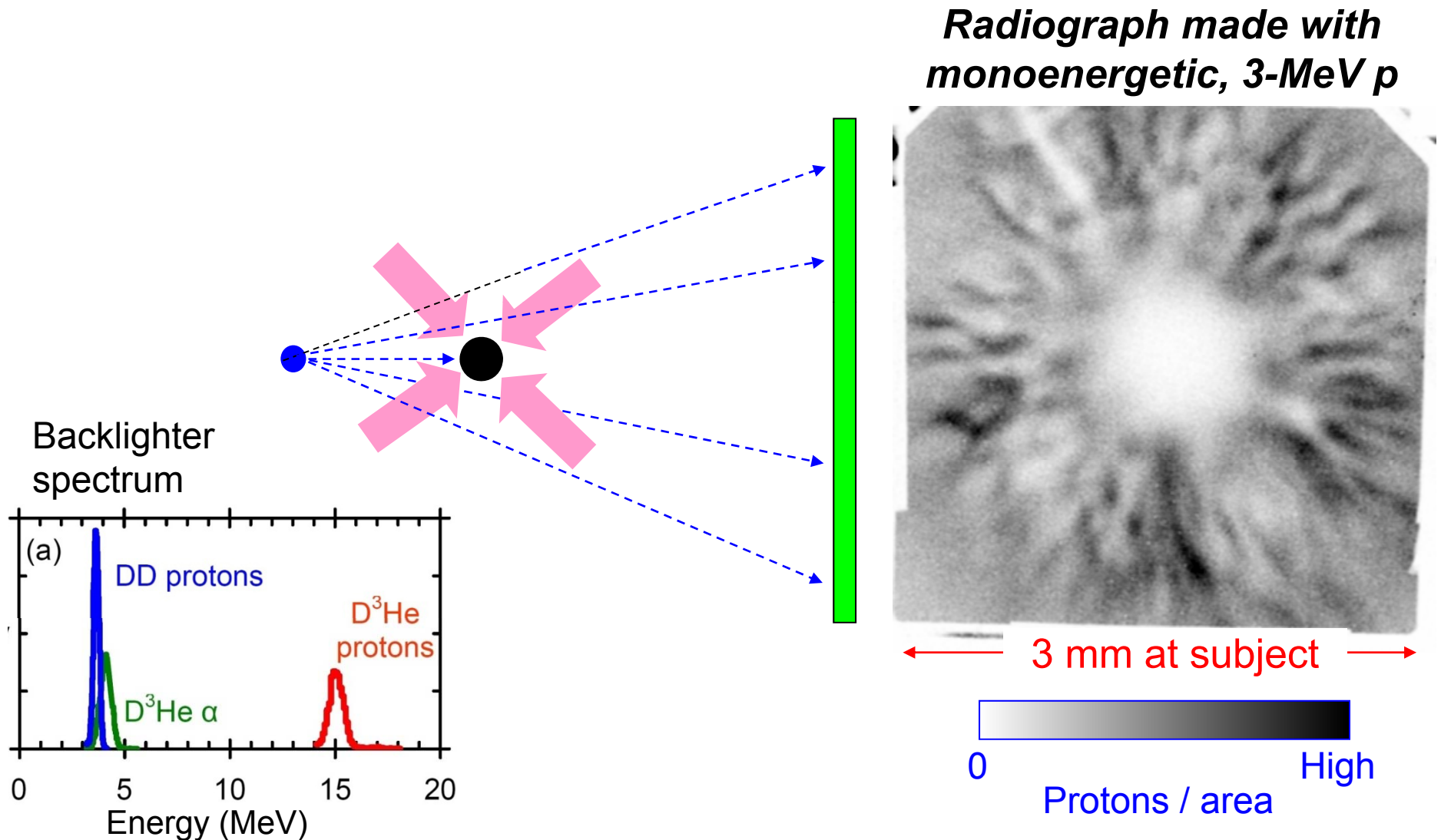
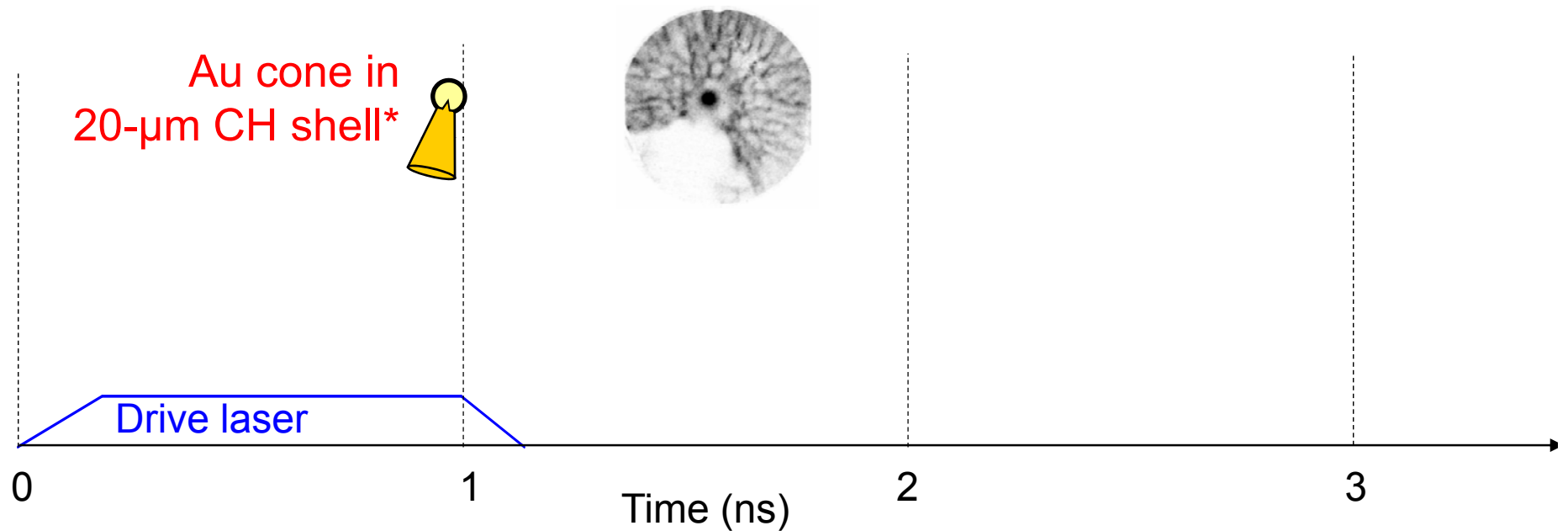


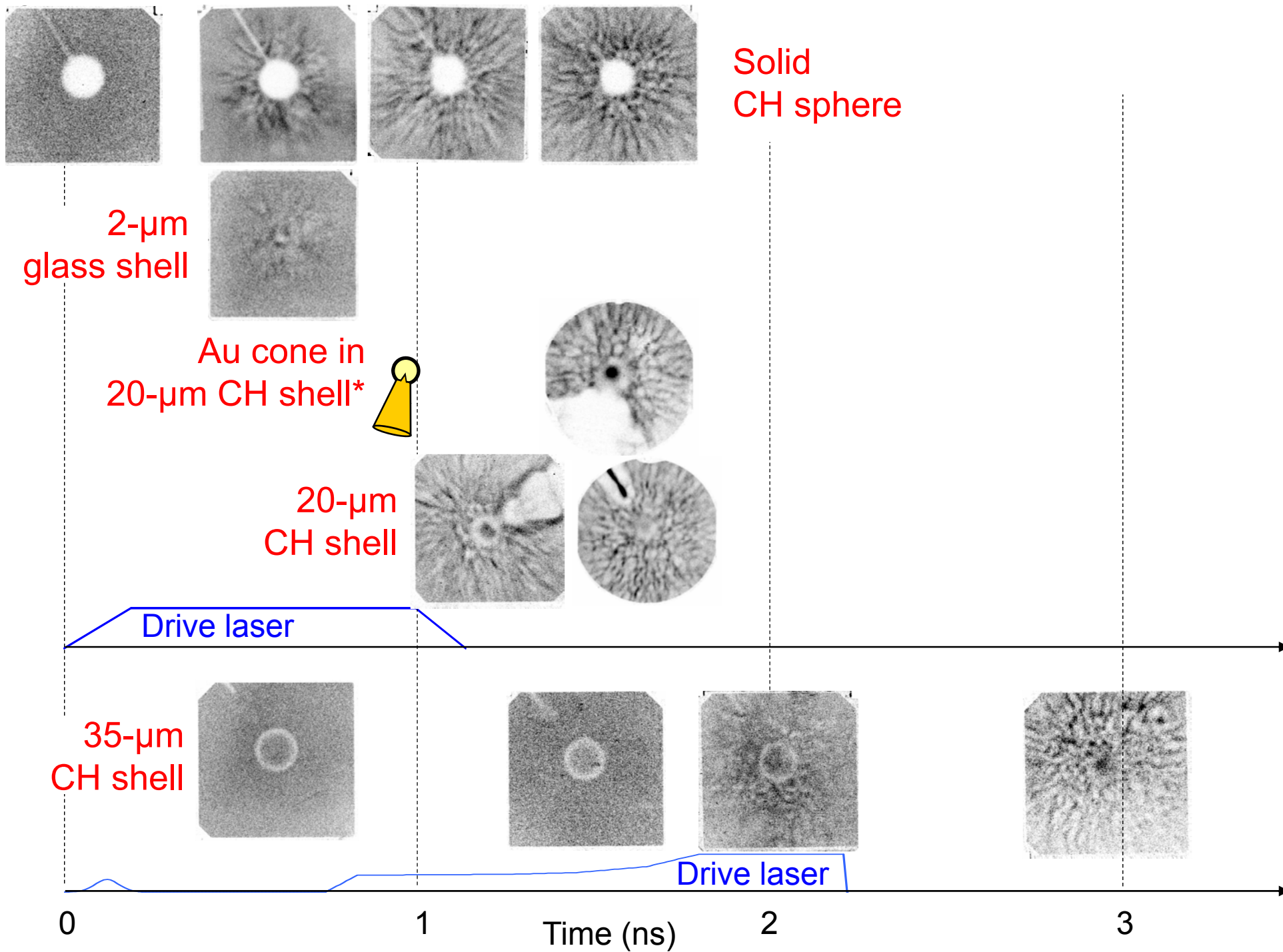
Observation and characterization of coronal filamentation in direct-drive ICF



15-MeV-proton radiography has demonstrated that filaments are endemic in direct-drive implosions



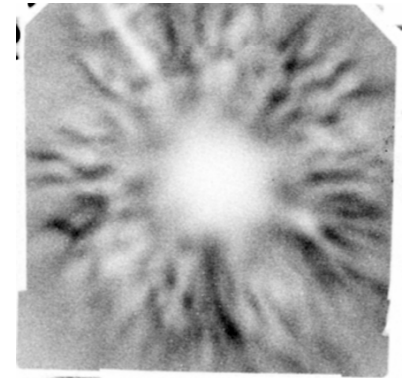
*Rygg *et. al*, *Science* (2008)



Summary:

1. Radial striations in the images correspond to

- Current filaments (~ 2 kA toward capsule)
- Surrounded by $B \sim 10$ T
- With separation near critical surface ≈ 100 μm



2. The filaments and fields

- Appear during the laser pulse – could affect drive efficiency
- Expand with the corona, with features frozen in for hundreds of ps

Collaborators

MIT

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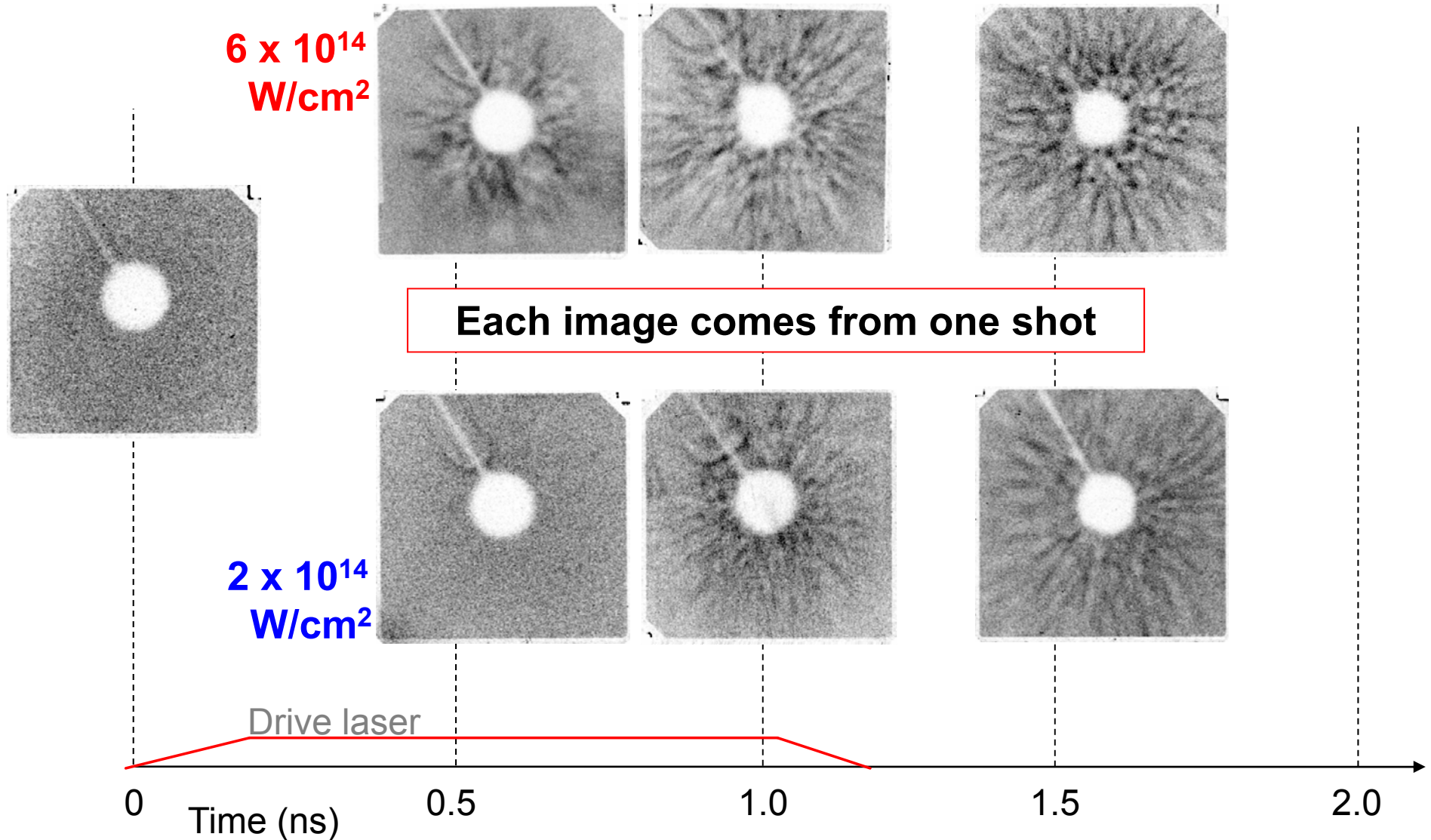
U of R – LLE

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D. Meyerhofer
V. Smalyuk

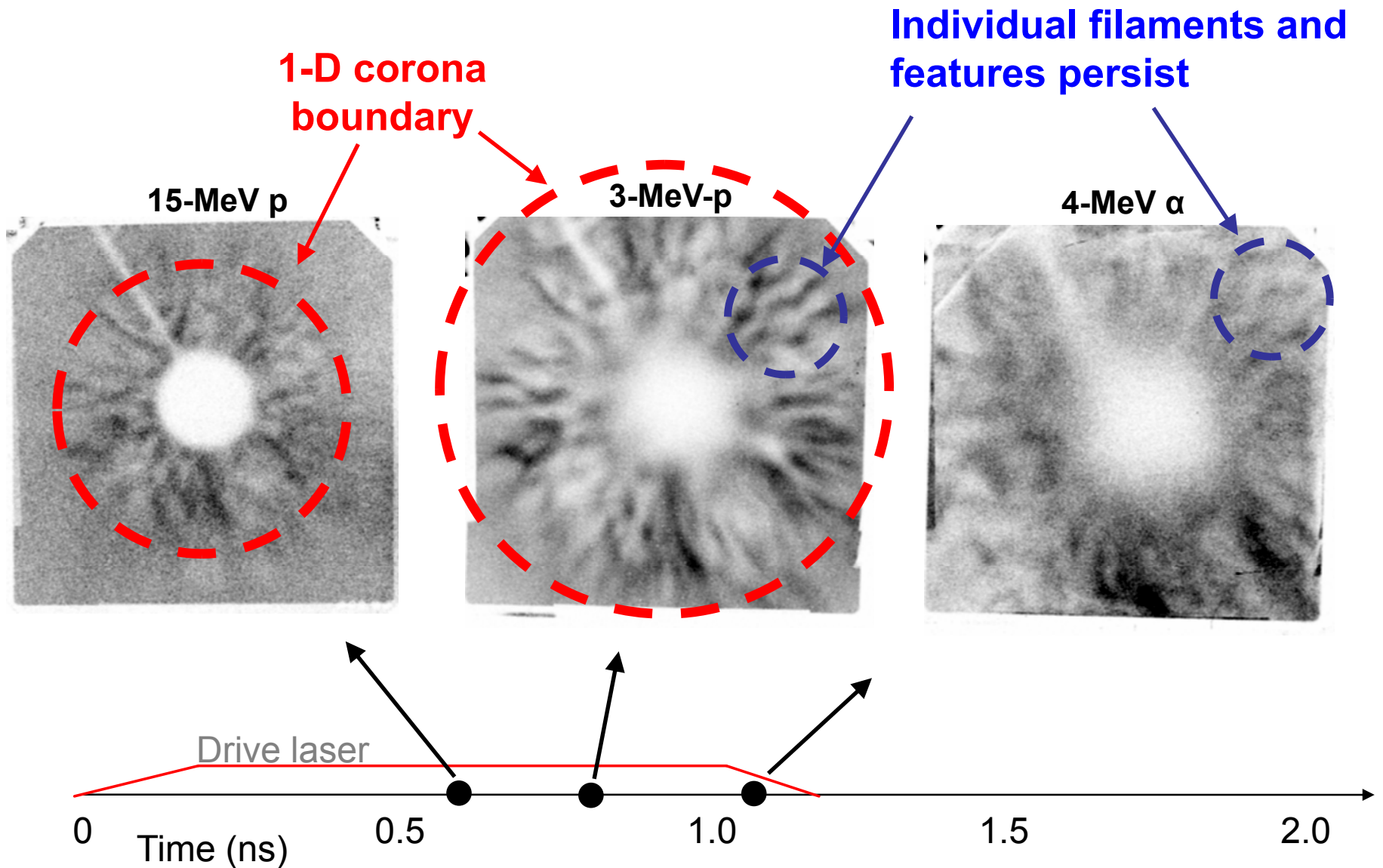
*Currently at *LLNL*

6 Radiographs were made of solid CH targets (860- μm diam.) driven with two different laser intensities

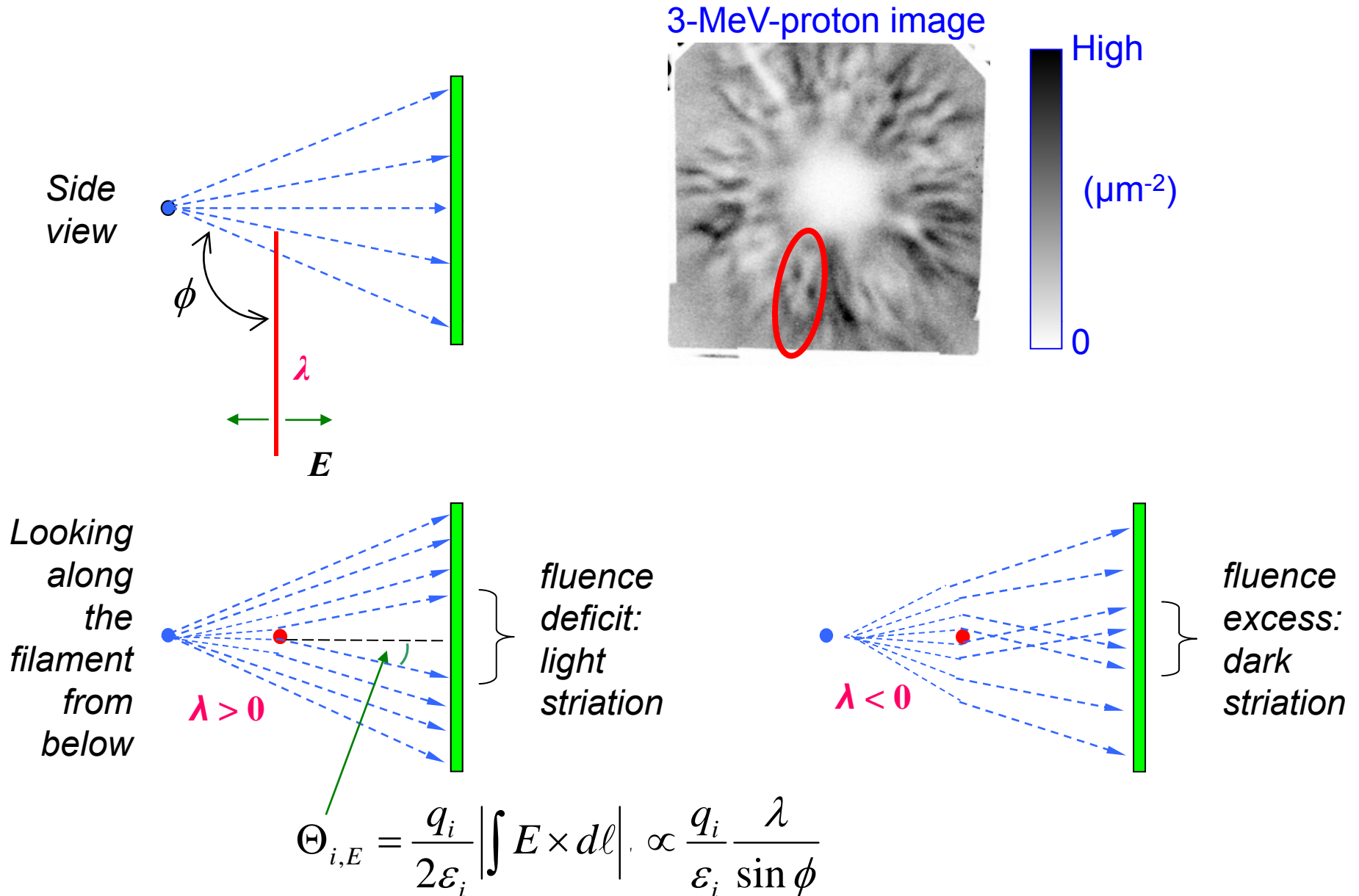
15-MeV p images



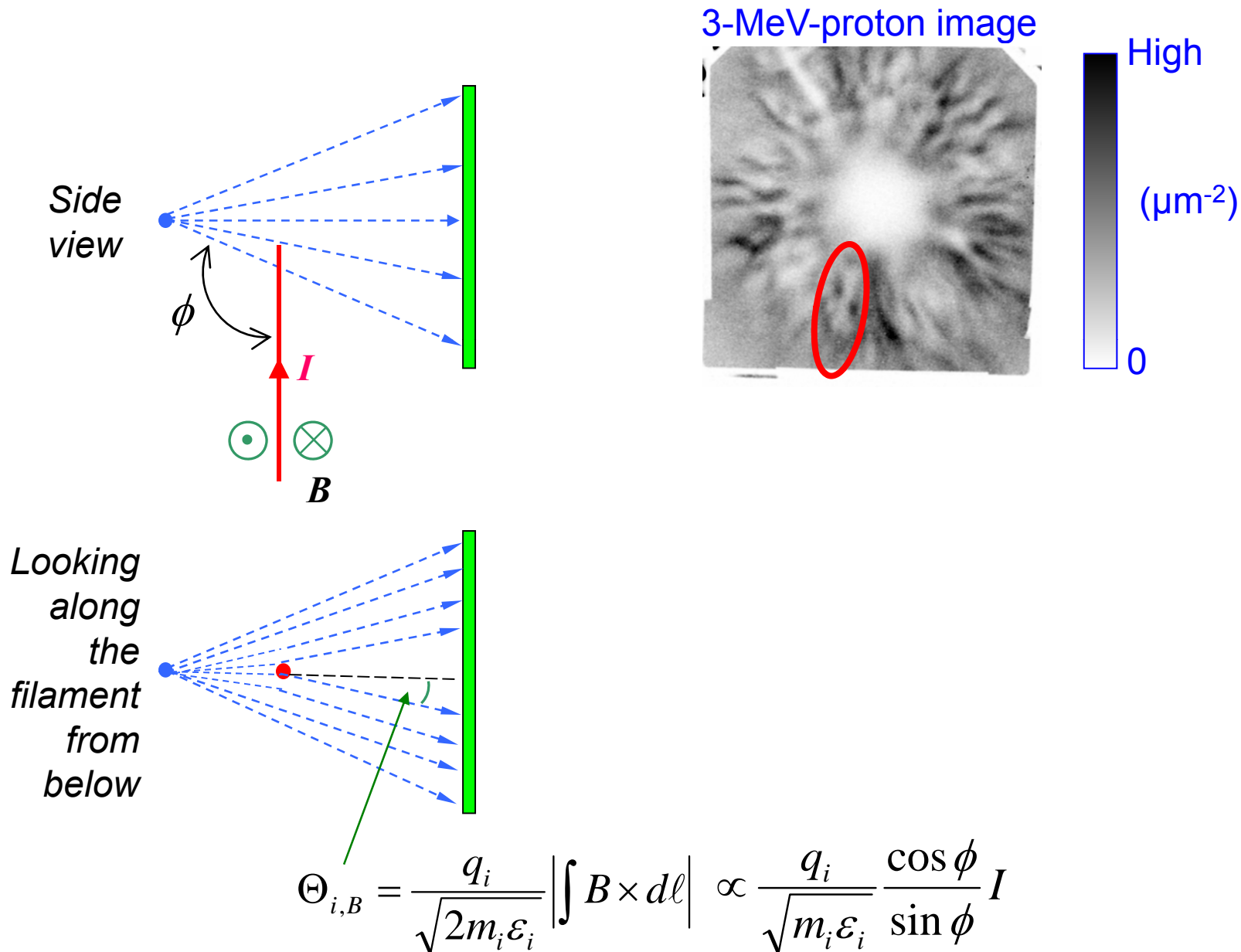
Each shot has images from different particles that sample fields at different times



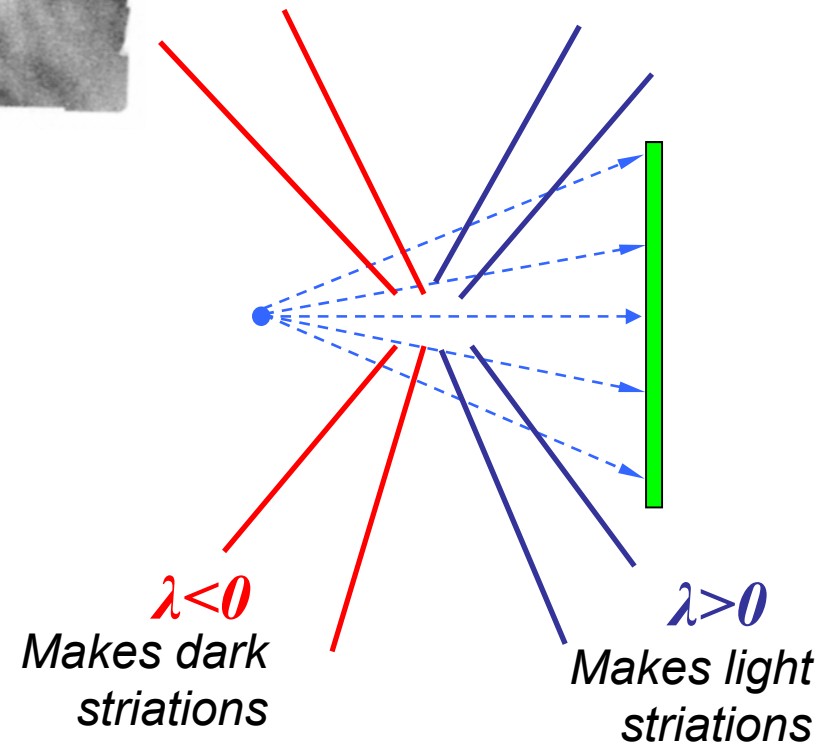
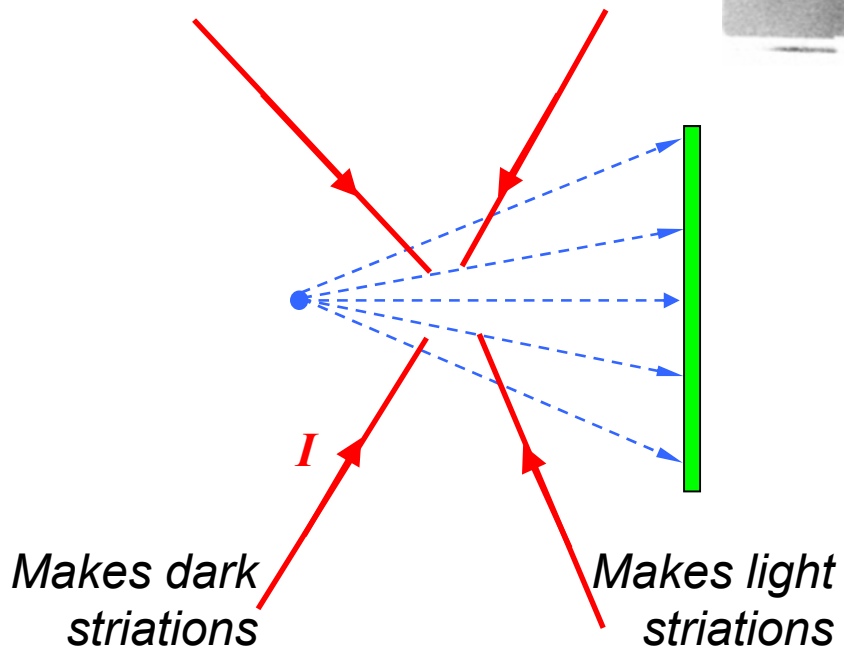
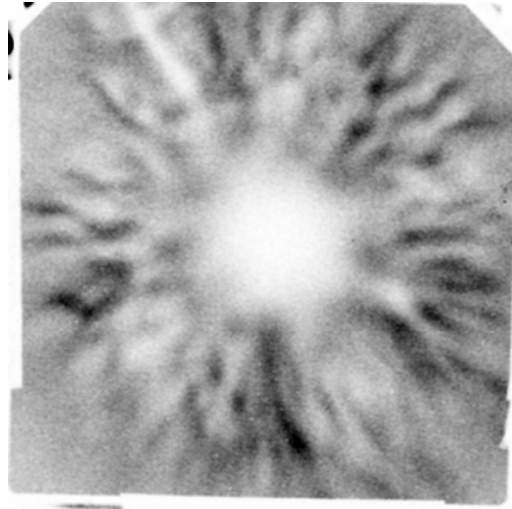
E fields could cause image striations



B fields could cause image striations



Light striations appearing “in front of” dark ones makes B more natural than E



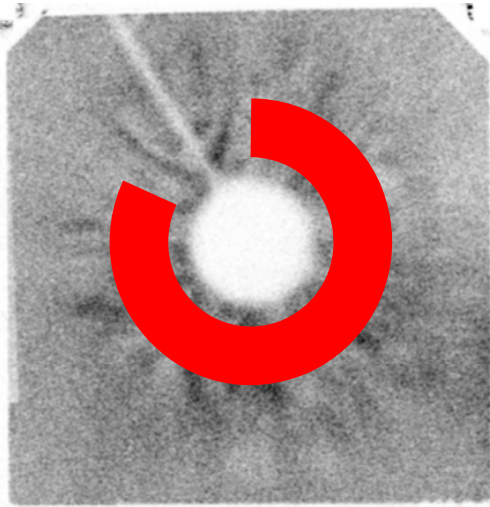
Deflections due to E and B fields scale differently with the energy, mass and charge of the imaging particle

$$\Theta_{i,E} \propto \frac{q_i}{\mathcal{E}_i}$$

$$\Theta_{i,B} \propto \frac{q_i}{\sqrt{m_i \mathcal{E}_i}}$$

- If we could accurately measure deflection angles for a filament with different imaging particles, we could see whether they scaled according to E or B .
- This is very difficult, because
 - The deflections are small and hard to quantify, especially for 15-MeV protons.

For small deflection angles, a similar scaling applies to measured image modulation due to filaments



Define an image region containing filaments but not stalk or capsule.

Define RMS image modulation due to filaments :

$$\Sigma \equiv \frac{\sqrt{\sigma_{total}^2 - \sigma_{stats}^2}}{\langle \text{fluence} \rangle}$$

If B:

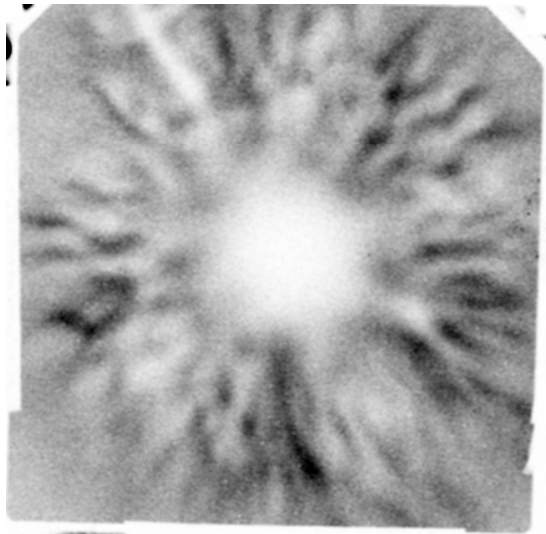
$$\Sigma_i \propto \frac{q_i}{\sqrt{m_i \varepsilon_i}} \sqrt{N_{fil}} |I_{fil}| \quad \Rightarrow \quad |I_{fil}| \propto \left[\frac{\sqrt{m_i \varepsilon_i}}{q_i} \right] \Sigma_i$$

If E:

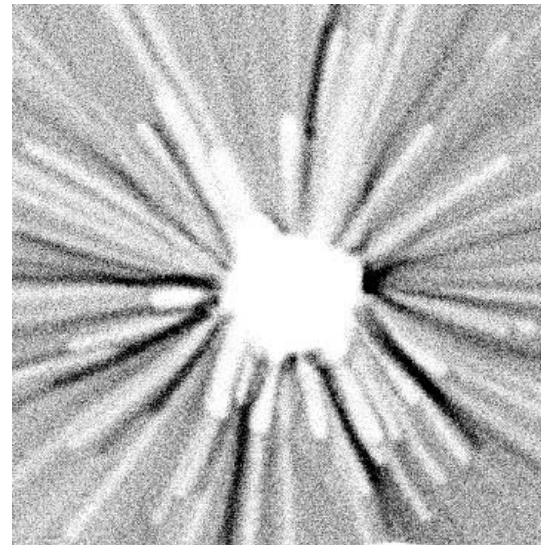
$$\Sigma_i \propto \frac{q_i}{\varepsilon_i} \sqrt{N_{fil}} |\lambda_{fil}| \quad \Rightarrow \quad |\lambda_{fil}| \propto \left[\frac{\varepsilon_i}{q_i} \right] \Sigma_i$$

The wavelength (filament-to-filament separation) can be estimated by comparison with simulations

Data

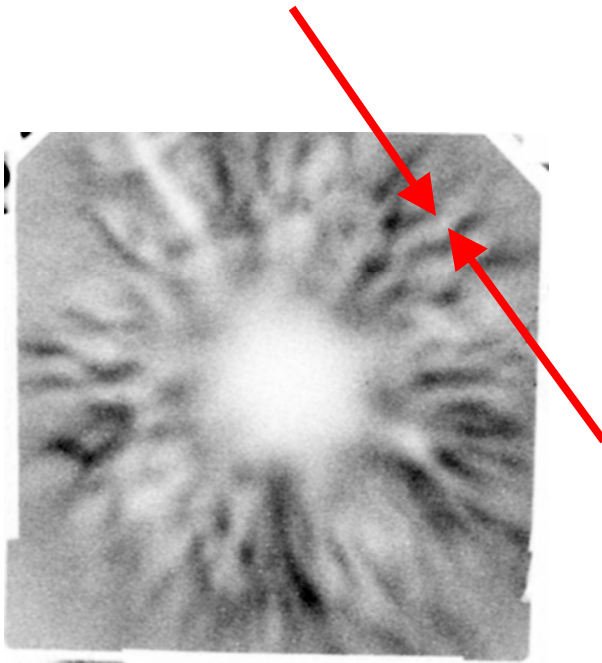


*Monte-Carlo simulation
for 200 filaments over 4π*



*~ 200 filaments distributed over 4π
means 100- μm separation at the capsule surface*

Field strength and filament current can be estimated from striation widths in the lowest-energy images



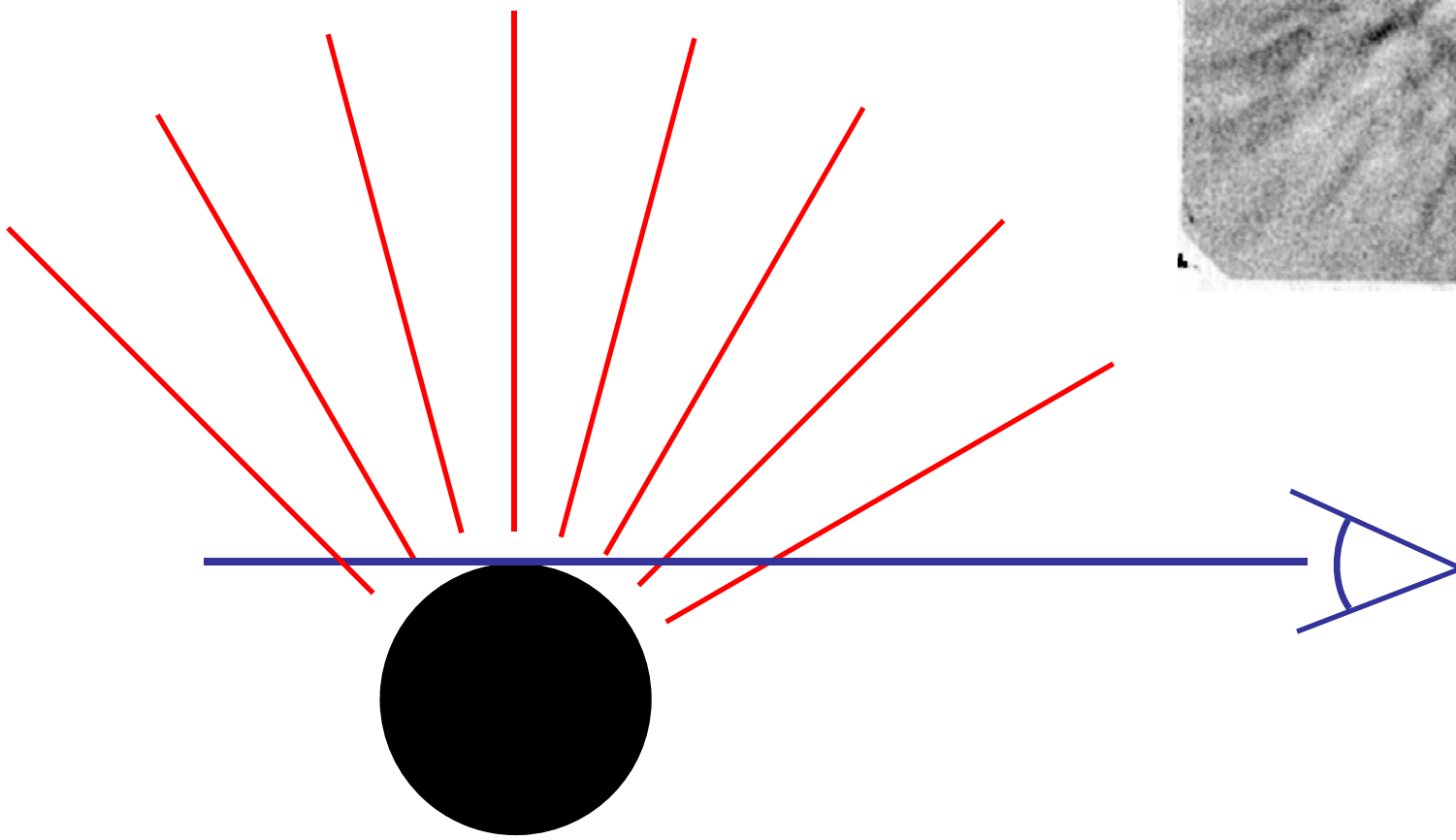
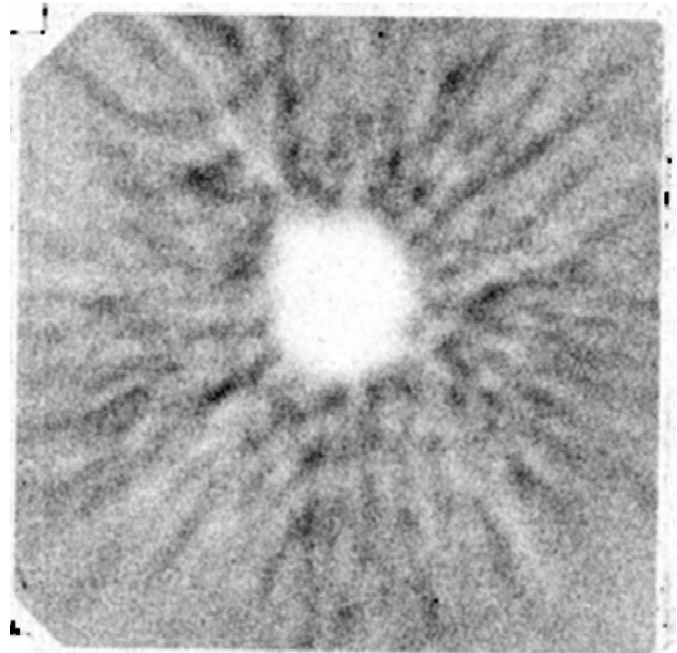
Half width gives a measure of deflection angle

$$\Theta_{i,B} = \frac{q_i}{\sqrt{2m_i\varepsilon_i}} \left| \int B \times d\ell \right|$$

$$B \sim 10 \text{ T}$$

Filaments appear to extend inward to the ablation surface

But we can't prove that
from the images



**This is the resulting picture of filament structure
at ~1 ns for laser intensity $6 \times 10^{14} \text{ W/cm}^2$**

