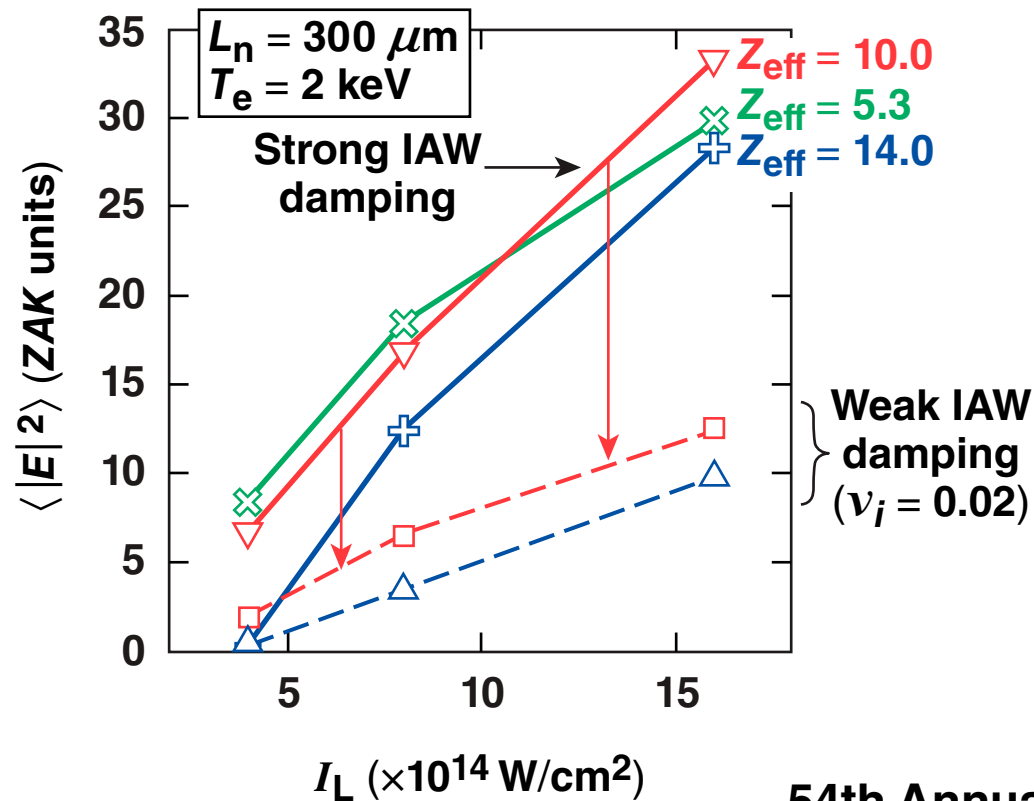


Mitigating Two-plasmon Decay Hot-electron Generation Through the Modification of Langmuir and Ion-acoustic Dissipation in Directly Driven Targets



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Two-plasmon–decay (TPD) preheat can be reduced through the manipulation of the collisional and Landau damping of Langmuir and ion-acoustic waves



- **Langmuir wave (LW) collisional damping has an impact on the growth rate/thresholds [in addition to the hydrodynamic variables ($\nabla n_e, T_e$)]**
 - importance increases with the scale length
- **Nonlinear saturation is sensitive to plasma composition ($\langle Z \rangle, T_i, \langle Z \rangle^2$) via the ion-acoustic damping rate**
 - predictions with the 2-D code **ZAK***
 - hot-electron predictions with the 2-D code **QZAK****
 - suggestions for planar experiments

* D. F. DuBois *et al.*, Phys. Rev. Lett. **74**, 3983 (1995);
D. A. Russell and D. F. DuBois, Phys. Rev. Lett. **86**, 428 (2001).
** D. A. Russell *et al.*, presented at the 42nd Annual Anomalous
Conference, Key West, FL, 25–29 June 2012.

Collaborators



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The ZAK/QZAK model of TPD is used to predict linear instability and nonlinear saturation caused by density fluctuations and quasilinear diffusion

- “Extended” Zakharov equations used in ZAK*

$$\nabla \cdot \left[D_{LW} - \omega_0^2 (\delta n + \delta N) / n_0 \right] E = \left(e / 4 m_c \right) \nabla \cdot \left[\nabla (E_0 \cdot \bar{E}) - E_0 \nabla \cdot \bar{E} \right] + S_E$$

Gradient threshold

$$D_{IAW} \delta n = \nabla^2 \left(|E|^2 + \frac{1}{4} |E_0|^2 \right) / (16 \pi M_i) + S_{\delta n}$$

TPD source term

Dispersion relations for LW's
and ion acoustic waves (IAW's)

$$D_{LW} = \left[2i\omega_{p0} (D_t + \nu_e^*) + 3\nu_e^2 \nabla^2 \right]$$

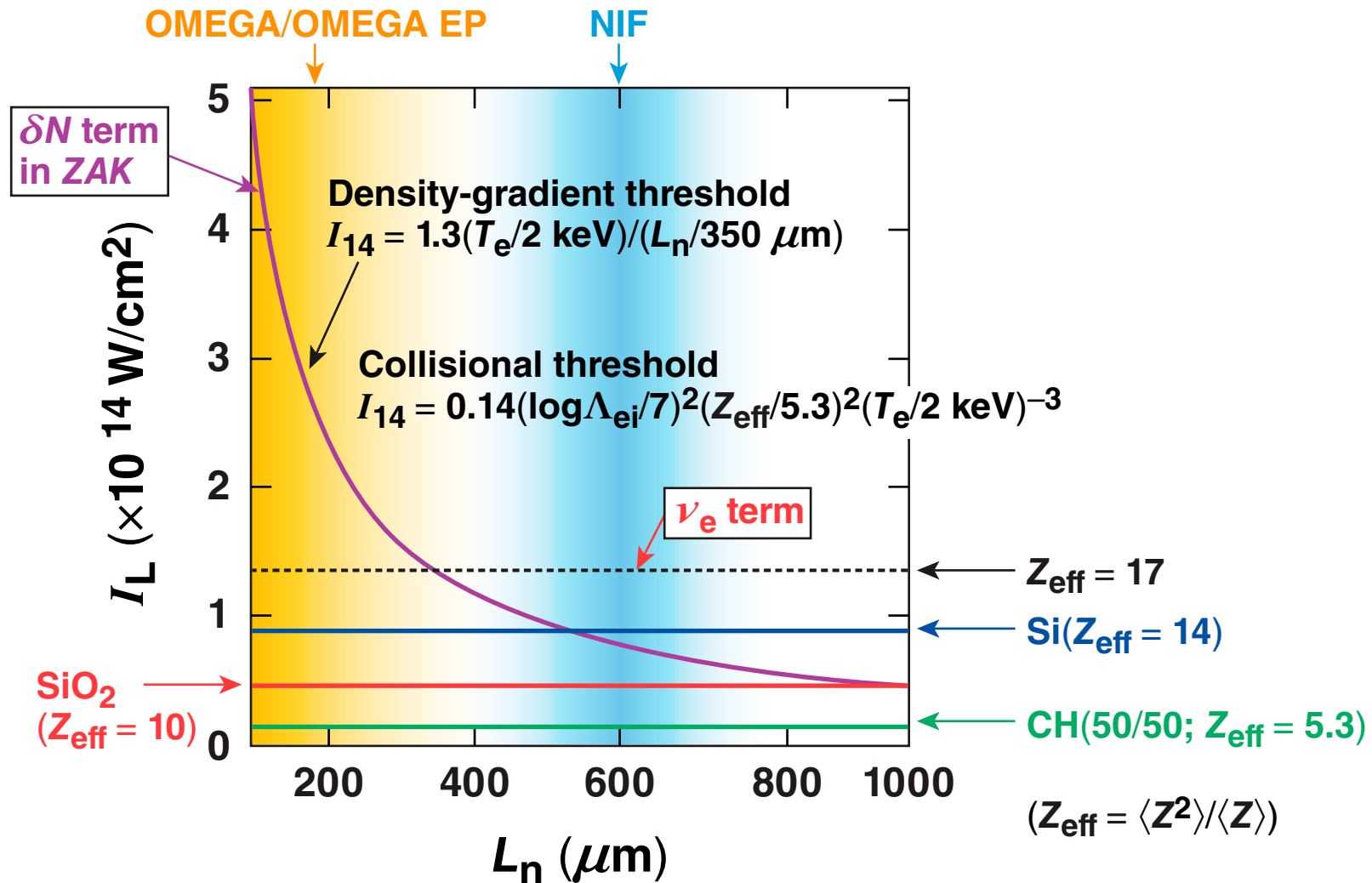
Collisional threshold

$$D_{IAW} = (D_t^2 + 2\nu_i^* D_t - c_s^2 \nabla^2)$$

- The diffusion equation $\frac{\partial \langle f_e \rangle}{\partial t} + \frac{\partial}{\partial \vec{v}} \cdot \left(D(\vec{v}) \cdot \frac{\partial \langle f_e \rangle}{\partial \vec{v}} \right) = \sigma (\langle f_e \rangle - f_M)$

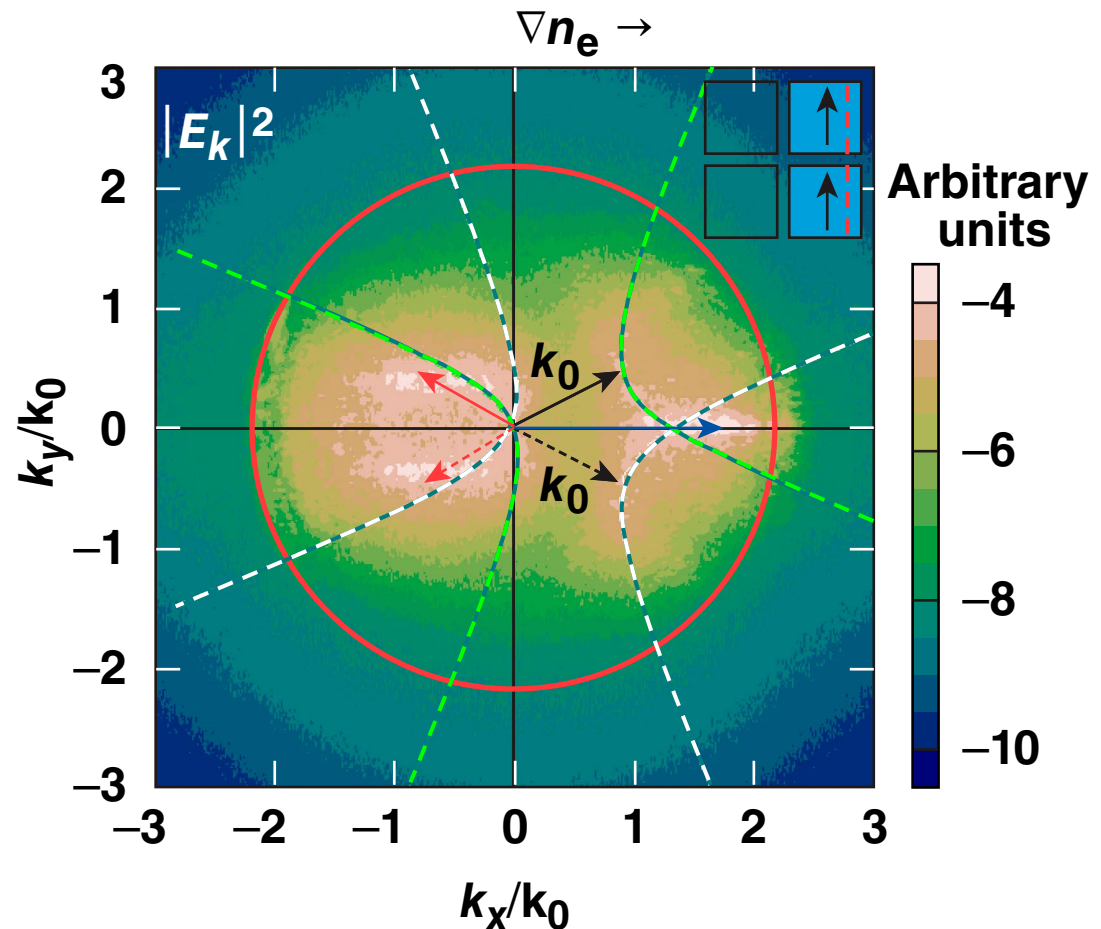
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The collisional threshold can be made to exceed the gradient threshold (for Si at ignition scale)



Two-dimensional calculations assume two overlapped plane electromagnetic (EM) waves* polarized in their plane of incidence

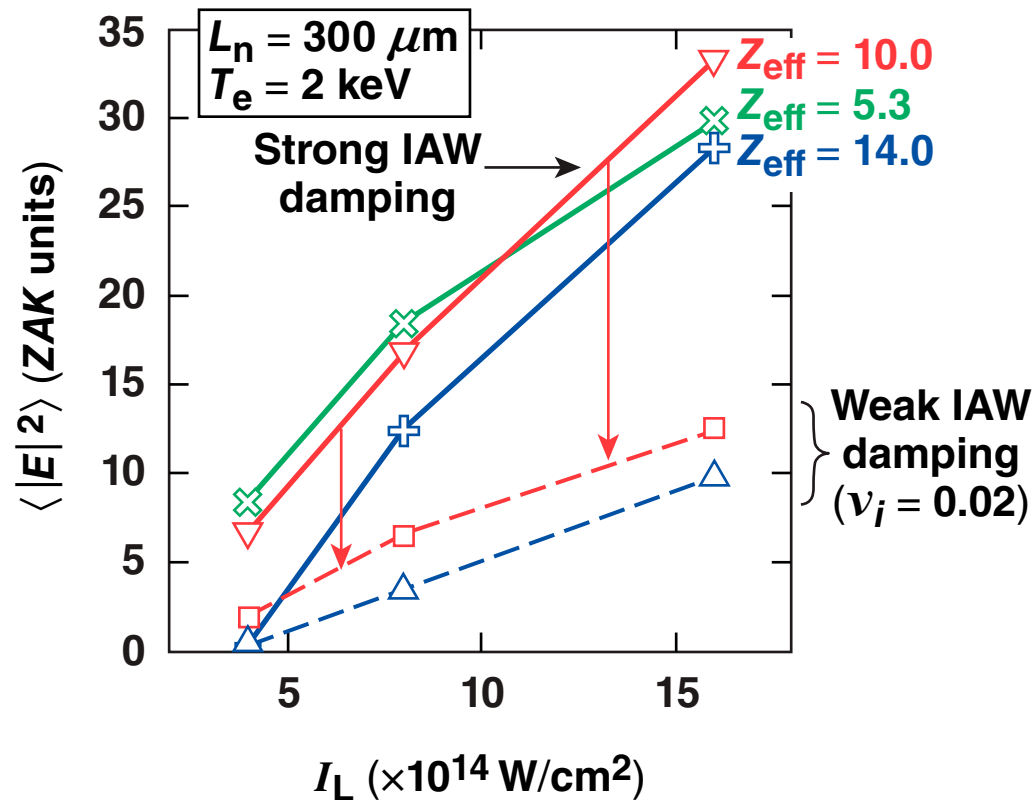
- This has the feature that TPD is driven unstable convectively even when the single-beam intensities are below threshold
- $\langle |E|^2 \rangle$ is computed at saturation (ZAK)
- Hot-electron distribution (QZAK)



*D. T. Michel *et al.*, presented at the 42nd Annual Anomalous Conference, Key West, FL, 25–29 June 2012.

Decreasing the IAW damping leads to a dramatic reduction in the level of electrostatic fluctuations

- Increased collisional damping of LWs leads to a lower saturated level of electrostatic fluctuations* for $Z_{\text{eff}} = 14$

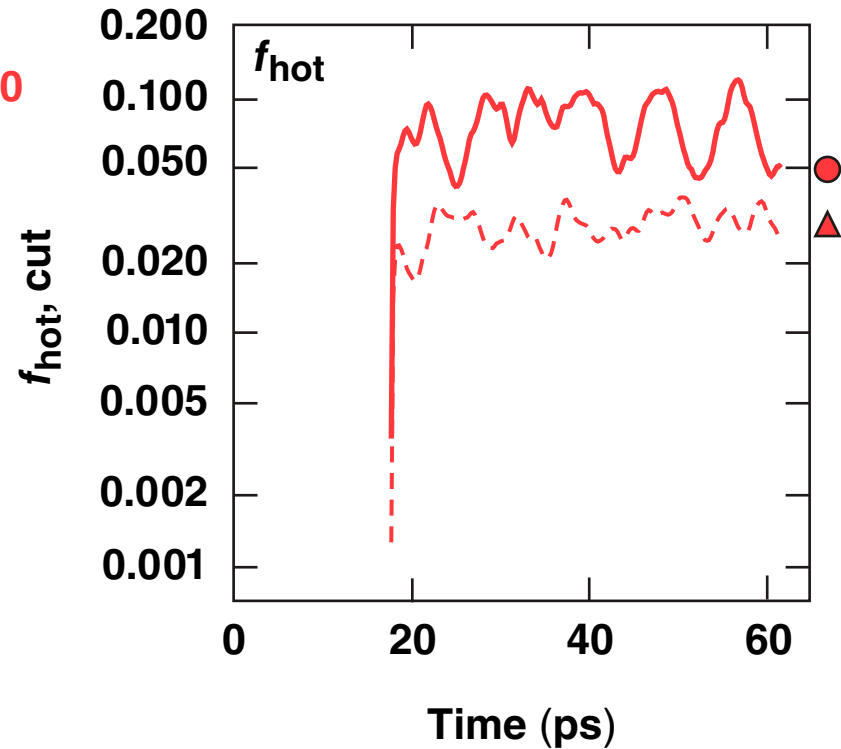
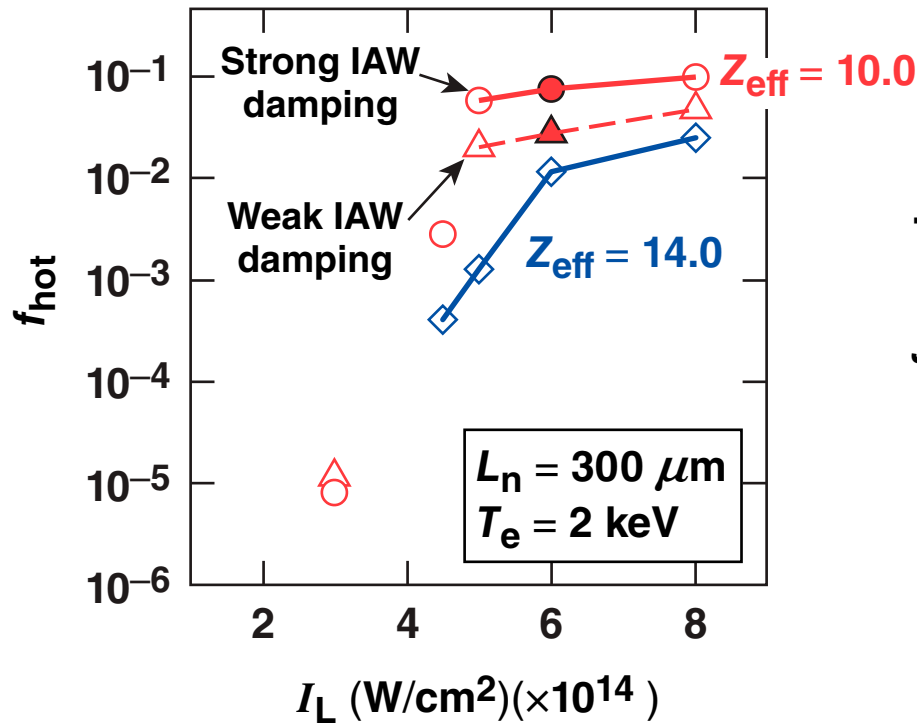


QZAK calculations

QZAK predicts less hot-electron production for a plasma with weakly damped ion-acoustic waves



A factor 2 to 3 reduction in f_{hot}

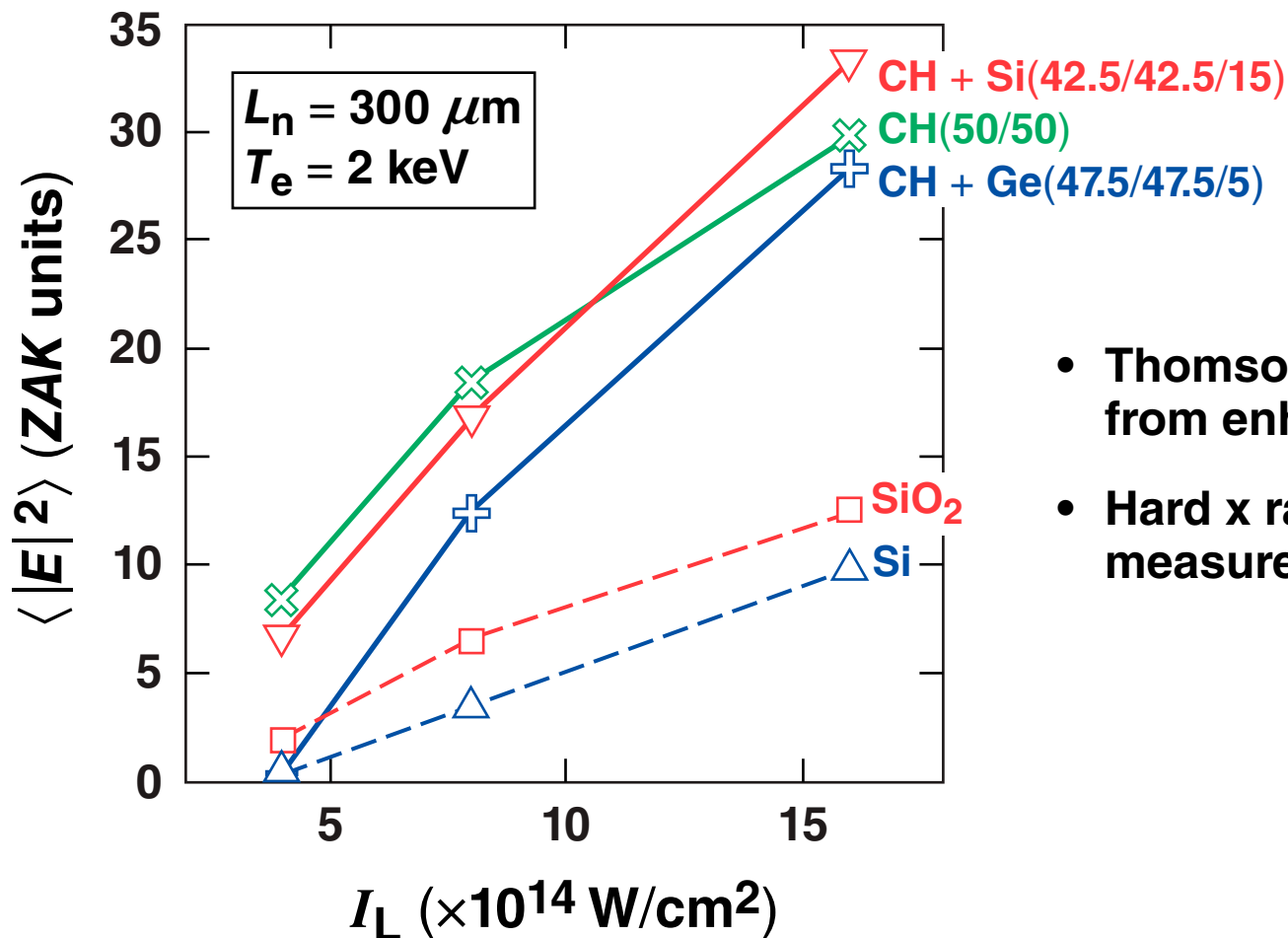


Ion-acoustic damping can be manipulated by modifying the plasma composition



- Ion Landau damping decreases with ZT_e / T_i ($c_s \gg v_{ti}$)
- Electron Landau damping in IAW's is always weak ($v_{te} \gg c_s$)
- Light ions (e.g., hydrogen) can increase the damping rate
- The multi-ion dispersion relation is solved by finding the most weakly damped mode*
- Ion-ion collisions complicate the matter
- Part of the spectrum will be collisionally damped because of ion viscosity ($k\lambda_{ij} < 10$)

These effects can be investigated experimentally in planar targets on OMEGA/OMEGA EP†



- Thomson scattering from enhanced IAW*
- Hard x rays can be measured

† S. X. Hu *et al.*, presented at the 42nd Annual Anomalous Conference, Key West, FL, 25–29 June 2012.

* R. K. Follet *et al.*, presented at the 42nd Annual Anomalous Conference, Key West, FL, 25–29 June 2012.

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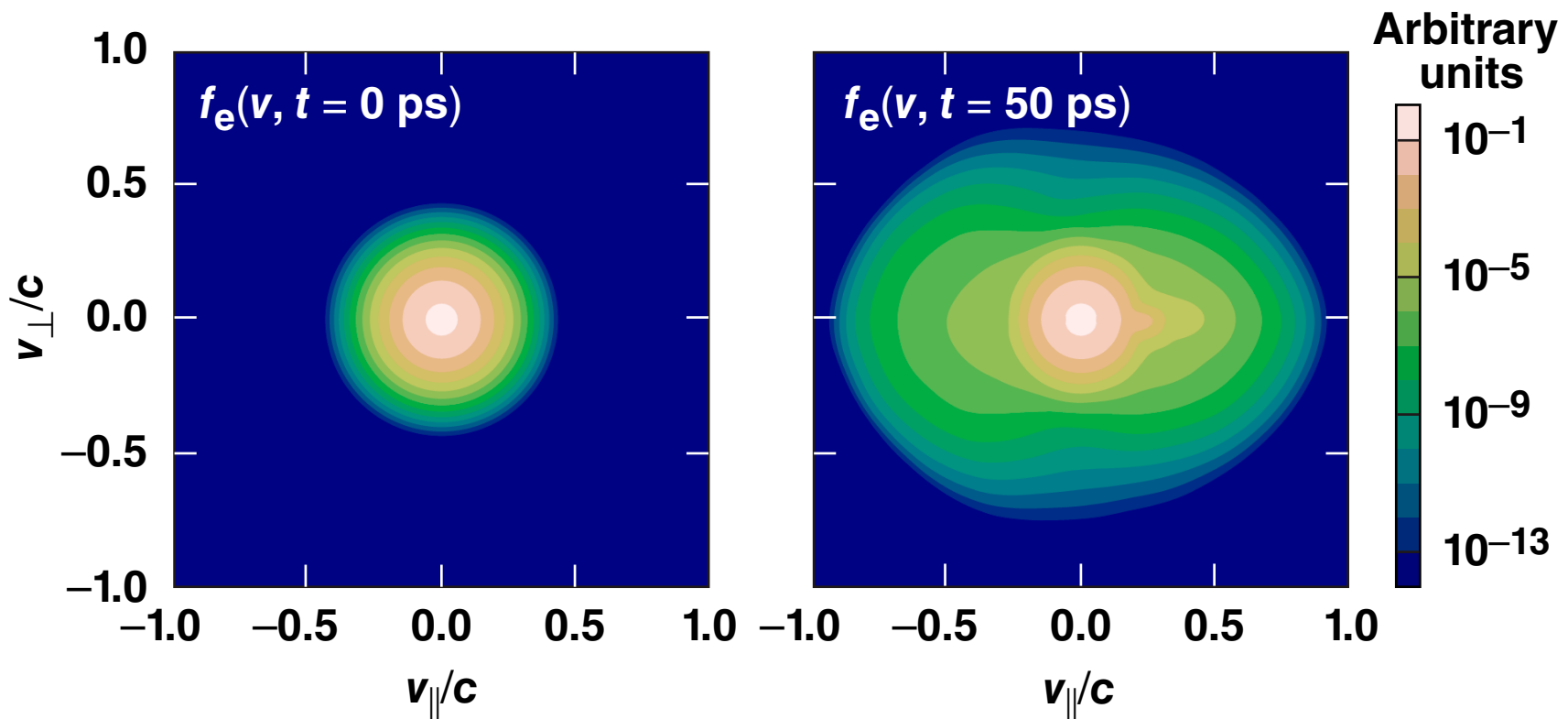
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QZAK* is an extension of ZAK to self-consistently include hot-electron generation in the quasilinear approximation



- The diffusion equation
$$\frac{\partial \langle f_e \rangle}{\partial t} + \frac{\partial}{\partial \vec{v}} \cdot \left(\mathbf{D}(\vec{v}) \cdot \frac{\partial \langle f_e \rangle}{\partial \vec{v}} \right) = \sigma (\langle f_e \rangle - f_M)$$



*D. A. Russell *et al.*, D. A. Russell *et al.*, presented at the 42nd Annual Anomalous Conference, Key West, FL, 25–29 June 2012.; H. X. Vu *et al.*, presented at the 42nd Annual Anomalous Conference, Key West, FL, 25–29 June 2012.