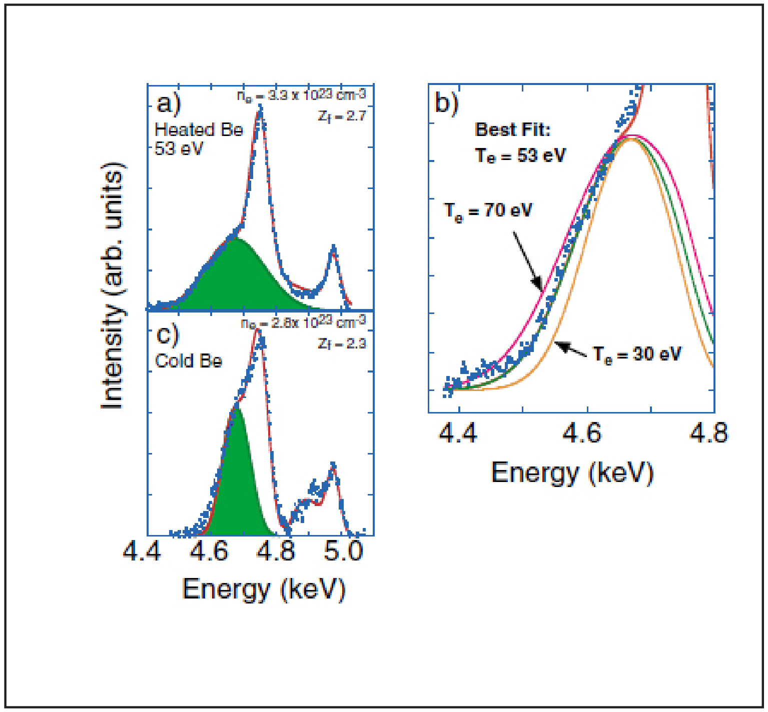
- Motivation
- XRTS
- **Experiments**

The first noncollective XRTS experiment diagnosed isochorically heated Be on OMEGA

Noncollective scattering from radiatively heated Be



Noncollective XRTS



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S. H. Glenzer *et al*., Phys. Rev. Lett. 90, 175002 (2003).

Collective XRTS was first demonstrated using isochorically heated Be on OMEGA

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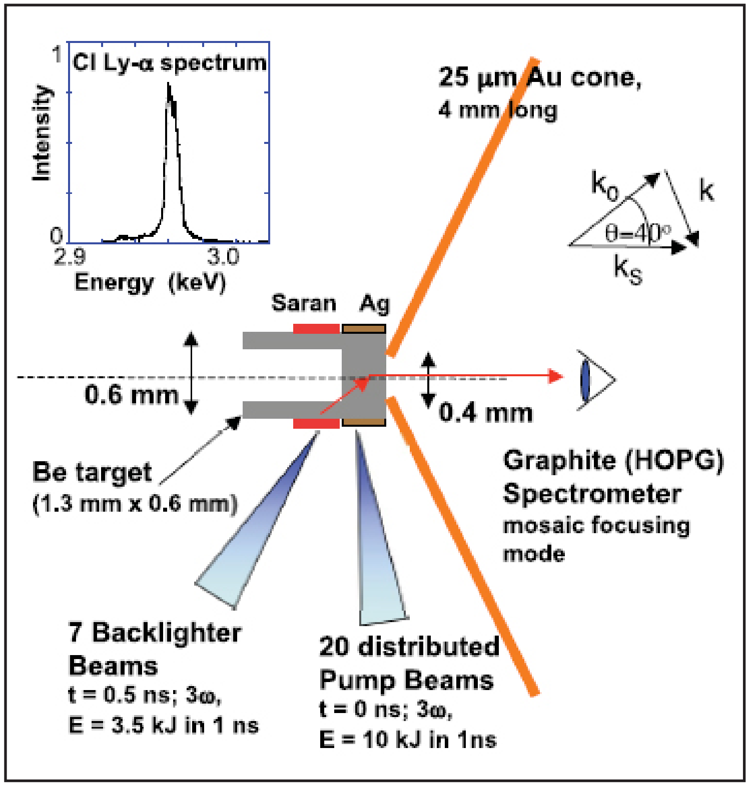
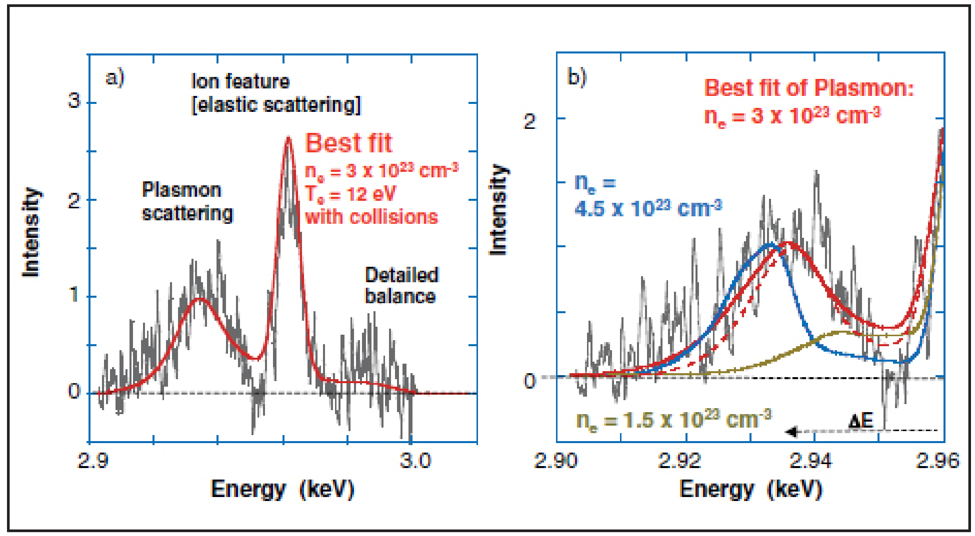
Collective scattering from radiatively heated Be



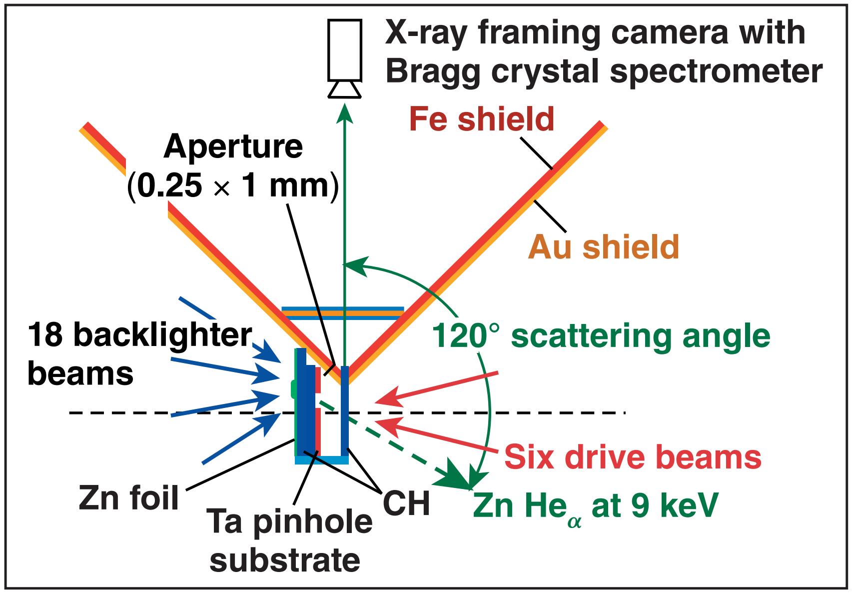
Collective XRTS



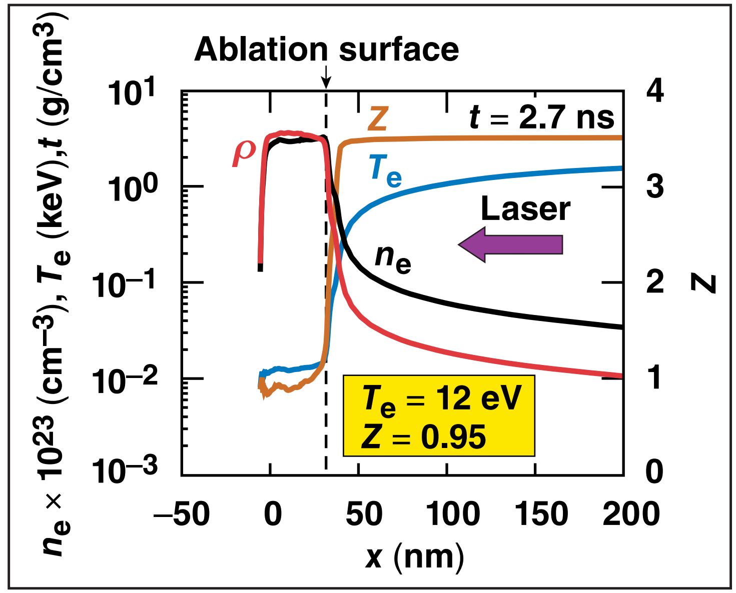
S. H. Glenzer *et al*., Phys. Rev. Lett. 98, 065002 (2007).

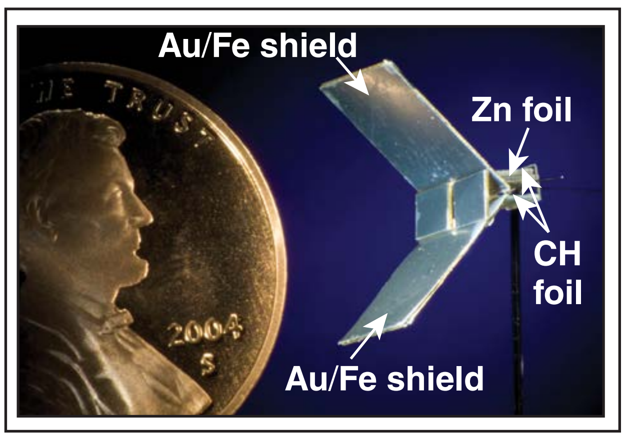
XRTS was applied to a laser-driven, shock-heated, and compressed CH foil on OMEGA

Shock heated and compressed CH foil



1-D simulation

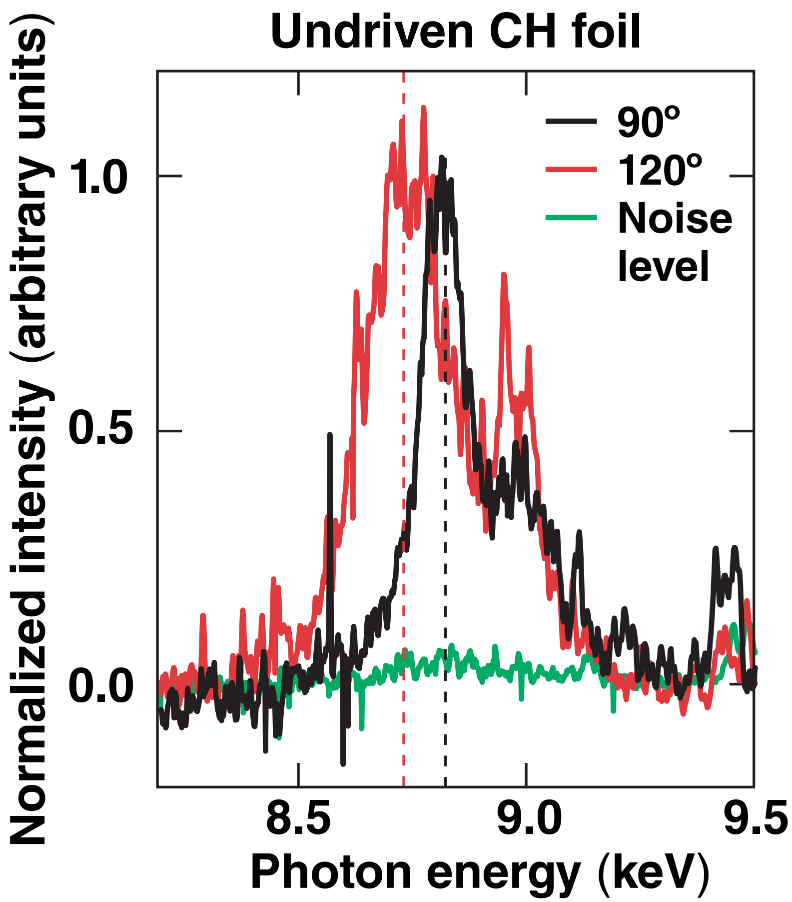




E21074

H. Sawada, S. P Regan *et al.*, Phys. Plasmas 14, 122703 (2007).

The Compton energy downshift is observed to be greater for the larger scattering angle



Compton downshifted energy (eV)

$$\Delta E_c = \frac{\hbar^2 k^2}{2m_e} \quad k = \frac{4\pi}{\lambda_0} \sin\left(\frac{\theta}{2}\right)$$

θ: scattering angle
λ_0: wavelength of probe

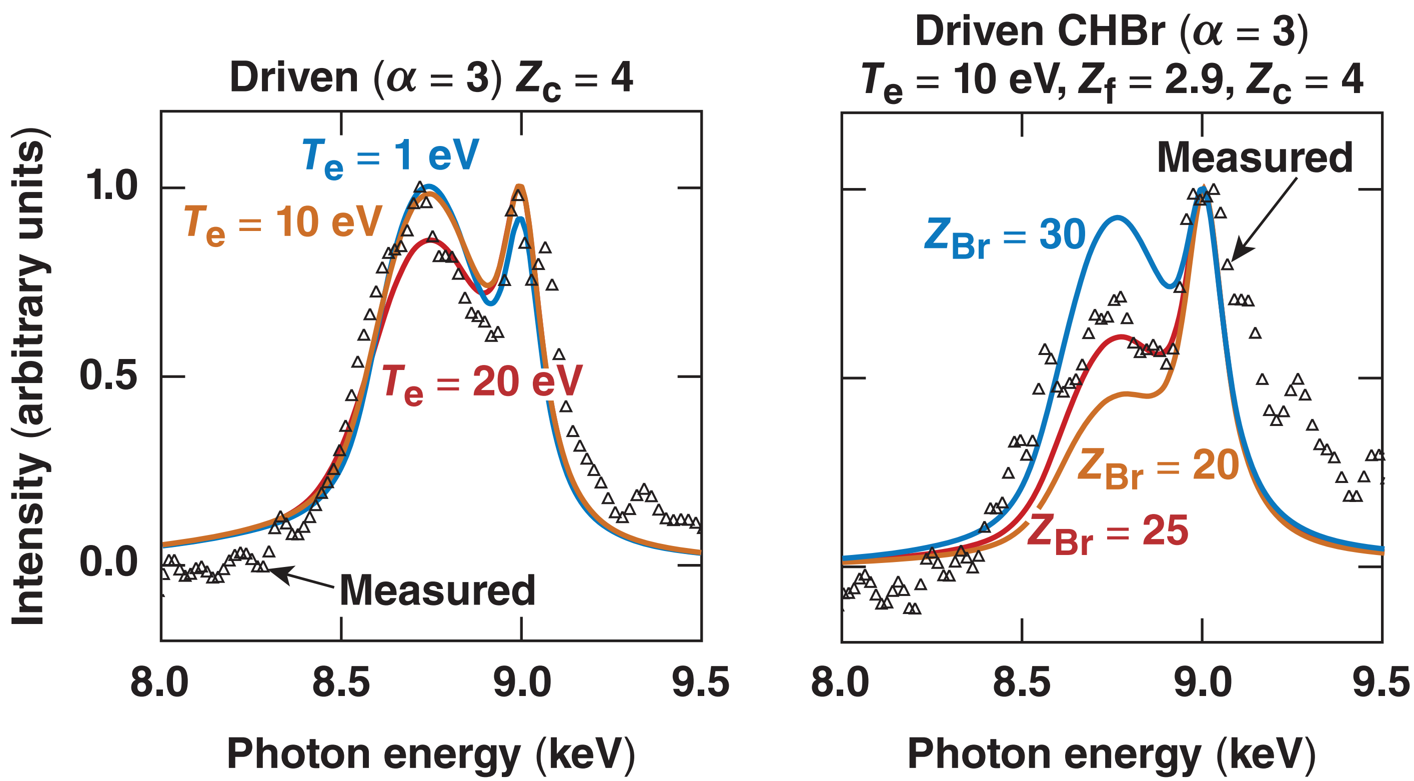
(Zn He$_\alpha$ ~ 1.3 Å ~ 9.0 keV)

ΔE_C = 178 eV (for 90°)
= 267 eV (for 120°)

E13475c

H. Sawada, S. P. Regan *et al.*, Phys. Plasmas 14, 122703 (2007).

An electron temperature of 10 to 20 eV was inferred from noncollective XRTS

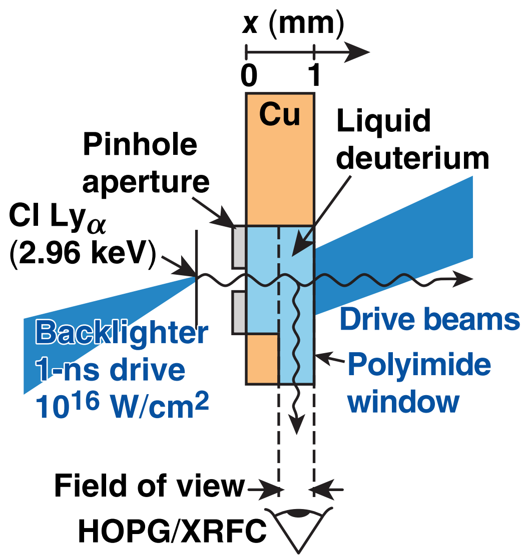


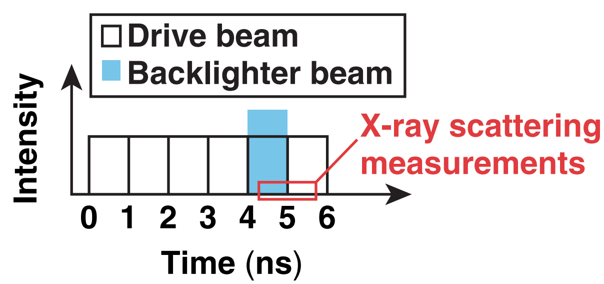
Br dopant in plastic increased the elastic scattering component.

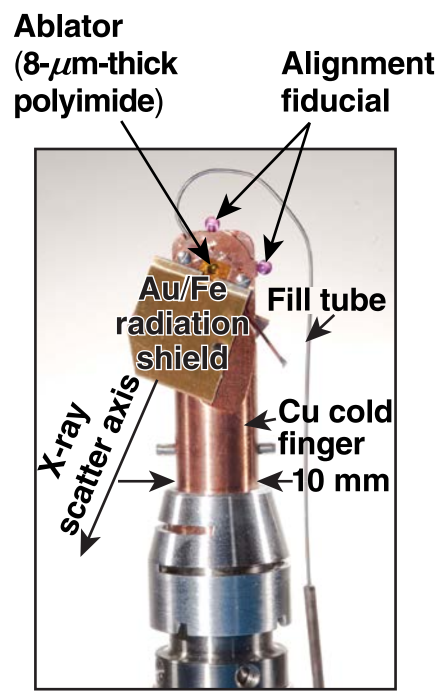
E15958

H. Sawada, S. P. Regan *et al.*, Phys. Plasmas 14, 122703 (2007).

A planar cryogenic target was developed to measure noncollective XRTS from shocked liquid deuterium





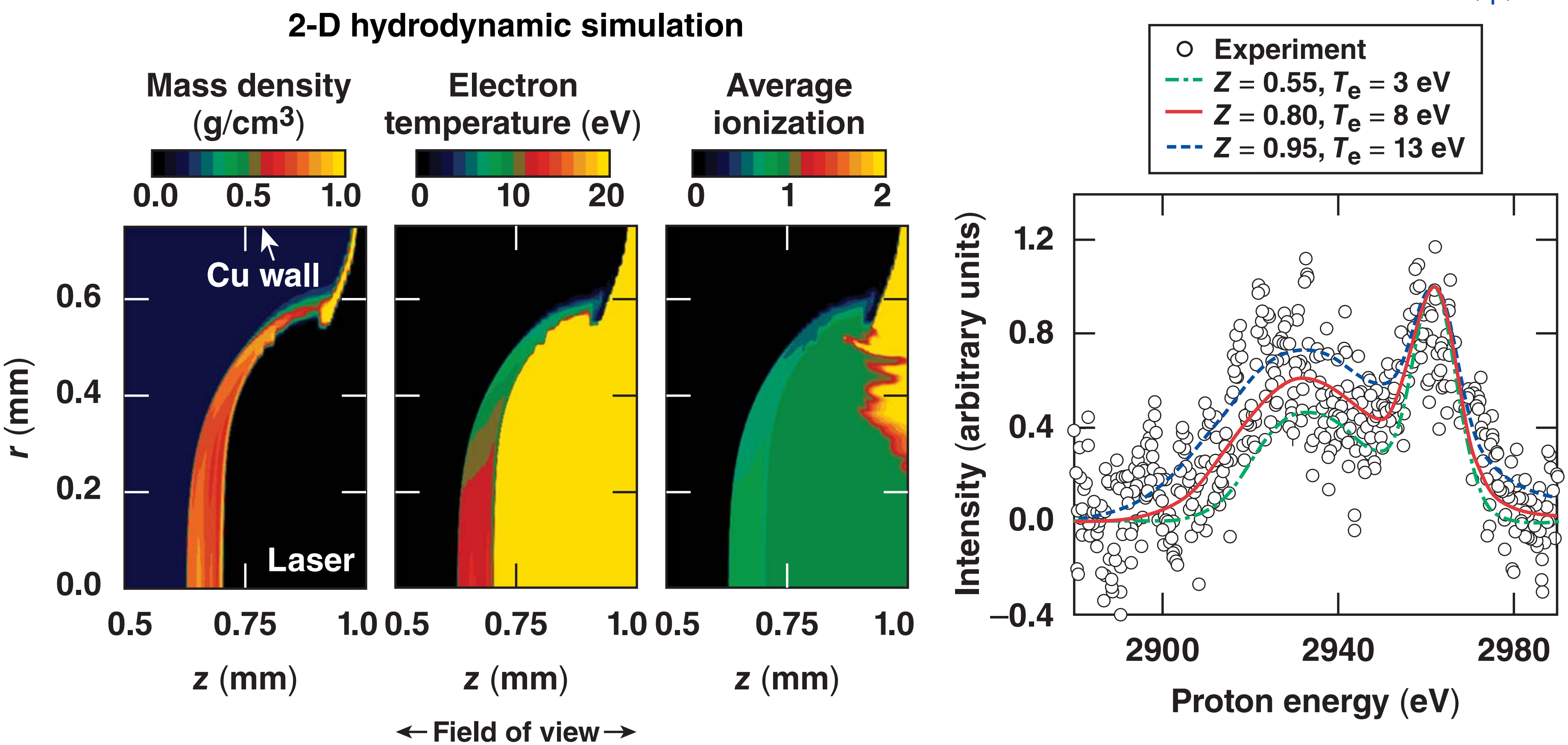


A 90° scattering geometry is used.

*S. P. Regan *et al.*, "Inelastic X-Ray Scattering from Shocked Liquid Deuterium," to be published in Physical Review Letters.

E19954b

2-D hydrodynamic simulations are consistent with the experimental results



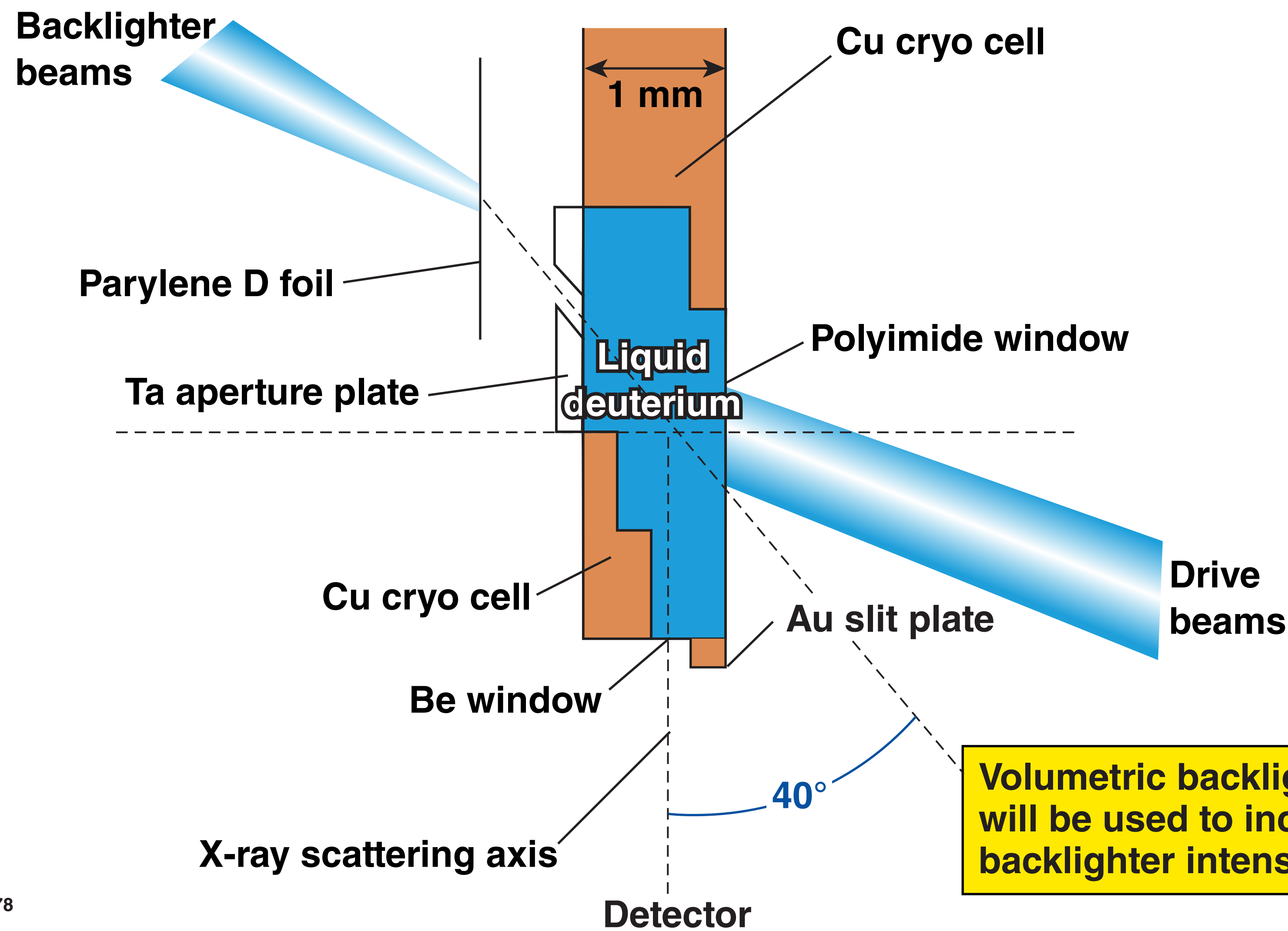
Spectral imaging would greatly benefit this experiment.

*S. P. Regan *et al.*, "Inelastic X-Ray Scattering from Shocked Liquid Deuterium," to be published in Physical Review Letters.

E18449a

Scattering angle will be decreased from 90° to 40° to measure collective XRTS and infer electron density

Shocked liquid deuterium experiment with 40° scattering geometry

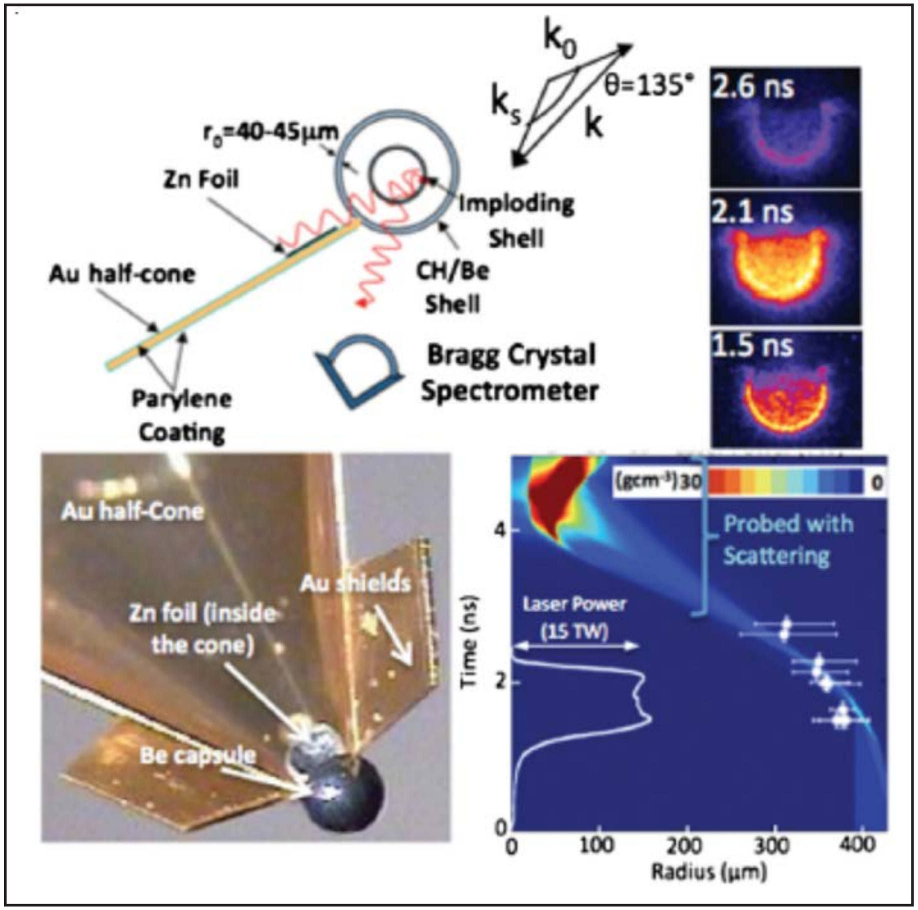


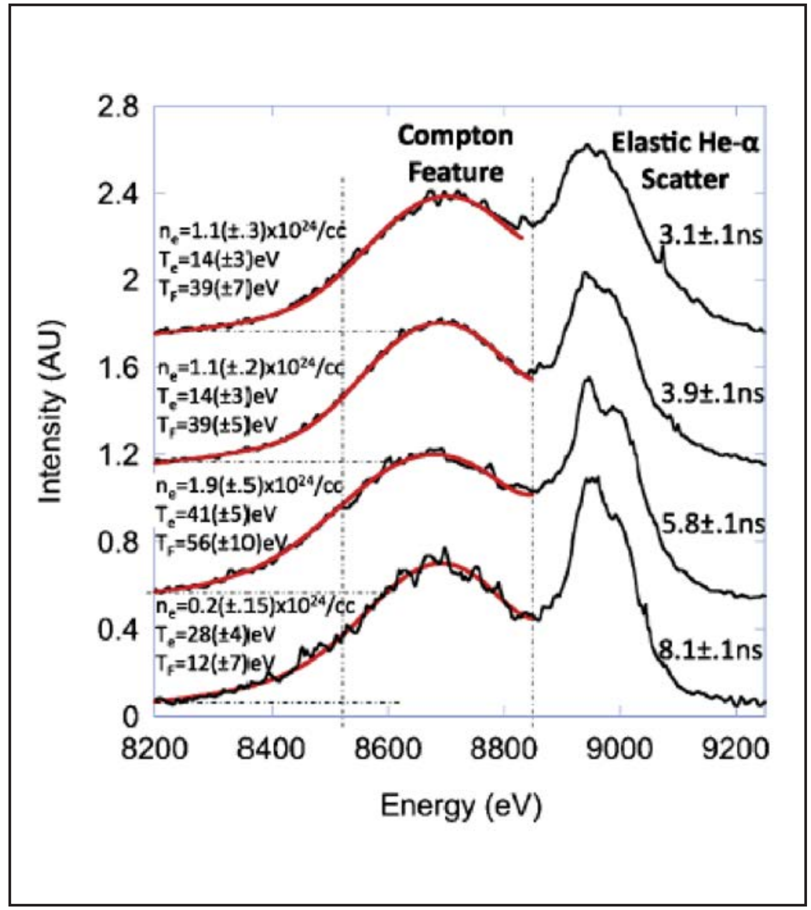
Volumetric backlighters will be used to increase backlighter intensity.

E21078

Compressed CH and Be shells were imploded on OMEGA and diagnosed with XRTS

Spherical direct-drive implosion



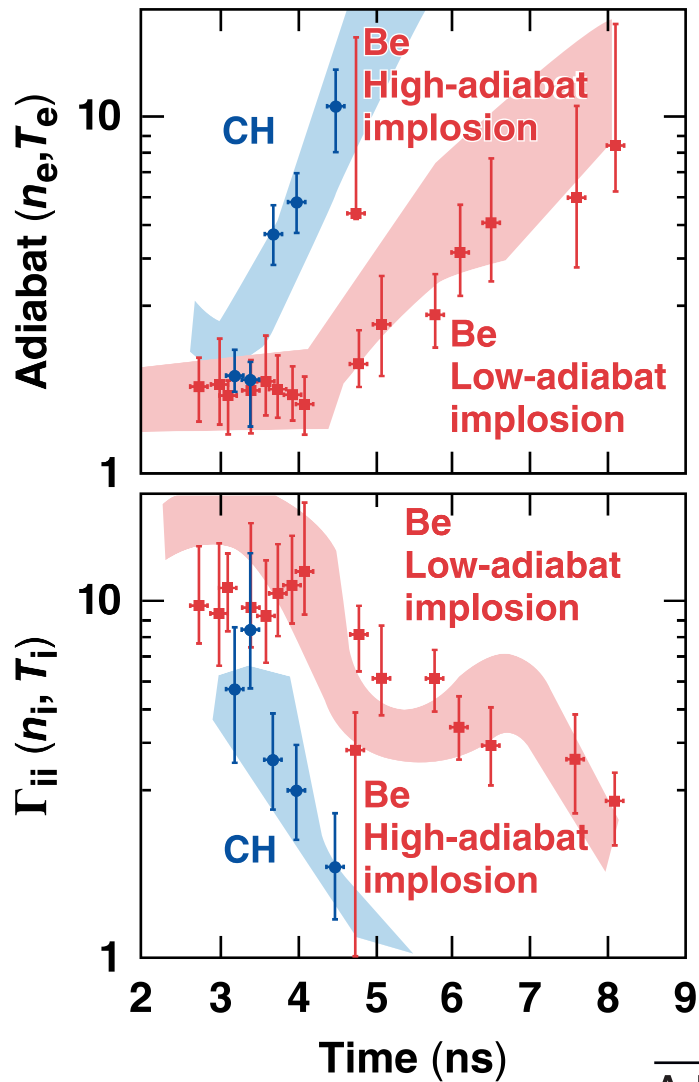


The adiabat (P_{fuel}/P_{fermi}) and the ion-ion coupling parameter were diagnosed with XRTS.

E21079

A. L. Kritcher *et al.*, Phys. Rev. Lett. 107, 015002 (2011).

The adiabat and the ion–ion coupling parameter were inferred from the XRTS measurements



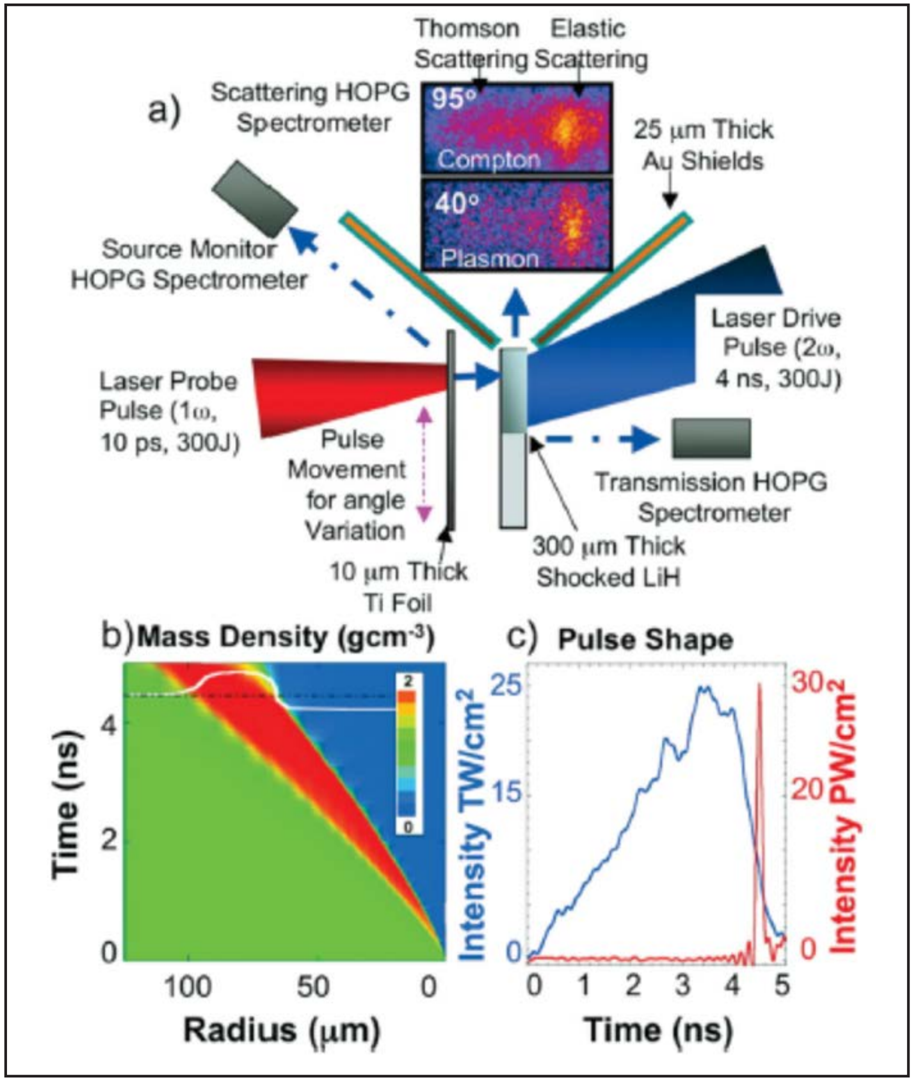
$$\text{adiabat} = \frac{P_{\text{final}}}{P_{\text{Fermi}}}$$

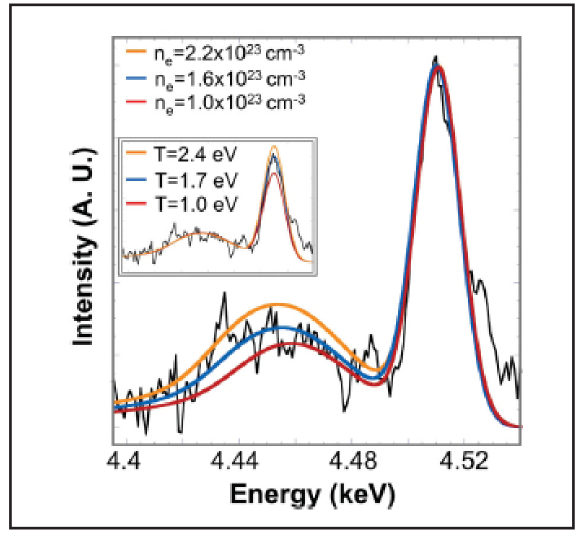
$$\Gamma_{ii} = (Ze)^2 / (4\pi\varepsilon_0 a\, k_B T_i)$$
$$a = \left(4/3\pi n_i\right)^{-1/3}$$

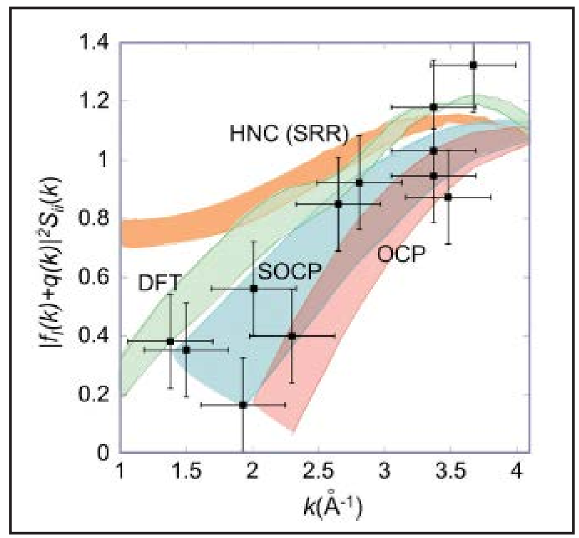
A. L. Kritcher *et al.*, Phys. Rev. Lett. 107, 015002 (2011).

E21080

Ultrafast x-ray thomson-scattering was used to probe ionic structure of shock compressed LiH on Titan



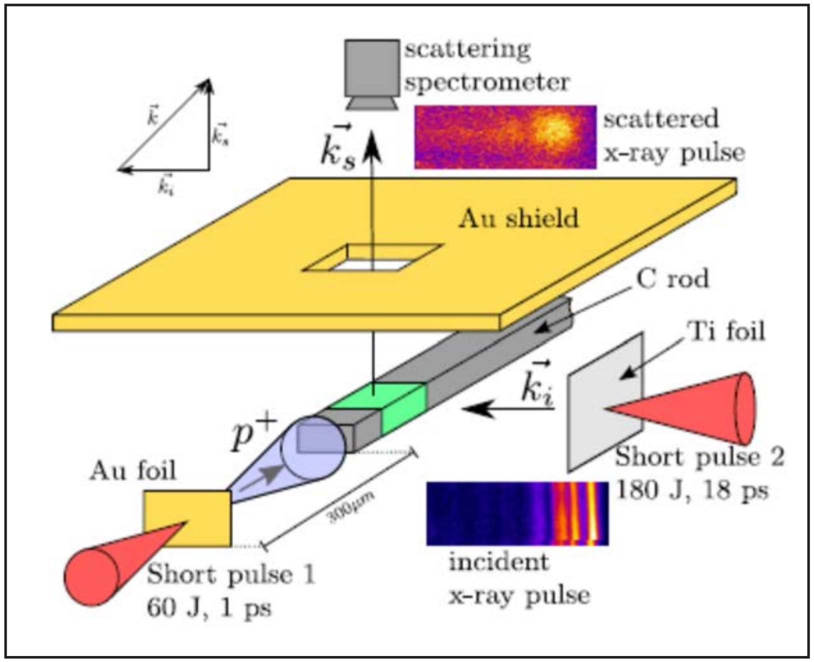


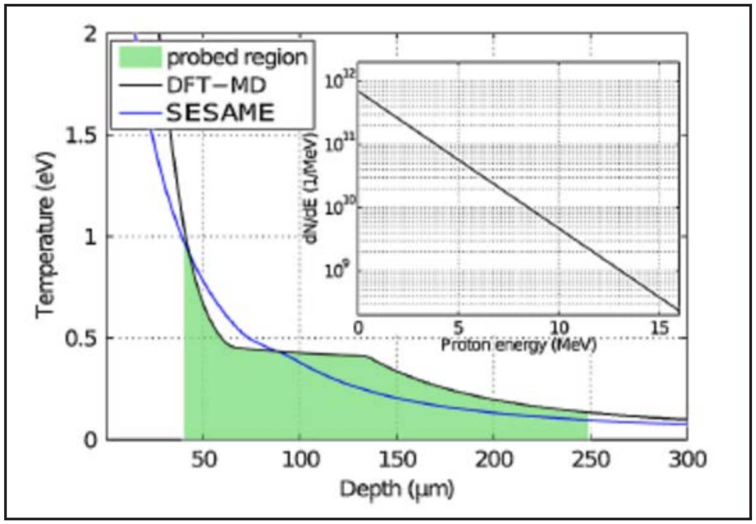


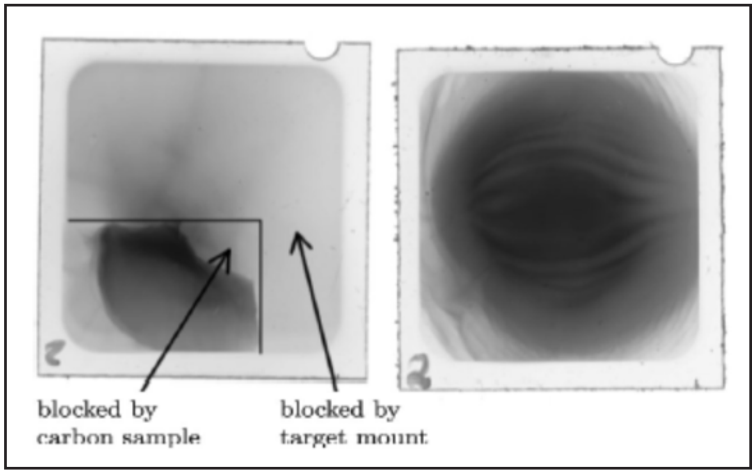
A. L. Kritcher *et al.*, Phys. Rev. Lett. 103, 25004 (2009).

The amount of liquid carbon was diagnosed with elastic x-ray scattering from a proton-heated target

Ultrafast proton-heating experiment



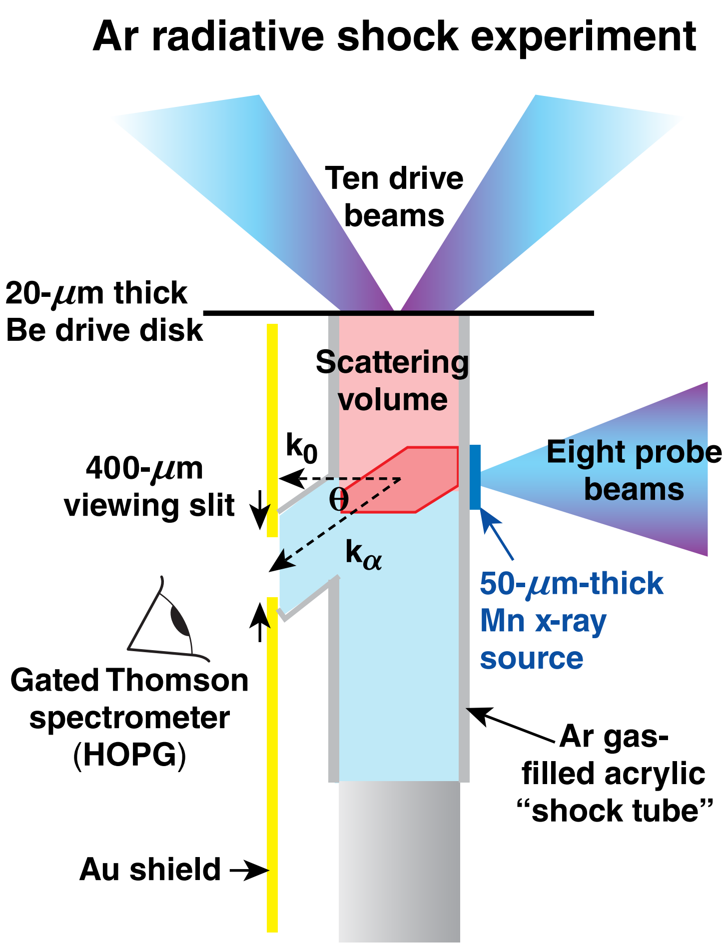


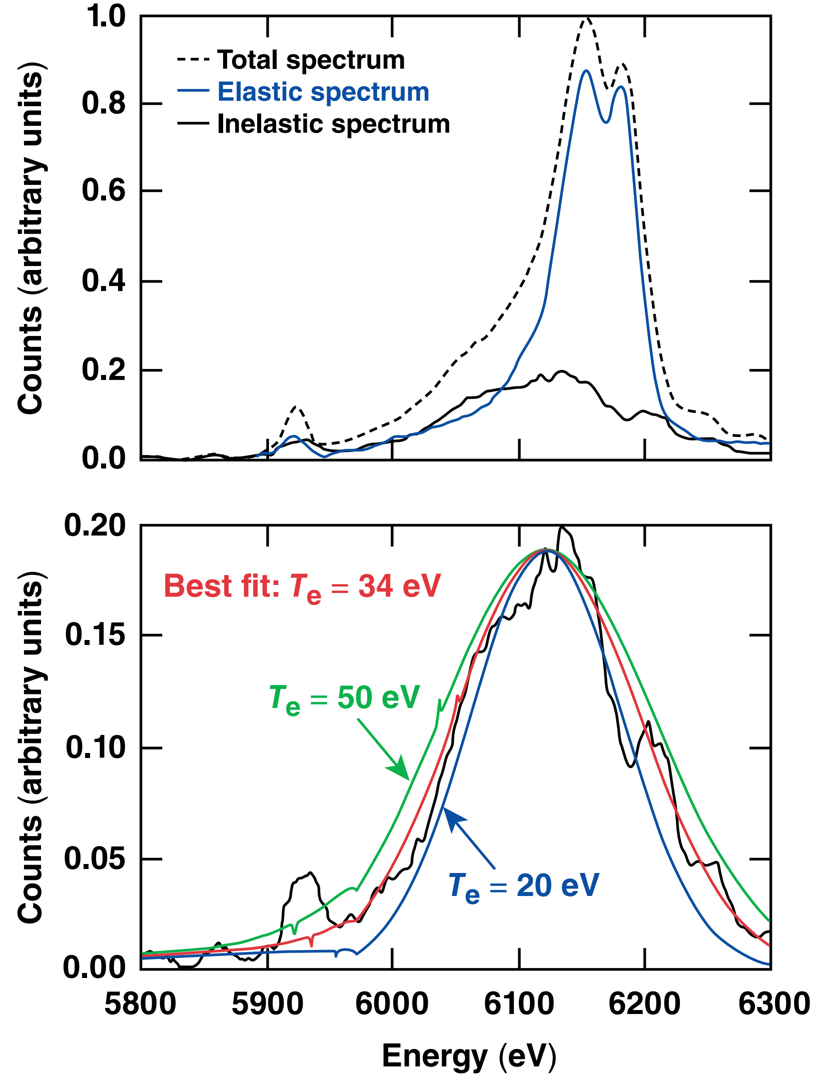


E21082

A. Pelka *et al.*, Phys. Rev. Lett. 105, 265701 (2010).

A radiative shock wave in Ar was diagnosed with XRTS on OMEGA

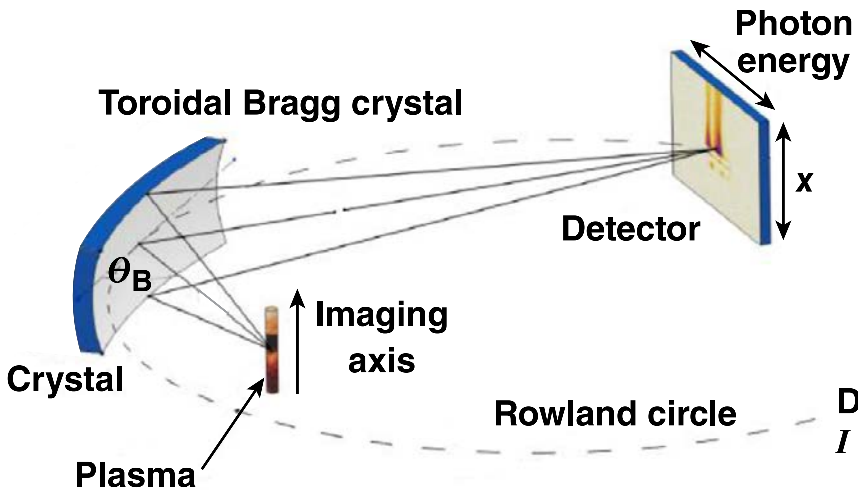


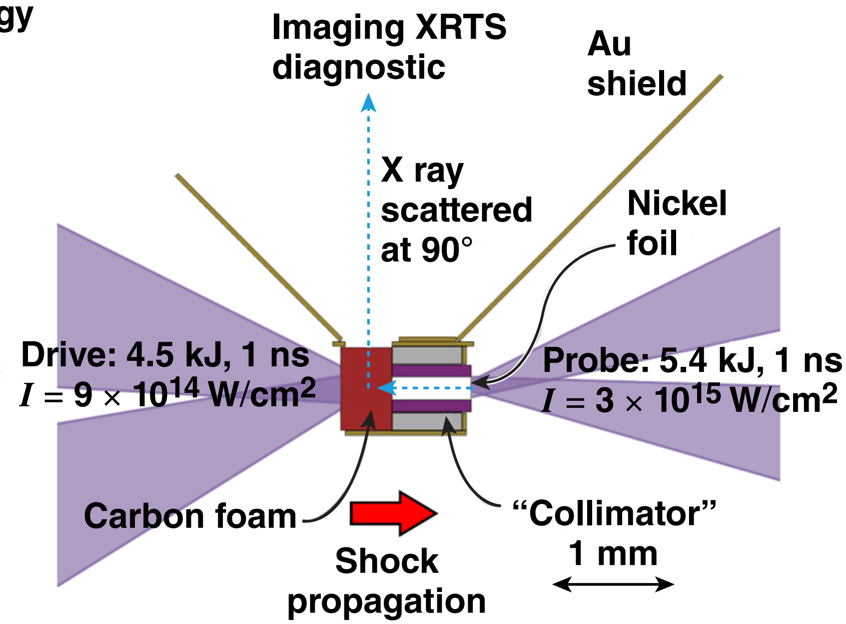


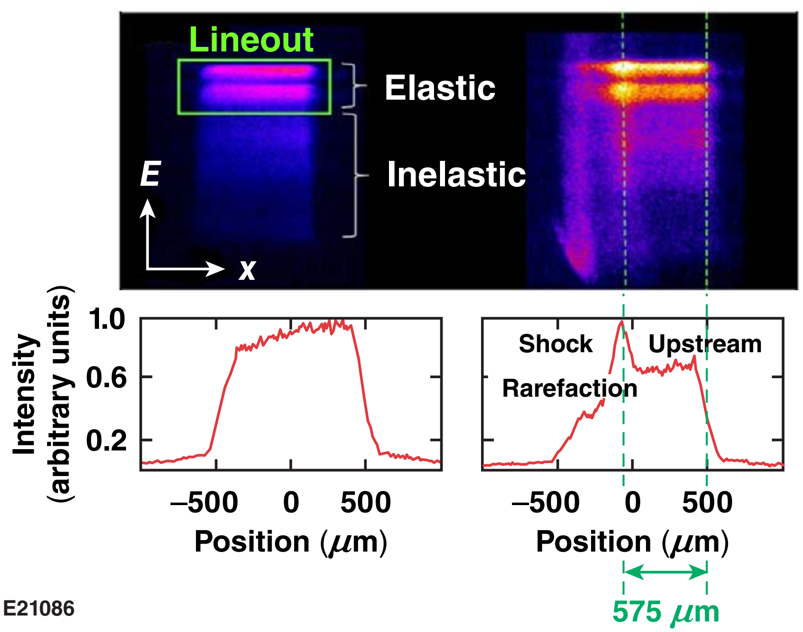
E21084

*A. J. Visco, *et al.,* PRL. 108, 145001 (2012).

The University of Michigan and LANL have deployed an imaging spectrometer for XRTS on OMEGA







Spatially resolved XRTS (dx ~ 50 μm) was recorded.

Gamboa *et al.*, to be presented at the 19th Topical Conference on High Temperature Plasma Diagnostics, Monterey, CA 6–10 May (2012).

E21086

Summary/Conclusions

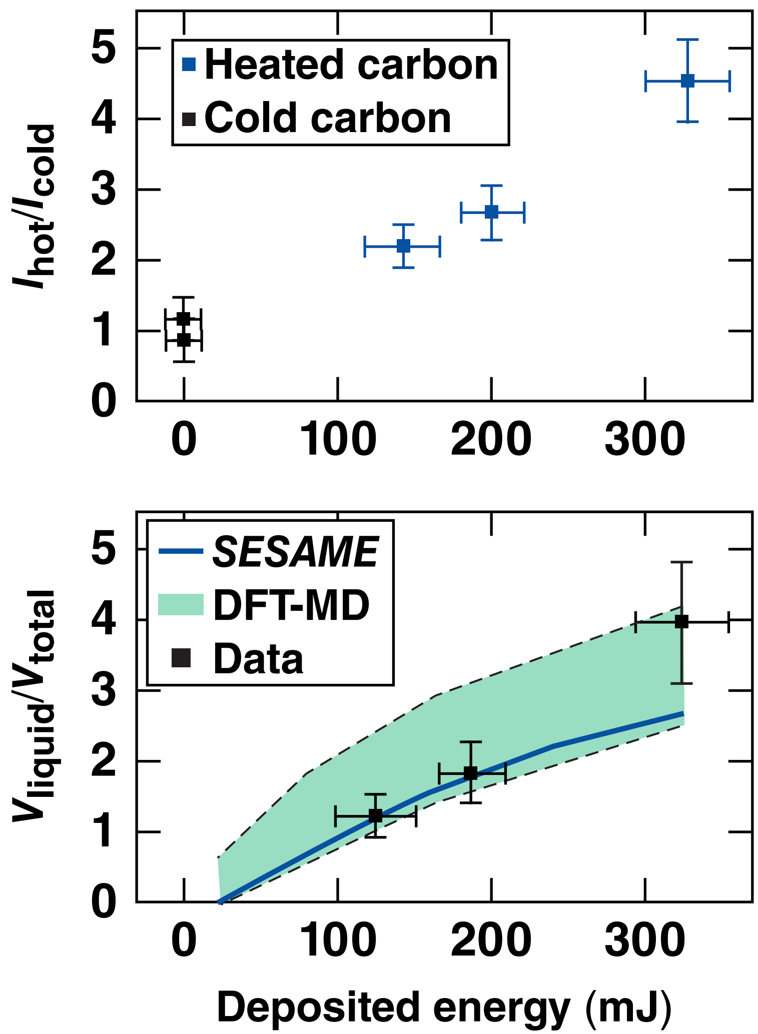
X-ray Thomson scattering* (XRTS) is a primary diagnostic for high-energy-density-physics (HEDP) experiments

- The conditions of dense plasmas are probed with XRTS, especially warm, dense matter with $T_e \sim T_F$ and the ratio of potential energy to kinetic energy of the ions greater than unity
- The elastic and inelastic x-ray scattering features are spectrally resolved
  - scattering from electrons (noncollective) → T_e, Z
  - scattering from plasmons (collective) → n_e
  - elastic scattering → structure of matter
- A review of XRTS experiments is presented
  - inertial confinement fusion (ICF)
  - radiation heated
  - shock heated and compressed
  - proton heated
  - laboratory astrophysics (radiative shocks and planetary interiors)

Many XRTS experiments need spatially resolved spectral measurements.

*S. H. Glenzer and R. Redmer, Rev. Mod. Phys. 81, 1625 (2008).
O. L. Landen *et al.*, J. Quant. Spectrosc. Radiat. Transf. 71, 465 (2001).
G. Gregori *et al.*, Phys. Rev. E 67, 026412(2003).

Experimental data on liquid fraction is used to validate equation-of-state models



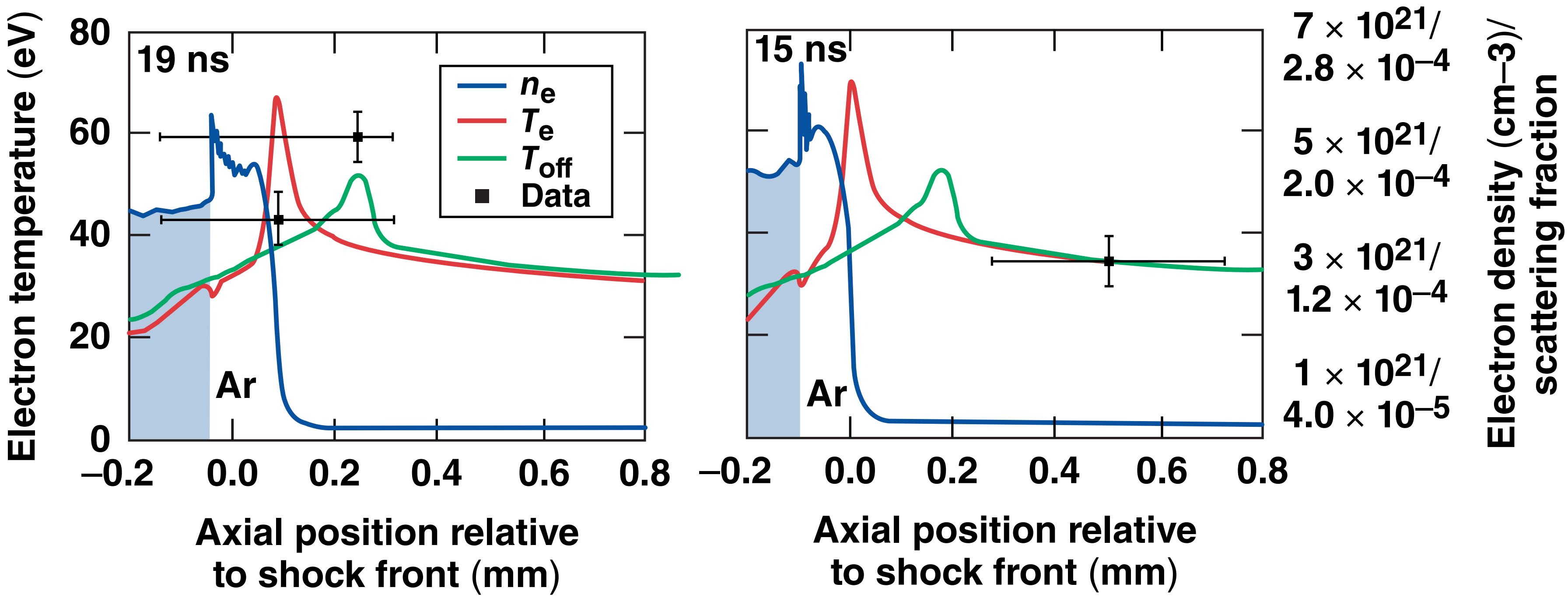
$$I_{sc} = C I_0 \left[V_s S_s^{tot}(k) + V_l S_l^{tot}(k) \right]$$

$$\frac{V_l}{V_{tot}} = \left(\frac{I_{hot}}{I_{cold}} - 1\right)\left(\frac{S_l^{tot}}{S_s^{tot}} - 1\right)$$

The correct modeling of the liquid content at high pressures in carbon-rich planets has implications on the predicted magnetic fields.

A. Pelka *et al*., Phys. Rev. Lett. 105, 265701 (2010).

The electron temperature in the shocked region is higher than in the radiative precursor



The predictive capability of radiative shock waves is validated in the laboratory with XRTS and x-ray radiography.

E21085

*A. J. Visco, *et al.,* PRL. 108, 145001 (2012).

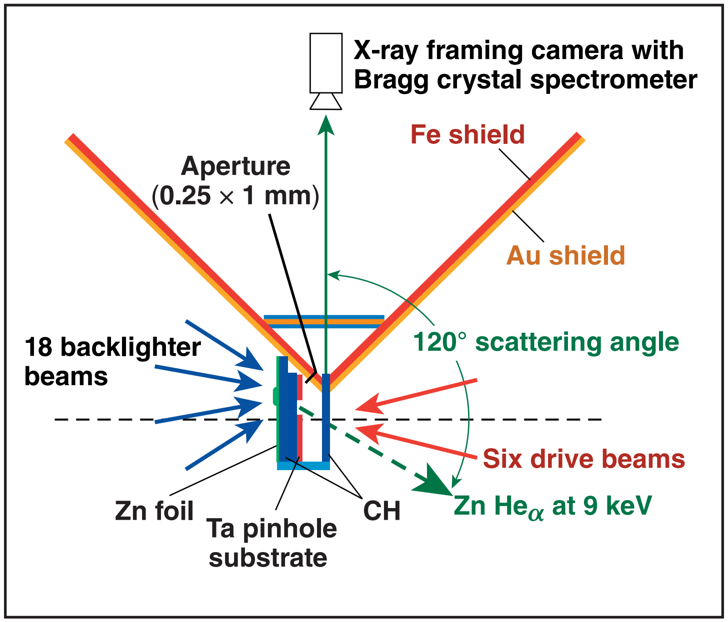
A spatial resolution of ~10 μm for spectral measurments would greatly benefit many XRTS experiments

- General-use spectral imaging diagnostic
- Development of brighter backlighters
- Improve accuracy of target alignment
  - install foreground lighting on OMEGA Target Viewing System

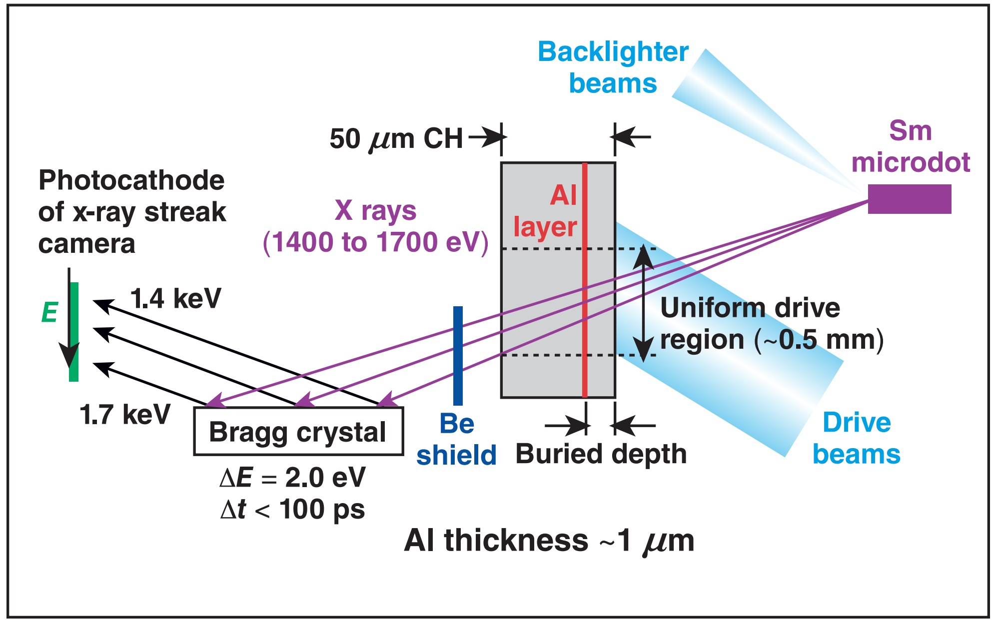
OLUG could help to develop specifications for next-generation XRTS measurements on OMEGA.

XRTS is a noninvasive probe

XRTS experiment[1]



X-ray absorption spectroscopy[2]



Compared to absorption spectroscopy, no tracer layers are required for XRTS.

[1]H. Sawada, S. P. Regan *et al.*, Phys. Plasmas 14, 122703 (2007).
[2]H. Sawada, S. P. Regan *et al.*, Phys. Plasmas 16, 052702 (2009).

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