

# A First-Principles Equation-of-State Table of Beryllium for High-Energy-Density Plasma Simulations



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## Summary

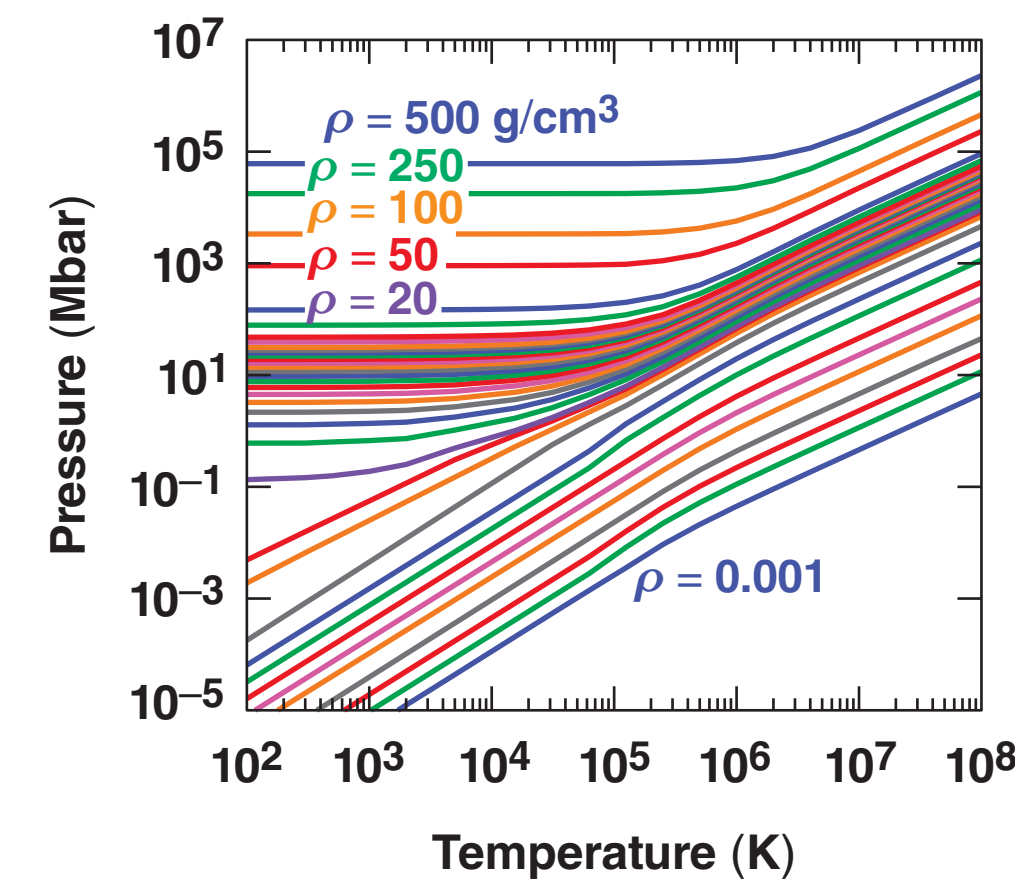
An accurate equation-of-state (EOS) table of beryllium has been built from first-principles calculations for inertial confinement fusion (ICF) and high-energy-density-physics (HEDP) applications



- Based on density-functional-theory (DFT) calculations [combining Kohn–Sham molecular dynamics (KSMD) and orbital-free molecular dynamics (OFMD)], we have established a wide-range beryllium EOS table of density  $\rho = 0.001$  to  $\rho = 500$  g/cm<sup>3</sup> and temperature  $T = 2000$  to  $10^8$  K
- The first-principles equation-of-state (FPEOS) table is in good agreement with the widely used *SESAME* EOS table (*SESAME* 2023)
- By implementing the FPEOS table into the 1-D radiation–hydrodynamics code *LILAC*, we studied the EOS effects on beryllium-shell target implosions; the FPEOS simulation predicts a higher neutron yield (~15%) compared to the simulation using the *SESAME* 2023 EOS table

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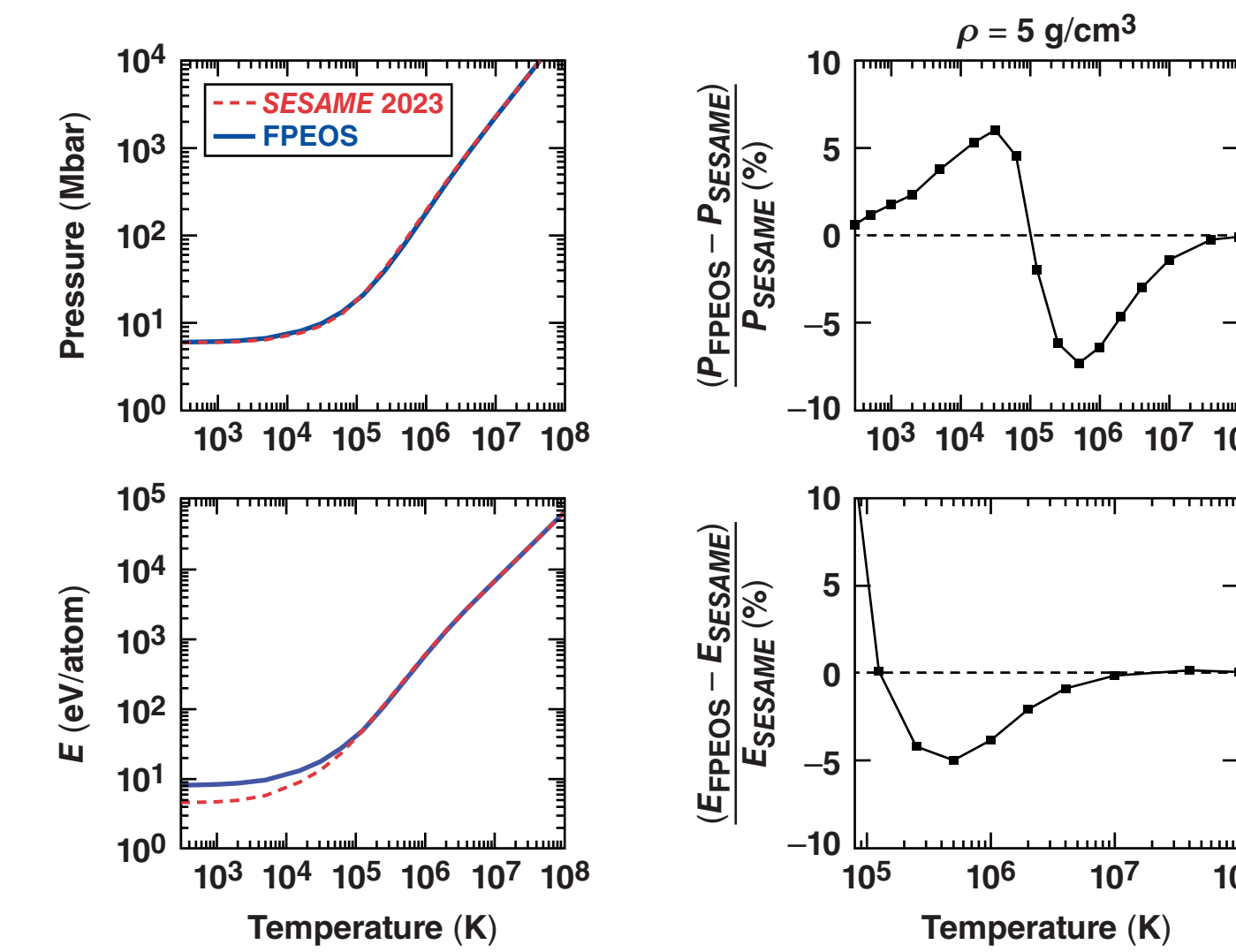
The FPEOS of beryllium has been calculated for densities and temperatures from  $\rho = 0.001$  g/cm<sup>3</sup> to  $\rho = 500$  g/cm<sup>3</sup> and  $T = 2,000$  K to  $10^8$  K



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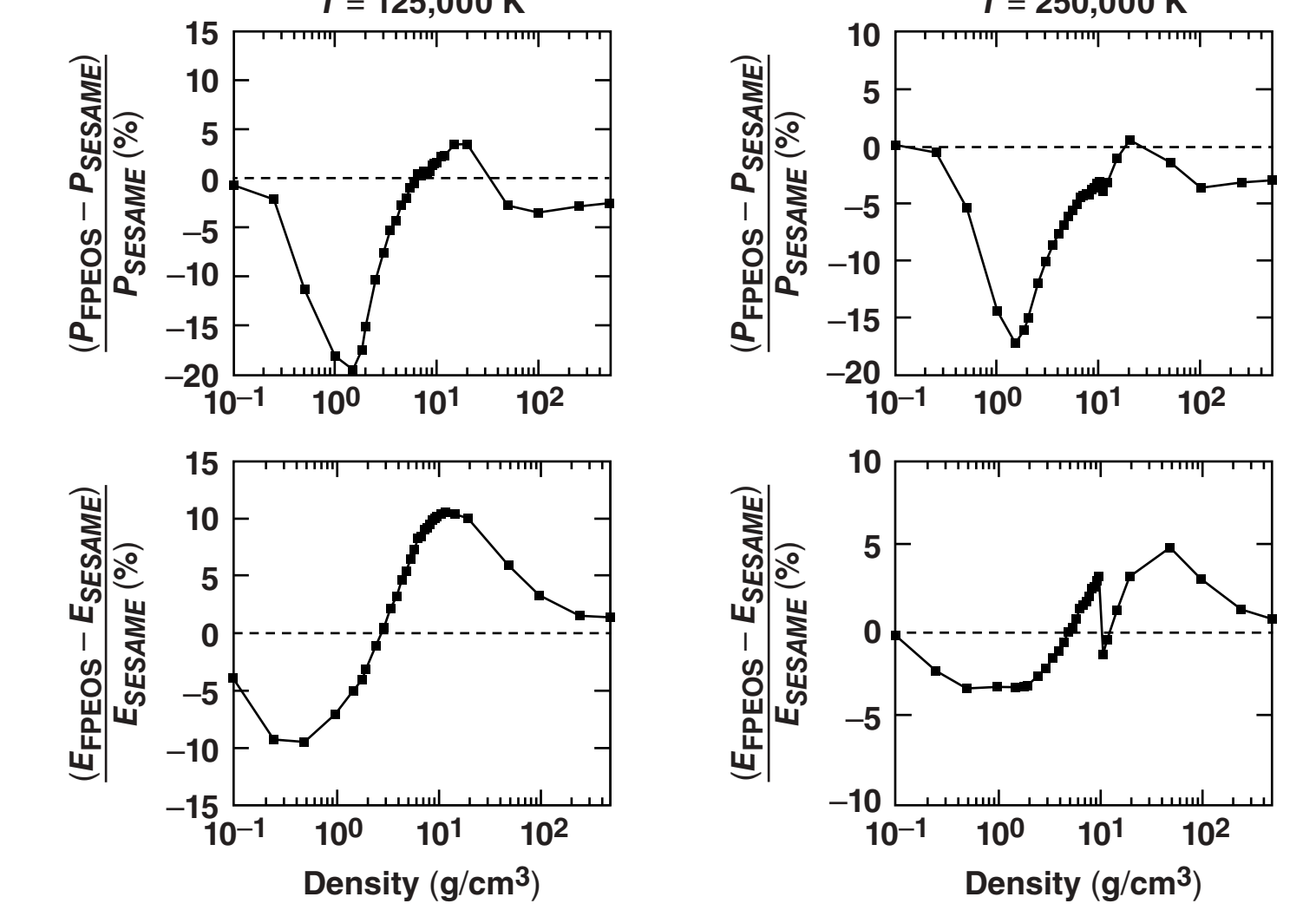
Y. H. Ding and S. X. Hu, Phys. Plasmas 24, 062702 (2017).

The off-Hugoniot equation-of-state comparison between FPEOS and *SESAME* at certain densities



TC13422

The off-Hugoniot equation-of-state comparison between FPEOS and *SESAME* at certain temperatures



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In the WDM regime the difference is most significant (20%).

It is of great importance to calculate the EOS of beryllium for a wide range of conditions by using reliable methods



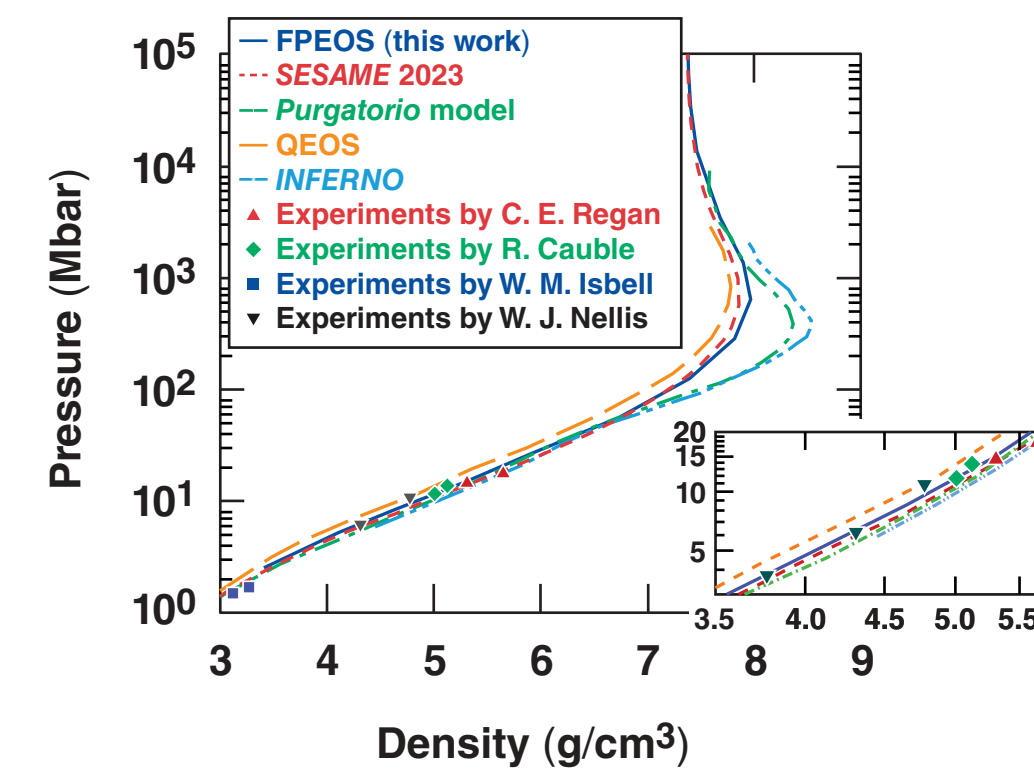
- In ICF applications, target designs are based mainly on the radiation–hydrodynamics simulations, where the materials will experience many different pressure (up to  $10^5$  Mbar) and temperature (up to  $10^8$  K) conditions; accurate properties of beryllium under such extreme conditions are essential for ICF applications
- Theoretical EOS models are particularly difficult to generate in the so-called warm-dense-matter (WDM) regime, where both strongly coupled ( $T > 1$ ) and degeneracy effects ( $T < 1$ ) are important\*
- Quantum molecular dynamics (QMD),\*\* based on DFT, has proven to be a reliable method for studying the many-body quantum systems of dense plasmas; QMD simulations have been shown to work well for EOS calculations\*\*\*

Combining orbital-based-DFT, KSMD, and OFMD, we can investigate wide-ranged EOS tables.

\*S. X. Hu, B. Militzer, V. N. Goncharov, and S. Skupsky, Phys. Rev. Lett. 104, 225003 (2010); Phys. Rev. B 84, 224109 (2011); S. X. Hu, T. R. Boehly, and L. A. Collins, Phys. Rev. E 89, 053104 (2014).  
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TC13417

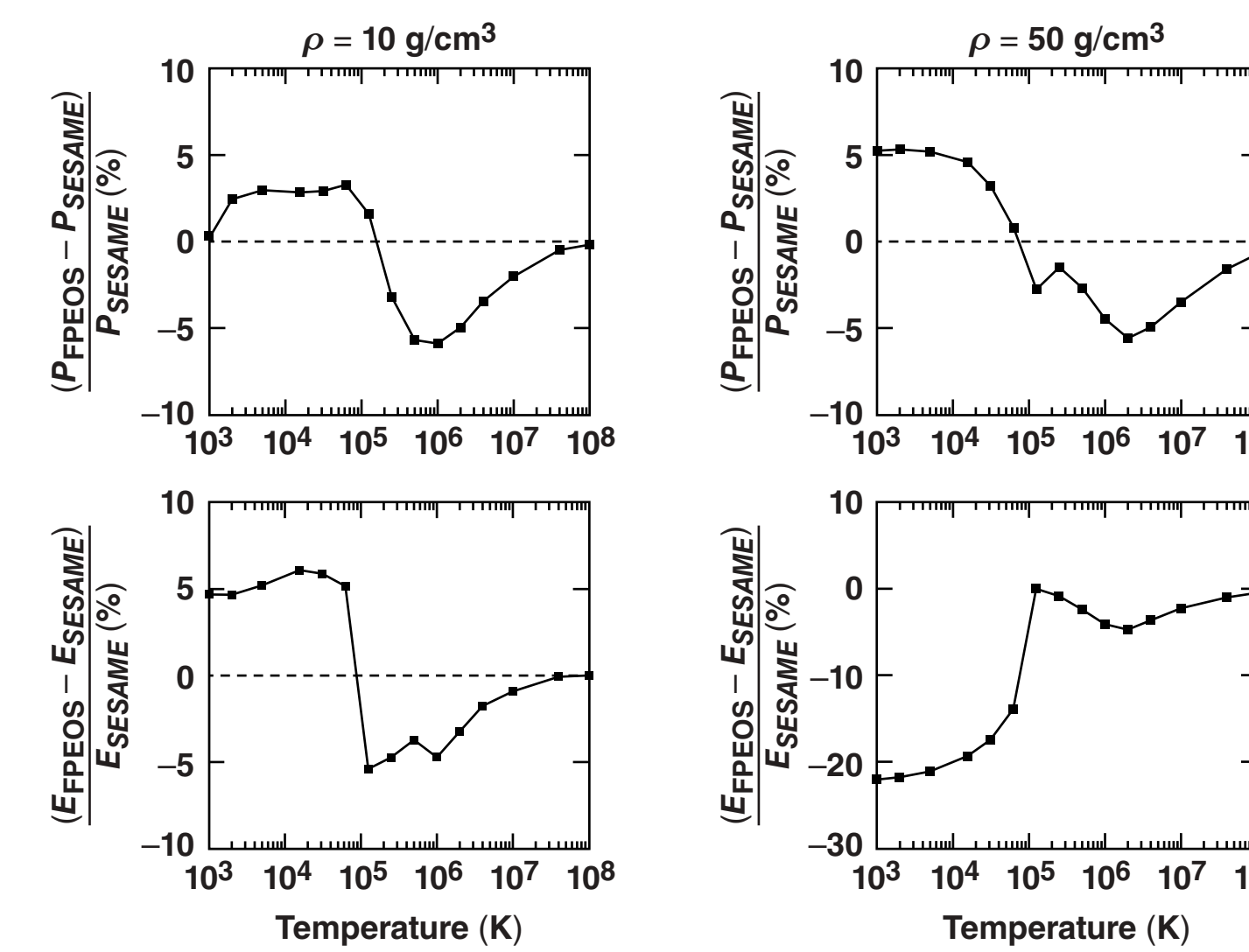
The calculated principal shock Hugoniot of beryllium from FPEOS has been compared with other theoretical models and experiments



The FPEOS Hugoniot pressure of beryllium is in good agreement (within 10%) with the widely used *SESAME* model (*SESAME* 2023) in the low-compression-ratio region; the pressure differences can be up to 30% in the high-compression region.

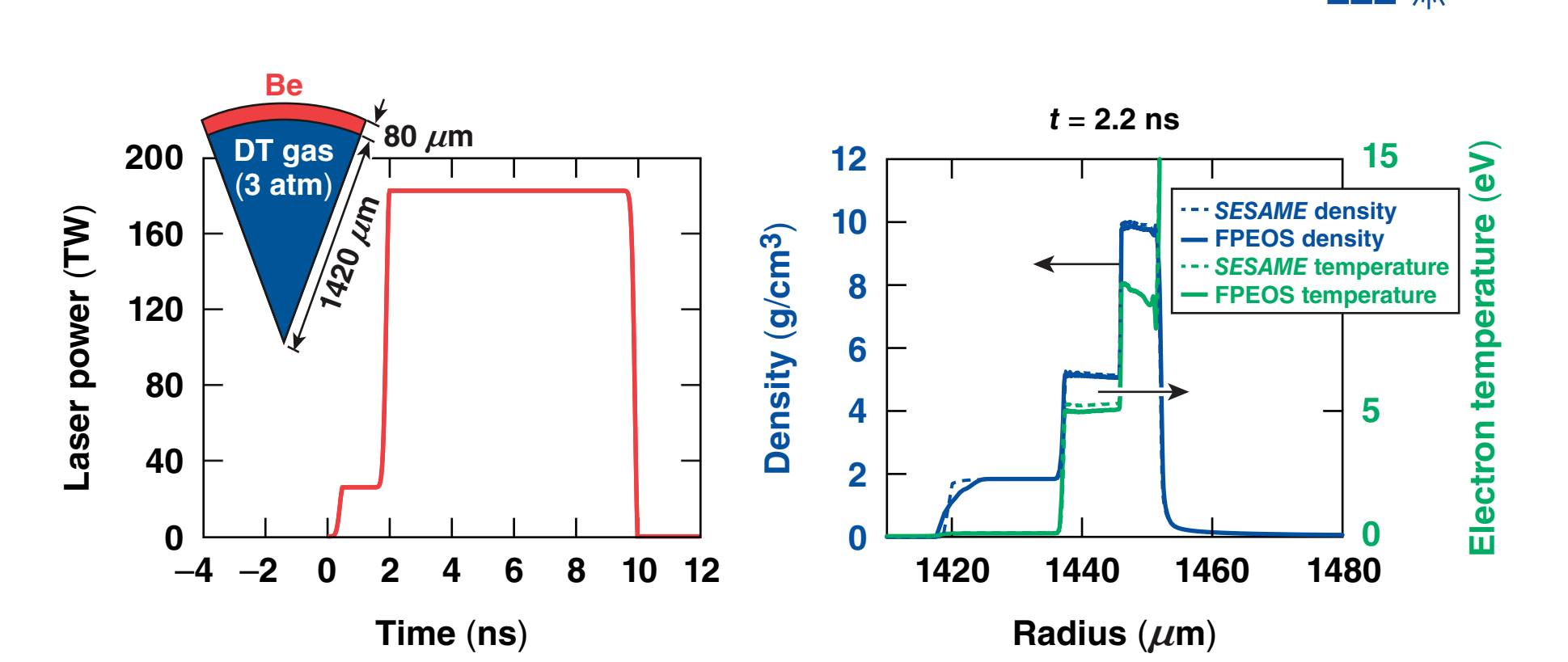
TC13419

The off-Hugoniot equation-of-state comparison between FPEOS and *SESAME* at certain densities



TC13423

The effects of beryllium FPEOS on high-energy-density (HED) plasma simulations



For DT, gas two simulations used the same FPEOS table and the first-principles opacity table.

TC13424

QMD calculations are based on DFT



Total energy in terms of density

$$E = -\frac{1}{2} \sum_n \int \phi_n^*(\vec{r}) \nabla^2 \phi_n(\vec{r}) d\vec{r} + \int n(\vec{r}) V_{ne}(\vec{r}) d\vec{r} + \frac{1}{2} \iint d\vec{r} d\vec{r}' \frac{n(\vec{r})n(\vec{r}')}{|\vec{r}-\vec{r}'|} + \int \mu_{xc}(n)n(\vec{r}) d\vec{r}$$

Minimize with respect to single-particle orbitals

$$\frac{\delta E}{\delta \phi_i} = \frac{\delta T}{\delta \phi_i} + \left[ \frac{\delta E_{en}}{\delta n_i} + \frac{\delta E_{ee}}{\delta n_i} + \frac{\delta E_{xc}}{\delta n_i} \right] \frac{\delta n_i(\vec{r})}{\delta \phi_i(\vec{r})} = \epsilon_i \phi_i(\vec{r})$$

The Kohn–Sham equations can be solved self-consistently

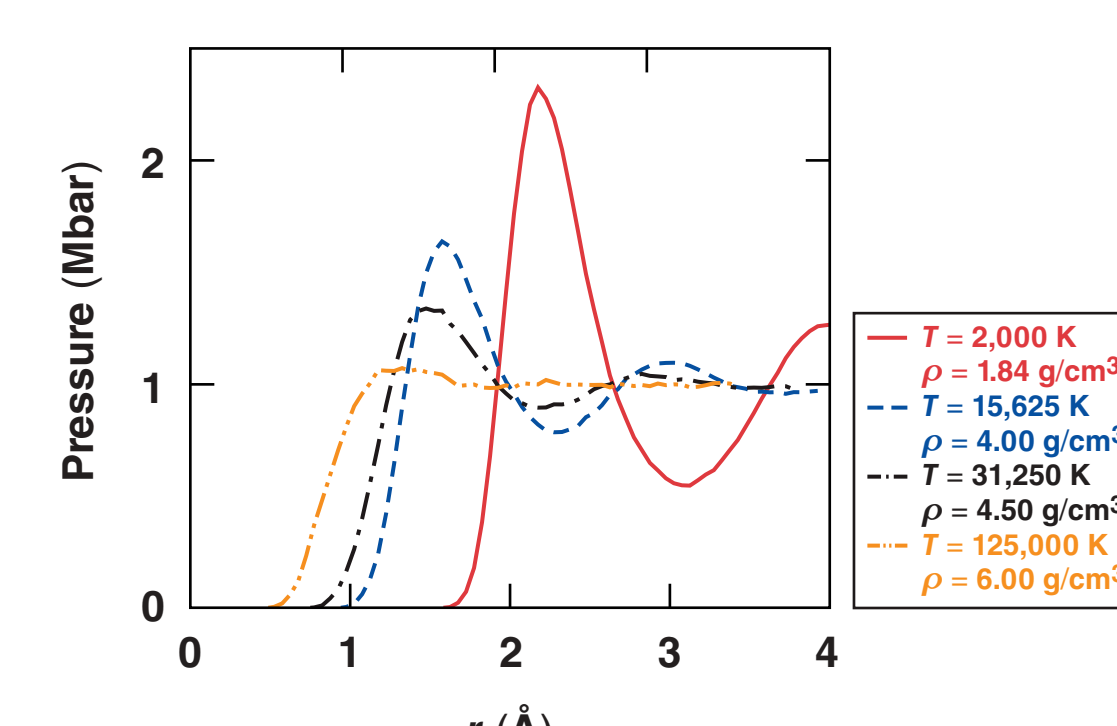
$$-\frac{1}{2} \nabla^2 \phi_i(\vec{r}) + \left[ V_{ne}(\vec{r}) + \int \frac{n(\vec{r}')}{|\vec{r}-\vec{r}'|} d\vec{r}' + \mu_{xc}(n) + n(\vec{r}) \frac{\delta \mu_{xc}(n)}{\delta n(\vec{r})} \right] \phi_i(\vec{r}) = \epsilon_i \phi_i(\vec{r})$$

$$(T + V_{eff}) \phi_i(\vec{r}) = \epsilon_i \phi_i(\vec{r})$$

Ion force for molecular dynamics  $F = -\nabla_{\vec{R}} \langle \psi | \hat{H} | \psi \rangle = -\frac{\partial}{\partial \vec{R}} (F_{el} + F_{ii})$

TC13526

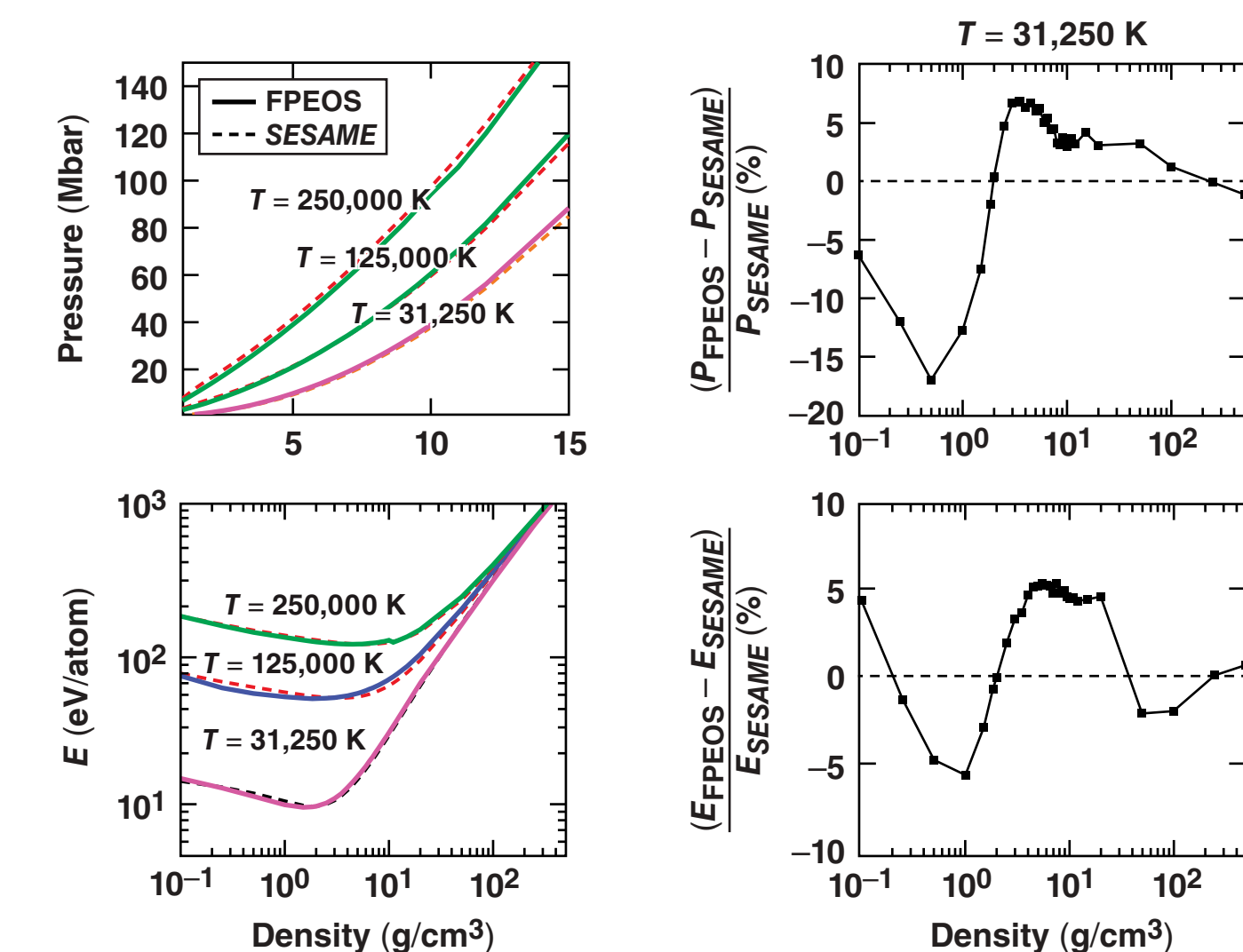
The ion–ion pair correlation function shows the structure change of Be along the principle Hugoniot curve



The peaks in the pair correlation  $g(r)$  gradually disappear along the principle Hugoniot curve, which indicates the structure change from solid to liquid state.

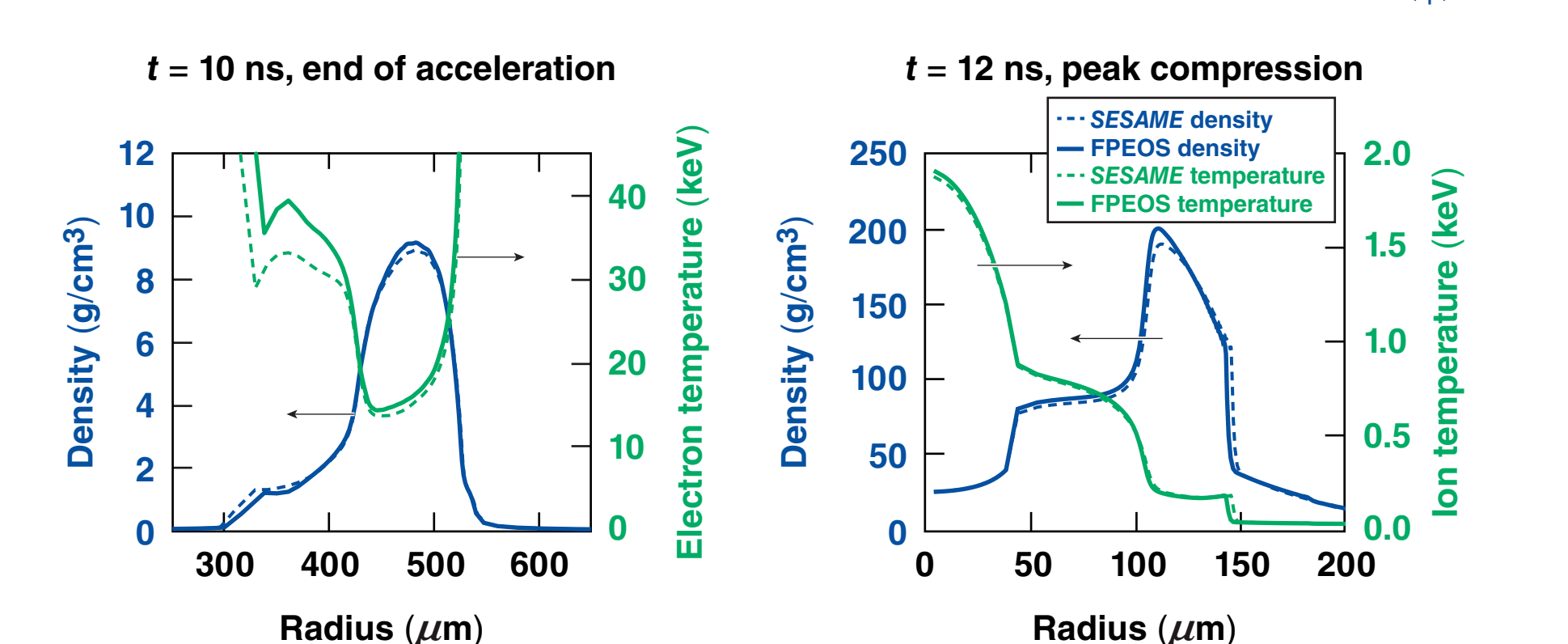
TC13527

The off-Hugoniot equation-of-state comparison between FPEOS and *SESAME* at certain temperatures



TC13420

The effects of beryllium FPEOS on HED plasma simulations



The FPEOS simulation also predicted a higher total neutron yield ( $y = 3.76 \times 10^{14}$ ), which is ~15% higher than the *SESAME* simulation, ( $y = 3.28 \times 10^{14}$ ).

TC13425

## Summary

# An accurate equation-of-state (EOS) table of beryllium has been built from first-principles calculations for inertial confinement fusion (ICF) and high-energy-density–physics (HEDP) applications



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# It is of great importance to calculate the EOS of beryllium for a wide range of conditions by using reliable methods



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**Combining orbital-based–DFT, KSMD, and OFMD, we can investigate wide-ranged EOS tables.**

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# QMD calculations are based on DFT

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- Minimize with respect to single-particle orbitals

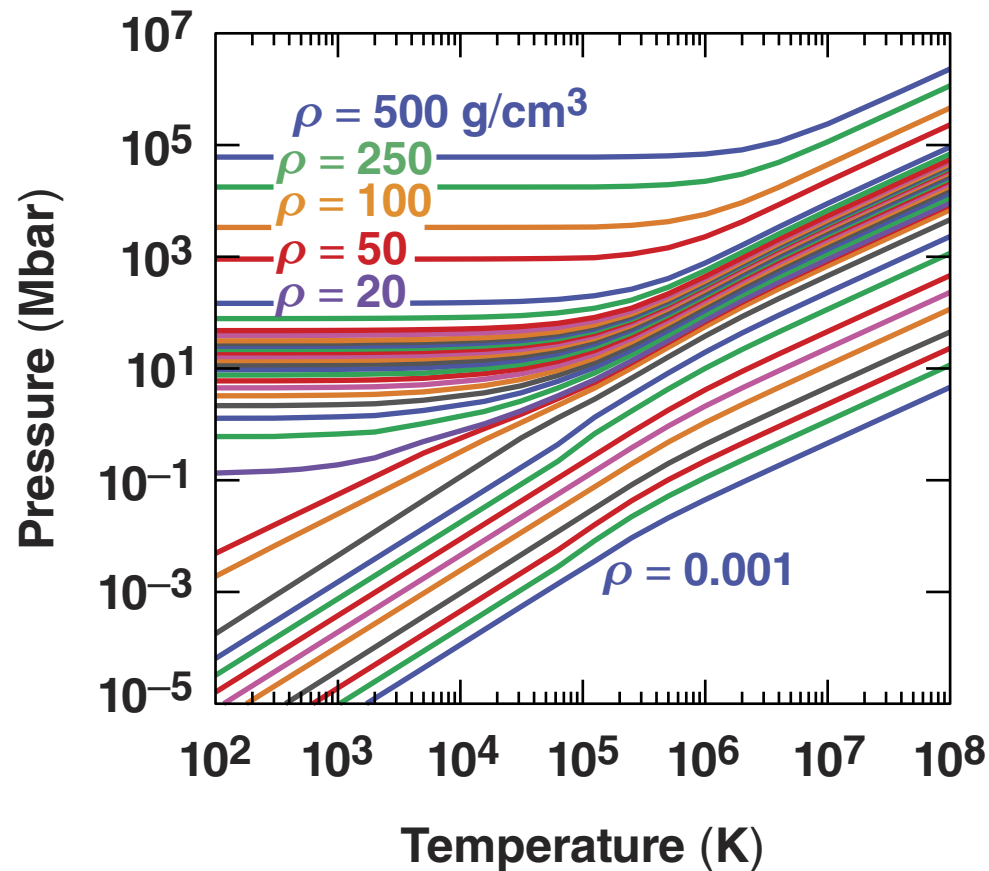
$$\frac{\delta E}{\delta \phi_i} = \frac{\delta T}{\delta \phi_i(\vec{r})} + \left[ \frac{\delta E_{\text{en}}}{\delta n_i(\vec{r})} + \frac{\delta E_{\text{ee}}}{\delta n_i(\vec{r})} + \frac{\delta E_{\text{xc}}}{\delta n_i(\vec{r})} \right] \frac{\delta n_i(\vec{r})}{\delta \phi_i(\vec{r})} = \epsilon_i \phi_i(\vec{r})$$

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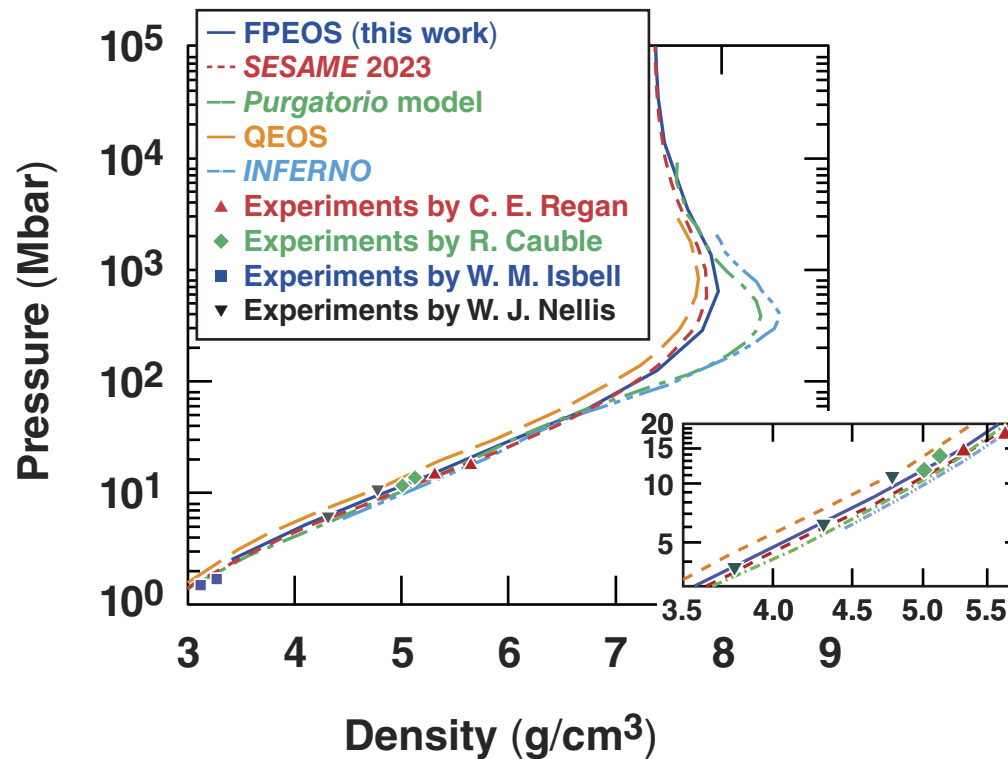
$$-\frac{1}{2} \nabla^2 \phi_i(\vec{r}) + \left[ V_{\text{ne}}(\vec{r}) + \int \frac{n(\vec{r}')}{|\vec{r} - \vec{r}'|} d\vec{r}' + \mu_{\text{xc}}(n) + n(\vec{r}) \frac{\delta \mu_{\text{xc}}(n)}{\delta n(\vec{r})} \right] \phi_i(\vec{r}) = \epsilon_i \phi_i(\vec{r})$$
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- Ion force for molecular dynamics  $\mathbf{F} = -\nabla_{\mathbf{R}} \langle \psi | \hat{H} | \psi \rangle = \frac{\partial}{\partial \mathbf{R}} (\mathbf{F}_{\text{ei}} + \mathbf{F}_{\text{ii}})$

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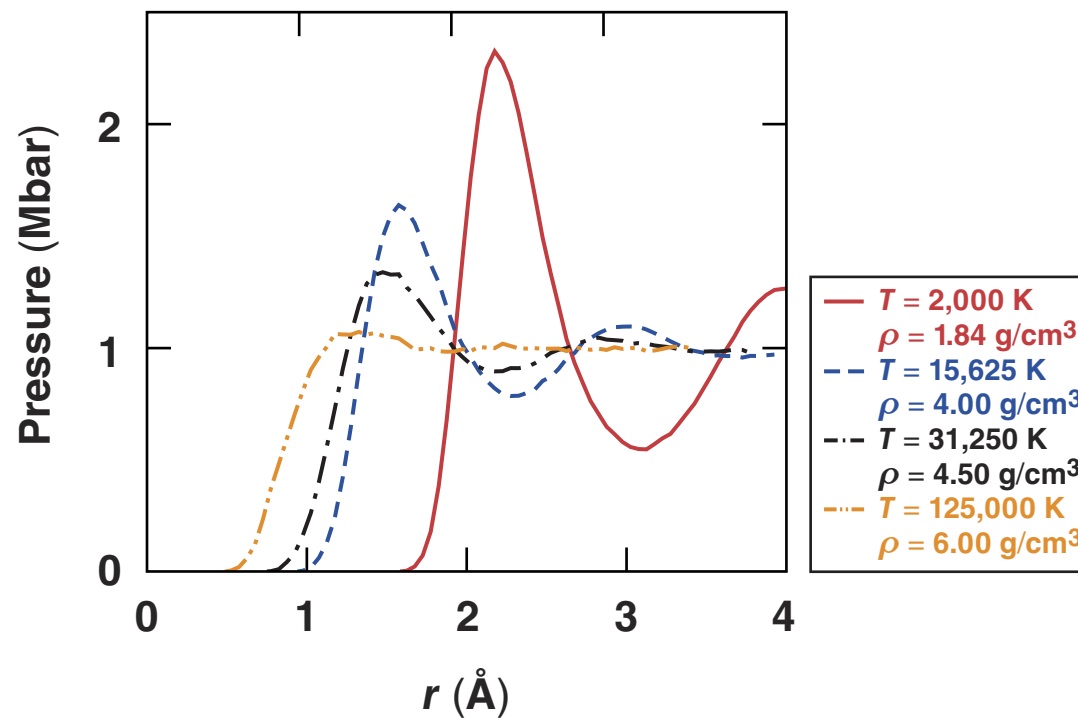


# The calculated principal shock Hugoniot of beryllium from FPEOS has been compared with other theoretical models and experiments



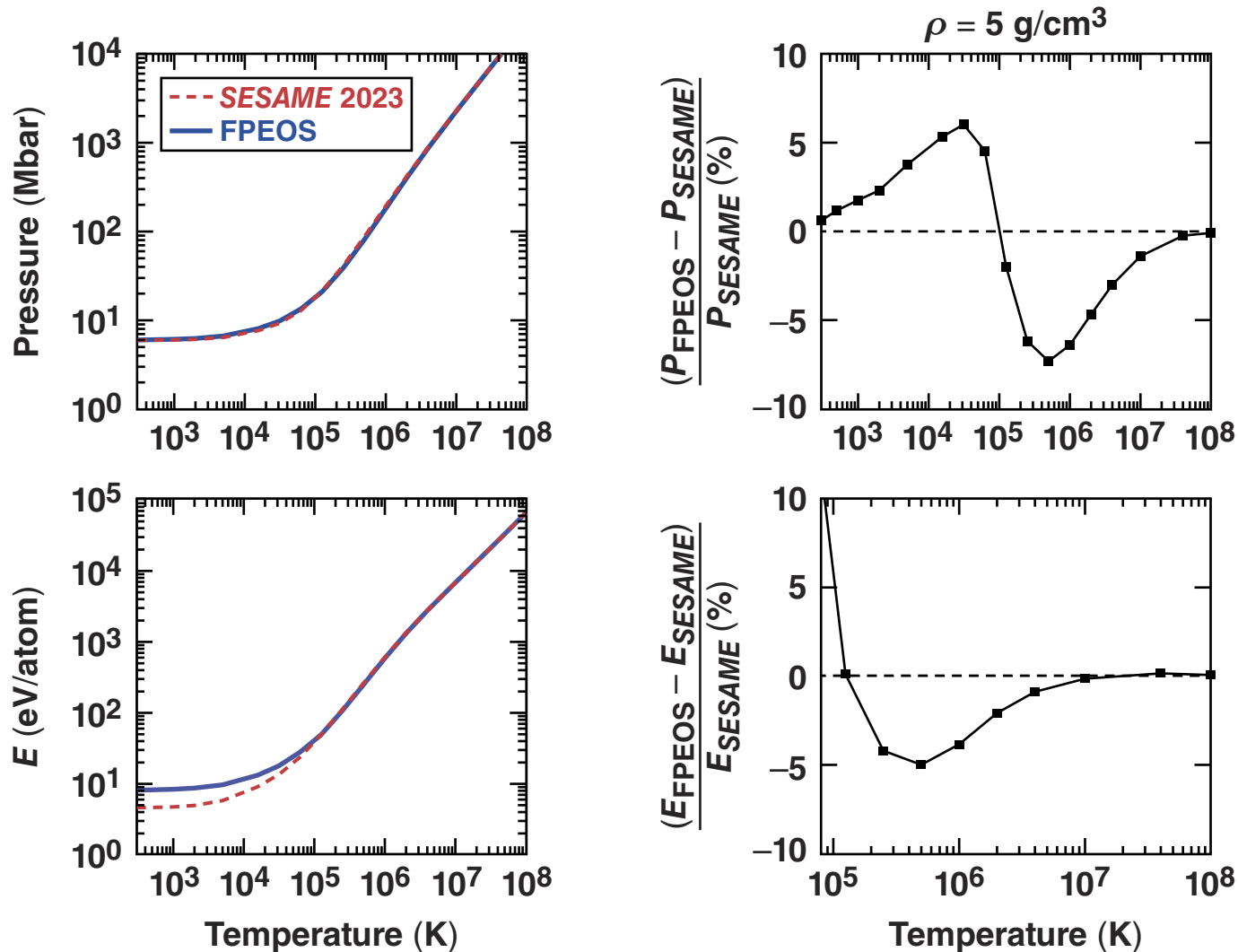
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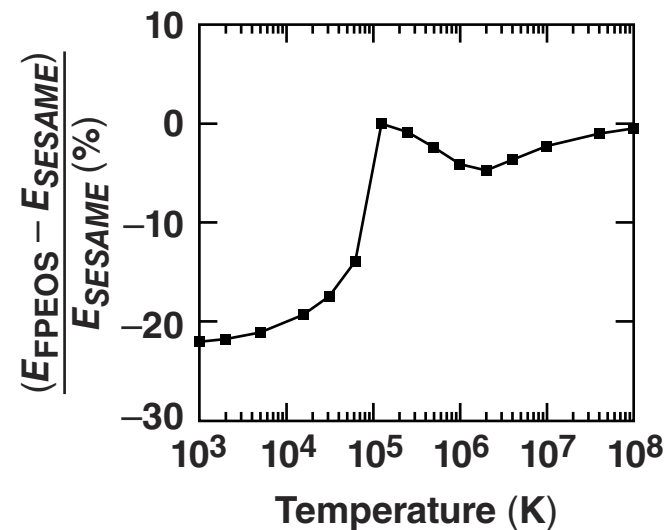
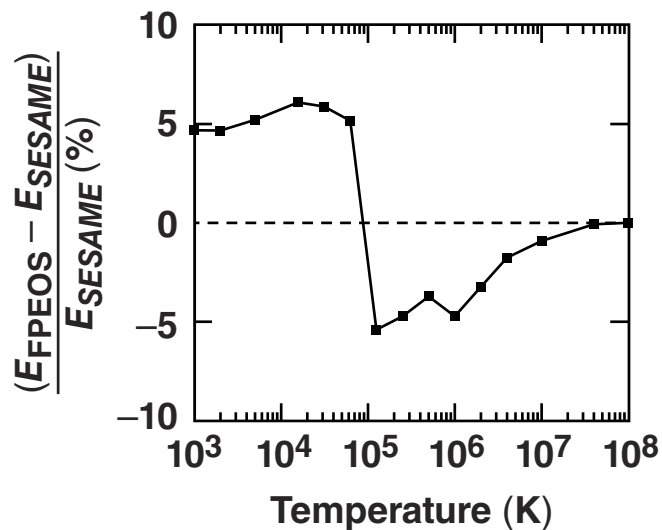
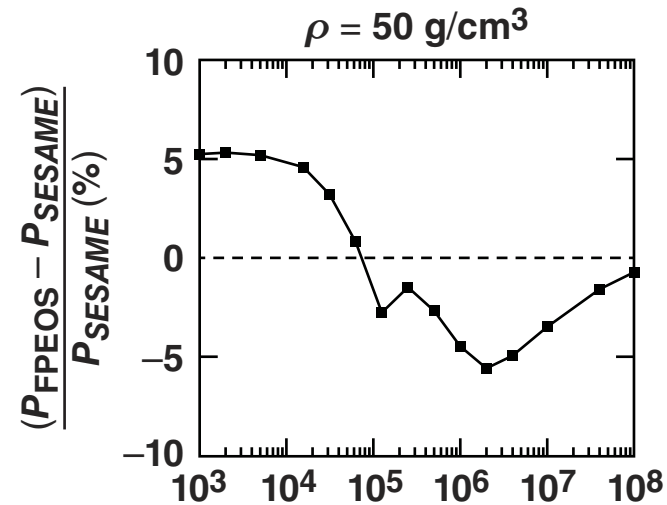
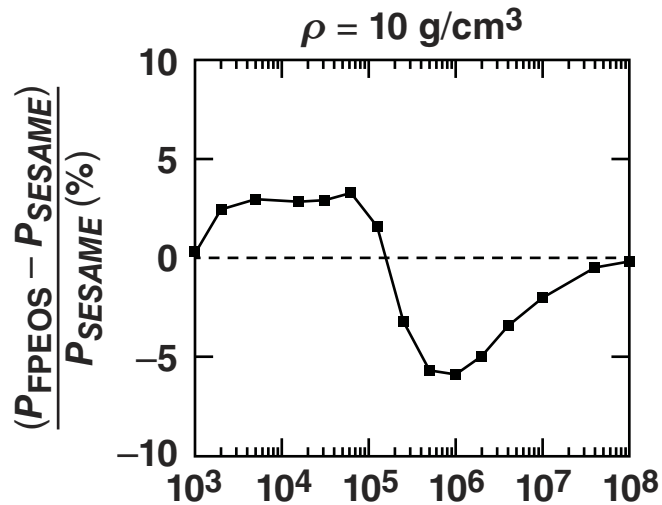
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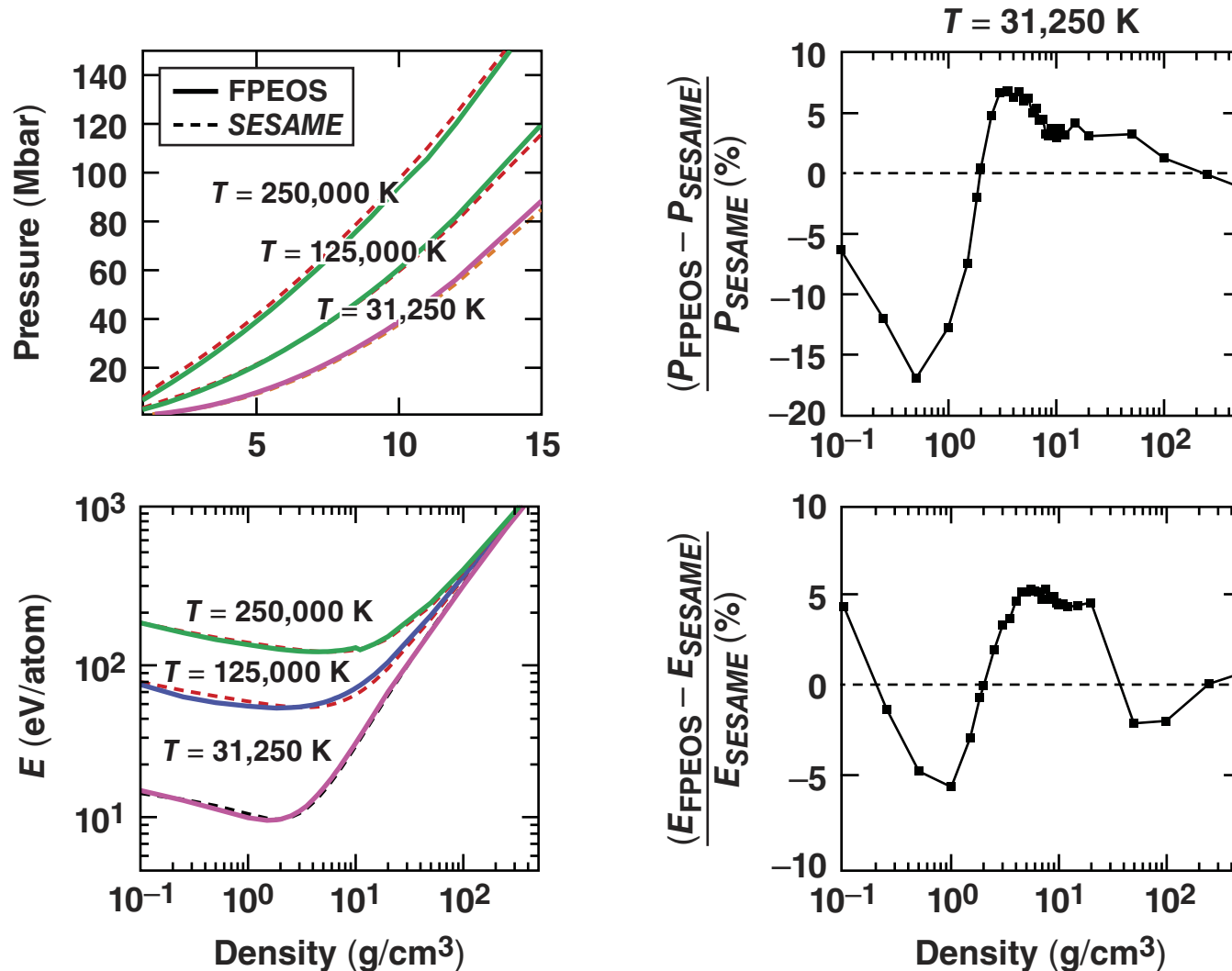




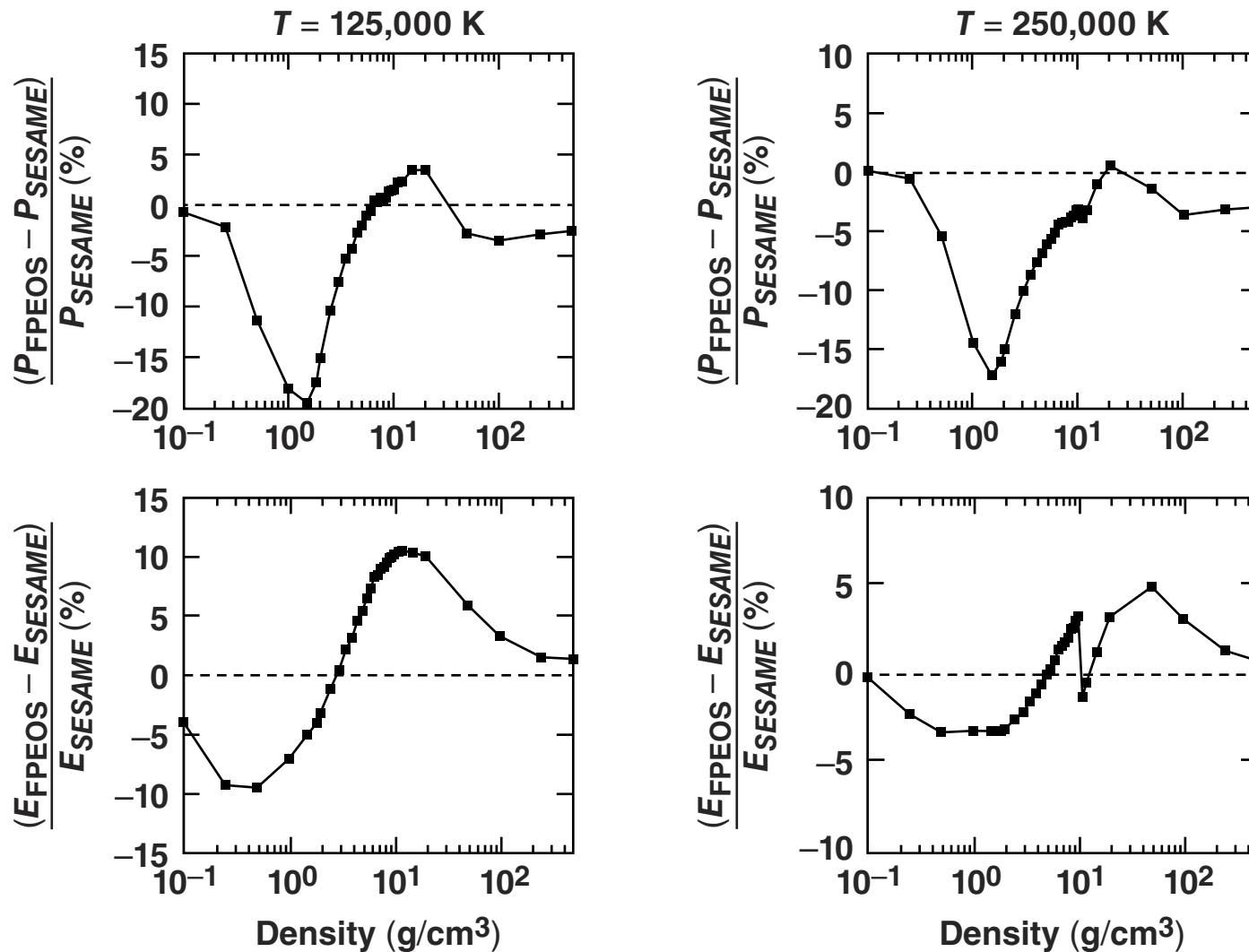
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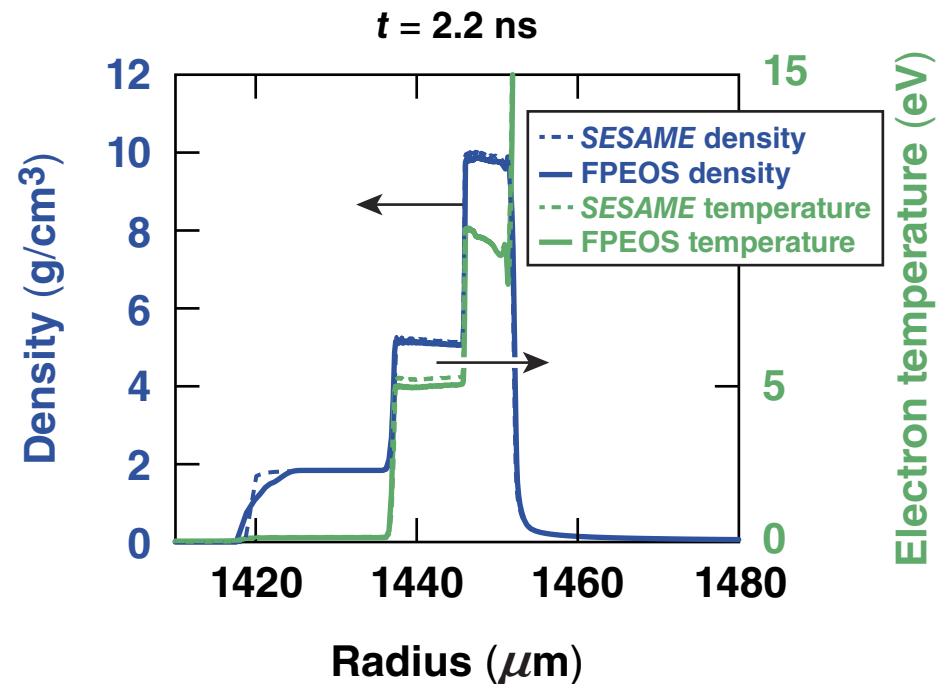
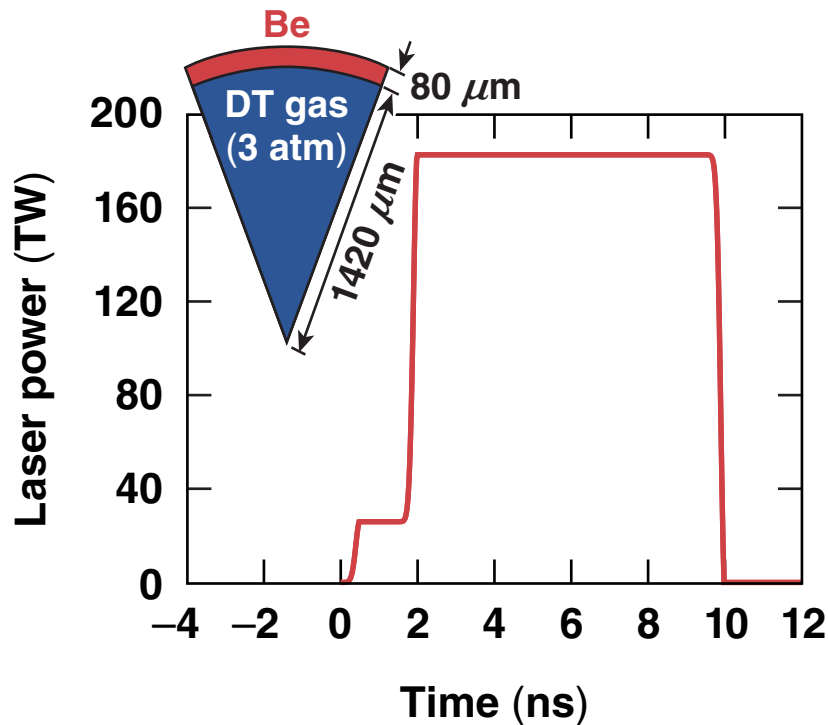


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In the WDM regime the difference is most significant (20%).

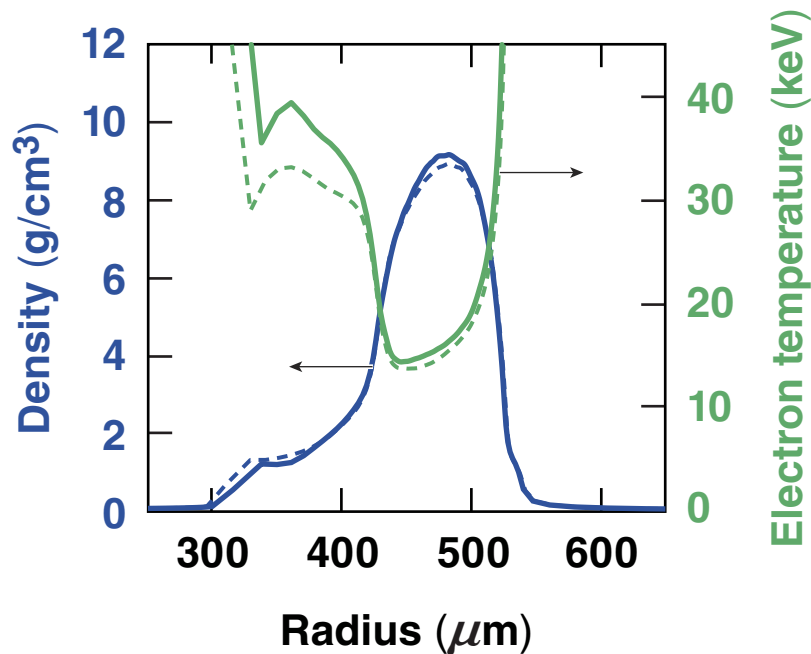
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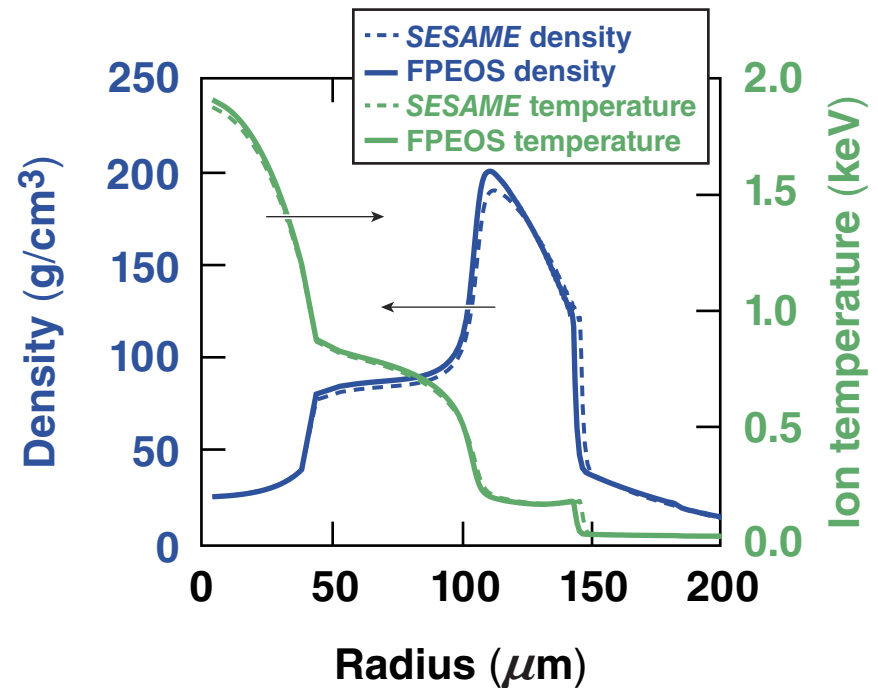
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# The effects of beryllium FPEOS on HED plasma simulations

$t = 10$  ns, end of acceleration



$t = 12$  ns, peak compression



The FPEOS simulation also predicted a higher total neutron yield ( $y = 3.76 \times 10^{14}$ ), which is  $\sim 15\%$  higher than the *SESAME* simulation, ( $y = 3.28 \times 10^{14}$ ).