

Theory & validation of electron-ion thermal relaxation in thermonuclear DT plasma for ICF

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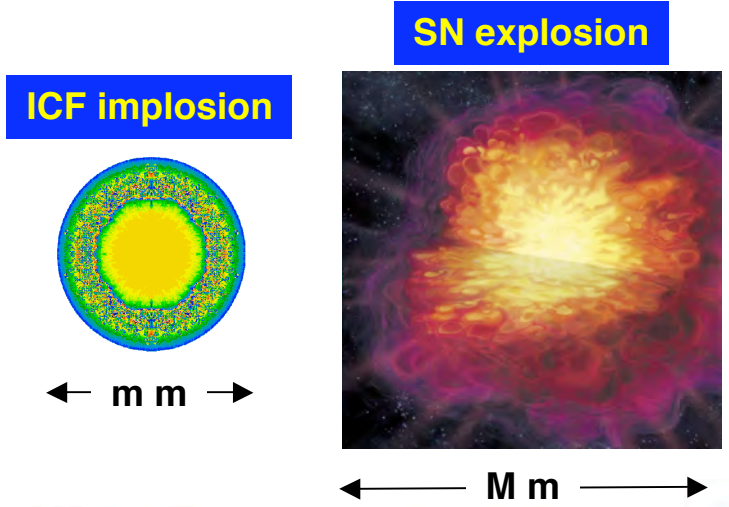
Project Overview

Thermal Relaxation in plasmas

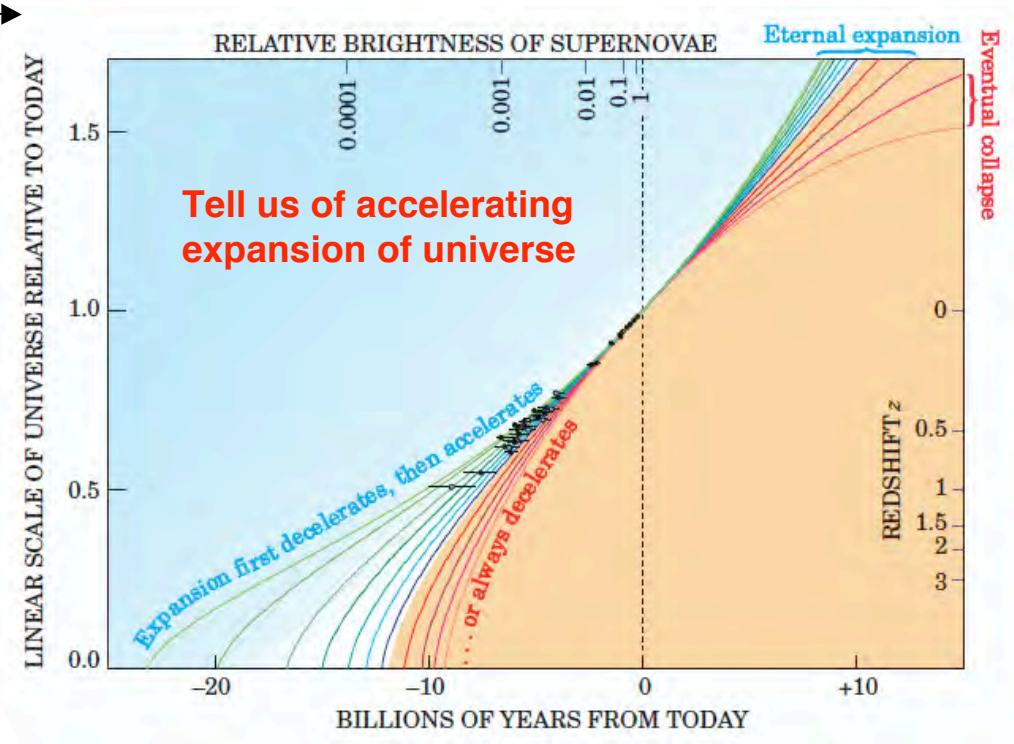
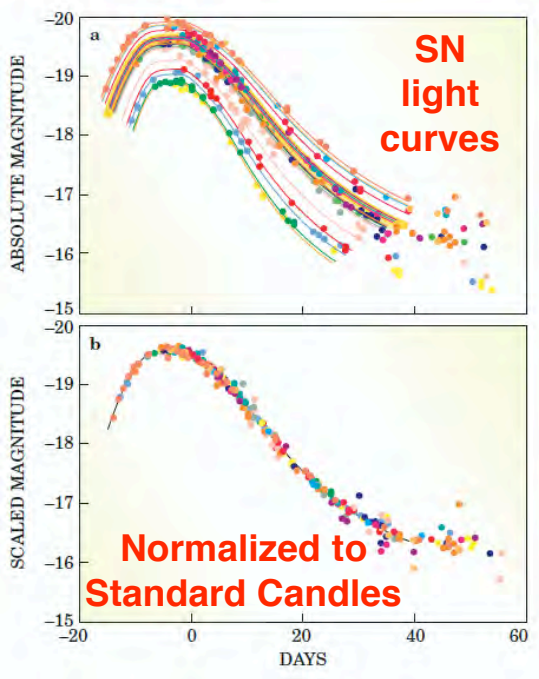
Thermonuclear burn & Ejecta

Opportunities @ LANL

LANL scientists strive to describe physical systems that depend on a wide variety of fundamental processes at various scales



- | | | |
|------------------------|---|------------------|
| Hydrodynamics | ⇒ | Complex EOS |
| Solid state properties | | Fluid turbulence |
| Warm-dense-matter | | Plasma physics |
| Radiation transport | | Opacity |
| Nuclear reactions | ⇒ | Hydrodynamics |



We use high-resolution, 3D single-physics numerical simulations to clarify the physics assumptions in our application codes

Multi-physics codes describe macroscopic (dis-) assembly of applications
Combine properties & transport rates for many different processes

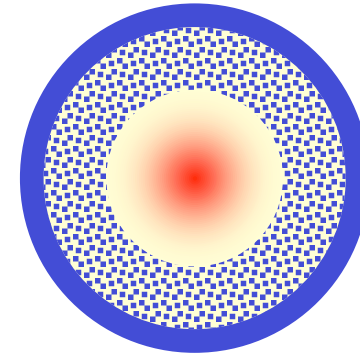
Hydrodynamic instabilities cause mix & reduce yield

Mix can be adjusted to obtain observed yield

But burn temperature smaller than observed

But He³/D results differ from pure gases

But ...



ICF capsule

Single-physics simulations clarify microscopic processes

Material properties & transport rates

Material strength & damage

Transport from turbulent hydrodynamics

Atomic mixing rates

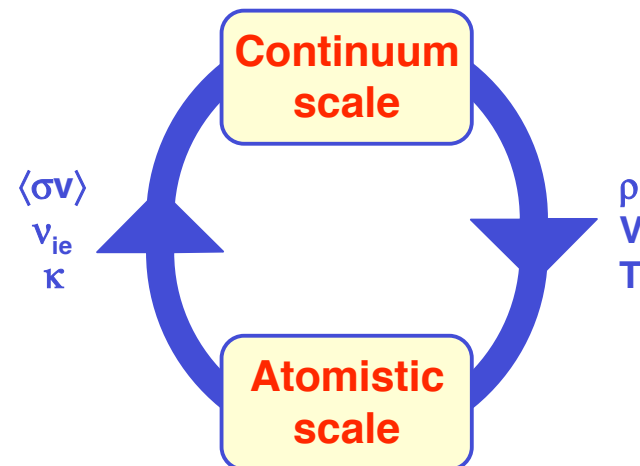
Electron-ion relaxation rates - ν_{ie}

Radiation transport in complex mixtures - κ

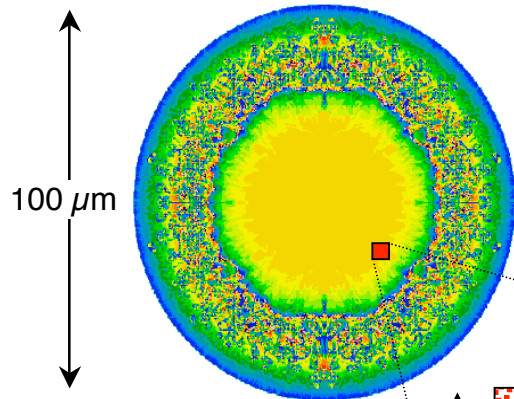
DT fusion rates - $\langle\sigma v\rangle$

Alpha energy deposition

...



Physics issues are coupled using multi-scale computing up to 1 PetaFlop/s on Roadrunner Computer (LANL)

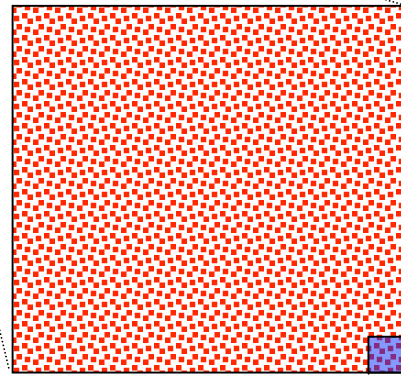


Continuum simulation of NIF implosion

Hydrodynamics, EOS, radiation transport, nuclear

Assumes Maxwellian fluids & Spitzer collisions

$$N_p^{1/2} \lambda_D \sim 10 \mu\text{m}$$



PIC simulation of TN plasma

Macro-particles with TN burn

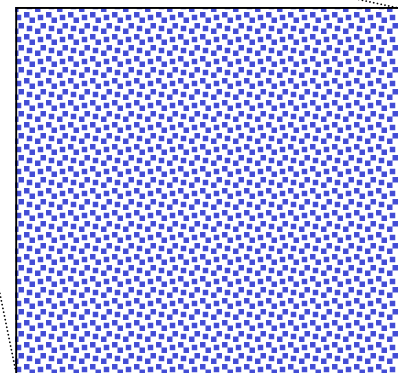
Kinetic & collective effects

Assumes Spitzer collisions

MD simulation of hot plasma

Electrons & ions with Coulomb potential

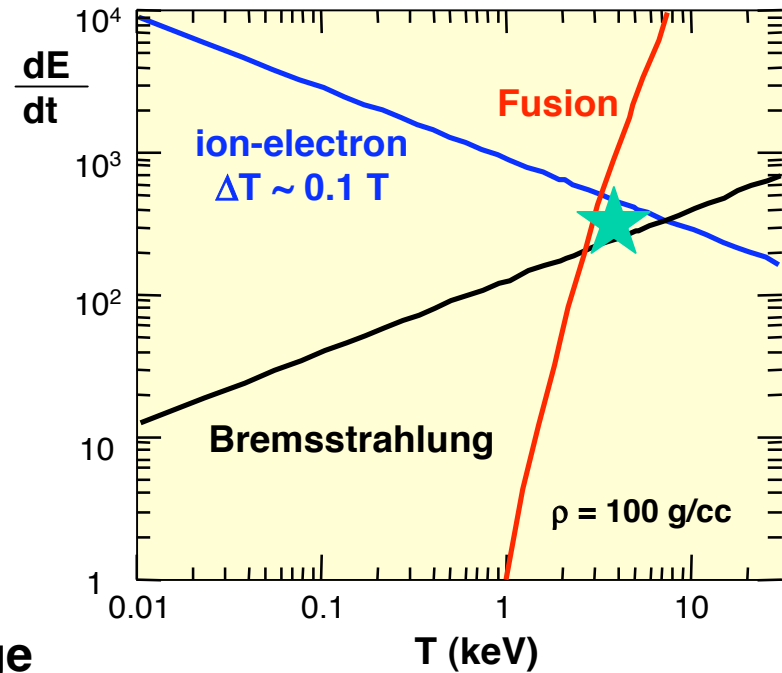
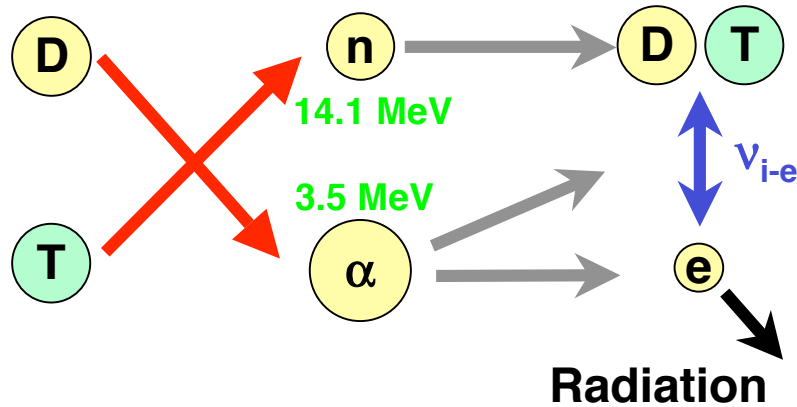
Validates new statistical collision theory



$$(N_p/4n)^{1/3} \sim 0.03 \mu\text{m}$$

Thermonuclear burn in ICF plasma involves many collisional transport processes whose rates can be improved

Thermonuclear processes in ICF



Transport rates depend on long-range Coulomb collisions

Maxwellian electrons & ions

Thermal relaxation

Thermal conductivity

Electrical conductivity

Radiative

Uncertainty in Coulomb log

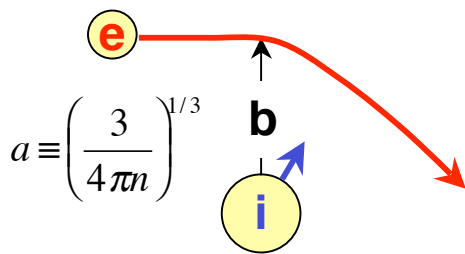
$$\frac{dT_i}{dt} \sim \nu_{ie} (T_e - T_i)$$

$$\nu_{process} \sim \nu_o \ln \Lambda_{process} \propto \frac{n}{T_e^{3/2}} \ln \left(\frac{b_{max}}{b_{min}} \right)$$

Temperature decay not exponential
because $\nu \propto T^{-3/2}$

Legacy calculations of temperature relaxation between electrons & ions diverge due to Coulomb force $\propto r^{-2}$

Particle collisions can be described by integrating Rutherford cross-section in Boltzman equation, but this **diverges** due to **distant encounters**



$$v_o = \frac{8}{3} \frac{\sqrt{2\pi m}}{M} \frac{e^2 e_i^2}{T^{3/2}} n$$

$$v_{ie} = v_o \int_0^{b_{\max}} \frac{b db}{b^2 + (Ze^2 / mv^2)^2}$$

$$v_{ie} = \frac{v_o}{2} \ln \left(1 + \left(3 \frac{b_{\max}}{b_{\min}} \right)^2 \right)$$

$$b_{\min} = \frac{Ze^2}{k_B T_e}$$

$$b_{\max} \Rightarrow \infty$$

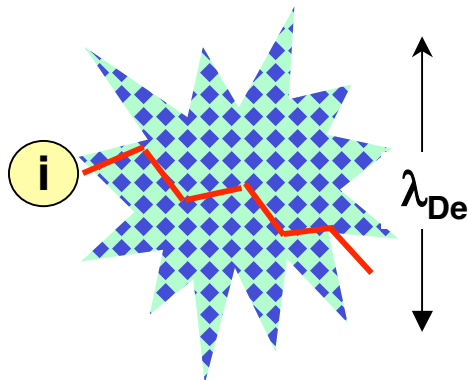
Screening parameter b_{\max}

Spitzer: $\lambda_{De} = \sqrt{\frac{k_B T}{4\pi n e^2}}$

Landau: $\lambda_{DH} = \frac{\lambda_{De}}{\sqrt{1+Z}}$

Chandrasekhar: $a = \left(\frac{4\pi n}{3} \right)^{-1/3}$

Plasma fluctuations ($\epsilon(k, \omega) \sim 0$) due to discrete electrons & ions are described by Lenard-Balescu, but this **diverges** due to **close encounters**



$$v_{ie} = v_o \int_0^{k_{\max}} \frac{k^3 dk}{|k \cdot \epsilon(k, 0) \cdot k|^2}$$

$$v_{ie} = v_o \int_0^{k_{\max}} \frac{dk}{k} \frac{k^4}{(k^2 + k_{De}^2)^2} = v_o \ln \left(\frac{\lambda_{De}}{b_{\min}} \right)$$

$$k_{\max} \sim \frac{1}{b_{\min}} \Rightarrow \infty$$

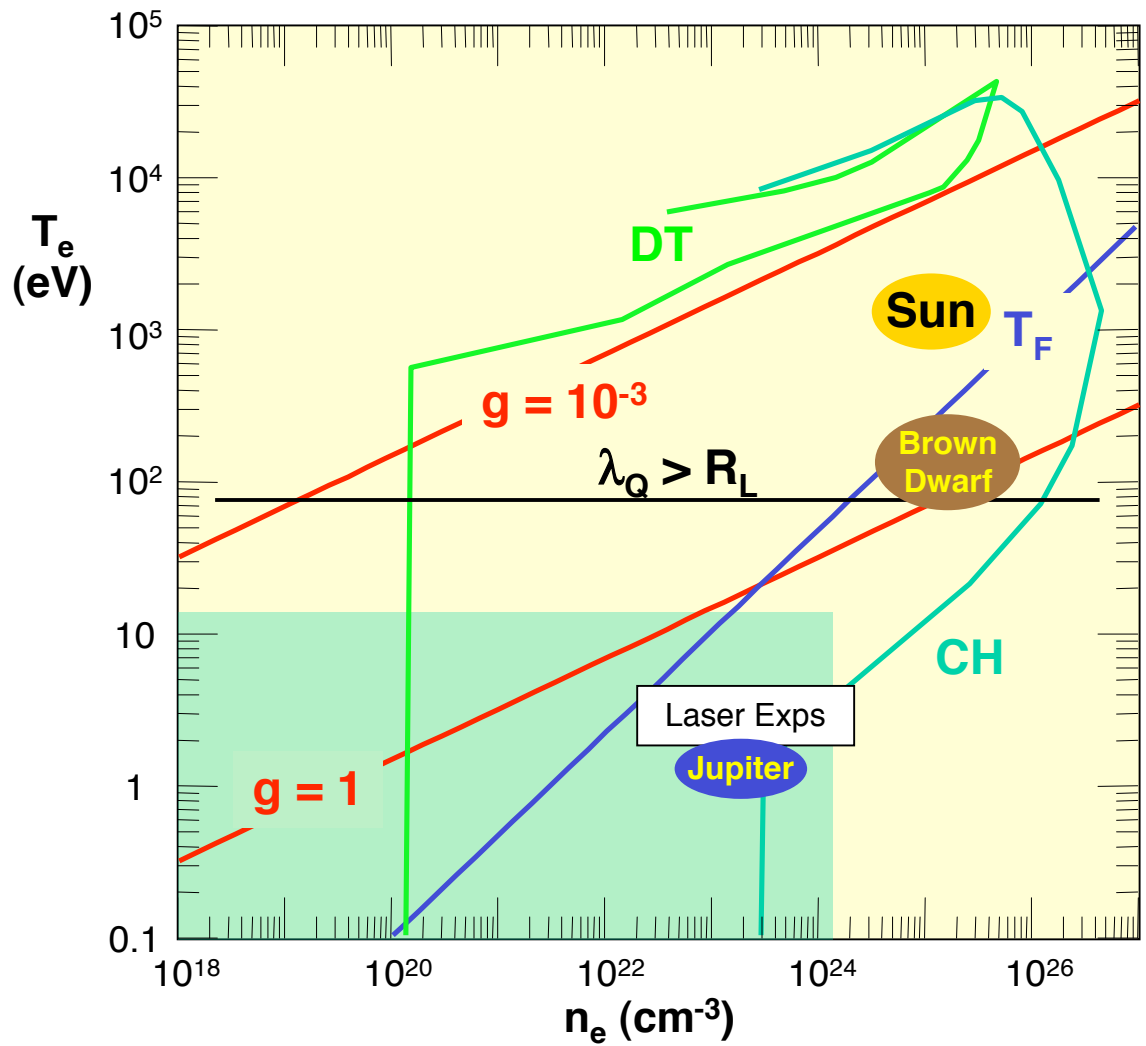
$b_{\min} \sim 1 / k_{\max}$

Landau: $R_L = \frac{Ze^2}{k_B T}$

deBroglie: $\lambda_Q = \frac{\hbar}{\sqrt{mk_B T}}$

Need to capture close & distant encounters self-consistently

NIF capsule traverses many difficult plasma regimes, but ignition plasma has three simplifying characteristics



Weakly coupled plasma

$$g \equiv \frac{Ze^2}{\lambda_{De} k_B T} = \frac{R_L}{\lambda_{De}} \ll 1$$

$$R_L = Ze^2 / k_B T_e$$

$$\lambda_{De} = \sqrt{k_B T_e / 4\pi n_e e^2}$$

Quantum diffraction

$$\lambda_Q = \hbar / \sqrt{m k_N T_e} > 1.67 R_L$$

$$T_e > 2.8 \frac{m e^4}{k_B \hbar^2} = 76 eV$$

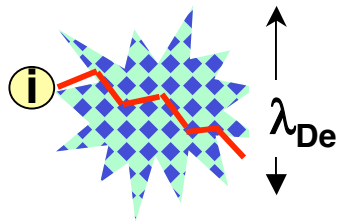
Maxwellian distribution

$$T_e > T_F = \frac{\hbar^2}{2 m k_B} (3\pi^2 n)^{2/3}$$

ICF codes should treat ALL regimes with unified theory that includes particle correlations (g) & quantum diffraction (λ_Q) & statistics (T_F)

Theory was developed to describe temperature relaxation that includes degeneracy & particle correlations self-consistently

Ion temperature in a spatially uniform, unmagnetized plasma depends on work done on ions (M, n_i, T_i) by electrons ($m, n_e = Z n_i, T_e$)



$$\frac{3}{2} n_i k_B \frac{dT_i}{dt} = \langle \delta \vec{j}_i \cdot \delta \vec{F}_{ie} \rangle$$

Jerome Daligault

Fourier (ω, k) components of fluctuating ion current δj_i & e-i force δF_{ie} are given by

Continuity equation: $\delta \vec{j}_i(k, \omega) = \omega \delta n_i(k, \omega) \vec{k} / k^2$

Gauss' law: $\delta \vec{F}_{ie}(k, \omega) = ik V_{ie}(k) \delta n_e(k, \omega)$

Fourier transform of interaction potential
 $V_{ie} \propto 1 / k^2$

δj_i & δF_{ie} depend on self-consistent density fluctuations δn_α & plasma susceptibility χ_α for each species α , and we assume a linear response

$$\delta n_\alpha(k, \omega) \equiv n_\alpha - \langle n_\alpha \rangle = \delta n_\alpha^s + \chi_\alpha \sum_{\beta=e,i} V_{\alpha\beta} \delta n_\beta \left[1 - G_{\alpha\beta}(k, \omega) \right]$$

Mean field assumes particles are allowed everywhere with equal probability
 \Rightarrow 'pair distribution' $g_{ie}(r) = 1$



'Local field correction' $G_{\alpha\beta}$ related to Fourier transform of $g_{ie}(r)$
 $\Rightarrow G_{\alpha\beta} = 1$ when particles excluded

Coupling between ion & electron subsystems is relatively slow due to small mass ratio m/M & this allows tractable solution

Linear response for density fluctuations introduces total dielectric $D(k, \omega)$

$$\frac{3}{2} n_i \frac{dT_i}{dt} = \pi \int \frac{d^3k}{(2\pi)^3} \int d\omega \frac{V_{ie}(k)}{|D(k, \omega)|^2} [T_i \text{Im} A_{ei}(k, \omega) \text{Im} \chi_i(k, \omega) - T_e \text{Im} A_{ie}(k, \omega) \text{Im} \chi_e(k, \omega)]$$

ABSORPTION - EMISSION of plasma fluctuations

where $A_{\alpha\beta} \equiv u_{\alpha\beta} \chi_\alpha (1 - u_{\beta\beta} \chi_\beta)$, $D \equiv (1 - u_{ee} \chi_e)(1 - u_{ii} \chi_i) - u_{ei} u_{ie} \chi_e \chi_i$, $u_{\alpha\beta} \equiv V_{\alpha\beta} (1 - G_{\alpha\beta})$

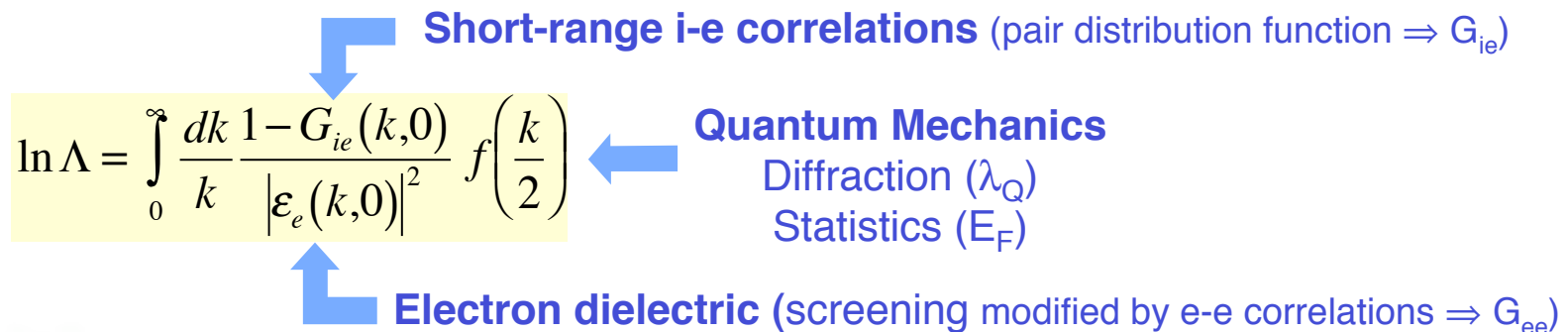
ω -integration is greatly simplified by the small mass ratio $m / M \ll 1$

Electron response near $\omega_{pi} \ll \omega_{pe} \Rightarrow \chi_e(k, \omega) \sim \chi_e(k, 0)$

Ion response evaluated using f-sum rule: $\int d\omega (f_{abs} - f_{ems}) \propto n_e$

Removes i - i interactions
 \Rightarrow NO ion screening
 Like Bethe stopping power &
 TRK sum-rule in spectroscopy

Result reduces to familiar relaxation rate $\nu_{ei} = \nu_o \ln \Lambda$ but with a generalized Coulomb log that includes **ALL physics self-consistently**



We used plasma MD simulations to test comprehensive model

MD simulations provide *ab-initio* transport rates with high accuracy

Classical e's & ions with Coulomb force

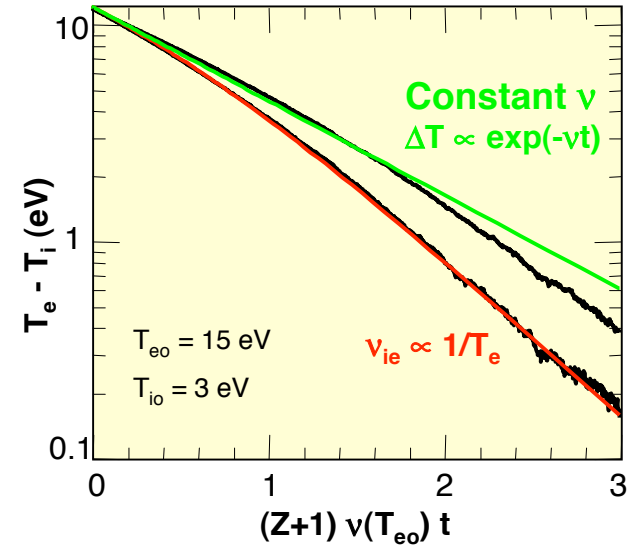
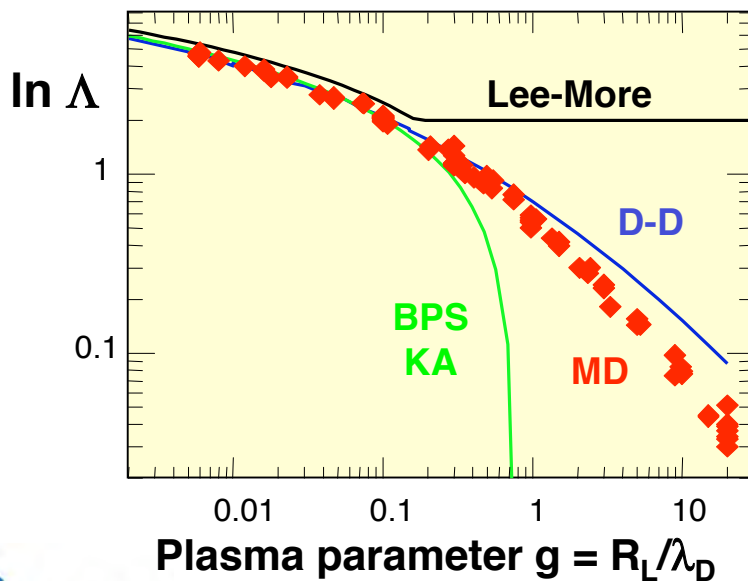
$N_{\text{particles}} \Rightarrow 10^6$ & $\delta t \Rightarrow 10^{-3}/\omega_{pe}$ for convergence

Can discern non-exponential decay due to $v(T_e)$

$$(Z+1) v_{e-i}(T_{eo}) t = 1 - T + \varepsilon \ln\left(\frac{\varepsilon - 1}{\varepsilon - T}\right) \quad \text{where} \quad \varepsilon \equiv \frac{Z + T_{io}/T_{eo}}{1 + Z}$$

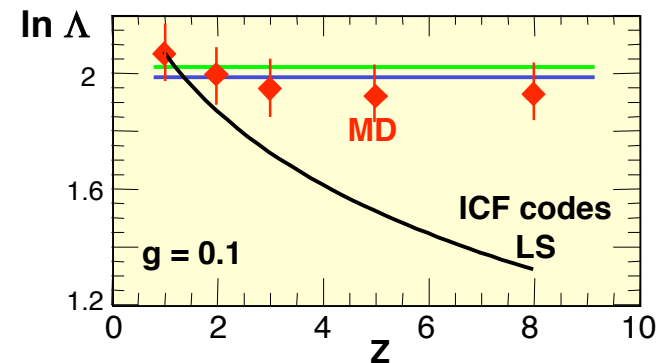
Coulomb log agrees with D & D theory

- Reduces to classical (KA '63, BPS '08) & quantum (Larkin '60) limits for weak coupling
- Extends results to finite coupling $g \Rightarrow 20$



NO D-H screening for v_{ie}

OK for electrical conductivity



Thermonuclear (TN) burn in DT plasma involves many complex processes of comparable time scales that compete with dis-assembly

Use VPIC code to simulate infinite D-T plasma @ $\rho = 100 \text{ g/cc}$, $T = 10 \text{ keV}$

Bowers et al., Phys Plasma **15**, 55703 (08)

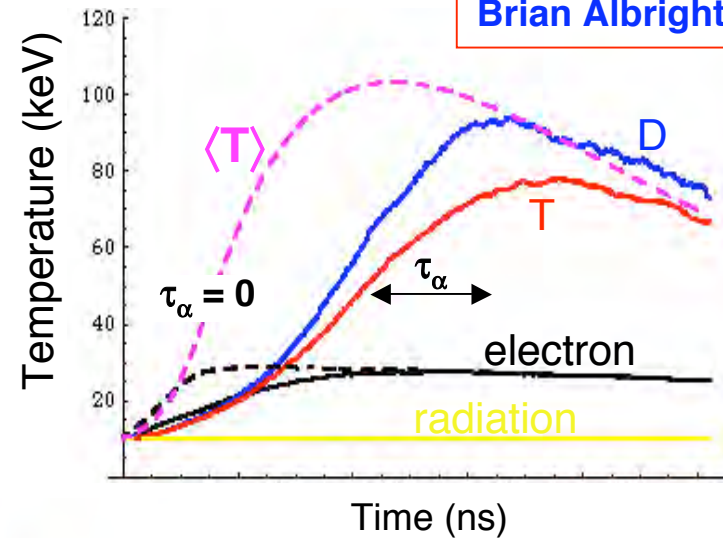
Kinetic ions, D, T, α 's

Fluid electrons

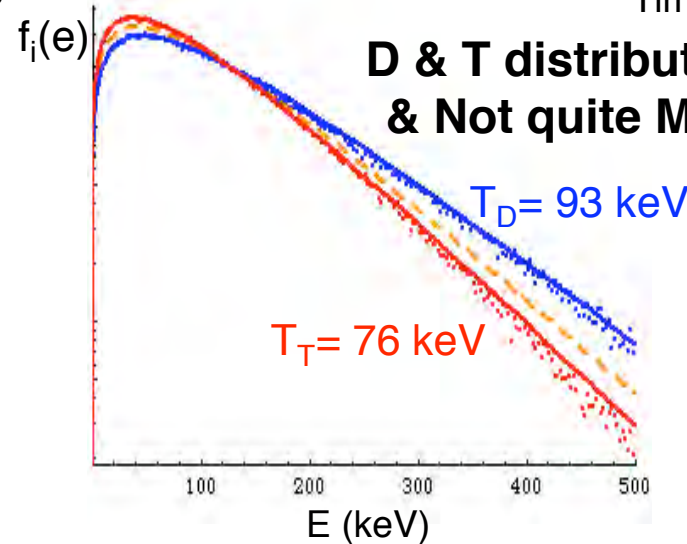
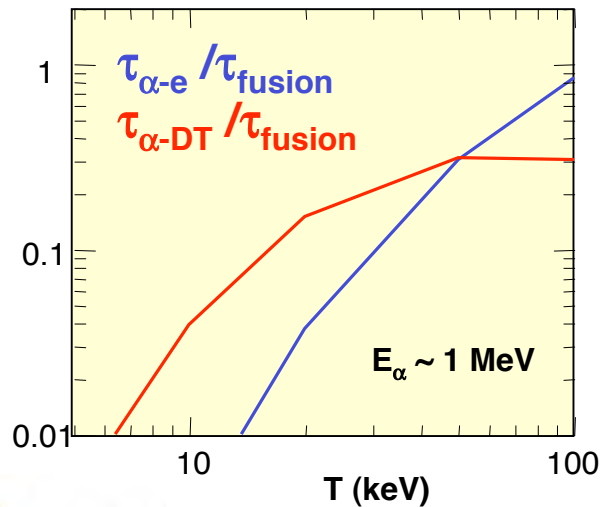
Radiation energy sink

α stopping moderates TN burn

Brian Albright



α 's couple first to electrons, then to ions as T increases

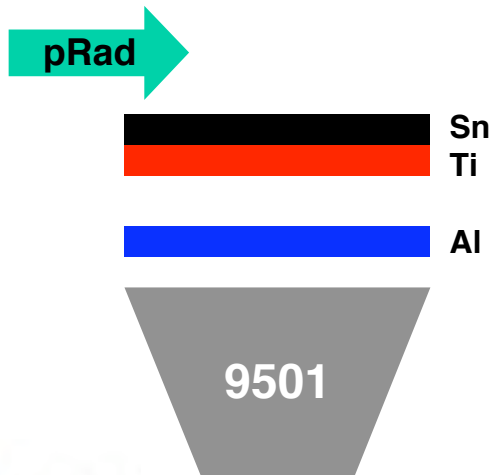
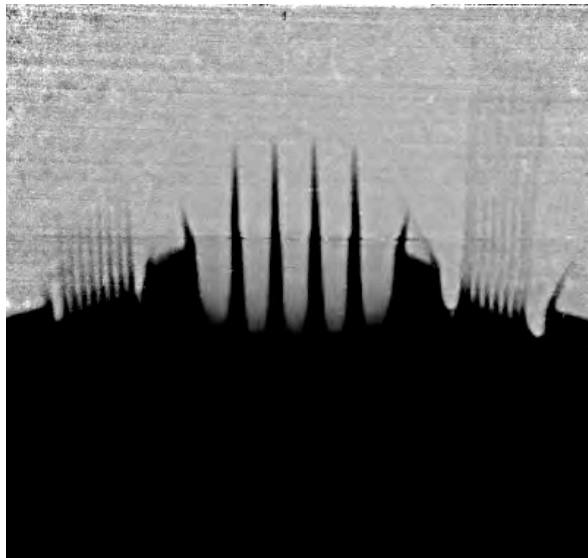


D & T distributions differ & Not quite Maxwellian

Need to couple results to ICF design code

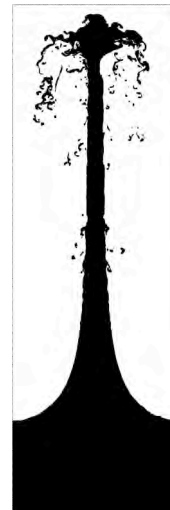
We couple hydrodynamics & molecular dynamics codes to study Richtmyer-Meshkov instability over variety of scales & conditions

Proton radiography exp's
Buttler et al.

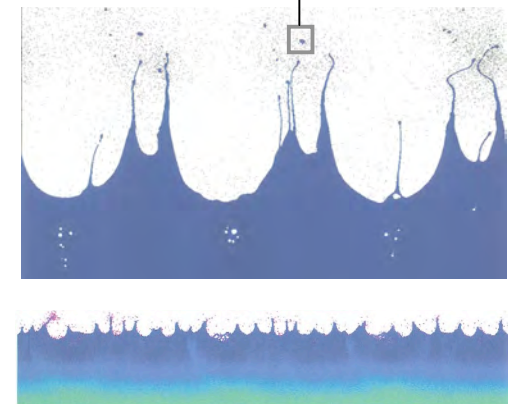
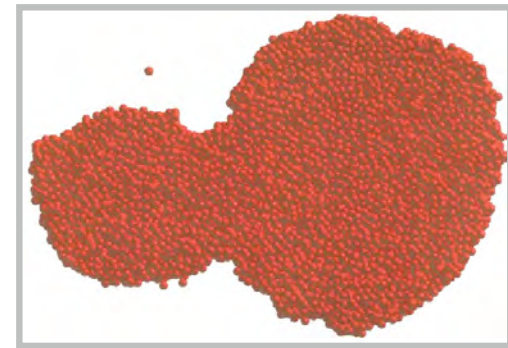


Hydro simulations with
perfect elastic-plastic,
Terrones

γ -fluids on
FLASH (UC)
Ramaprabhu



MD simulations to
atomic scale
Germann & Cherne



We solve complex physical problems for National Security

Multi-physics codes describe complex physical phenomena in HED regime

Defense, energy, astrophysics, climate, biology ...

High-resolution, single-physics simulations clarify fundamental properties

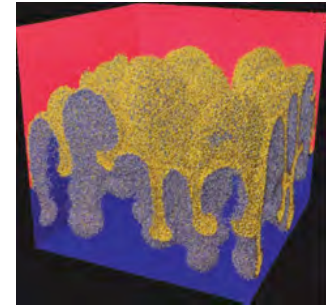
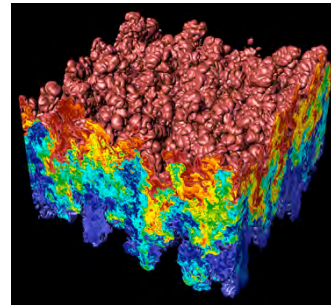
Turbulent hydrodynamics

Particle-in-cell (collective plasmas)

Molecular dynamics (HED properties)

Monte Carlo methods (transport)

Theory

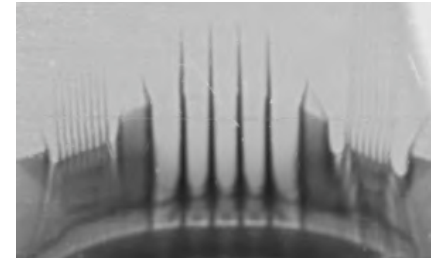


World-class computing & experimental facilities

Roadrunner @ 1 PetaFlop/s

Proton-Radiography

NIF (LLNL) & Ω (UR)



LANL jobs are challenging, relevant, stable & justly compensated

Annualized salaries



Undergraduates \$ 20 - 40 k

Graduates \$ 45 - 60 k

Post-docs \$ 70 - 85 k

Scientist 3 ~ \$ 112 k

Scientist 5 ~ \$ 167 k

Management i = 1, 6 + 25 %



Range ~ ± 25 %

