The Effects of Beam Geometry and Polarization on Two-Plasmon Decay Driven by Multiple Laser Beams



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For multibeam two-plasmon decay (TPD), the absolute instability usually dominates

- The TPD threshold is sensitive to the number, orientation, and polarization of the beams
- For two beams polarized out of their common plane a collective absolute mode near k = 0 dominates TPD
- For two beams polarized in their common plane there are two absolute TPD modes; the dominant one depends on angle of incidence
- The thresholds of multi-beam absolute modes decrease with larger incidence angles and increased polarization components in the plane of the common wave

The temporal growth rate for single-beam TPD is maximized on a hyperbola in k space for a single beam



- The hyperbola lies in the plane of polarization
- Different points on the hyperbola correspond to decays occurring at different densities; larger wavev ectors → smaller densities

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A single beam shows maximal convective gain along the hyperbola



The expected convective gain enhancement is seen for two pump beams polarized in their common plane



When the beams are polarized out of their common plane, enhanced gain is seen near the origin



The presence of enhanced gain near the origin raises the possibility of absolute instability there



For a single beam, the absolute TPD threshold (Simon *et al.*) is lower than the Rosenbluth convective threshold

- The Simon threshold is $\eta \equiv \frac{I_{14}L_{\mu}}{233T_{keV}} > 1$
- The Rosenbluth convective gain is $G_R = \frac{2\pi\gamma_0^2}{\kappa' V_1 V_2} = \frac{I_{14}L_{\mu}}{53.6T_{keV}} \approx 4.35 \eta$

- The nominal convective threshold is $G_R > 2\pi$, or $\eta > \frac{2\pi}{4.35} \approx 1.44$
- Therefore, the absolute instability appears below the convective instability threshold; this, in general, remains true for multiple beams

Fourier analysis of the time-dependent TPD equations results in a set of first-order linear equations that can be integrated numerically

$$\frac{du}{dk_{x}} = \frac{1}{2} e^{i\alpha\beta^{1/2}k_{01x}(k_{x}-k_{xr})^{2}} \left(\frac{k^{2}-k_{d1}^{2}}{kk_{d1}}\right) \alpha_{1}(\hat{\varepsilon}_{1}\cdot k) u_{d1}$$
$$+ \frac{1}{2} e^{i\alpha\beta^{1/2}k_{02x}(k_{x}-k_{xr})^{2}} \left(\frac{k^{2}-k_{d2}^{2}}{kk_{d2}}\right) \alpha_{2}(\hat{\varepsilon}_{2}\cdot k) u_{d2}$$

$$\frac{\mathrm{d}u_{\mathrm{d}1}}{\mathrm{d}k_{\mathrm{x}}} = -\frac{1}{2} \mathrm{e}^{-i\alpha\beta^{1/2}k_{01\mathrm{x}}(k_{\mathrm{x}}-k_{\mathrm{x}r})^{2}} \left(\frac{k^{2}-k_{\mathrm{d}1}^{2}}{k_{\mathrm{d}1}}\right) \alpha_{1}\left(\hat{\varepsilon}_{1}\cdot k\right) u$$

$$\frac{\mathrm{d}u_{\mathrm{d}2}}{\mathrm{d}k_{\mathrm{x}}} = -\frac{1}{2} \mathrm{e}^{-i\alpha\beta^{1/2}k_{02\mathrm{x}}(k_{\mathrm{x}}-k_{\mathrm{xr}})^{2}} \left(\frac{k^{2}-k_{\mathrm{d}2}^{2}}{kk_{\mathrm{d}2}}\right) \alpha_{2} \left(\hat{\varepsilon}_{2} \cdot k\right) u$$

where
$$\alpha_{i} = \frac{4k_{0} |\upsilon_{0i}|}{\omega_{0}} k_{0}L$$
 and $\beta_{i} = \frac{9\upsilon_{e}^{4}k_{0}^{2}}{|\upsilon_{0i}|^{2}\omega_{0}^{2}}$

Numerical integration of these equations gives spatial gain; divergent gain indicates absolute threshold



The absolute threshold for TPD depends on the angle of incidence and polarization





• As a result, the threshold is reduced by $\sim \cos^2 \theta$

For two *p*-polarized beams, the interaction regions are separated, reducing synergy



• For two s-polarized beams the separation is much smaller

For two *p*-polarized beams, an on-axis absolute mode with $k_v = 0$ has the lowest threshold at larger angles



At larger angles, the on-axis mode is closer to the hyperbolas than the off-axis modes





 $\theta = 40^{\circ}$

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With more beams, the absolute TPD threshold for the on-axis mode is quite sensitive to the cone angle



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