Planar Laser–Plasma Interaction Experiments at Direct-Drive Ignition-Relevant Scale Lengths at the National Ignition Facility





Summary

Planar experiments at the National Ignition Facility (NIF) have investigated laser-plasma interaction (LPI) hot-electron production at direct-drive ignition-relevant coronal conditions



- NIF planar experiments achieve ignition-relevant scale lengths $(L_n \sim 400 \text{ to } 700 \ \mu\text{m})$ and electron temperatures $(T_e \sim 4 \text{ to } 5 \text{ keV})$
- The fraction of laser energy converted to hot electrons increased with laser intensity from $f_{hot} \sim 0.5\%$ to 2.3%—from 6 to 15×10^{14} W/cm²—while T_{hot} was ~50 keV
- Stimulated Raman scattering (SRS) is inferred to be the dominant hot-electron source at these conditions
- The use of Si ablators reduces the observed SRS, f_{hot} , and T_{hot} relative to CH, using small angle beams

These results indicate a viable ignition-design space for direct drive.





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- Motivation for direct-drive planar LPI experiments on the NIF and platform development
- Hot-electron and scattered-light results: dominance of SRS
- LPI/hot-electron preheat mitigation strategies and future work







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Motivation

The National Direct-Drive Program includes OMEGA and NIF experiments to study direct-drive physics



Laser coupling, preheat, imprint, and hydrodynamically scaled implosions

Laser coupling, preheat, and imprint at the MJ scale



TC10256w

Motivation

Hot-electron preheat is a potential concern for direct-drive-ignition designs



Limit of ~0.15% laser energy into fuel preheat + angular divergence \rightarrow limit of ~0.7% laser energy into hot electrons generated.

*B. Yaakobi et al., Phys. Plasmas 20, 092706 (2013).



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Motivation

Direct-drive (DD)-ignition designs predict long density scale lengths and high electron temperatures under which LPI may occur



One-dimensional simulated plasma conditions for an igniting direct-drive design



Experiments must be performed at these conditions to understand LPI and assess hot-electron levels at the NIF/ignition scale.

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Currently, ignition-relevant coronal plasma conditions can only be achieved in NIF planar experiments

Two-dimensional DRACO simulated plasma conditions at $n_c/4$

| | NIF ignition designs | Ongoing NIF planar experiments | Ongoing NIF implosions | OMEGA implosions |
|----------------------|------------------------------------|--------------------------------------|---------------------------|------------------------------------|
| L_{n} (μ m) | 600 | 400 to 700 | 360 | 150 |
| T _e (keV) | 3.5 to 5 | 3 to 5 | 3.2 | 2.8 |
| $I_{L} (W/cm^2)$ | $(6 \text{ to } 8) \times 10^{14}$ | $(6 \text{ to } 15) \times 10^{14}$ | $5 	imes 10^{14}$ | $(5 \text{ to } 7) \times 10^{14}$ |

Note: incident laser intensity is $\sim 2 \times$ larger than intensity at $n_c/4$ at ignition-relevant conditions because of absorption



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Two initial planar experiments were performed on the NIF to constrain plasma conditions



Cross-beam energy transfer does not have a strong influence on conditions at $n_c/4$.

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The isoelectronic ratio^{*} of the Mn/Co K-shell emission lines is used to infer $T_e = 4.6 \pm 1.1$ keV at $n_c/4$



Based on modeling, discrepancy can be partially explained by self-heating of the microdot.

*R. Marjoribanks et al., Phys. Rev. A <u>46</u>, R1747 (1992).



TC12489g

In subsequent experiments at higher laser intensity, the wavelength of $\omega/2$ emission was used to infer $T_{\rm e} \sim 4.5$ keV at $n_{\rm c}/4$



These measurements match DRACO predictions of ignition-relevant $T_e = 4.5$ keV.

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W. Seka et al., Phys. Fluids 28, 2570 (1985).



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Three experiments explored the scaling of hot-electron properties with laser intensity at ~500- μ m scale lengths and ~4-keV temperatures





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^{*}T. J. B. Collins *et al.*, Phys. Plasmas <u>19</u>, 056308 (2012).

^{**}A. A. Solodov, this conference.

Considering overlapped laser intensities, these experiments are well above threshold for two-plasmon decay (TPD) and SRS

• Absolute instability thresholds for a single beam at normal incidence

$$\frac{\text{TPD}^{*}}{\text{TPD}^{*}} \qquad \begin{array}{l} I_{14,\text{thr},\text{TPD}} = 230 \ T_{e,\text{keV}} / L_{n,\mu\text{m}} \\ \rightarrow \eta_{\text{TPD}} = I_{\text{overlapped}} / I_{\text{thr},\text{TPD}} \sim 3 \text{ to } 8 \end{array}$$

$$\underline{SRS^{**}} \qquad I_{14, \text{thr, TPD}} = 2377 / (L_{n, \mu m})^{4/3} \\ \rightarrow \eta_{SRS} = I_{\text{overlapped}} / I_{14, \text{thr, SRS}} \sim 10 \text{ to } 25$$

• These experiments overlapped 64 beams (16 NIF quads)

**C. S. Liu, M. N. Rosenbluth, and R. B. White, Phys. Fluids <u>17</u>, 1211 (1974).



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^{*}A. Simon *et al.*, Phys. Fluids <u>26</u>, 3107 (1983).

Time-integrated hard x-ray data show f_{hot} (E_{hot}/E_{laser}) increases with laser intensity, while T_{hot} is constant



*f*_{hot} results are encouraging for direct-drive–ignition designs and constrain design space.

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Scattered-light measurements to identify the hot-electron source were optimized by orienting the target normal to the optical diagnostics



as have been observed previously on OMEGA.*

*W. Seka et al., Phys. Rev. Lett. <u>112</u>, 145001 (2014). **Near-backscatter imager



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The optical spectrum indicates a sharp, red-shifted $\omega/2$ feature as well as SRS at shorter wavelengths





The $\omega/2$ feature observed on the NIF is in contrast to that observed on OMEGA, which showed both blue- and red-shifted $\omega/2$, and is attributed to TPD



On the NIF, the observed $\omega/2$ emission is attributed to absolute SRS, although the presence of TPD cannot yet be ruled out.



W. Seka et al., WeO-5, this conference.

The dominance of SRS at the NIF scale is explained by evaluating the absolute thresholds of SRS* versus TPD**



*A. Simon et al., Phys. Fluids <u>26</u>, 3107 (1983).

**C. S. Liu, M. N. Rosenbluth, and R. B. White, Phys. Fluids <u>17</u>, 1211 (1974).



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Further evidence that SRS is the hot-electron source can be obtained by inferring the total SRS produced



Diagnostic views are limited: absolute energy measurements are available at two locations

Extrapolating to the total SRS requires understanding the SRS mechanism.



Observation of SRS at multiple locations from different drive beams provides strong evidence of SRS sidescattering



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This observation is explained by tangential SRS sidescatter, which allows for SRS observation at large angles and wavelength independent of drive-beam angle



Knowledge of SRS mechanisms—absolute SRS ($\omega/2$) and sidescattered SRS—allows for extrapolation to the total SRS generated



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W. Seka et al., WeO-5, this conference.

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The use of a Si layer in the ablator has been proposed as a means of reducing hot-electron generation from near-quarter-critical LPI



Si ablators produce shorter scale lengths, higher electron temperatures, and more collisional damping in order to reduce LPI near $n_c/4$.

*V. N. Goncharov et al., Phys. Plasmas 21, 056315 (2014).



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Use of a Si ablator in planar experiments causes a reduction in observed SRS driven by small-angle beams, relative to CH





In comparison to the CH target, the Si target produced an ~30% lower f_{hot} and an ~15 keV lower T_{hot} for small-angle-beam drive



Future Work

The new optical Thomson-scattering (OTS) diagnostic on the NIF will be used to probe $3\omega/2$ emission and to measure plasma conditions



These experiments will assess the presence of TPD, confirm plasma conditions, and develop a platform for eventual use of 5ω Thomson scattering (TS) to probe plasma waves.



Future Work

Coupling of hot electrons to an implosion will be assessed by measuring their angular divergence using buried Mo layers



These NIF experiments will determine the relationship between hot-electron generation and the expected level of preheat in ignition-relevant implosions.



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^{*}B. Yaakobi et al., Phys. Plasmas 20, 092706 (2013).

Summary/Conclusions

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Appendix





Despite varying plasma conditions in a NIF experiment using a ramped laser pulse, a similar SRS spectrum is obtained



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SRS depends strongly on inner-beam laser intensity in CH ablators



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The use of an Si ablator in planar experiments causes a reduction in observed SRS driven by small-angle beams, relative to CH



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For an Si ablator, the SRS intensity at 23° depends strongly on inner-beam laser intensity, but SRS observed at 50° is still minimal





Compared to the CH target, the Si target produced an ~40% lower f_{hot} and an ~15-keV lower T_{hot} for small-angle-beam drive





Long-scale-length (>500- μ m), high-temperature (>3-keV) coronal plasma conditions are predicted by 2-D DRACO simulations



*A. A. Solodov *et al.*, presented at the Ninth International Conference on Inertial Fusion Sciences and Applications (IFSA 2015), Seattle, WA, 20–25 September 2015

**A. Simon et al., Phys. Fluids 26, 3107 (1983);

C. S. Liu, M. N. Rosenbluth, and R. B. White, Phys. Fluids 17, 1211 (1974).





Hard x-ray and Mo K $_{\alpha}$ emission caused by LPI-generated hot electrons were observed





Time-integrated hard x-ray spectra indicate $T_{hot} \sim 45\pm5$ keV, $f_{hot} \sim 1\%$ for both experiments



The beam angle of incidence did not have a strong effect on f_{hot} and T_{hot} .

*A. A. Solodov *et al.*, presented at the Ninth International Conference on Inertial Fusion Sciences and Applications (IFSA 2015), Seattle, WA, 20–25 September 2015.

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This observation is explained by tangential SRS sidescatter, which allows for SRS observation at large angles and wavelength independent of drive-beam angle



Tangential sidescatter exit angle does not depend on the incidence angle



P. A. Michel *et al.*, WeO-4, this conference. *Full-aperture backscatter station