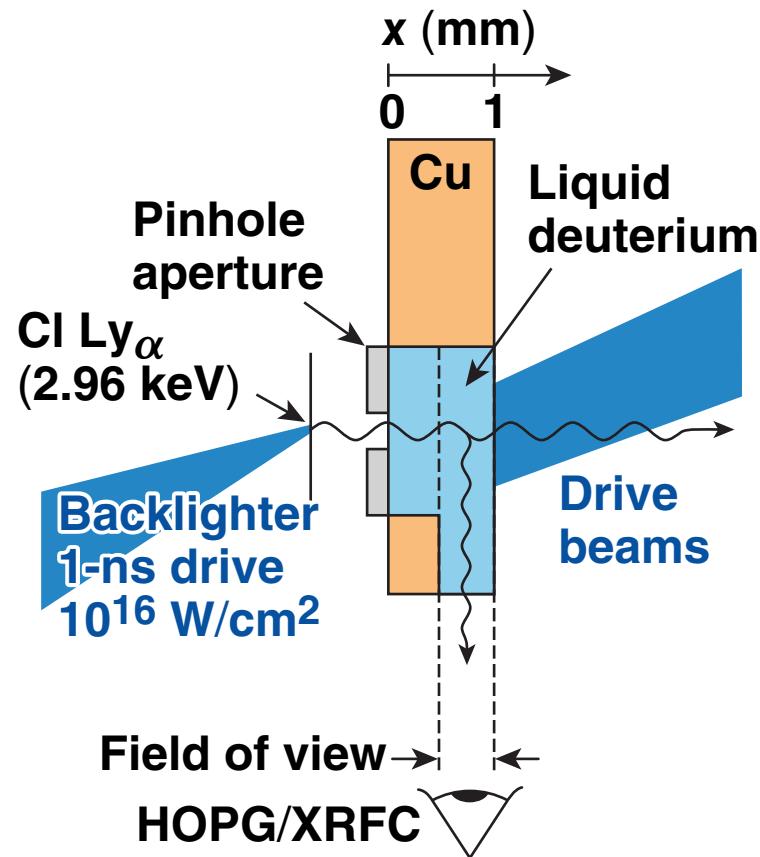


X-Ray Thomson-Scattering: Incisive Probe for Warm, Dense Matter



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Omega Laser Facility
Users' Group Workshop
Rochester, NY
25–27 April 2012

Summary

X-ray Thomson scattering* (XRTS) is a primary diagnostic for high-energy-density-physics (HEDP) experiments



- The conditions of dense plasmas are probed with XRTS, especially warm, dense matter with $T_e \sim T_F$ and the ratio of potential energy to kinetic energy of the ions greater than unity
- The elastic and inelastic x-ray scattering features are spectrally resolved
 - scattering from electrons (noncollective) $\rightarrow T_e, Z$
 - scattering from plasmons (collective) $\rightarrow n_e$
 - elastic scattering \rightarrow structure of matter
- A review of XRTS experiments is presented
 - inertial confinement fusion (ICF)
 - radiation heated
 - shock heated and compressed
 - proton heated
 - laboratory astrophysics (radiative shocks and planetary interiors)

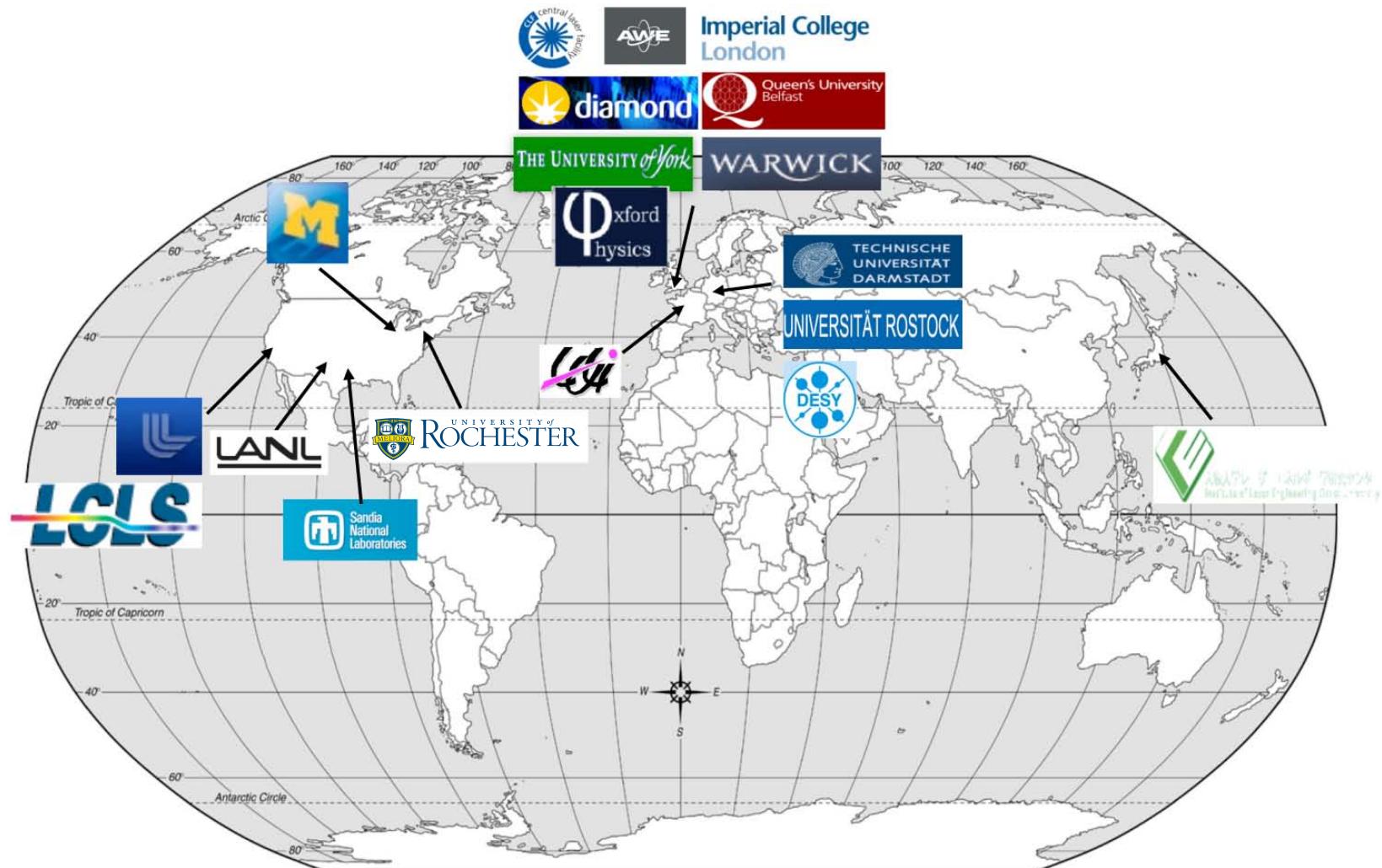
Many XRTS experiments need spatially resolved spectral measurements.

*S. H. Glenzer and R. Redmer, Rev. Mod. Phys. 81, 1625 (2008).

O. L. Landen *et al.*, J. Quant. Spectrosc. Radiat. Transf. 71, 465 (2001).

G. Gregori *et al.*, Phys. Rev. E 67, 026412(2003).

XRTS research involves many international collaborations



Collaborators for XRTS experiment of shocked liquid deuterium



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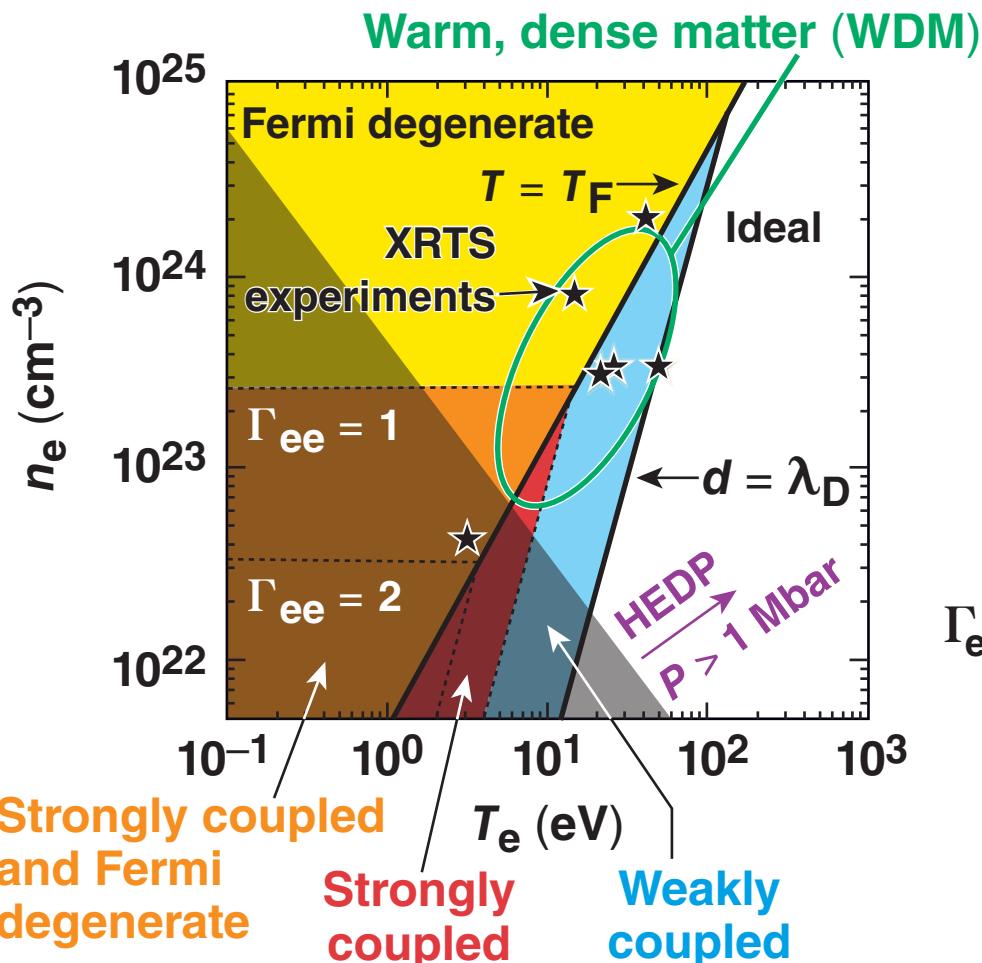
***also, Depts. of Mechanical Engineering and Physics and Astronomy,
University of Rochester**

Outline



- **Motivation**
- **XRTS**
- **Experiments**
- **Future**

X-ray Thomson scattering* (XRTS) is a primary diagnostic for HEDP experiments



Degeneracy parameter

$$\theta = \frac{k_B T_e}{E_F}, \quad E_F = \frac{\hbar^2}{2m_e} (3\pi^2 n_e)^{2/3}$$

Coupling parameter

$$\Gamma_{ee} = \frac{e^2}{4\pi\epsilon_0 dk_B T_e}, \quad d = \left(\frac{4\pi n_e}{3}\right)^{-1/3}$$

Penetrating x rays are needed to probe dense plasmas.

*S. H. Glenzer and R. Redmer, Rev. Mod. Phys. 81, 1625 (2008).

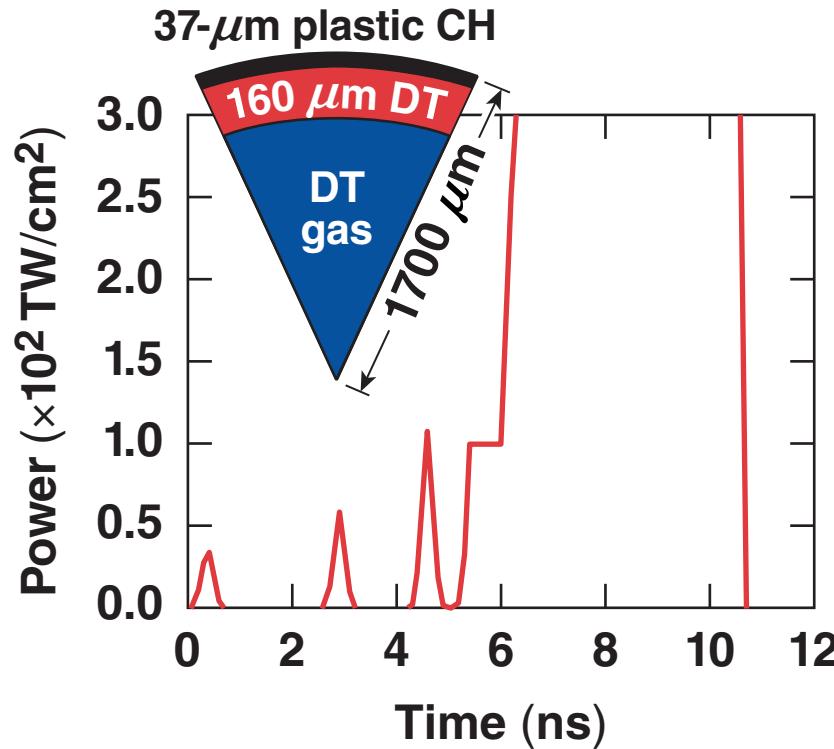
O. L. Landen et al., J. Quant. Spectrosc. Radiat. Transf. 71, 465 (2001).

G. Gregori et al., Phys. Rev. E 67, 026412(2003).

Diagnosing the shock heated and compressed-shell conditions in an ICF target is an ideal XRTS application



Triple-picket, direct-drive design for the NIF¹



$$E_{\min} \sim \frac{\alpha^{1.8}}{V_{\text{imp}}^{5.8}}$$

(minimum laser energy for ignition)^{2,3}

$$\text{Minimize} \rightarrow \alpha = \frac{P_{\text{fuel}}}{P_{\text{Fermi}}}$$

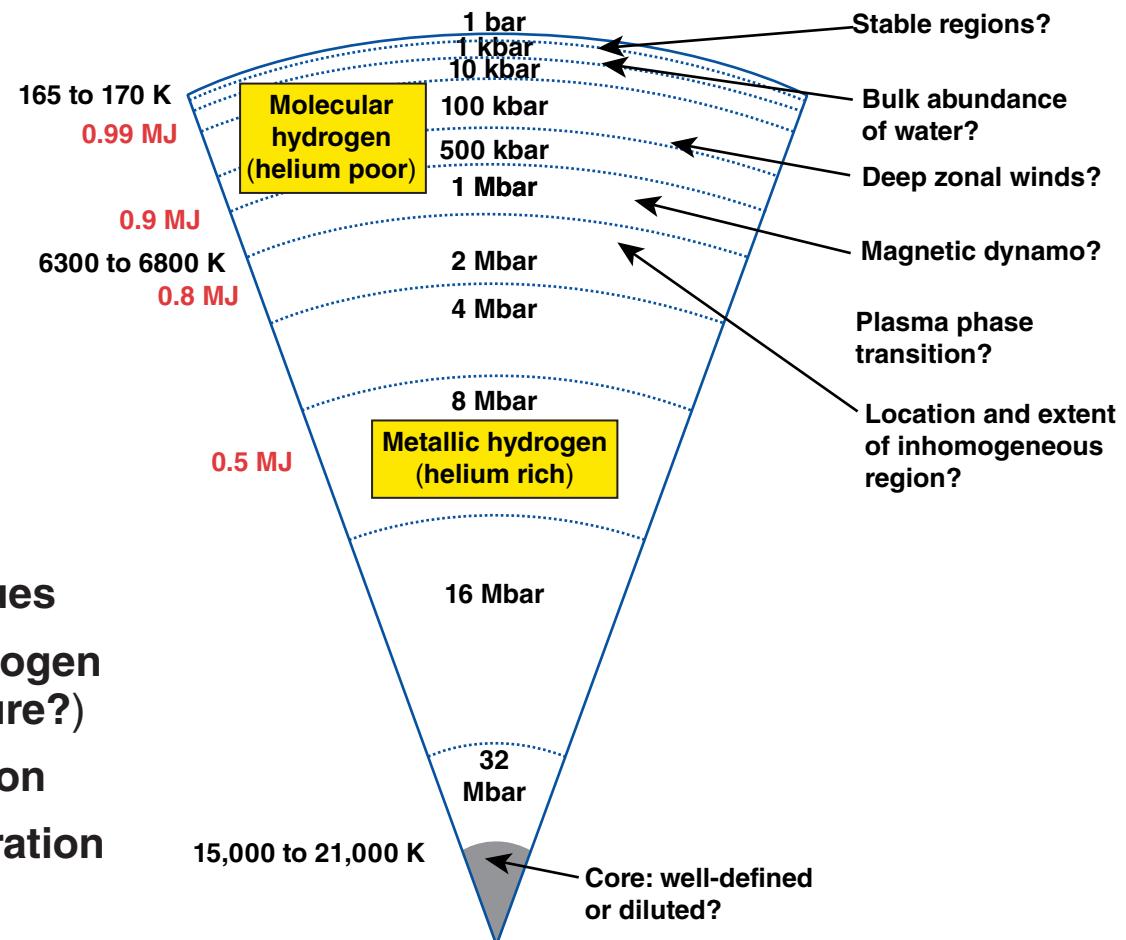
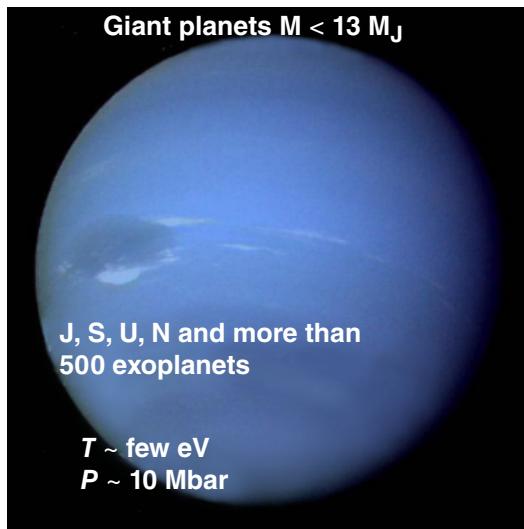
The minimum energy needed for ignition depends on the plasma conditions in the fuel layer (P_{fuel}), which can be diagnosed with XRTS.

¹V. N. Goncharov et al., Phys. Rev. Lett. **104**, 165001 (2010).

²M. C. Hermann, M. Tabak, and J. Lindl, Nuc. Fusion **41**, 99 (2001).

³R. Betti et al., Phys. Plasmas **9**, 2277 (2002).

XRTS can provide accurate measurements of conditions in HEDP plasmas designed for planetary science



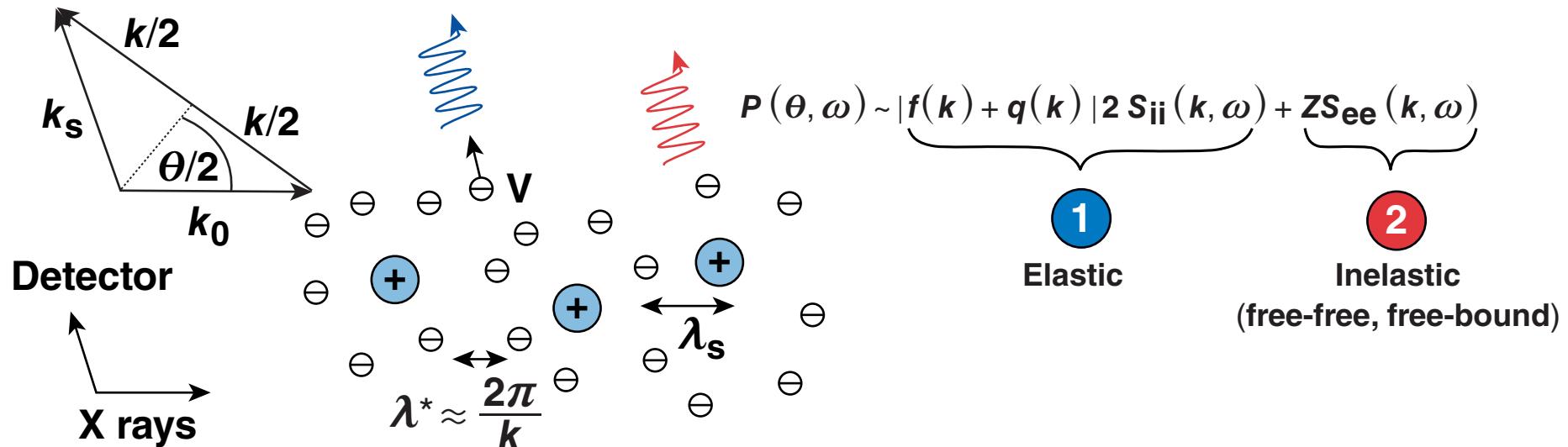
- Outstanding physics issues
 - metallization of hydrogen (pressure, temperature?)
 - H/He phase separation
 - magnetic-field generation

Outline



- Motivation
- XRTS
- Experiments
- Future

The properties of dense plasmas can be diagnosed with XRTS



1 Ion–ion correlations

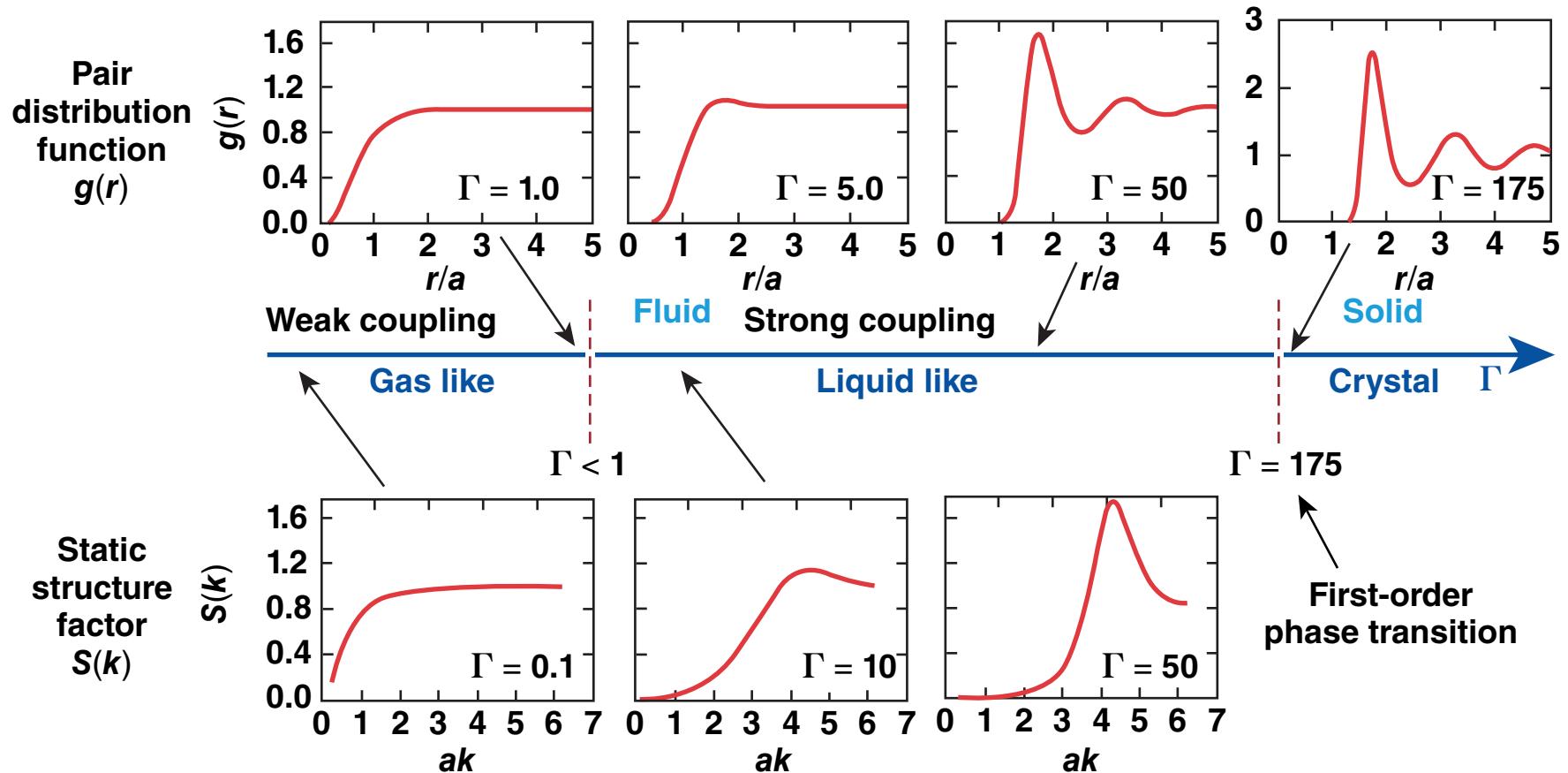
- scattering as a function of the scattering angle exhibit peaks that are representative of the structure (diffraction)

2 Electron–electron correlations

- scattering as a function of energy describes either single-particle dynamics or correlated plasmons

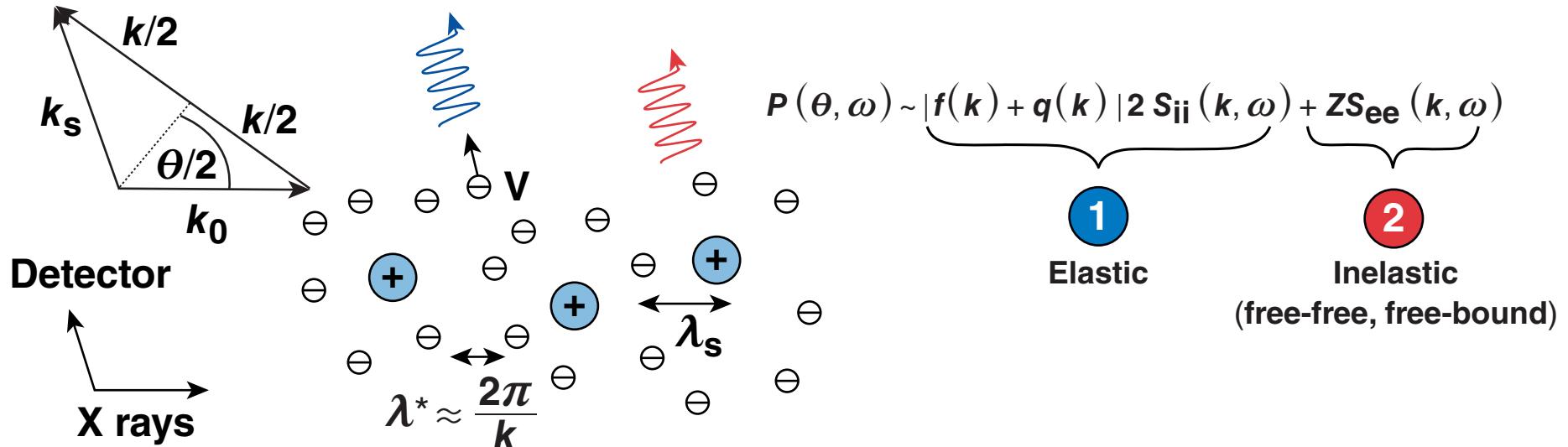
S. H. Glenzer and R. Redmer, Rev. Mod. Phys. 81, 1625 (2008).
G. Gregori et al., Phys. Rev. E 67, 26412 (2003).

Elastic scattering (diffraction) is used to test the degree of correlation between ions



- As the plasma gets denser (or cooler) correlations among ions emerge with the formation of crystalline structure

The microscopic characterization of warm, dense matter is inferred from inelastic scattering



1 Ion–ion correlations

- scattering as a function of the scattering angle exhibit peaks that are representative of the structure (diffraction)

2 Electron–electron correlations

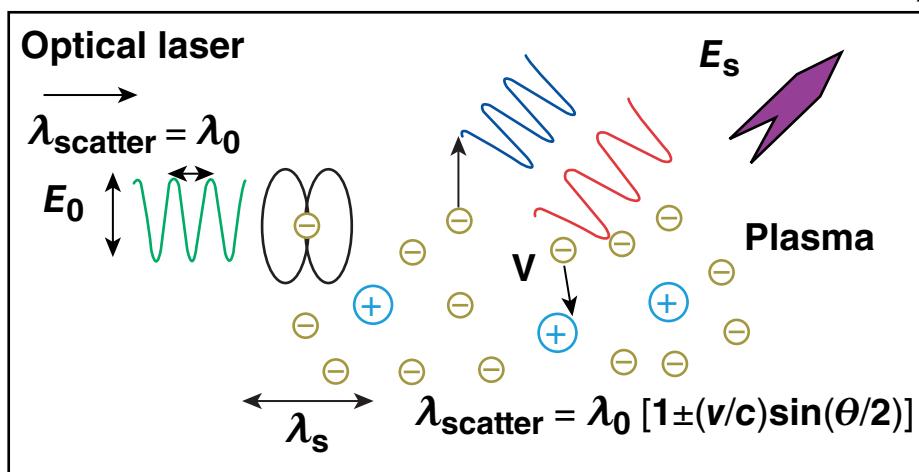
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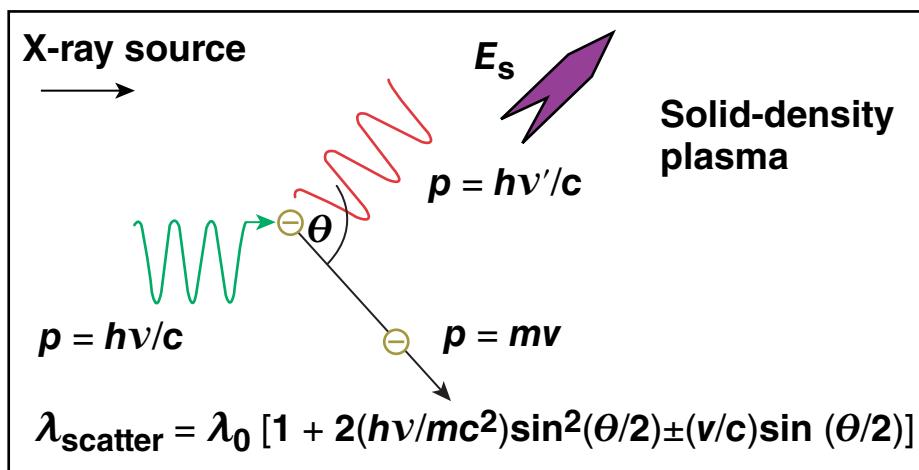
Inelastic x-ray scattering probes the electron-velocity distribution function



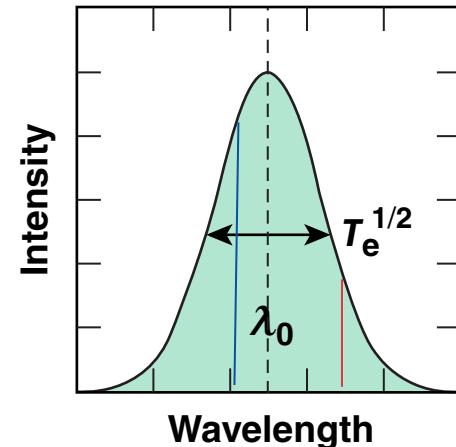
Noncollective Thomson scattering ($\lambda_{\text{scatter}} < \lambda_s$)



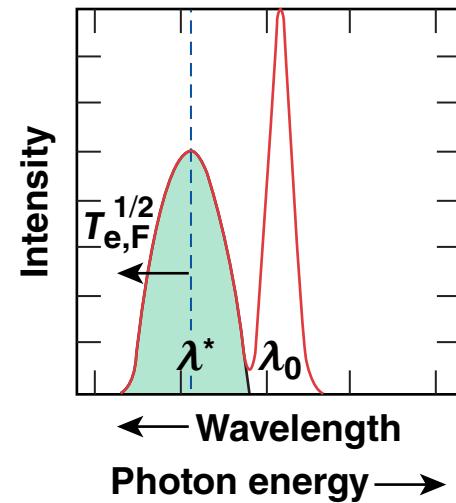
X-ray Compton scattering



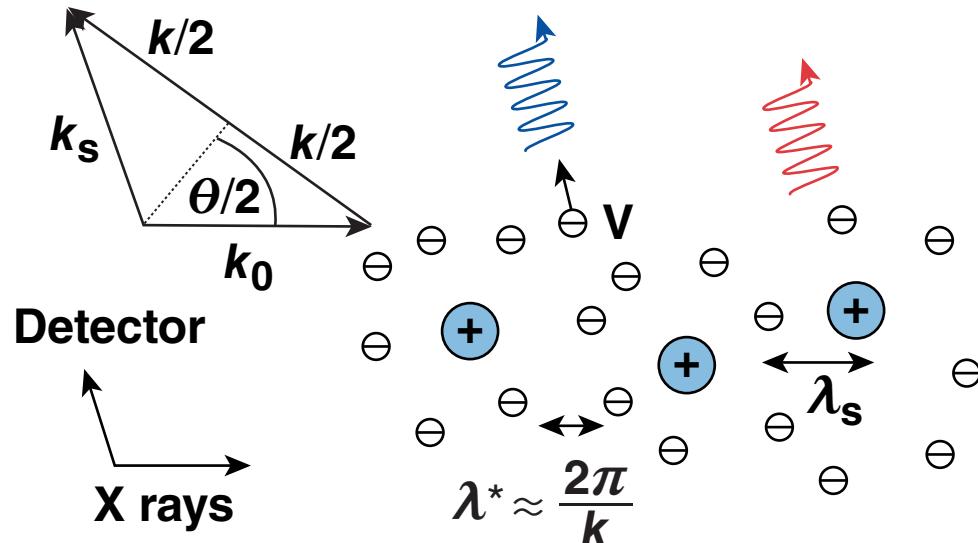
Boltzmann distribution



Compton peak Raleigh peak



Collective inelastic x-ray scattering is used to diagnose the electron density



Scattering parameter

$$\alpha = \frac{\lambda^*}{2\pi \lambda_s} = \frac{1}{k \lambda_s}$$

$$k = |\mathbf{k}| = 4\pi \frac{E_0}{hc} \sin(\theta/2)$$

- $\alpha < 1$: scattering from electrons (noncollective) $\rightarrow T_e, Z$
- $\alpha > 1$: scattering from plasmons (collective) $\rightarrow n_e$

$$\omega_{\text{plasmon}} \sim \omega_{\text{plasma}} = \sqrt{n_e e^2 / \epsilon_0 m_e}$$

Outline

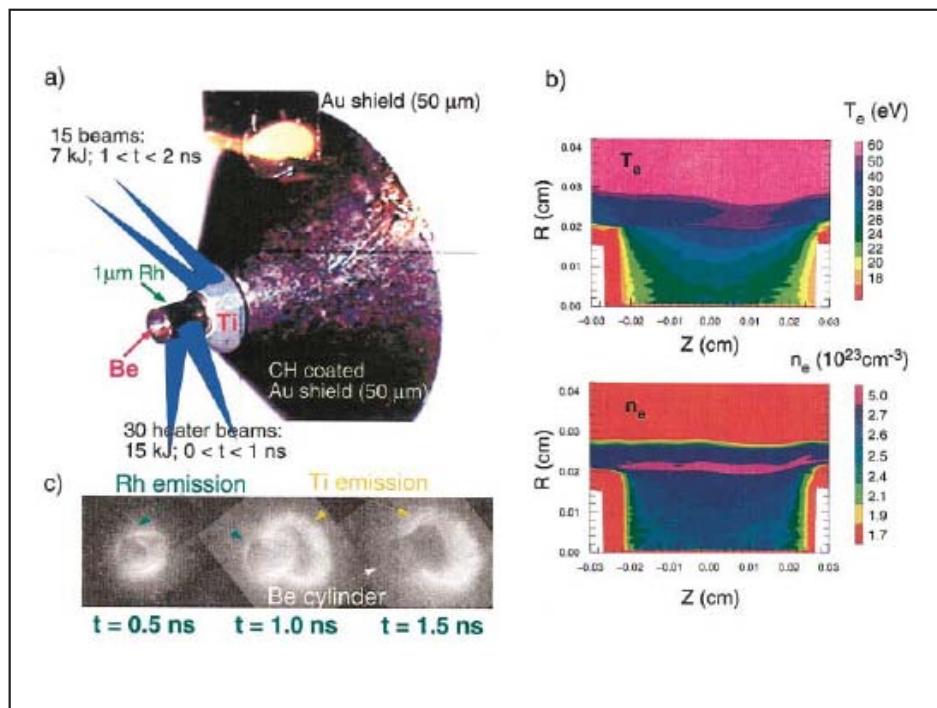


- Motivation
- XRTS
- **Experiments**
- Future

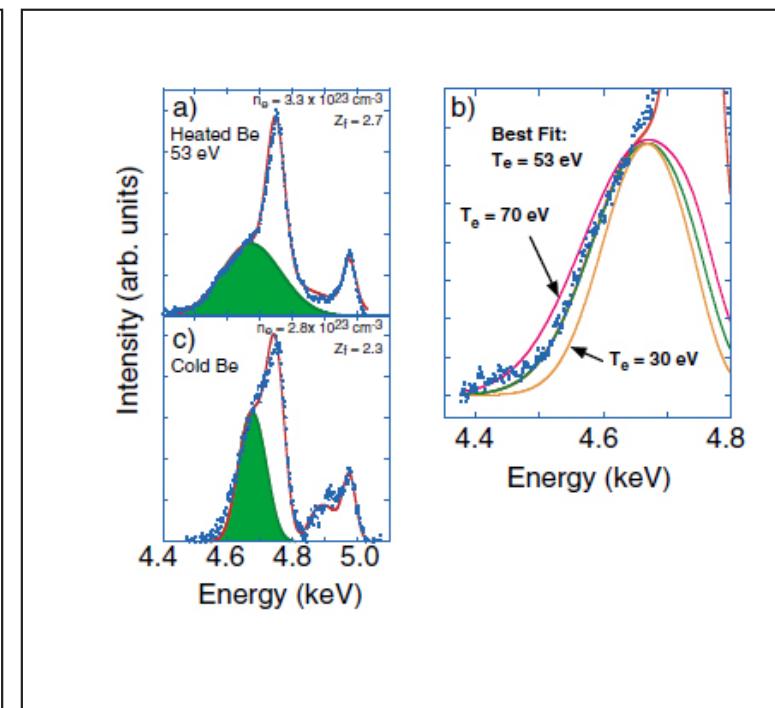
The first noncollective XRTS experiment diagnosed isochorically heated Be on OMEGA



Noncollective scattering from radiatively heated Be



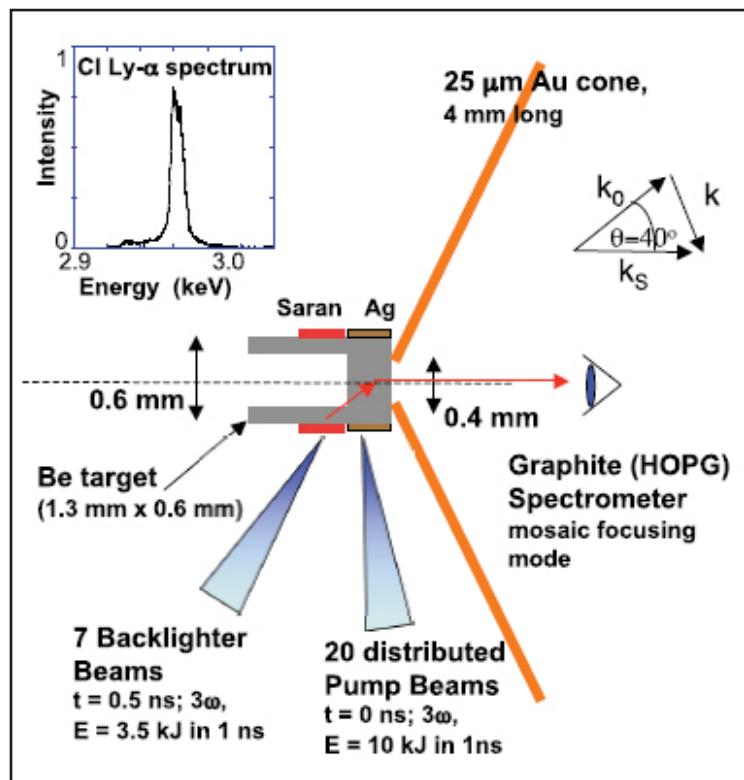
Noncollective XRTS



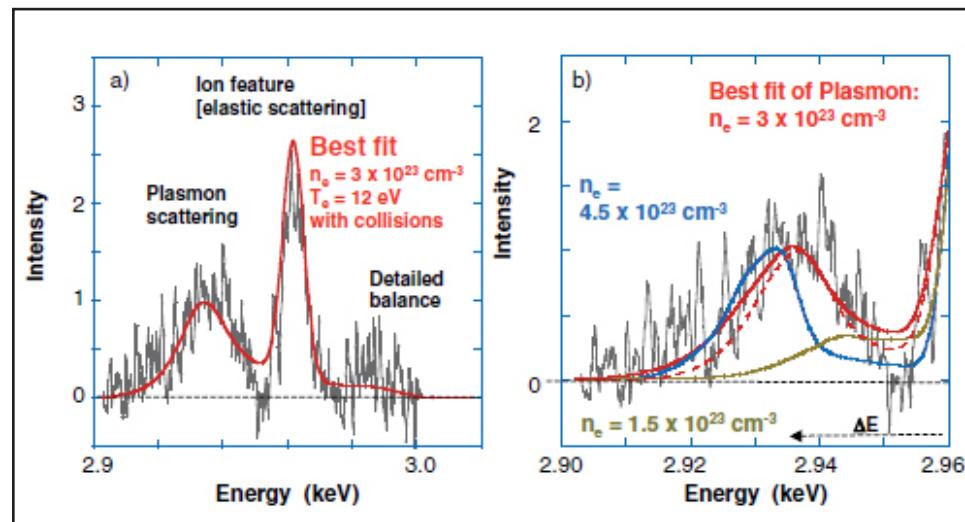
Collective XRTS was first demonstrated using isochorically heated Be on OMEGA



Collective scattering from radiatively heated Be



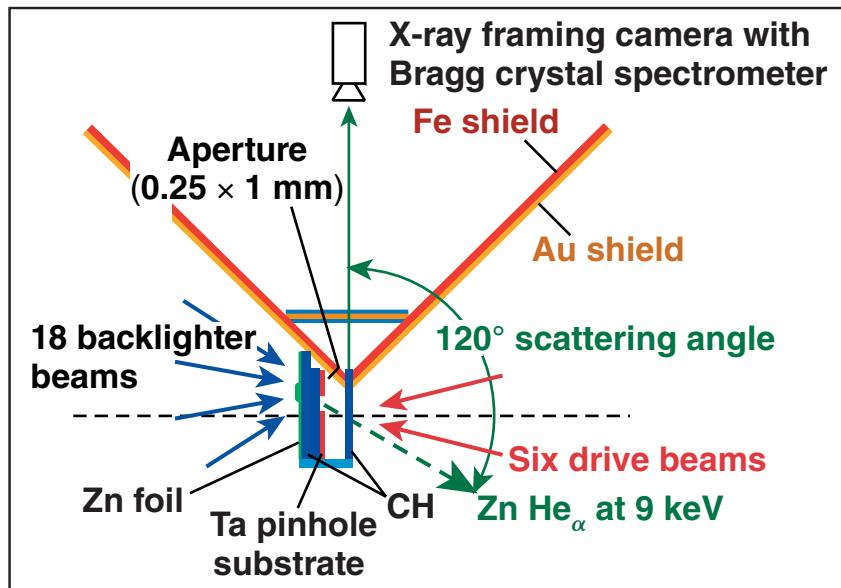
Collective XRTS



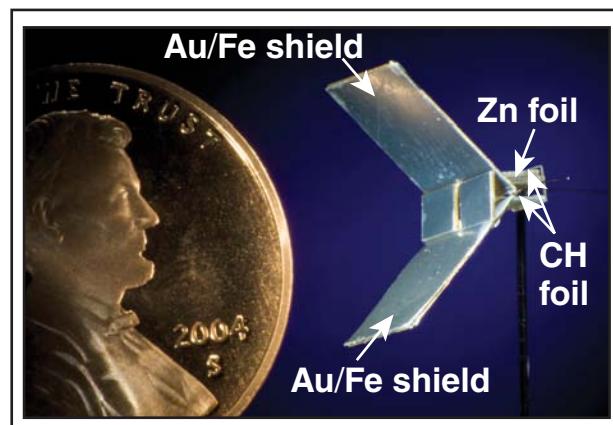
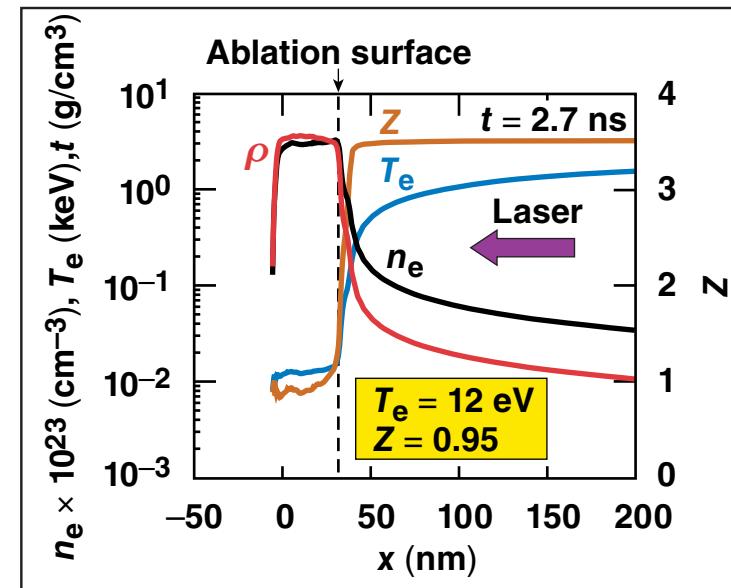
XRTS was applied to a laser-driven, shock-heated, and compressed CH foil on OMEGA



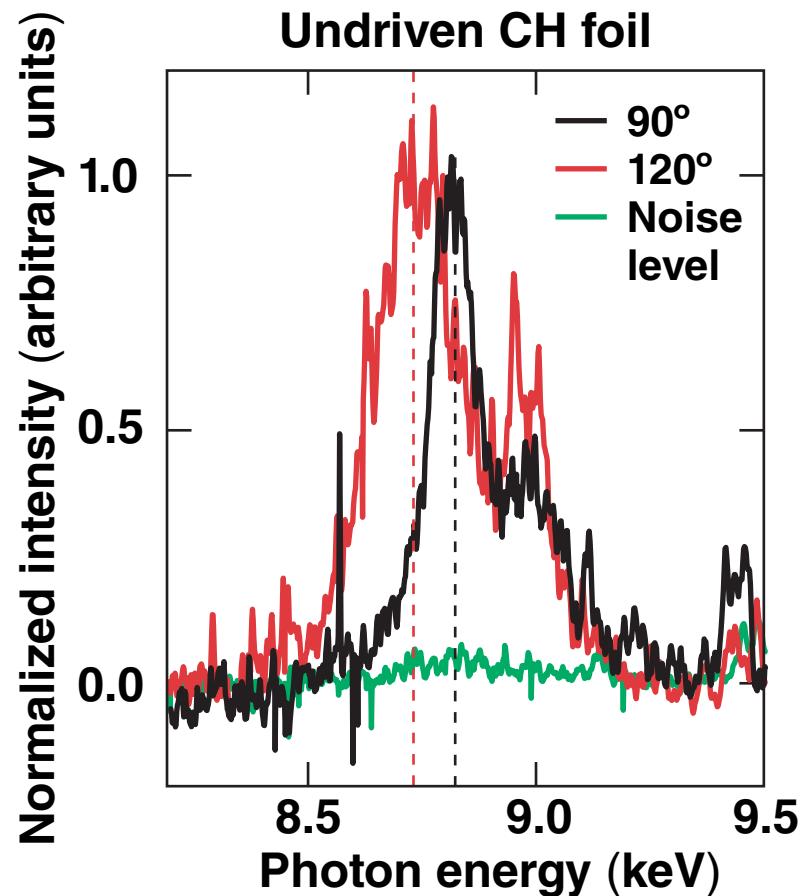
Shock heated and compressed CH foil



1-D simulation



The Compton energy downshift is observed to be greater for the larger scattering angle



Compton downshifted energy (eV)

$$\Delta E_C = \frac{\hbar^2 k^2}{2m_e} \quad k = \frac{4\pi}{\lambda_0} \sin\left(\frac{\theta}{2}\right)$$

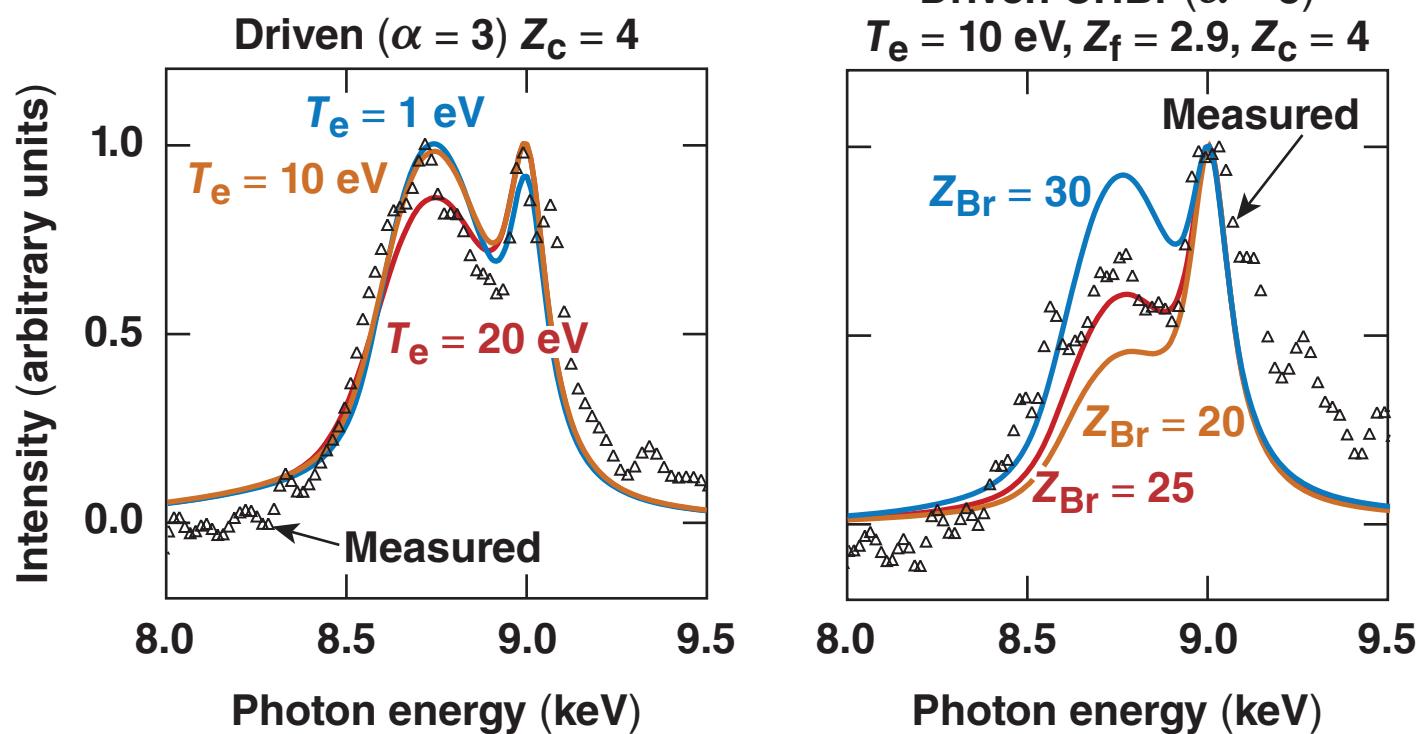
θ : scattering angle

λ_0 : wavelength of probe

(Zn He α $\sim 1.3 \text{ \AA} \sim 9.0 \text{ keV}$)

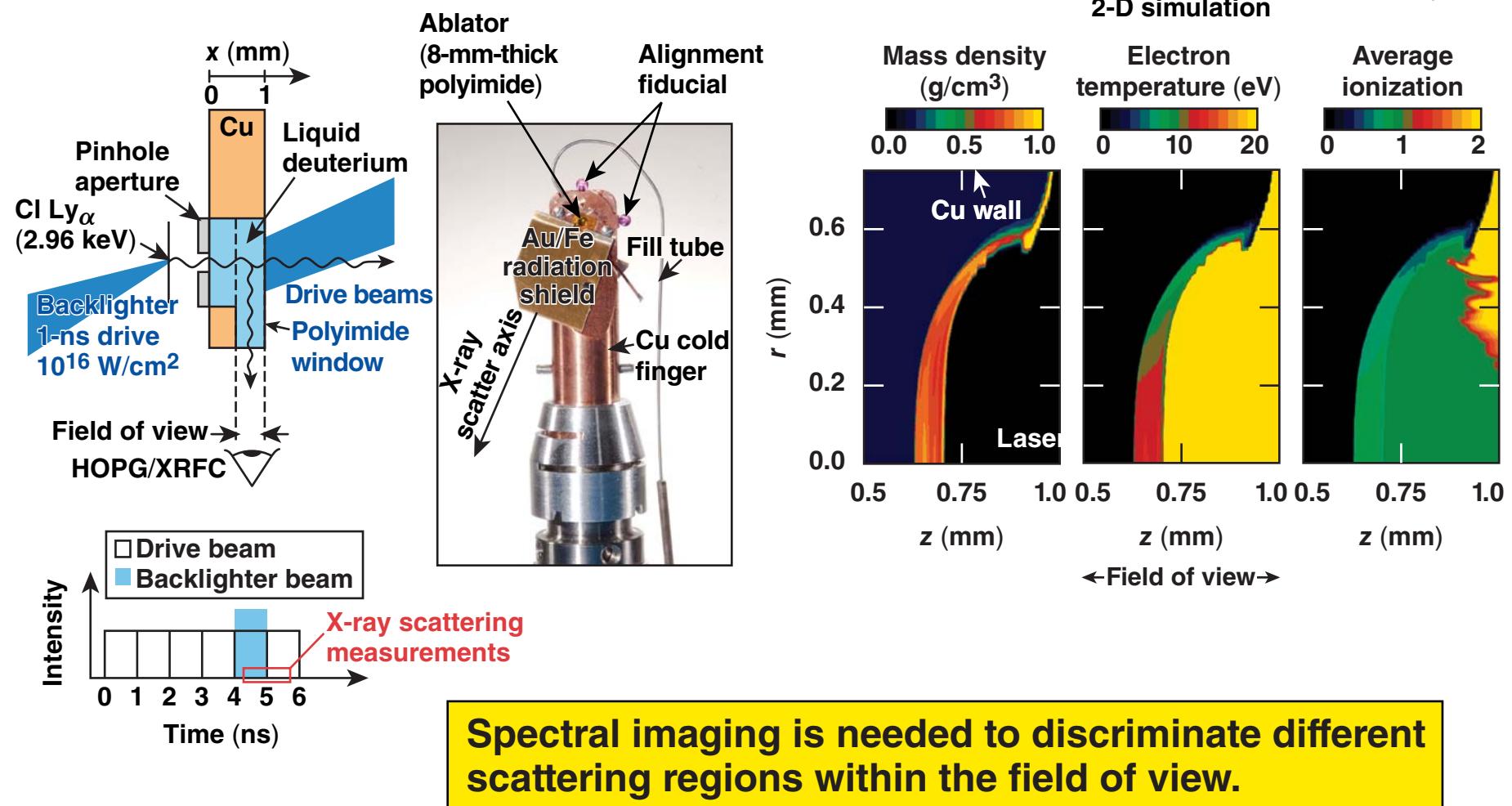
$\Delta E_C = 178 \text{ eV (for } 90^\circ)$
 $= 267 \text{ eV (for } 120^\circ)$

An electron temperature of 10 to 20 eV was inferred from noncollective XRTS

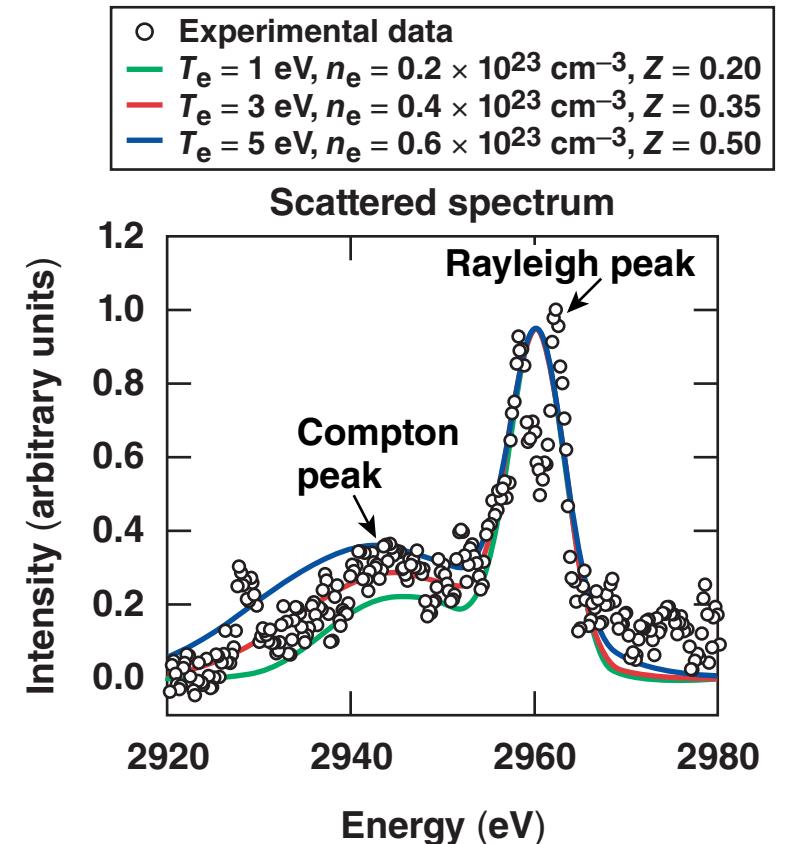
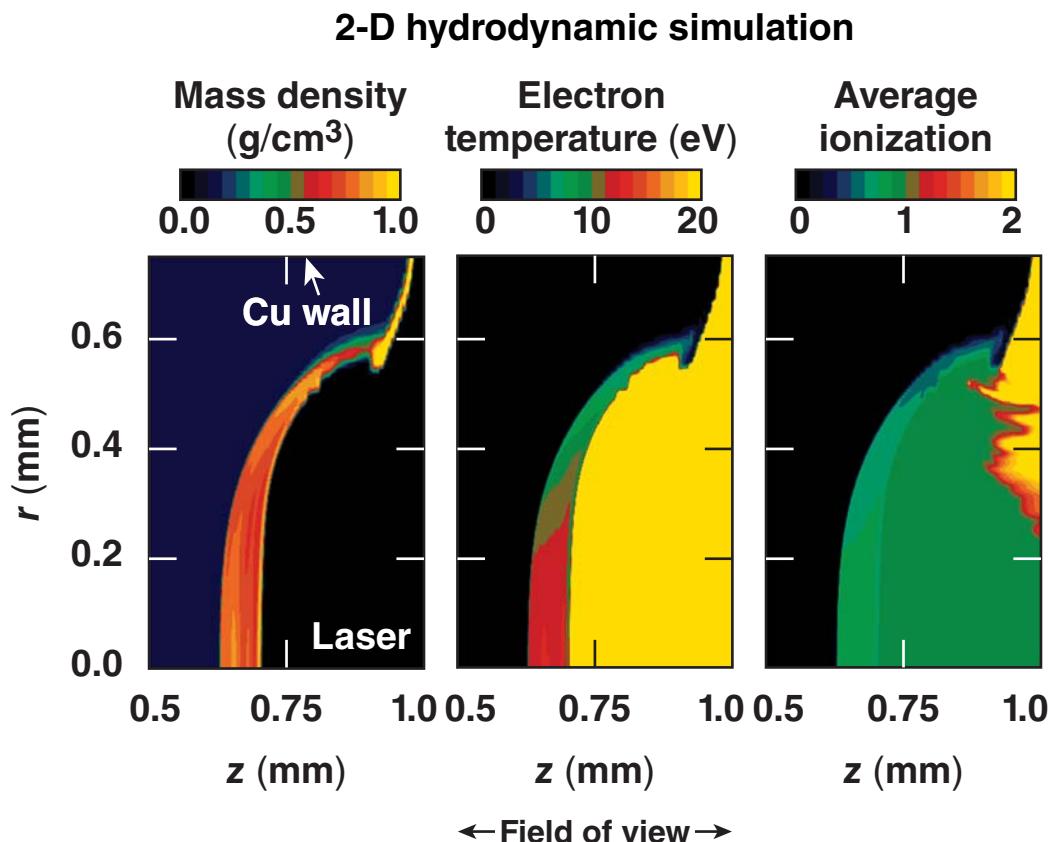


Br dopant in plastic increased the elastic scattering component.

Shocked liquid deuterium is being studied with noncollective XRTS to infer the electron temperature



Shocked liquid deuterium is being studied with noncollective XRTS to infer the electron temperature

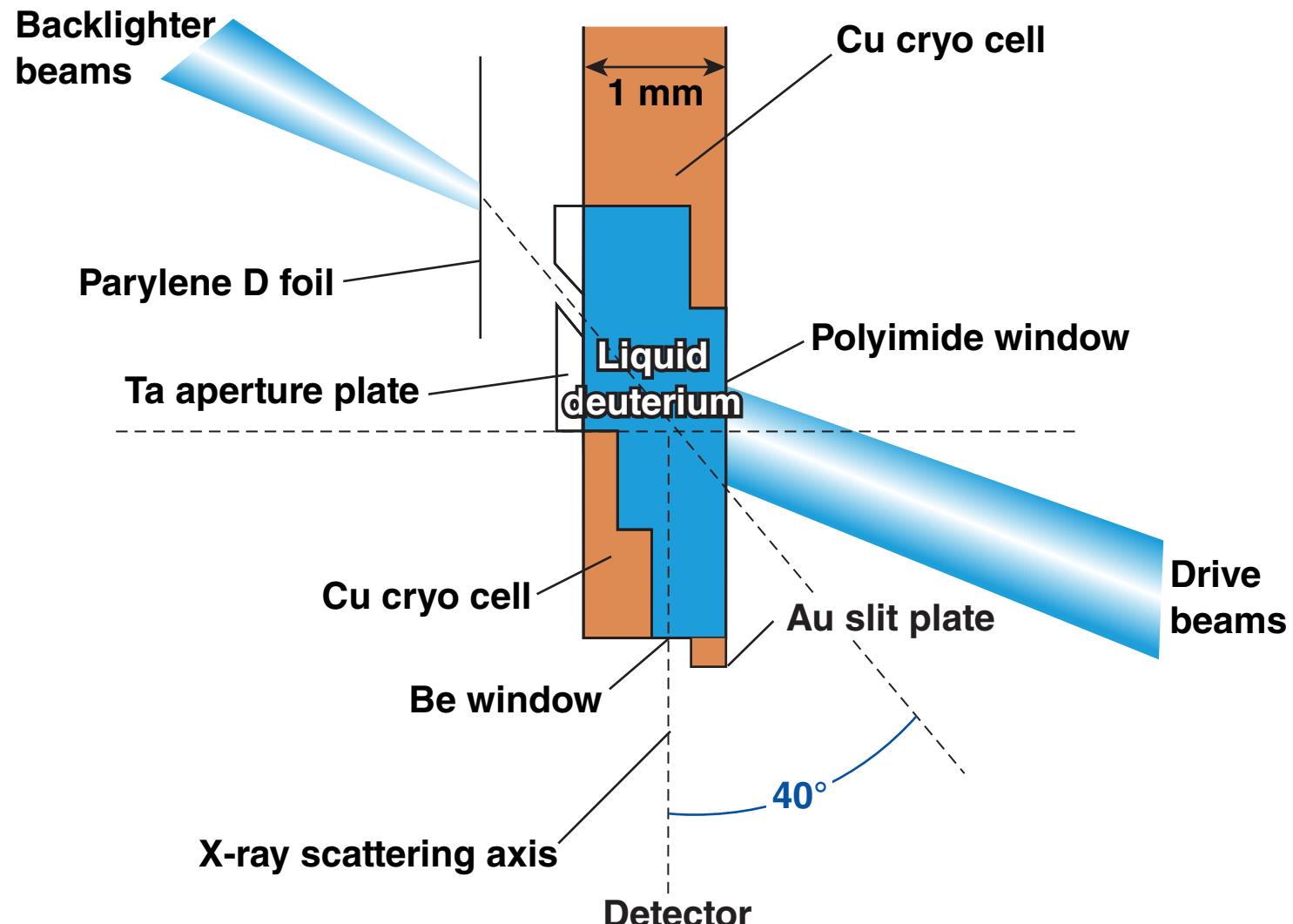


Spectral imaging is needed to discriminate scattering from shocked and unshocked matter.

Scattering angle will be decreased from 90° to 40° to measure collective XRTS and infer electron density



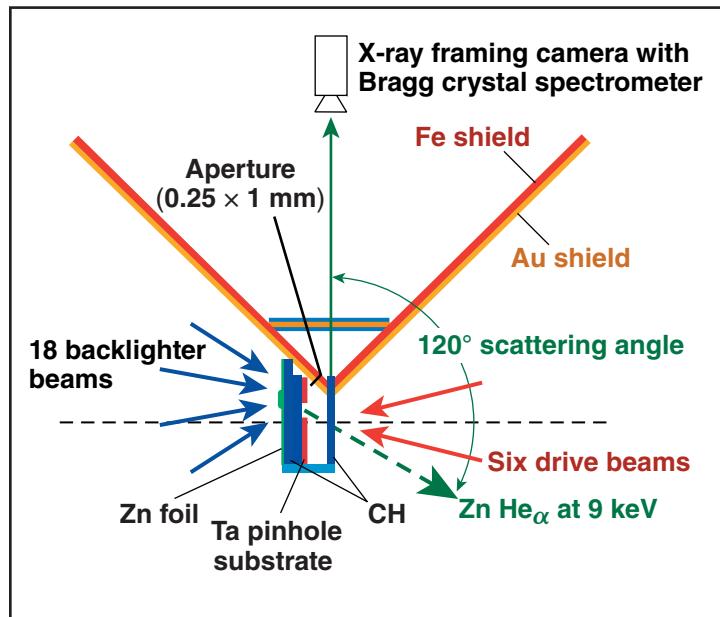
Shocked liquid deuterium experiment with 40° scattering geometry



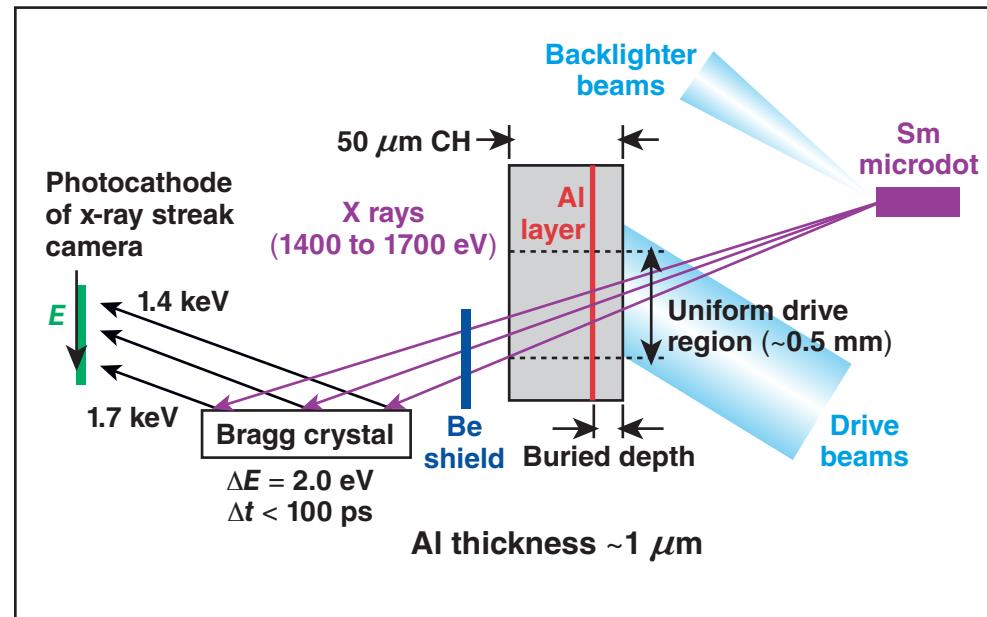
XRTS is a noninvasive probe



XRTS experiment¹



X-ray absorption spectroscopy²



Compared to absorption spectroscopy,
no tracer layers are required for XRTS.

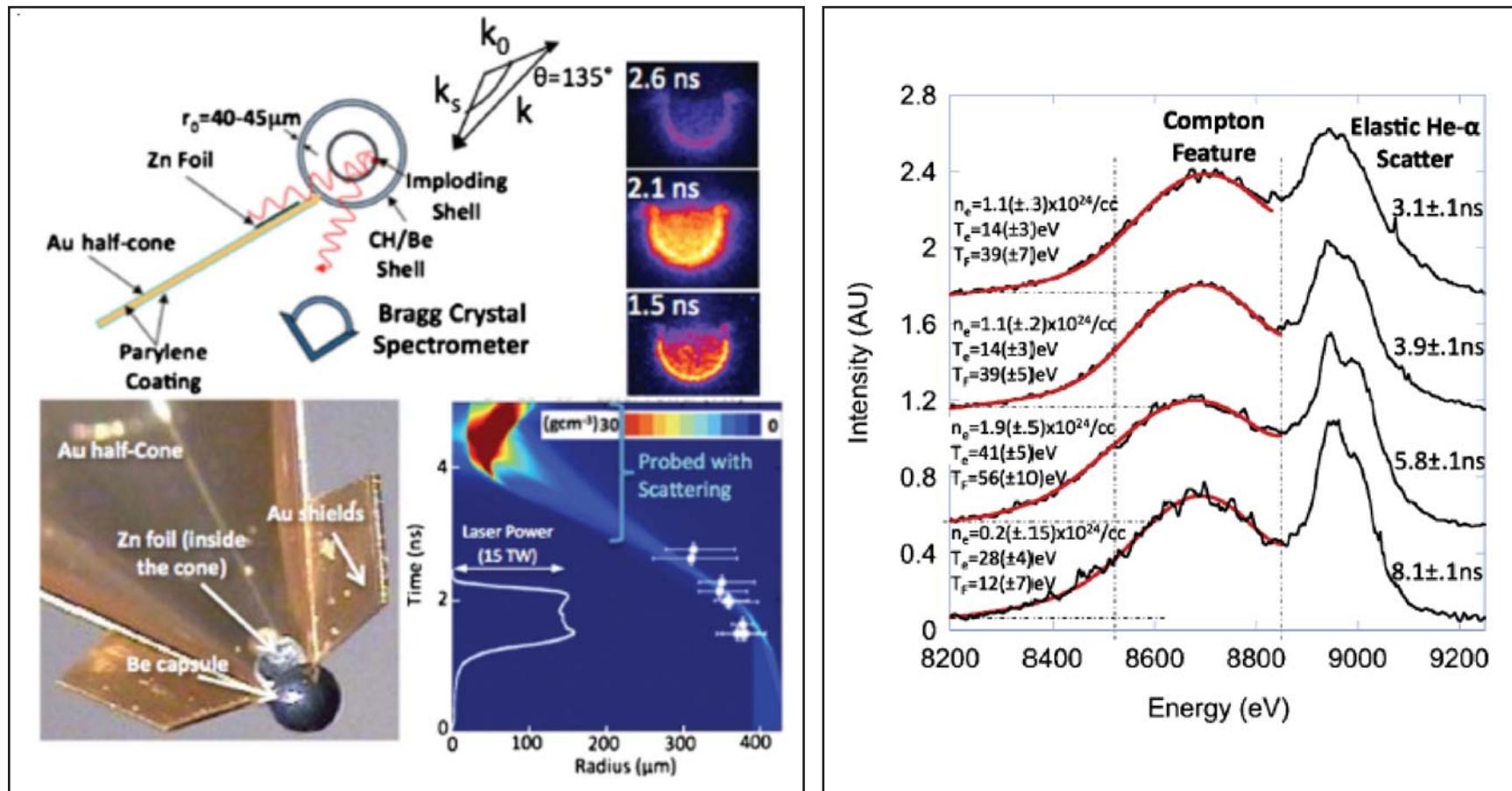
¹H. Sawada, S. P. Regan et al., Phys. Plasmas **14**, 122703 (2007).

²H. Sawada, S. P. Regan et al., Phys. Plasmas **16**, 052702 (2009).

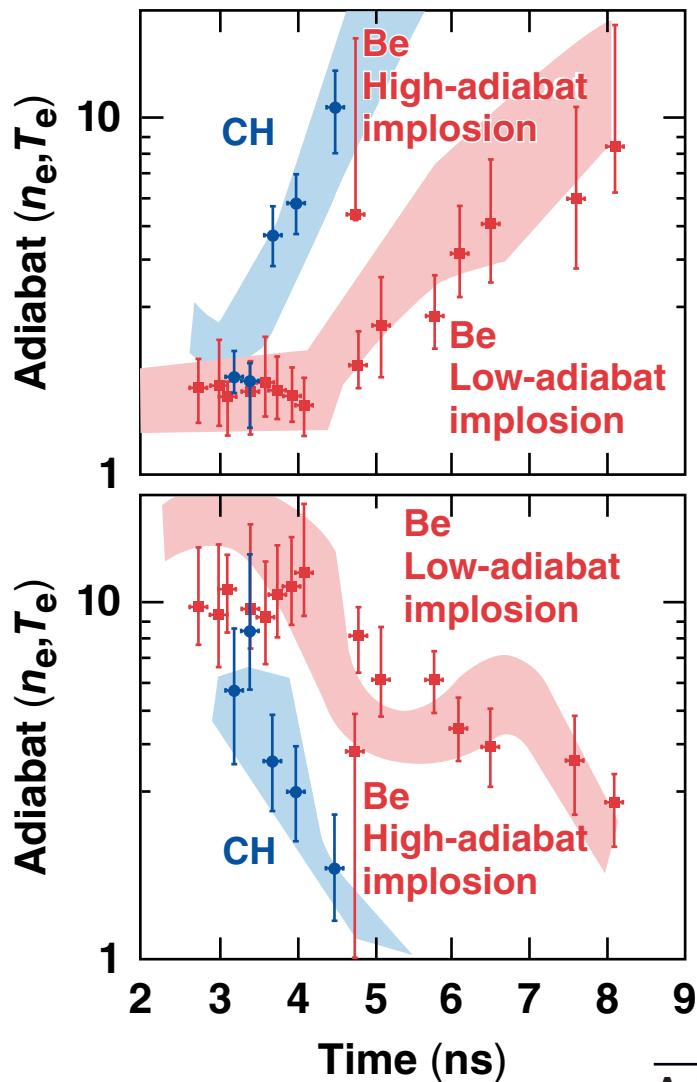
Compressed CH and Be shells were imploded on OMEGA and diagnosed with XRTS



Spherical direct-drive implosion



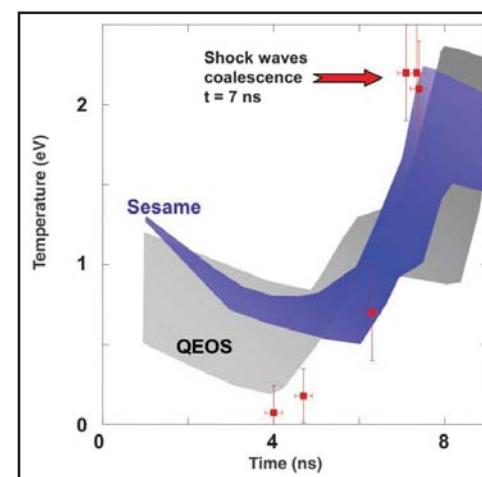
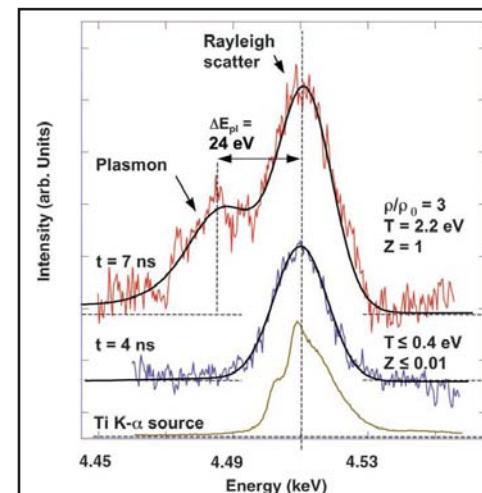
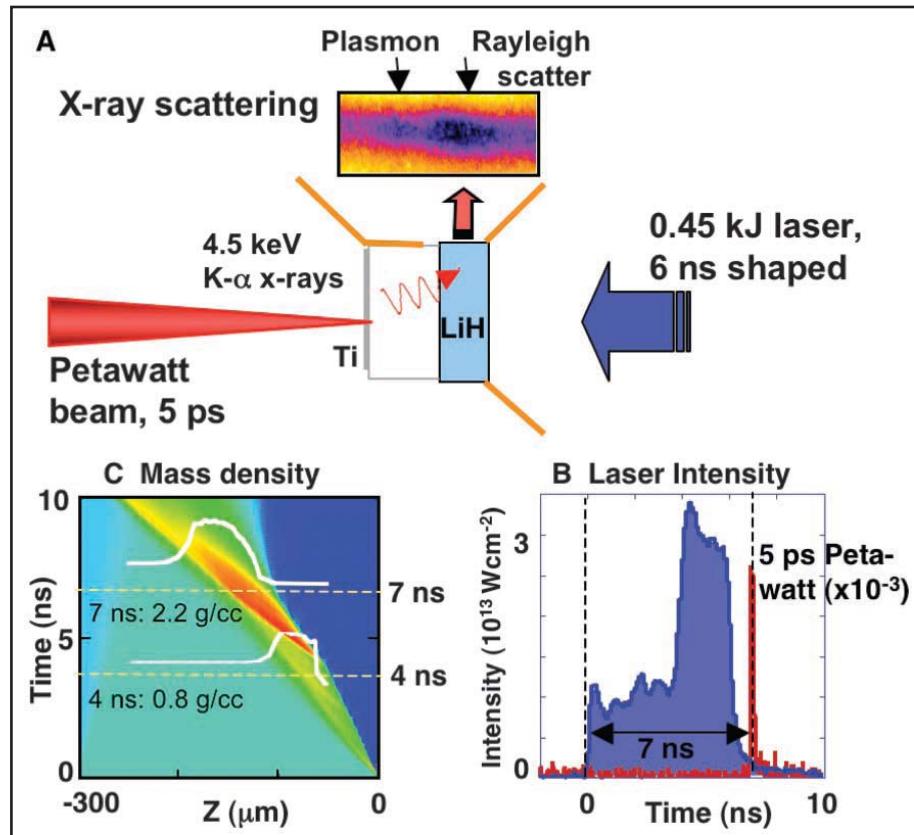
The adiabat and the ion–ion coupling parameter were inferred from the XRTS measurements



$$\alpha = \frac{P_{\text{final}}}{P_{\text{Fermi}}}$$

$$\Gamma_{ii} = (Ze)^2 / (4\pi\epsilon_0 a k_B T_i)$$
$$a = (4/3\pi n_i)^{-1/3}$$

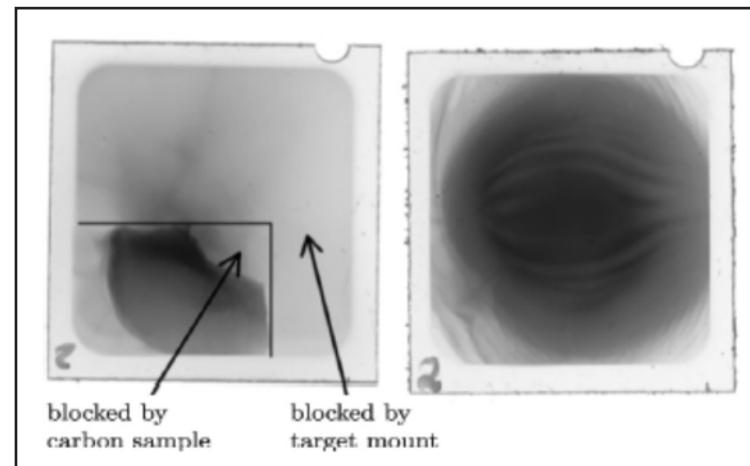
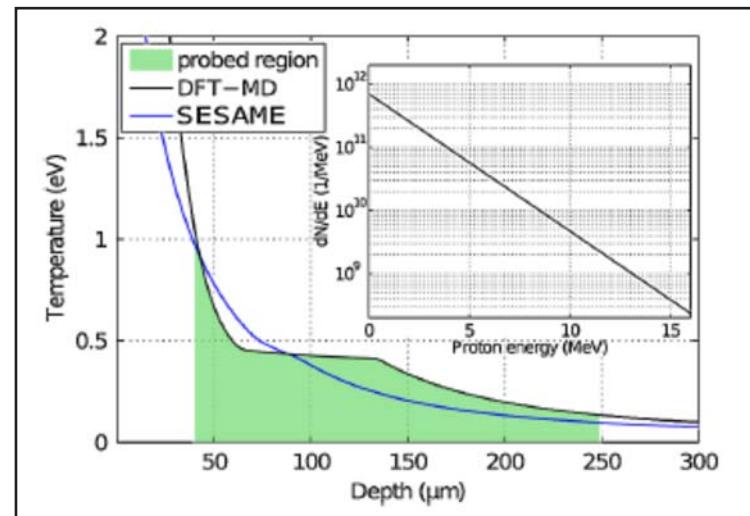
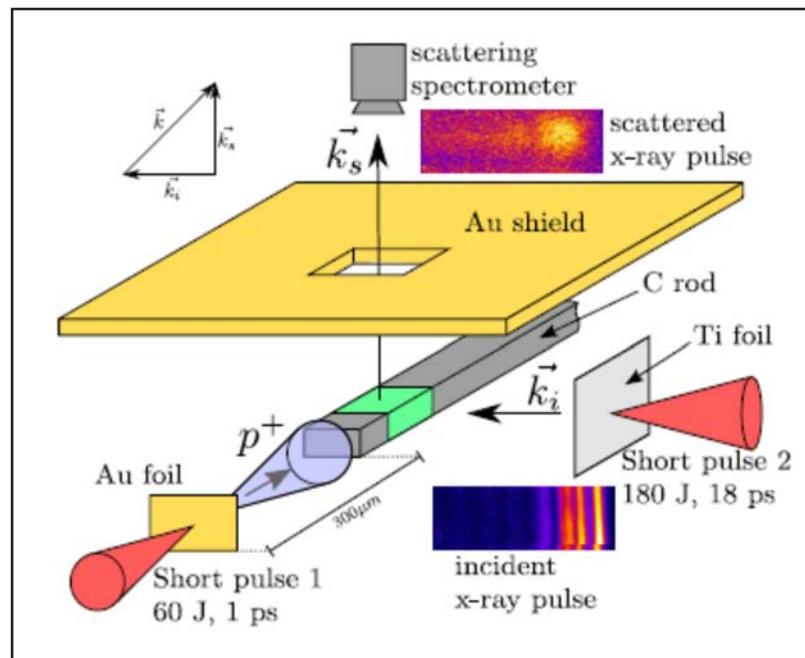
Ultrafast XRTS was used to probe shock-compressed matter



XRTS with 10-ps temporal resolution was achieved.

The amount of liquid carbon was diagnosed with elastic x-ray scattering from a proton-heated target

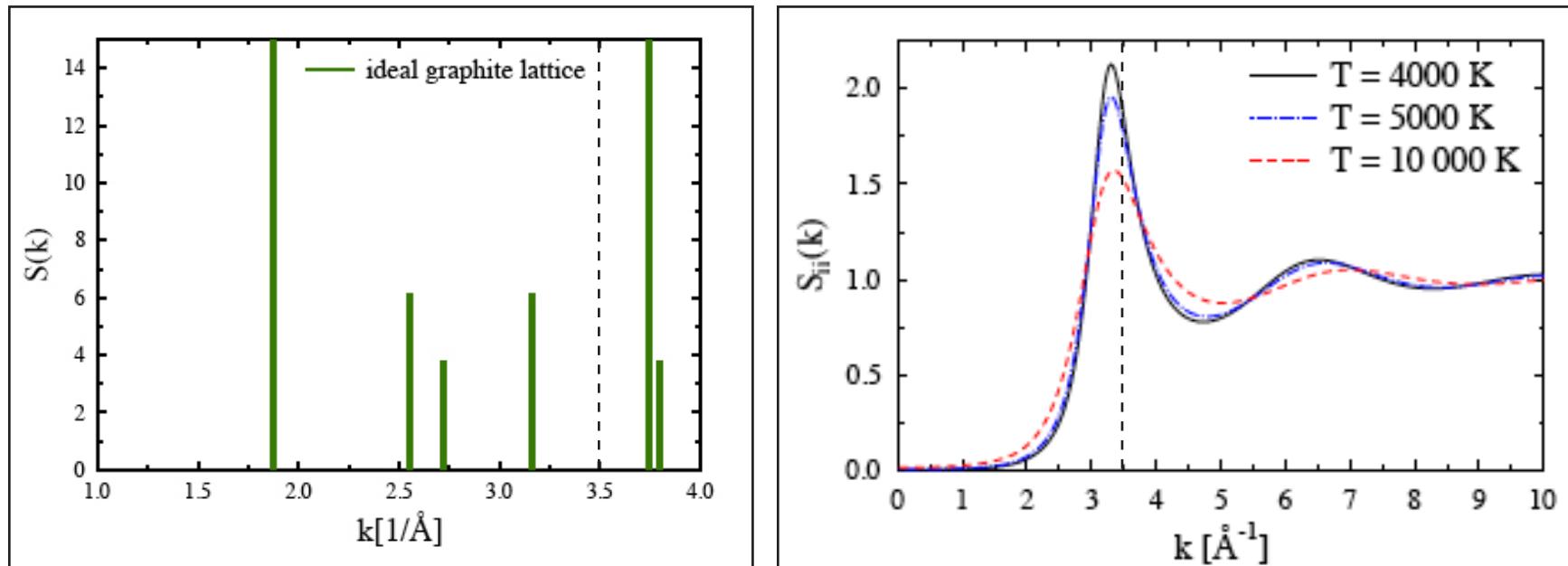
Ultrafast proton-heating experiment



DFT-MD simulations of heated carbon show that scattering comes only from liquid phase

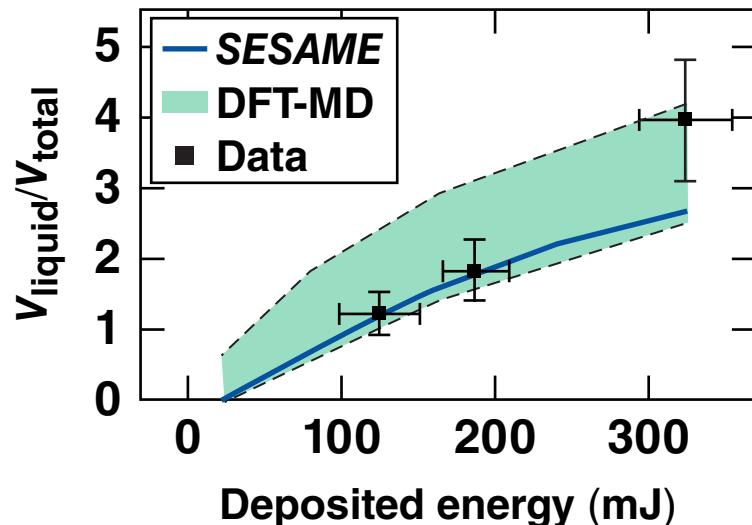
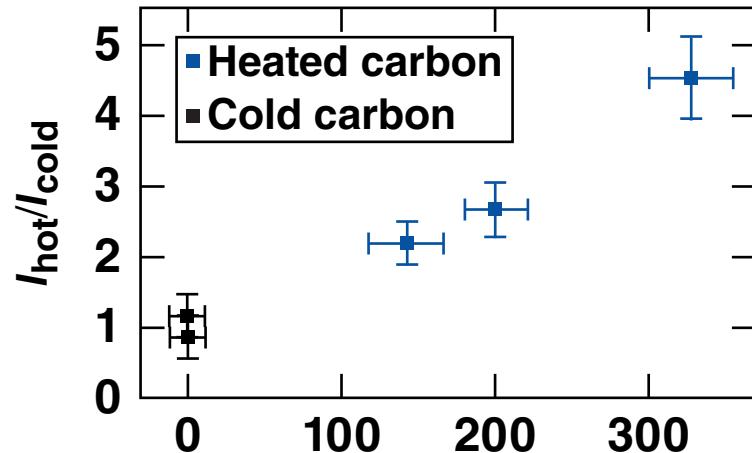


Density functional theory molecular dynamics (DFT-MD)



- Experiment probes $k = 3.5 \text{ \AA}^{-1}$

Experimental data on liquid fraction is used to validate equation-of-state models

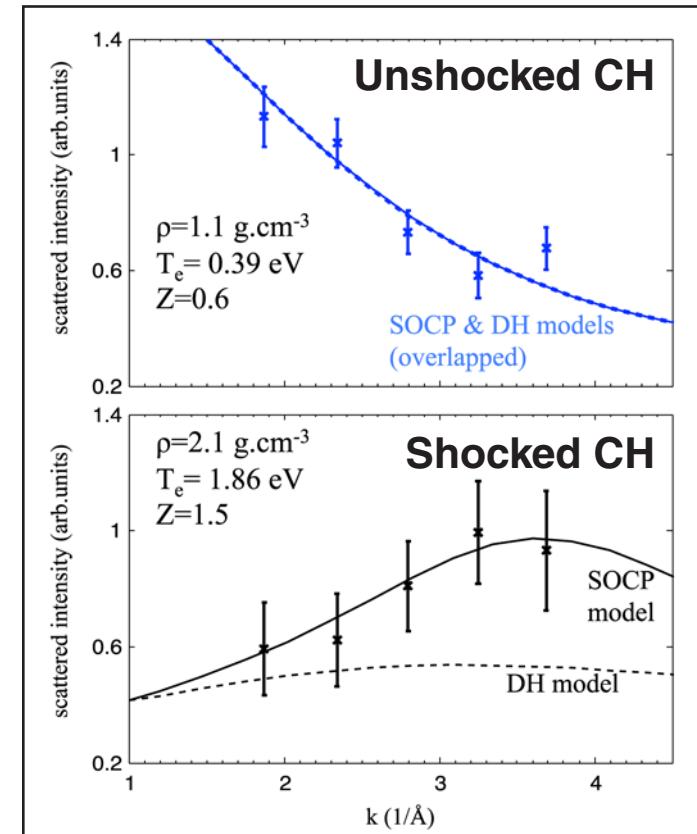
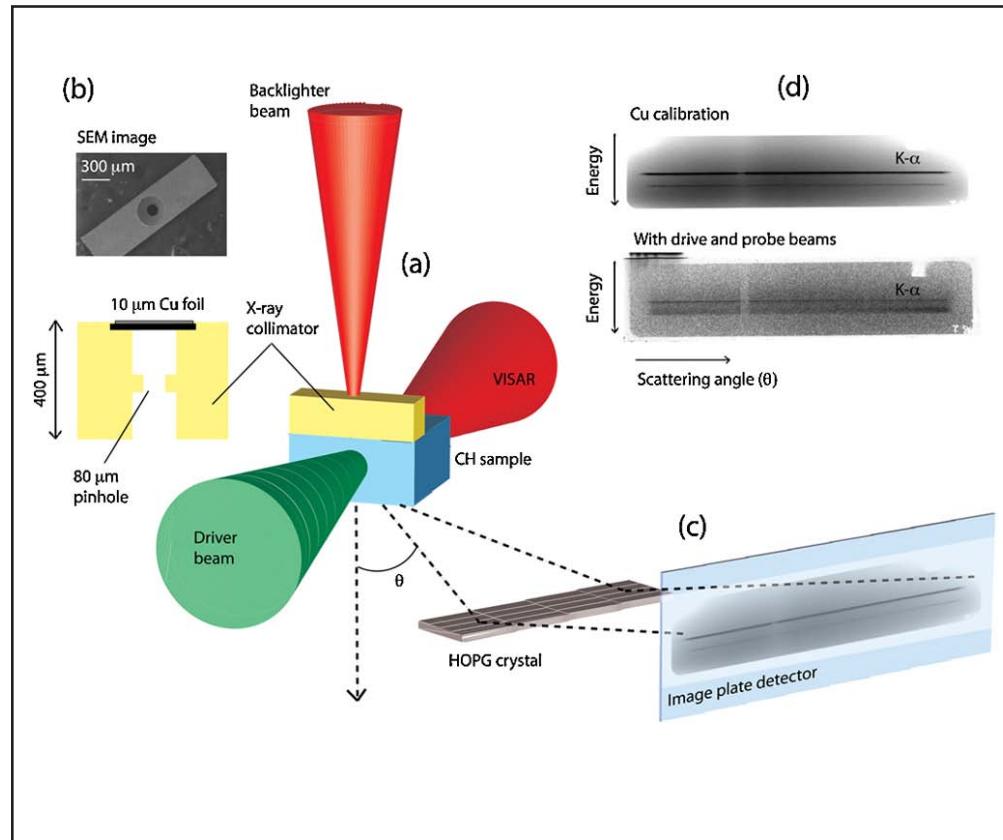


$$I_{\text{sc}} = CI_0 [V_s S_s^{\text{tot}}(k) + V_l S_l^{\text{tot}}(k)]$$

$$\frac{V_l}{V_{\text{tot}}} = \left(\frac{I_{\text{hot}}}{I_{\text{cold}}} - 1 \right) \left(\frac{S_l^{\text{tot}}}{S_s^{\text{tot}}} - 1 \right)$$

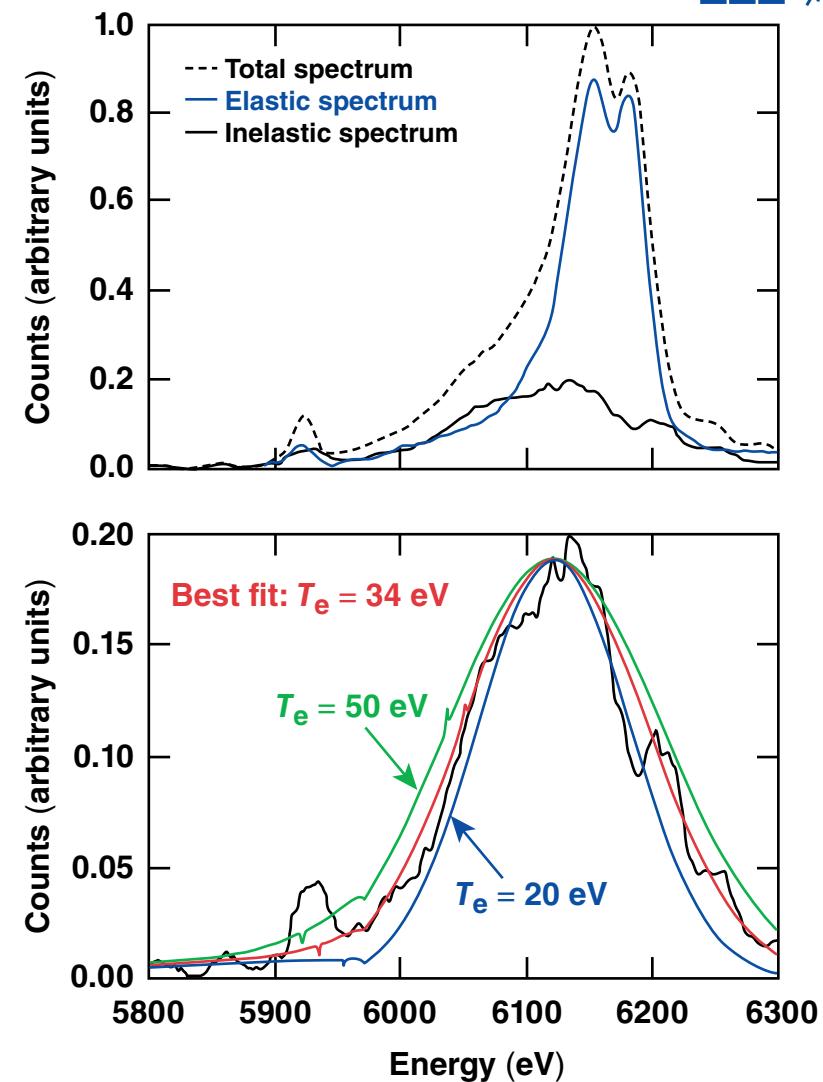
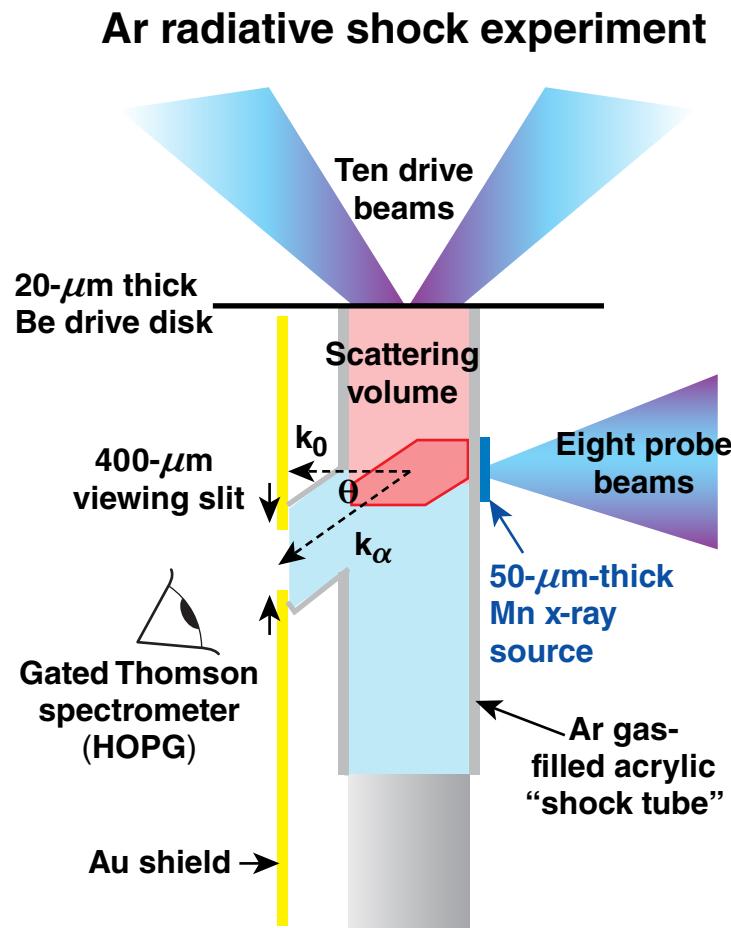
The correct modeling of the liquid content at high pressures in carbon-rich planets has implications on the predicted magnetic fields.

Time-resolved XRTS shows change in CH structure as shock wave enters the line of sight

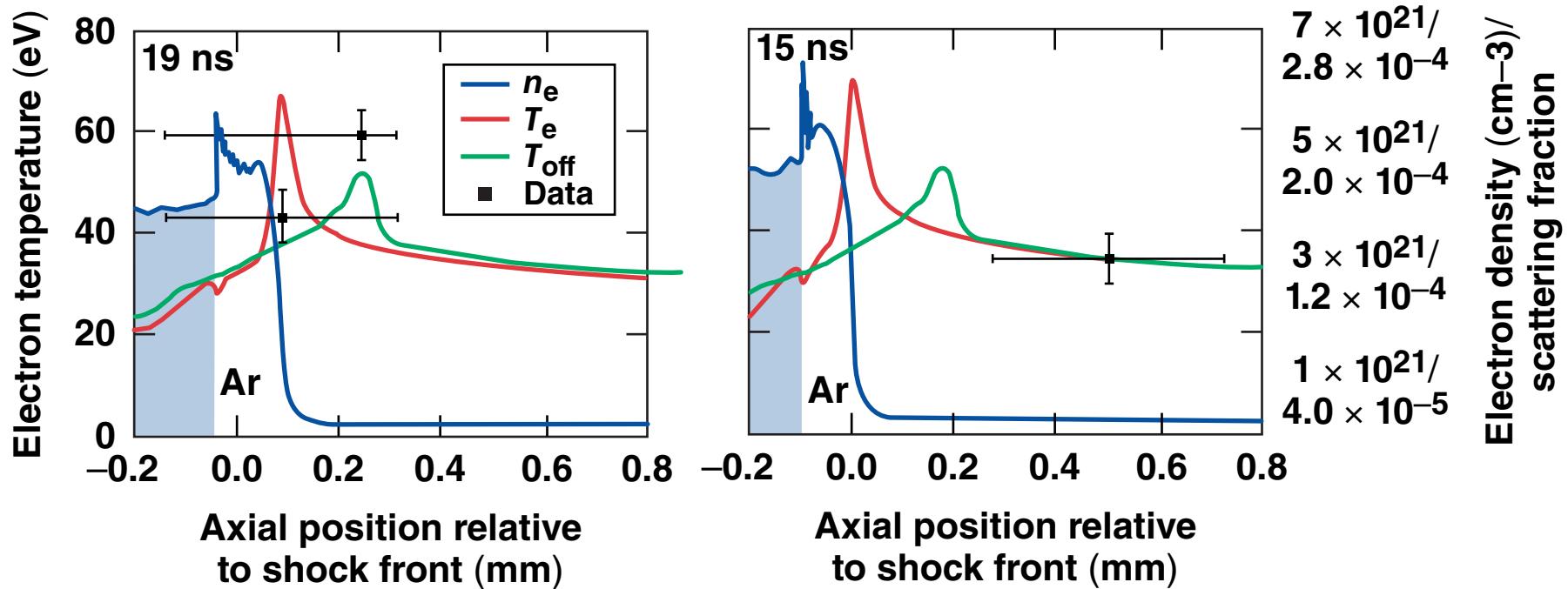


The shocked CH has a peak at 3.5 \AA^{-1} , which can be modeled when ion screening is included.

A radiative shock wave in Ar was diagnosed with XRTS on OMEGA

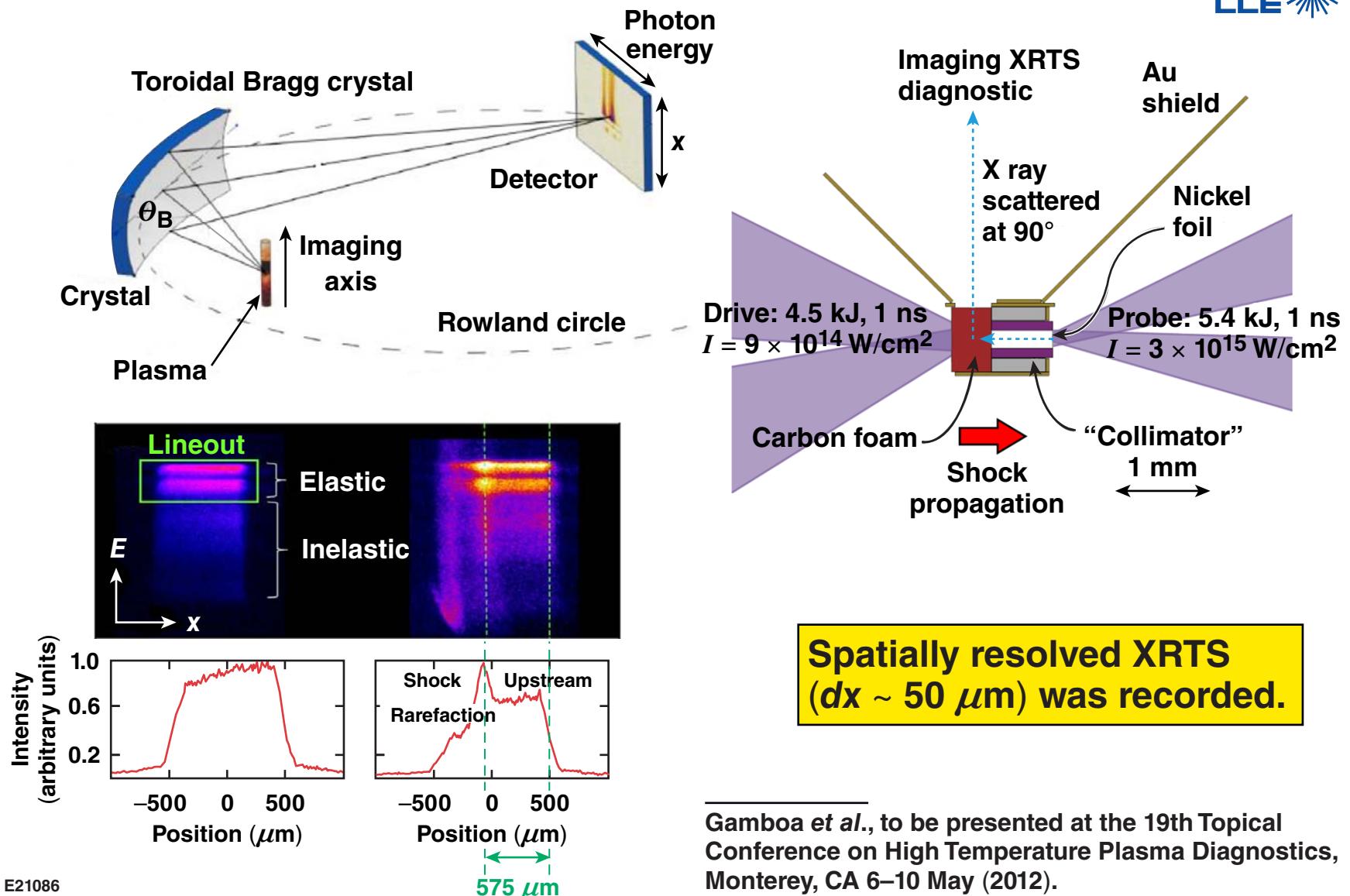


The electron temperature in the shocked region is higher than in the radiative precursor



The predictive capability of radiative shock waves is validated in the laboratory with XRTS and x-ray radiography.

The University of Michigan and LANL have deployed an imaging spectrometer for XRTS on OMEGA



Outline



- Motivation
- XRTS
- Experiments
- Future

A spatial resolution of $\sim 10 \mu\text{m}$ for spectral measurements would greatly benefit many XRTS experiments



- General-use spectral imaging diagnostic
- Development of brighter backlights
- Improve accuracy of target alignment
 - install foreground lighting on OMEGA Target Viewing System

OLUG could help to develop specifications for next-generation XRTS measurements on OMEGA.

Summary/Conclusions

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- A review of XRTS experiments is presented
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Many XRTS experiments need spatially resolved spectral measurements.

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