

# Comparisons Between Ray- and Wave-Based Calculations of Cross-Beam Energy Transfer

R. K. FOLLETT, D. H. EDGELL, D. H. FROULA, V. N. GONCHAROV, I. V. IGUMENSHCHEV, J. G. SHAW, and J. F. MYATT

University of Rochester, Laboratory for Laser Energetics

J. W. BATES, K. OBENSCHAIN, and J. L. WEAVER

Naval Research Laboratory

## Summary

A 3-D wave-based model has been developed to understand the physics of cross-beam energy transfer (CBET) in an inhomogeneous plasma



- Detailed CBET calculations are used to test ray-based CBET models that are implemented in hydrodynamics codes
- The comparisons generally highlight the accuracy of ray-based models
- Discrepancies between the models are found related to beam speckle and polarization smoothing

LPSE solves the time-envelope Maxwell's equations coupled to a linearized plasma response



Maxwell's equations (time envelope)  $\vec{E} = \Re\{\vec{E}(\vec{x}, t)\exp(-i\omega_0 t)\}$

$$\frac{2i\omega_0}{c^2} \frac{\partial}{\partial t} \vec{E} + \nabla^2 \vec{E} - \nabla(\nabla \cdot \vec{E}) + \frac{\omega_0^2}{c^2} \epsilon(\omega; \vec{x}, t) \vec{E} = 0$$

$$E(\omega_0; \vec{x}, t) = 1 - \frac{\omega_{pe}^2(\vec{x}, t)}{\omega_0(\omega_0 + i\nu_{ei})}$$

$$\omega_{pe}^2 = \frac{4\pi e^2 n_e(\vec{x}, t)}{m_e}$$

Plasma response

$$\left[ \partial_t + \vec{U}_0(\vec{x}) \cdot \nabla \right] \left( \frac{\delta n}{n_0} \right) = -W$$

$$\left[ \partial_t + \vec{U}_0(\vec{x}) \cdot \nabla + 2\nu_{IAW} \right] W = -\nabla^2 \left[ c_s^2 \left( \frac{\delta n}{n_0} \right) + \frac{e^2}{4m_e \omega_0^2} |\vec{E}|^2 \right]$$

$$W \equiv \nabla \cdot \delta \vec{U}$$

$$n_e \equiv n_0(\vec{x}) + \delta n(\vec{x}, t)$$

$$\vec{U} \equiv \vec{U}_0(\vec{x}) + \delta \vec{U}(\vec{x}, t)$$

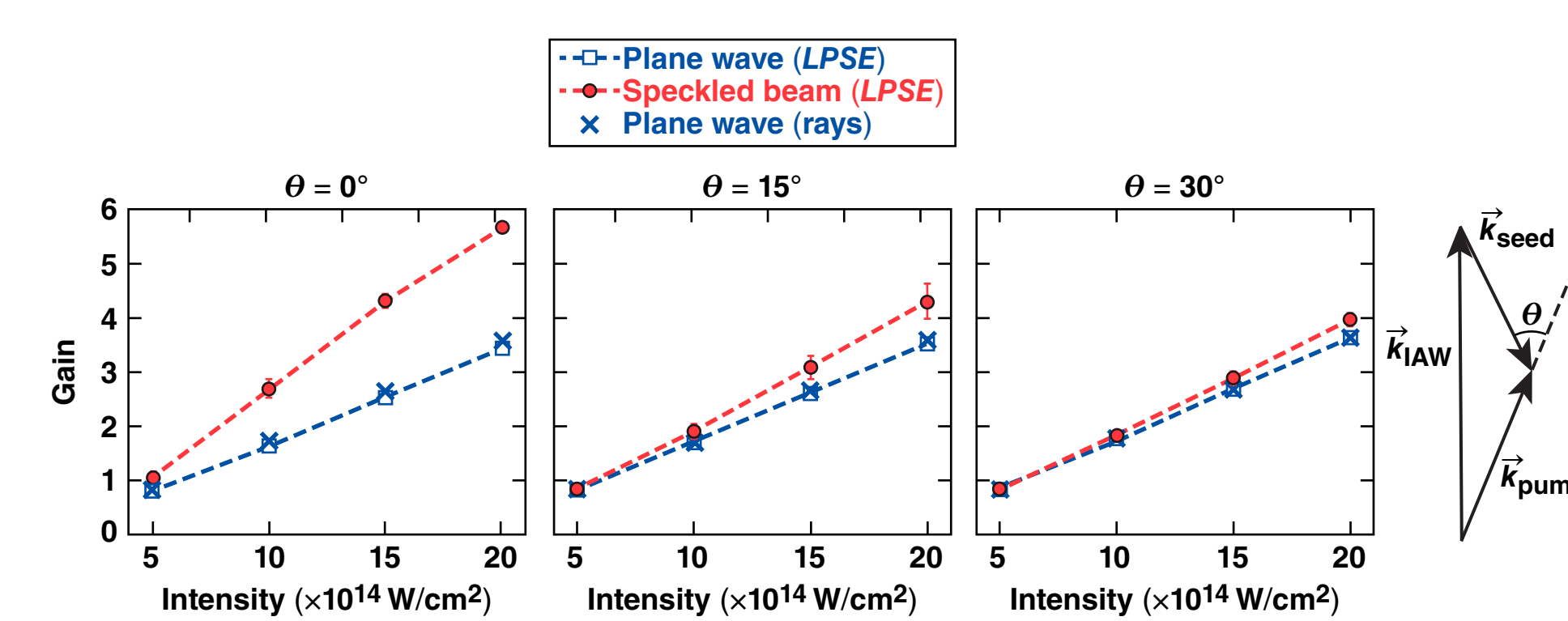
E26192

The CBET gain is sensitive to beam speckle for gains greater than ~1 and relative beam angles of less than ~30°



$$\text{Gain} \equiv \log \left( \frac{\text{Seed energy out}}{\text{Seed energy in}} \right)$$

CBET gain versus pump intensity for various beam-relative beam angles

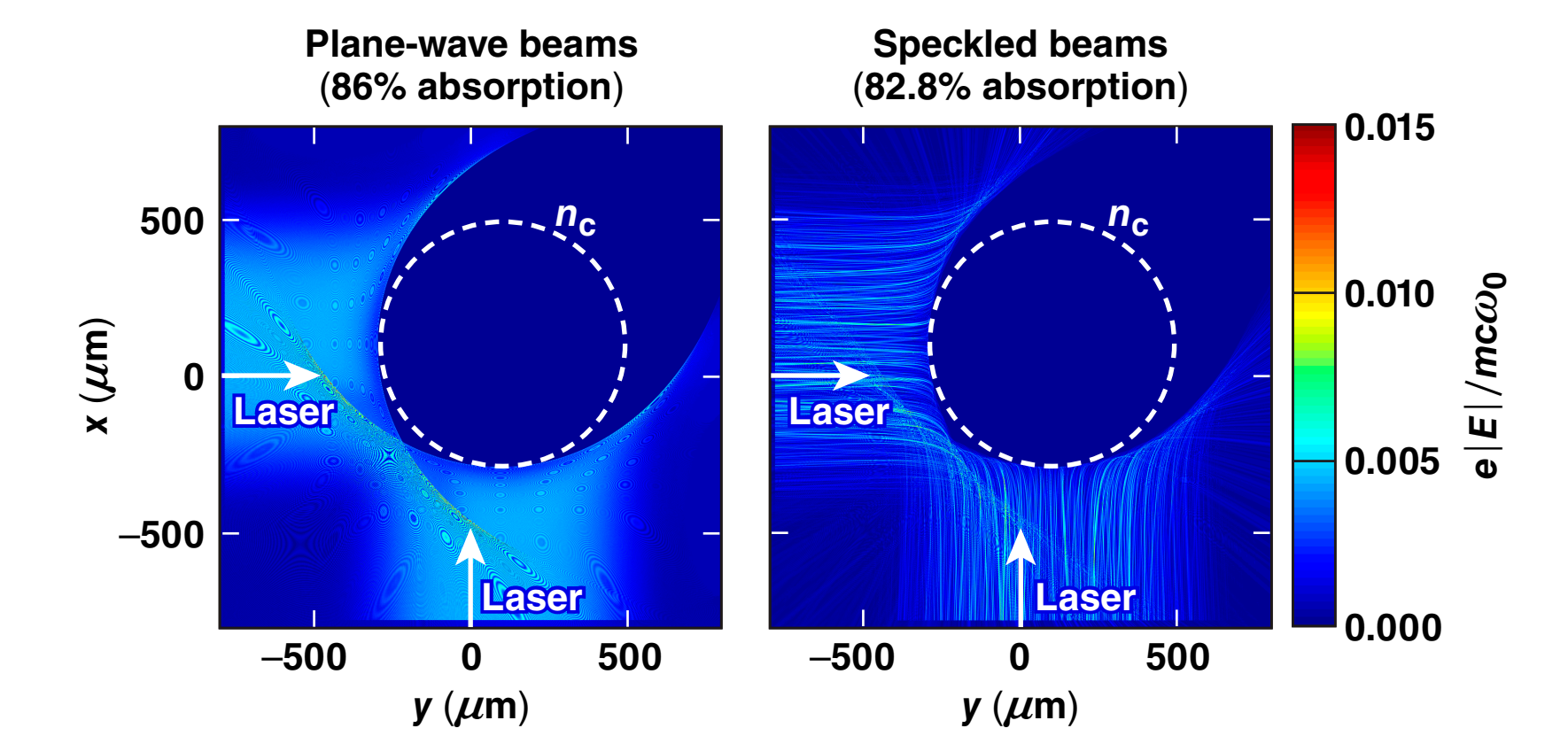


E26196

Speckled beams result in a modest decrease in laser absorption in OMEGA-scale, two-beam LPSE simulations at ICF-relevant plasma conditions



OMEGA-scale LPSE simulations ( $I_{\text{single beam}} = 2 \times 10^{14} \text{ W/cm}^2$ ) in a LILAC plasma profile (1.6 billion grid cells)

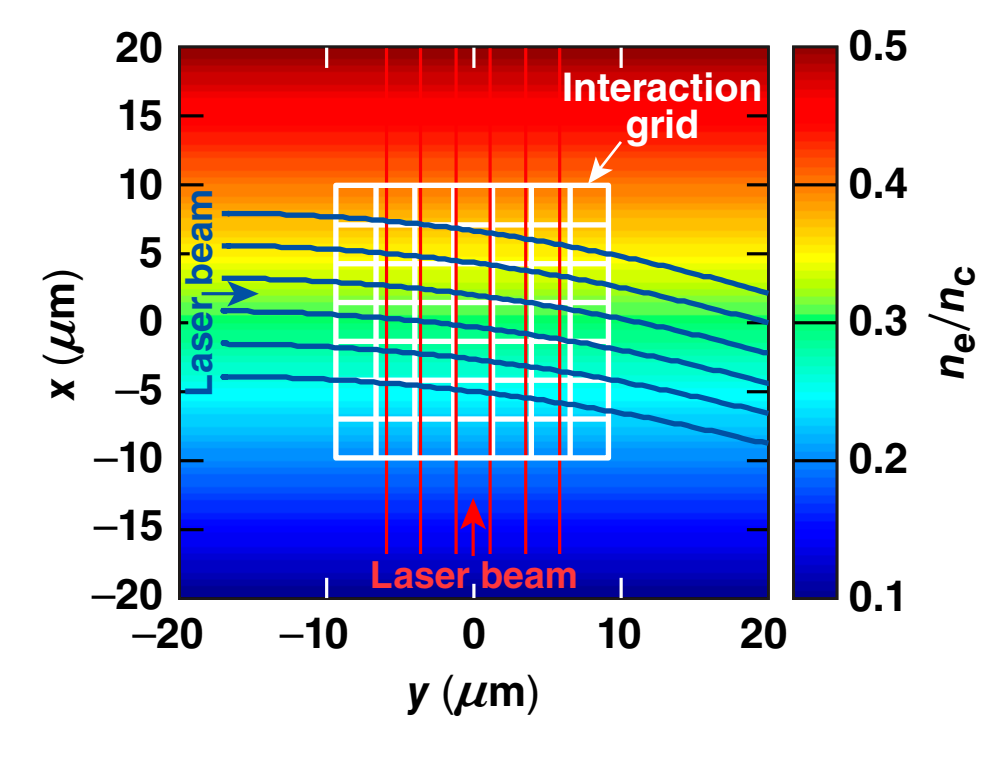


E26199

Ray-based models calculate CBET by considering pairwise interactions between rays



Schematic of ray-interaction calculation



CBET interaction between plane waves in a homogeneous plasma

$$\frac{dI_i}{ds} = \sum_j I_j L_{ij}^{-1}$$

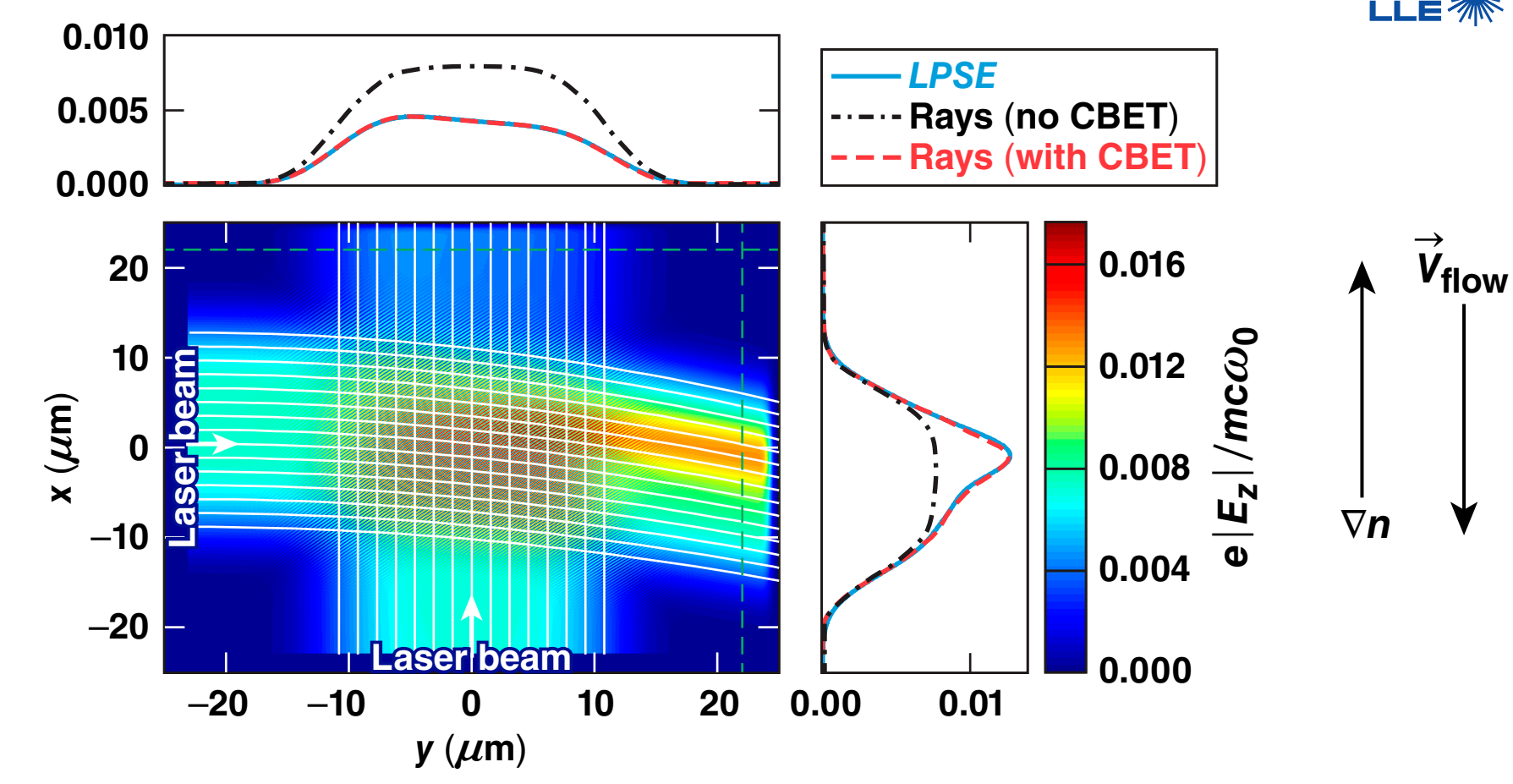
$$L_{ij}^{-1} = 5.88 \times 10^{-2} \frac{I_j \lambda}{T_e (1 + 3T_e / ZT_e)} \frac{n_s}{n_c} \frac{\omega_s}{\nu_i} P(\eta_{ij})$$

$$P(\eta) = \frac{\nu_i^2 \eta}{(\eta^2 - 1)^2 + \nu_i^2 \eta^2}$$

$$\eta_{ij} = \frac{\omega_j - \omega_i - (\vec{k}_j - \vec{k}_i) \cdot \vec{u}}{\omega_s}$$

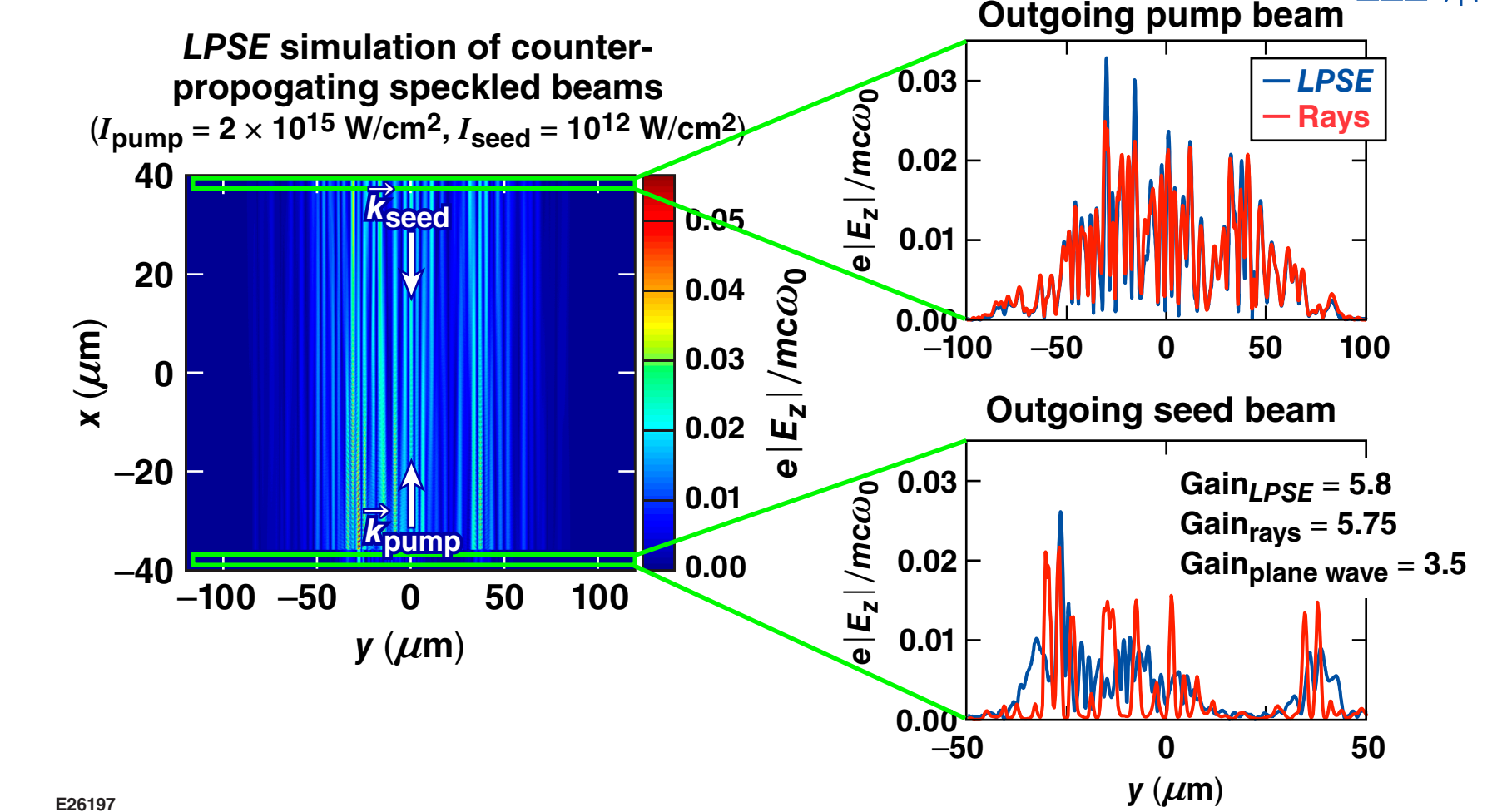
E26176

Ray- and wave-based CBET models give the same result in simple interaction geometries (plane-wave beams, no caustics)



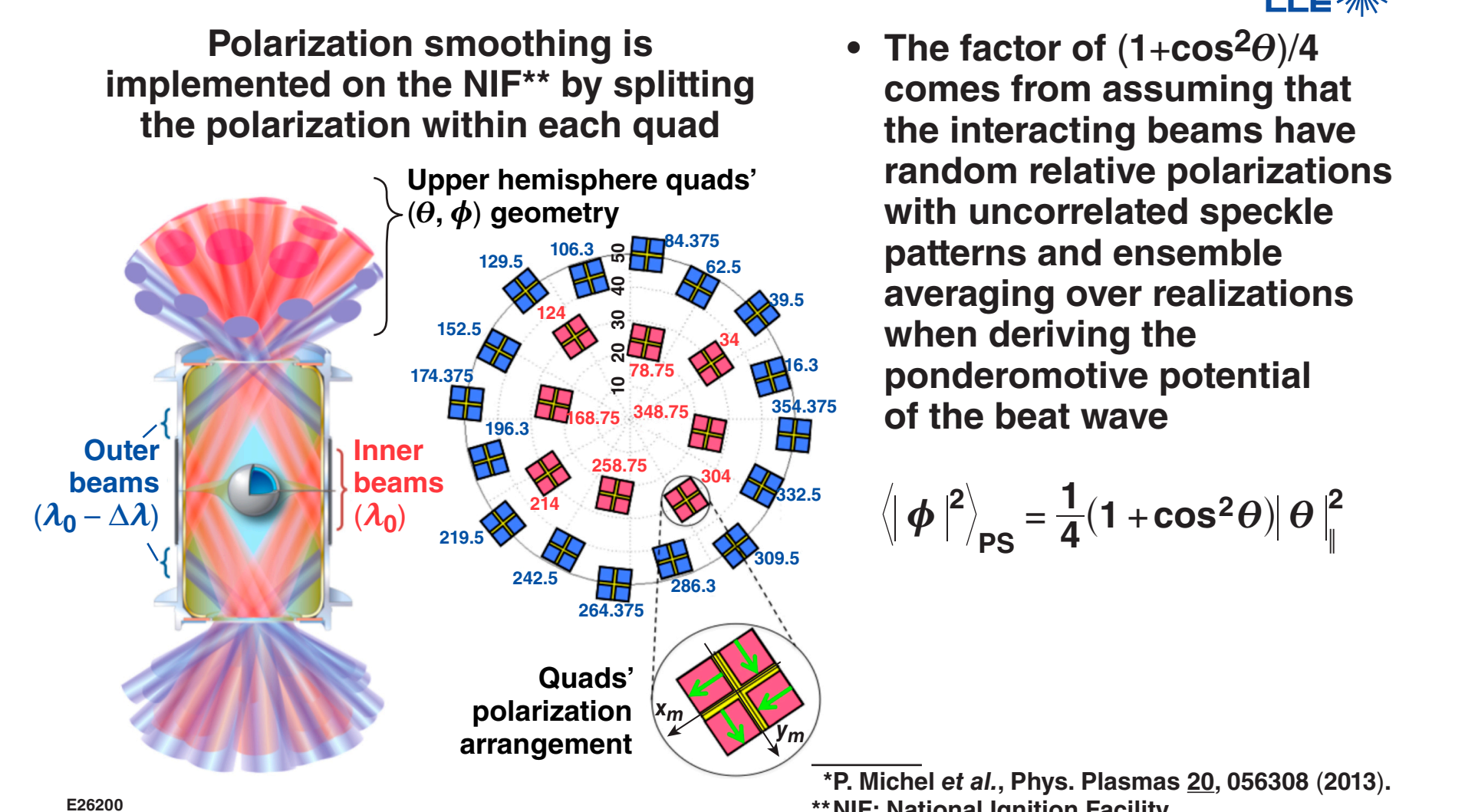
E26184

A good approximation to the average CBET between speckled beams can be obtained by using the linearity of Maxwell's equations to solve for the correct unperturbed field amplitudes in the ray-based calculation



E26187

Polarization smoothing is accounted for in ray-based CBET models by multiplying the gain coefficient by a factor of  $(1 + \cos^2\theta)/4$ \*



E26200

Ray-based CBET models make several approximations that are not always satisfied in inertial confinement fusion (ICF) applications

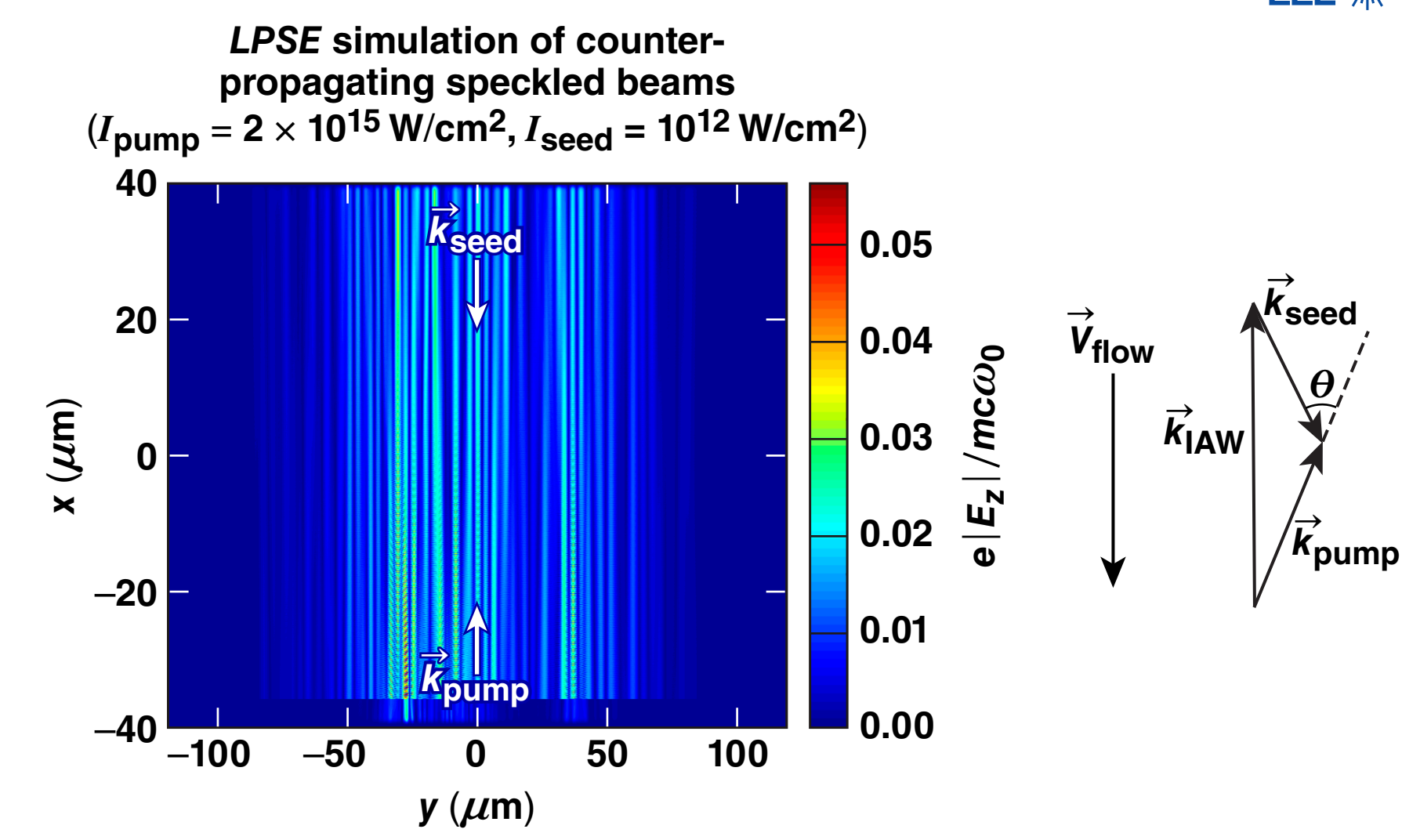


- Ion-acoustic waves (IAW's) ( $\delta n/n \ll 1$ )
- Steady-state convective gain
- Polarization-averaged coupling constant
- Pairwise coupling between beams
- Local plane-wave approximation
  - not valid for speckled beams or at caustics
- Eikonal approximation [Wentzel-Kramers-Brillouin (WKB), envelope]
  - breaks down at caustics

Approximations that are not made in LPSE

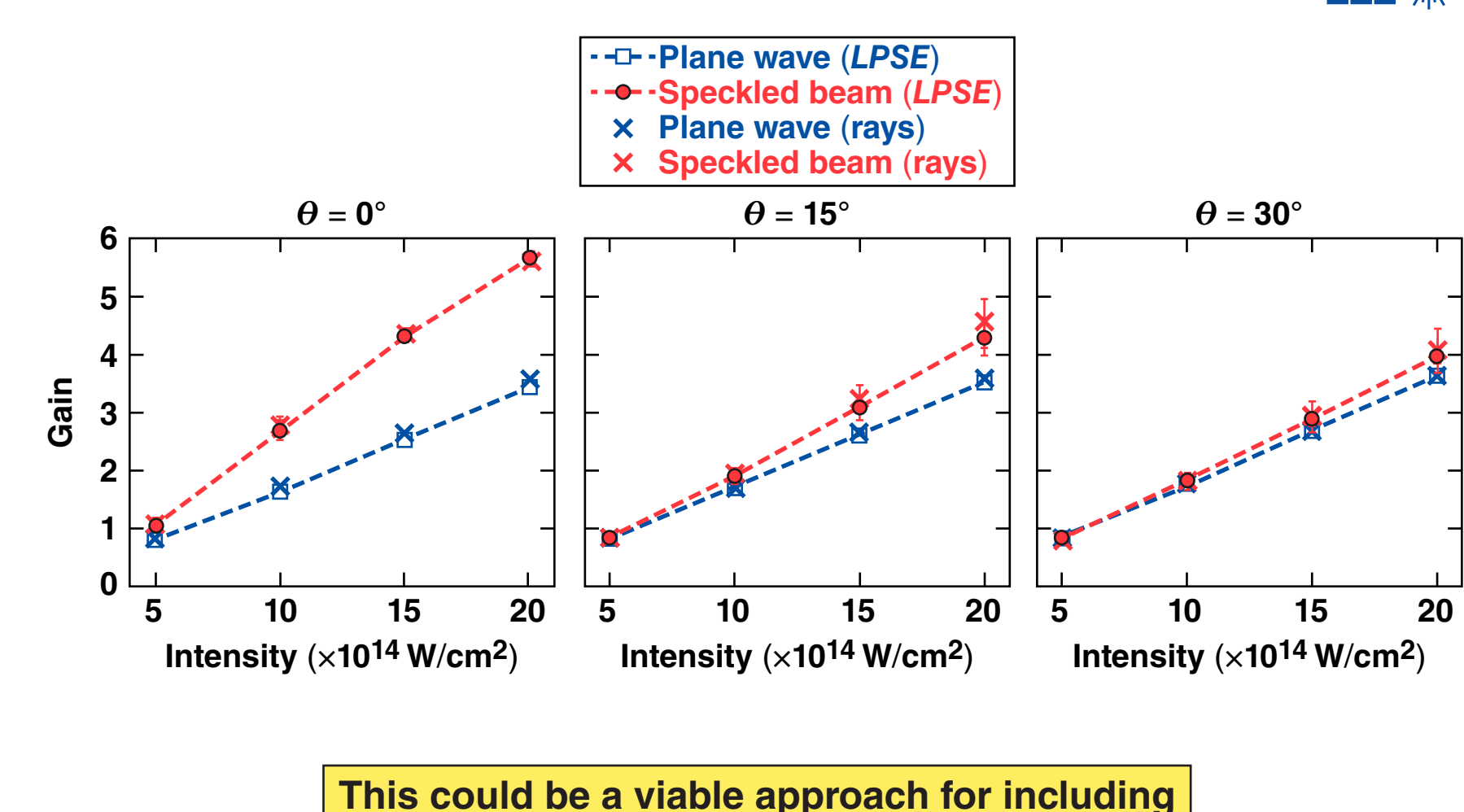
E26191

Speckled beams can transfer more energy than plane-wave beams with the same average intensity



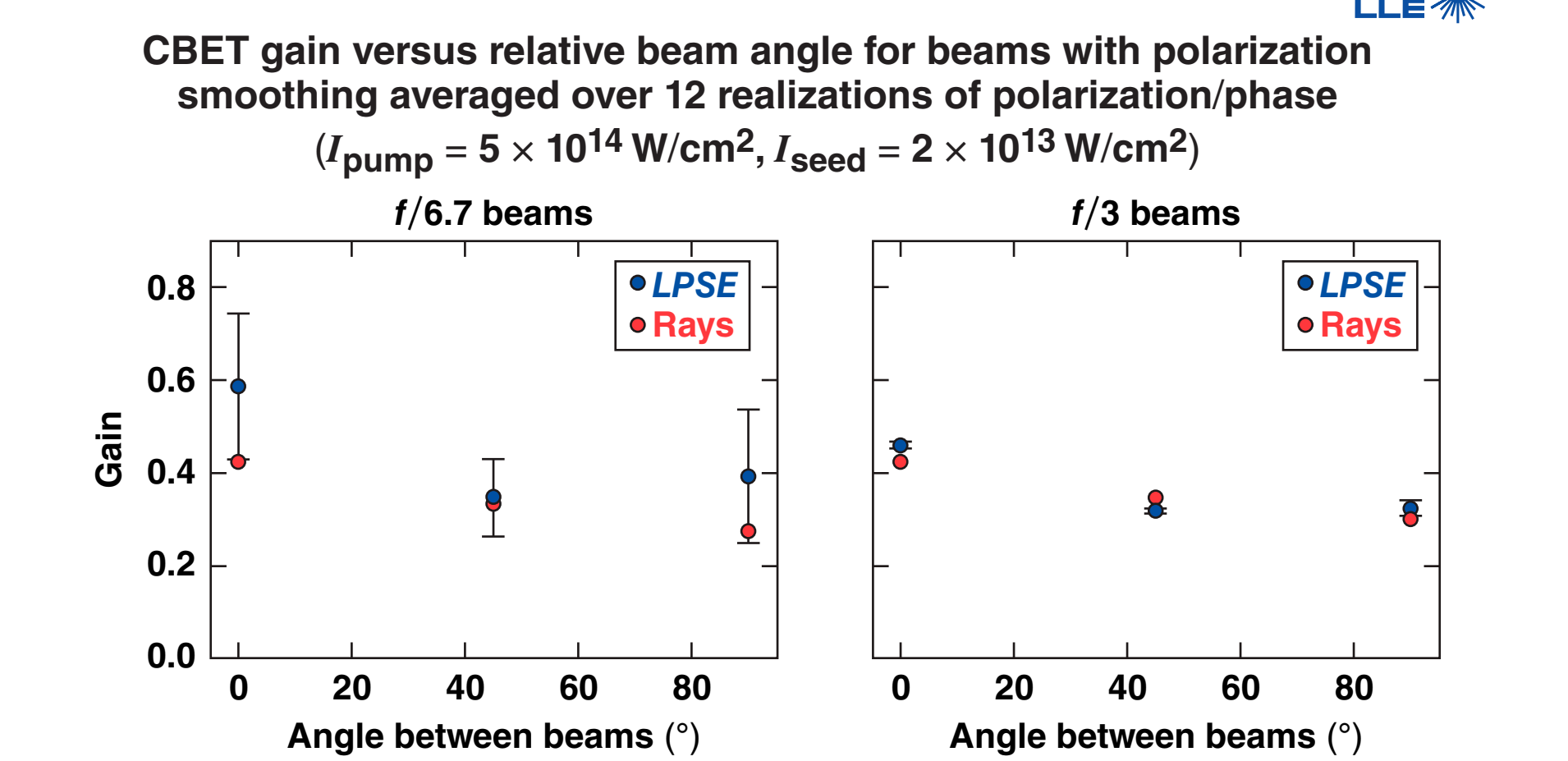
E26195

The speckled ray approach reproduces LPSE results to within one standard deviation of the average over realizations



E26198

The factor of  $(1 + \cos^2\theta)/4$  used to account for the modification of the CBET gain between beams with polarization smoothing is valid only when the speckle length is shorter than the interaction region\*



E26201

## Summary

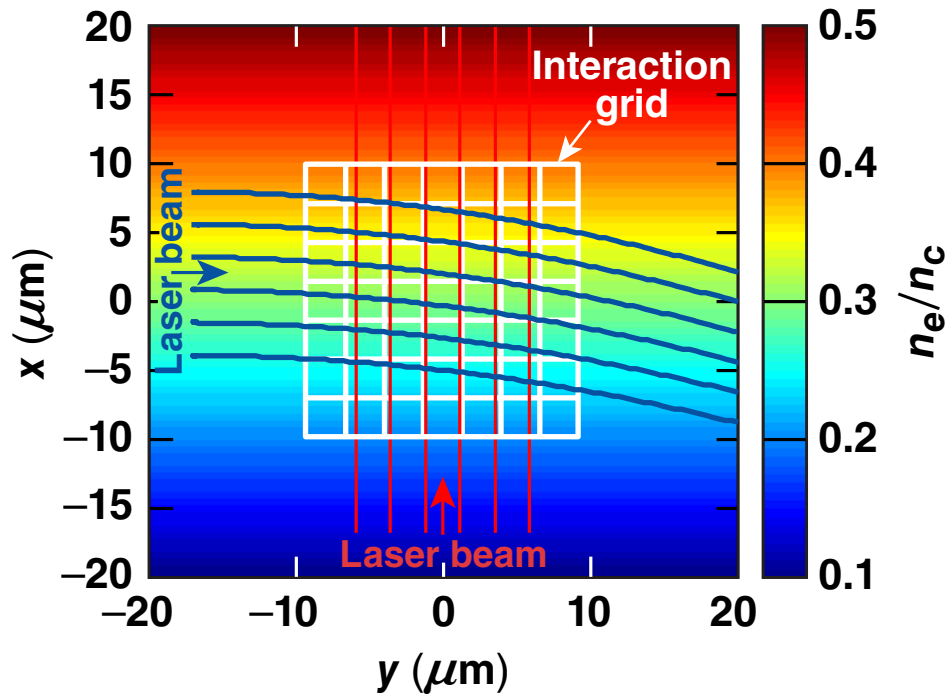
# A 3-D wave-based model has been developed to understand the physics of cross-beam energy transfer (CBET) in an inhomogeneous plasma



- Detailed CBET calculations are used to test ray-based CBET models that are implemented in hydrodynamics codes
- The comparisons generally highlight the accuracy of ray-based models
- Discrepancies between the models are found related to beam speckle and polarization smoothing

# Ray-based models calculate CBET by considering pairwise interactions between rays

Schematic of ray-interaction calculation



CBET interaction between plane waves in a homogeneous plasma

$$\frac{dI_i}{ds} = \sum_j I_j L_{ij}^{-1}$$

$$L_{ij}^{-1} = 5.88 \times 10^{-2} \frac{I_j \lambda}{T_e (1 + 3T_i / ZT_e)} \frac{n_e}{n_c} \frac{\omega_s}{\nu_i} P(\eta_{ij})$$

$$P(\eta) = \frac{\nu_i^2 \eta}{(\eta^2 - 1)^2 + \nu_i^2 \eta^2}$$

$$\eta_{ij} = \frac{\omega_j - \omega_i - (\vec{k}_j - \vec{k}_i) \cdot \vec{u}}{\omega_s}$$

# Ray-based CBET models make several approximations that are not always satisfied in inertial confinement fusion (ICF) applications

- Ion-acoustic waves (IAW's) ( $\delta_n/n \ll 1$ )
- Steady-state convective gain
- Polarization-averaged coupling constant
- Pairwise coupling between beams
- Local plane-wave approximation
  - not valid for speckled beams or at caustics
- Eikonal approximation  
[Wentzel–Kramers–Brillouin (WKB), envelope]
  - breaks down at caustics

Approximations that  
are not made in *LPSE*

# LPSE solves the time-enveloped Maxwell's equations coupled to a linearized plasma response



- Maxwell's equations (time enveloped)  $\vec{E} = \Re \{ \vec{E}(\vec{x}, t) \exp(-i\omega_0 t) \}$

$$\frac{2i\omega_0}{c^2} \frac{\partial}{\partial t} \vec{E} + \nabla^2 \vec{E} - \nabla(\nabla \cdot \vec{E}) + \frac{\omega_0^2}{c^2} \epsilon(\omega; \vec{x}, t) \vec{E} = 0$$

$$E(\omega_0; \vec{x}, t) = 1 - \frac{\omega_{pe}^2(\vec{x}, t)}{\omega_0(\omega_0 + i\nu_{ei})}$$

$$\omega_{pe}^2 = \frac{4\pi e^2 n_e(\vec{x}, t)}{m_e}$$

- Plasma response

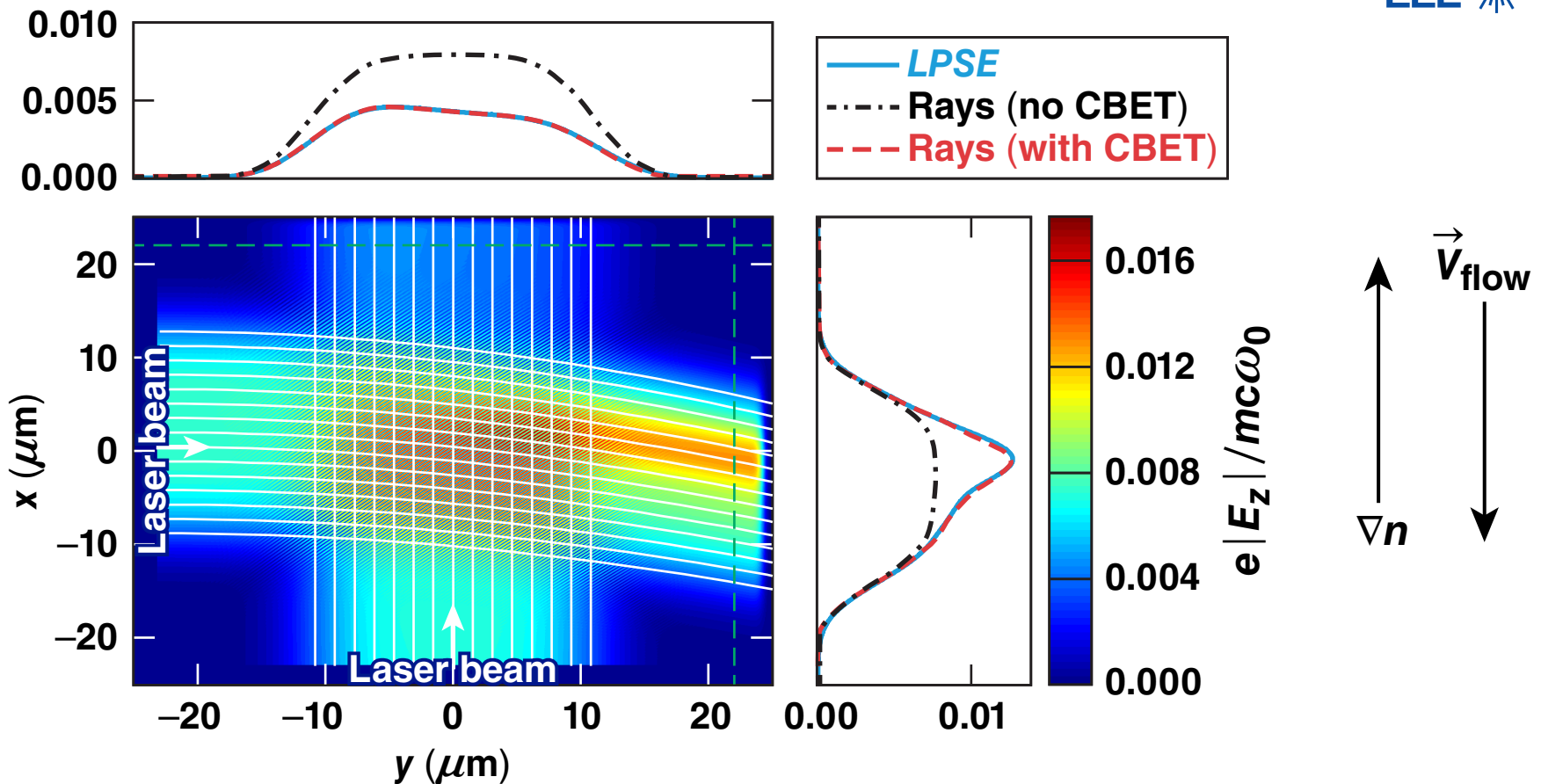
$$\begin{aligned} [\partial_t + \vec{U}_0(\vec{x}) \cdot \nabla] \left( \frac{\delta n}{n_0} \right) &= -W \\ [\partial_t + \vec{U}_0(\vec{x}) \cdot \nabla + 2\hat{\nu}_{IAW}] W &= -\nabla^2 \left[ c_s^2 \left( \frac{\delta n}{n_0} \right) + \frac{e^2}{4m_e \omega_0^2} |\vec{E}|^2 \right] \end{aligned}$$

$$W \equiv \nabla \cdot \delta \vec{U}$$

$$n_e \equiv n_0(\vec{x}) + \delta n(\vec{x}, t)$$

$$\vec{U} = \vec{U}_0(\vec{x}) + \delta \vec{U}(\vec{x}, t)$$

# Ray- and wave-based CBET models give the same result in simple interaction geometries (plane-wave beams, no caustics)

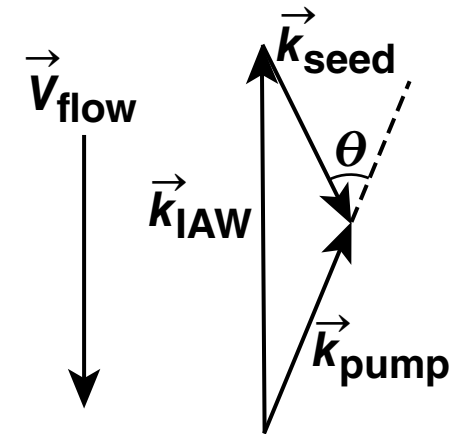
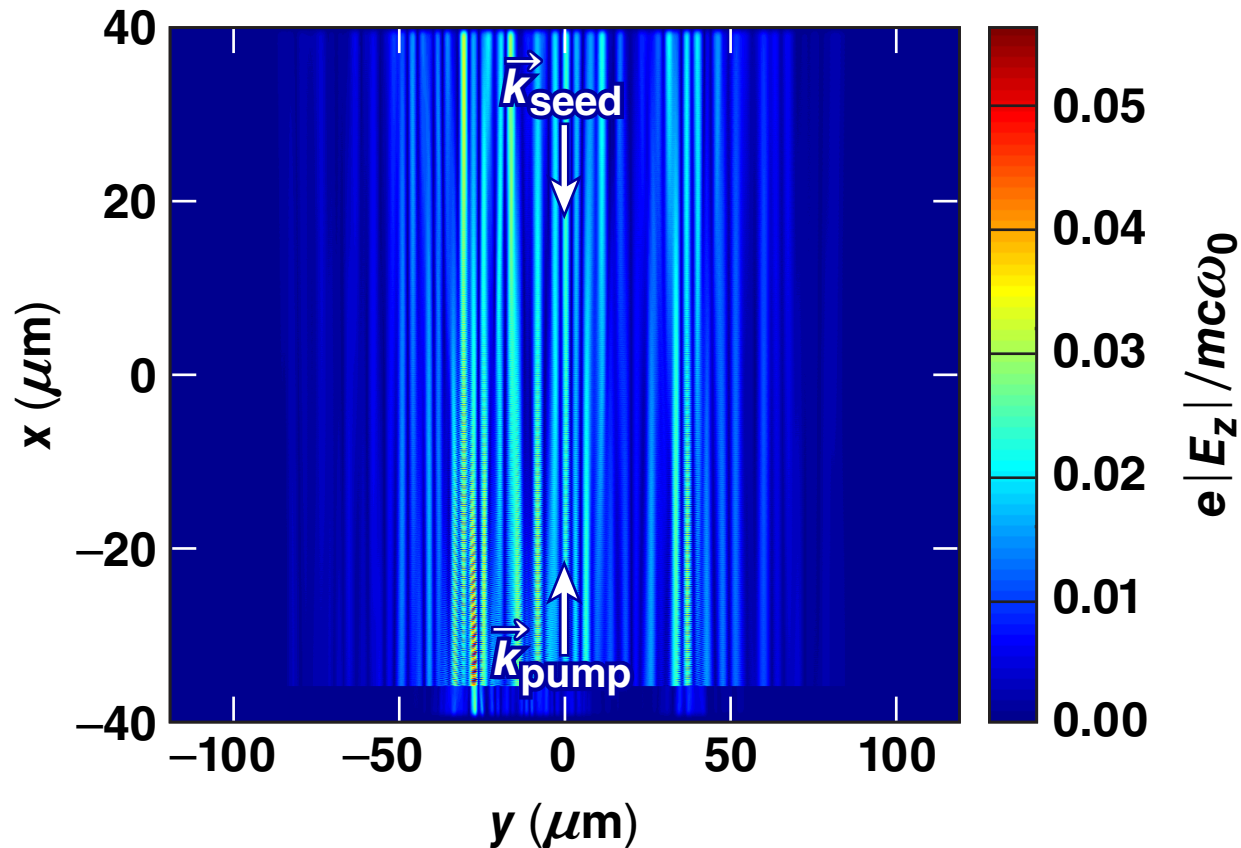


All of the approximations made in the ray model are satisfied in this configuration.

# Speckled beams can transfer more energy than plane-wave beams with the same average intensity

**LPSE simulation of counter-propagating speckled beams**

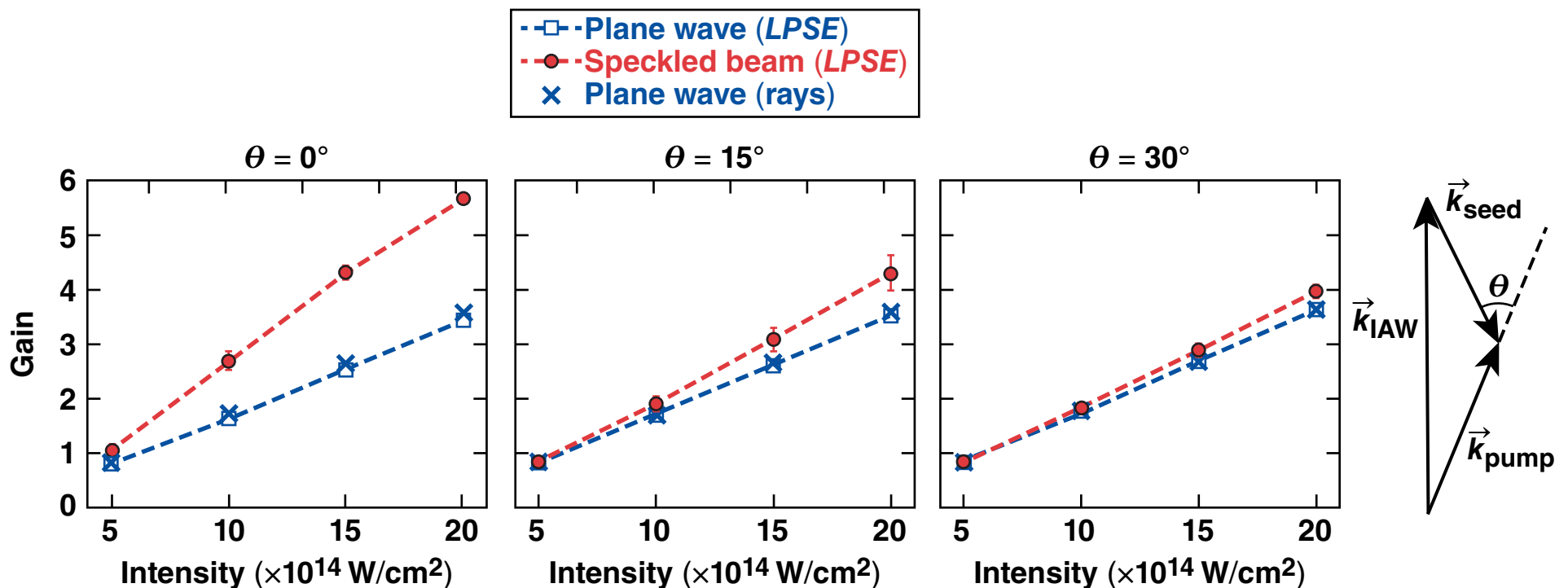
$(I_{\text{pump}} = 2 \times 10^{15} \text{ W/cm}^2, I_{\text{seed}} = 10^{12} \text{ W/cm}^2)$



# The CBET gain is sensitive to beam speckle for gains greater than $\sim 1$ and relative beam angles of less than $\sim 30^\circ$

$$\text{Gain} \equiv \log\left(\frac{\text{Seed energy out}}{\text{Seed energy in}}\right)$$

CBET gain versus pump intensity for various beam-relative beam angles

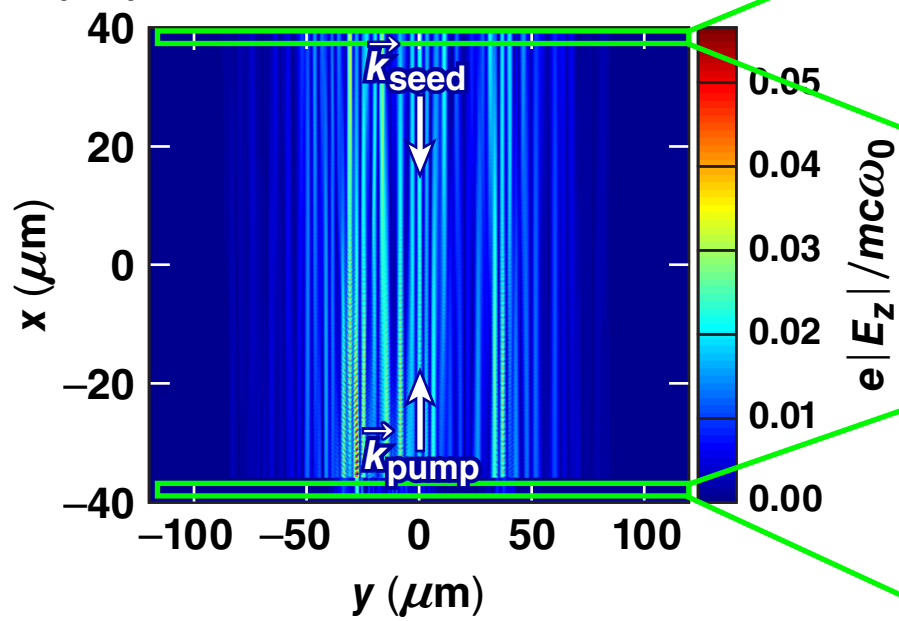




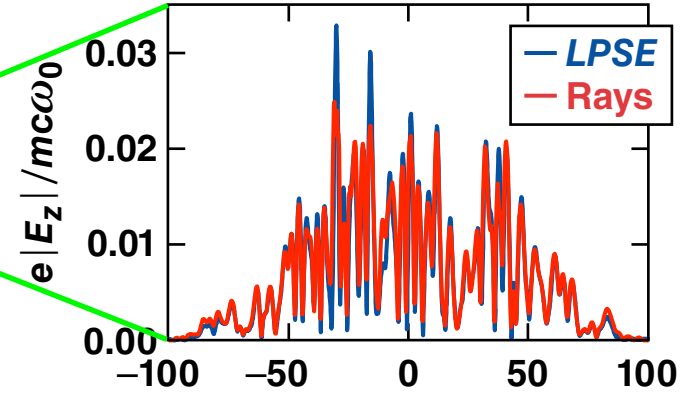
# A good approximation to the average CBET between speckled beams can be obtained by using the linearity of Maxwell's equations to solve for the correct unperturbed field amplitudes in the ray-based calculation



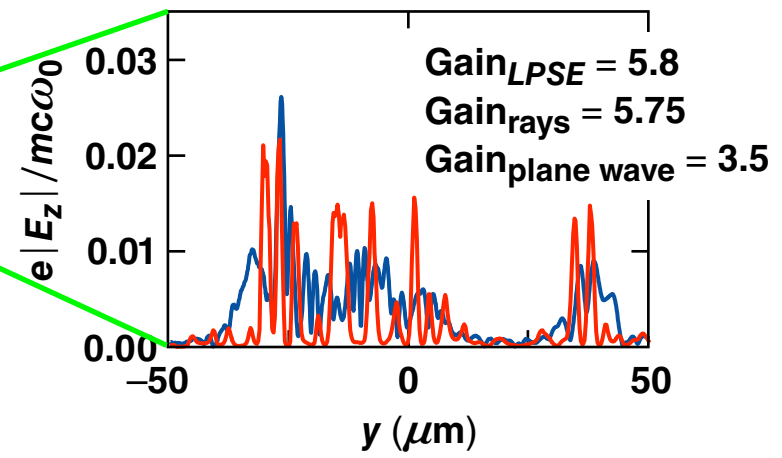
**LPSE simulation of counter-propagating speckled beams**  
 ( $I_{\text{pump}} = 2 \times 10^{15} \text{ W/cm}^2$ ,  $I_{\text{seed}} = 10^{12} \text{ W/cm}^2$ )



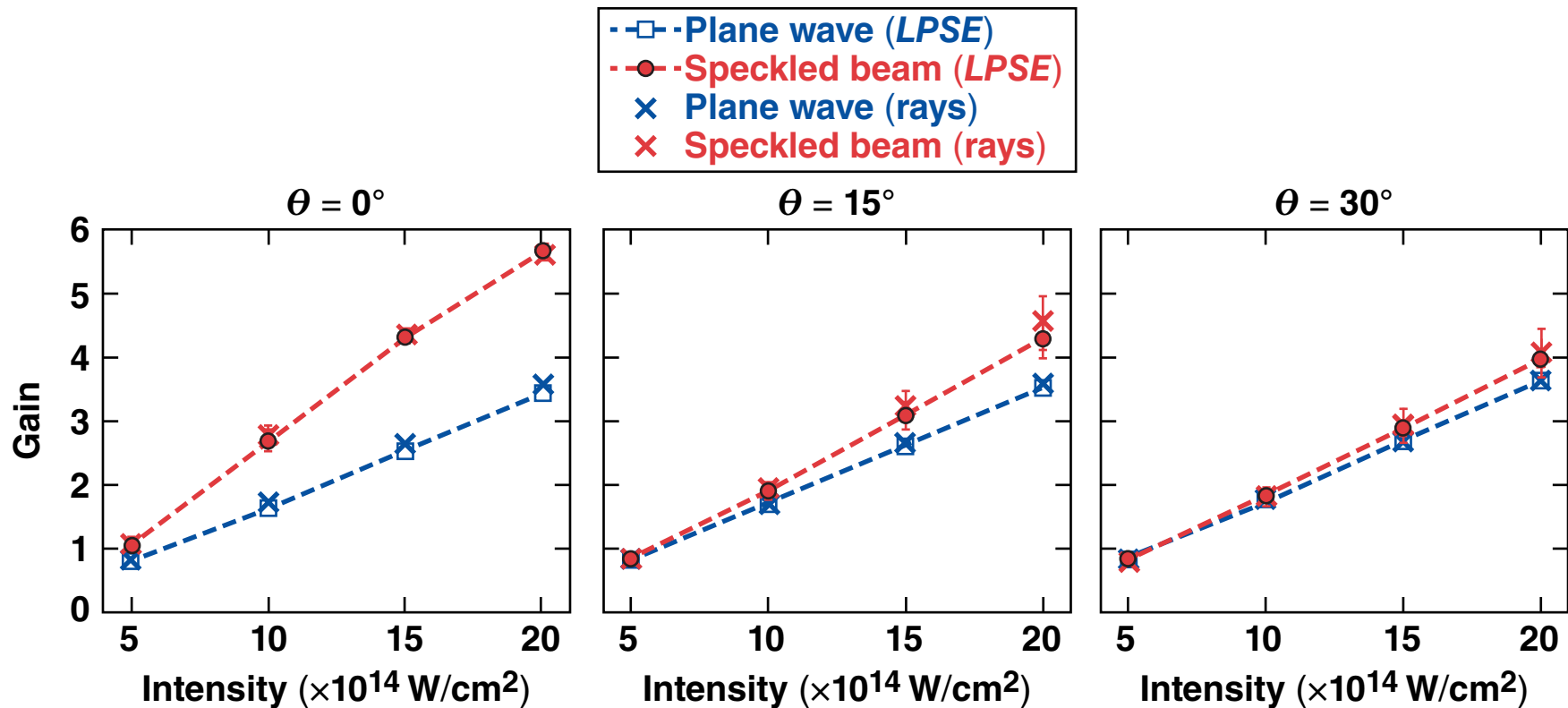
**Outgoing pump beam**



**Outgoing seed beam**



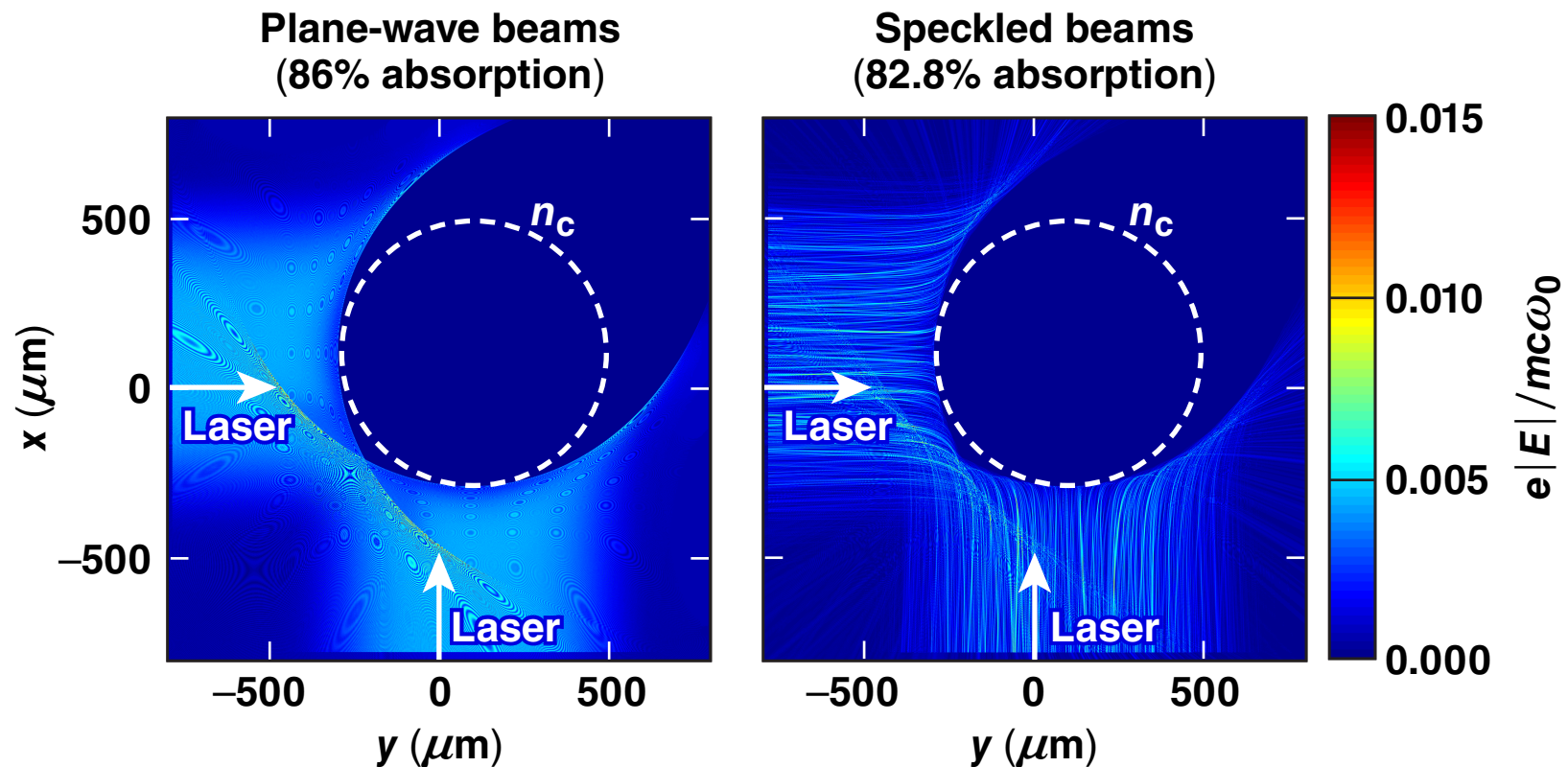
# The speckled ray approach reproduces *LPSE* results to within one standard deviation of the average over realizations



This could be a viable approach for including speckle effects in ray-based CBET models.

# Speckled beams result in a modest decrease in laser absorption in OMEGA-scale, two-beam *LPSE* simulations at ICF-relevant plasma conditions

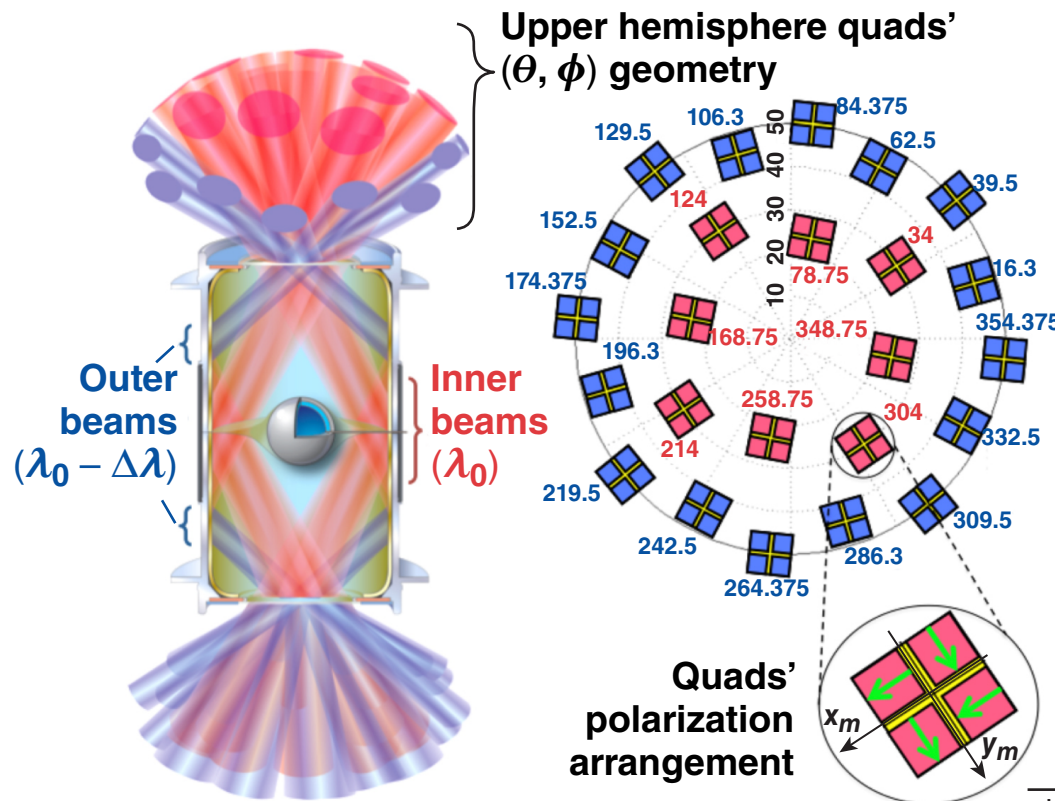
OMEGA-scale *LPSE* simulations ( $I_{\text{single beam}} = 2 \times 10^{14} \text{ W/cm}^2$ )  
in a *LILAC* plasma profile (1.6 billion grid cells)



# Polarization smoothing is accounted for in ray-based CBET models by multiplying the gain coefficient by a factor of $(1 + \cos^2\theta)/4^*$

Polarization smoothing is implemented on the NIF\*\* by splitting the polarization within each quad

- The factor of  $(1 + \cos^2\theta)/4$  comes from assuming that the interacting beams have random relative polarizations with uncorrelated speckle patterns and ensemble averaging over realizations when deriving the ponderomotive potential of the beat wave



$$\langle |\phi|^2 \rangle_{PS} = \frac{1}{4} (1 + \cos^2\theta) |\theta|_{\parallel}^2$$

\*P. Michel *et al.*, Phys. Plasmas 20, 056308 (2013).  
\*\*NIF: National Ignition Facility

The factor of  $(1+\cos^2\theta)/4$  used to account for the modification of the CBET gain between beams with polarization smoothing is valid only when the speckle length is shorter than the interaction region\*



CBET gain versus relative beam angle for beams with polarization smoothing averaged over 12 realizations of polarization/phase

$$(I_{\text{pump}} = 5 \times 10^{14} \text{ W/cm}^2, I_{\text{seed}} = 2 \times 10^{13} \text{ W/cm}^2)$$

